# How well can satellite altimetry and firn models resolve Antarctic firn thickness variations?

Maria T. Kappelsberger<sup>1</sup>, Martin Horwath<sup>1</sup>, Eric Buchta<sup>1</sup>, Matthias O. Willen<sup>1</sup>, Ludwig Schröder<sup>1,b</sup>, Sanne B.M. Veldhuijsen<sup>2</sup>, Peter Kuipers Munneke<sup>2</sup>, and Michiel R. van den Broeke<sup>2</sup>

**Correspondence:** Maria T. Kappelsberger (maria.kappelsberger@tu-dresden.de)

**Abstract.** Elevation changes of the Antarctic Ice Sheet (AIS) related to surface mass balance (SMB) and firn processes vary strongly in space and time. Their short-term sub-decadal natural variability is large and hampers the detection of long-term climate trends. Firn models or satellite altimetry observations are typically used to investigate such firn thickness changes. However, there is a large spread among firm models. Further, they do not fully explain observed firm thickness changes, especially on smaller temporal and spatial scales. Reconciled firn thickness variations will facilitate the detection of long-term trends from satellite altimetry, the resolution of the spatial patterns of such trends and, hence, their attribution to the underlying mechanisms. This study has two objectives: First, we quantify interannual Antarctic firn thickness variations on a 10 km grid scale. Second, we characterise errors in both the altimetry products and firn models. To achieve this, we jointly analyse satellite altimetry and firn modelling results in time and space. We use the timing of firn thickness variations from firn models and the satellite-observed amplitude of these variations to generate a combined product ('adjusted firn thickness variations') over the AIS for 1992–2017. The combined product characterises spatially resolved variations better than either firn models alone or altimetry alone. We detect highest absolute differences between the adjusted and modelled variations at lower elevations near the AIS margins, probably influenced by the lower resolution, more blurred and the less precise spatial distribution of the modelled variations. In a relative sense, the largest mismatch between the adjusted and modelled variations is found in the dry interior of the East Antarctic Ice Sheet (EAIS), in particular across large megadune fields. Here, the low signal-to-noise ratio poses a challenge for both models and altimetry to resolve firn thickness variations. The altimetric residuals still contain a A large part of the altimetry variance and include firm variance in the altimetric time series is not explained by the adjusted firm thickness variations. Analysis of the altimetric residuals indicate that they contain both firn model errors, such as firn signals not captured by the models, and altimetry errors. Apart from, such as time-variable penetration effects of radar altimetry signals, the residuals disclose patterns indicating uncertainties radar penetration effects but also errors in intermission calibration.

## 1 Introduction

The global mean sea level rose by  $3.05 \pm 0.24 \, \mathrm{mm \, yr^{-1}}$  during the period 1993–2016 (Horwath et al., 2022). Ice-mass loss from Antarctica contributed  $\sim 6\%$  to this rise (Horwath et al., 2022), and is likely to continue (Horwath et al., 2022; IPCC, 2021)

<sup>&</sup>lt;sup>1</sup>Institut für Planetare Geodäsie, Technische Universität Dresden, Dresden, Germany

<sup>&</sup>lt;sup>2</sup>Institute for Marine and Atmospheric Research Utrecht (IMAU), Utrecht University, Utrecht, The Netherlands

<sup>&</sup>lt;sup>b</sup>now at: Bundesamt für Kartographie und Geodäsie, Leipzig, Germany

(IPCC, 2021). The evolution of the Antarctic Ice Sheet (AIS) is of critical concern because the AIS contains the world's largest reservoir of frozen freshwater, the equivalent of ~58 m in global mean sea level (Fretwell et al., 2013), and projections of Antarctica's future contribution to the sea-level rise exhibit a large spread (Schlegel et al., 2018). Relative to 1995–2014, by 2100, Antarctica is expected to contribute 0.03 to 0.27 m and 0.03 to 0.34 m (likely ranges) to the global mean sea level rise under the low and very high greenhouse gas emissions scenario, respectively (Fox-Kemper et al., 2021) –(Schlegel et al., 2018; Fox-Kemper et al., 2021). In order to narrow the range of future sea-level rise projections, this spread we need to better understand the ice-sheet processes and, for this, improve through improved models and observational constraints to quantify the associated volume and mass changes with higher accuracy.

## 1.1 Antarctic mass balance and the role of SMB variations

The mass balance of a grounded ice sheet is commonly separated into three processes components: surface mass balance (SMB), ice discharge and basal mass balance, and ice discharge. SMB comprises total precipitation (snowfall, rainfall), total sublimation (from surface and drifting snow), drifting snow erosion and meltwater runoff (van den Broeke et al., 2016; van Wessem et al., 2018). It refers to processes occurring on the surface of the ice sheet in the snow and firn layer. Snow refers to the seasonal snow cover, i.e. it is less than a year old. Firn refers to multiyear snow and is defined as the transition from snow to glacier ice (van den Broeke, 2008). In the following, we refer to both snow and firn by the term firn layer. Ice discharge is the ice flow across the grounding lineand is linked to processes occurring in the ice layer (Willen et al., 2021). Basal mass balance is thought to be small (Otosaka et al., 2023a), and not considered here.

The current overall mass balance mass loss of the AIS is dominated by an increase in mass loss through ice discharge resulting from an acceleration of glacier flow, primarily from ice discharge from outlet glaciers of the West Antarctic Ice Sheet (WAIS) (Velicogna et al., 2020; Rignot et al., 2019) (Otosaka et al., 2023b). However, uncertainties in the long-term SMB limit the attribution of mass balance components when evaluating satellite data (Willen et al., 2021). On interannual to decadal timescales, variations in SMB (dominated by precipitation) control the variability of the Antarctic mass balance on interannual to decadal timescales (Rignot et al., 2019; Davison et al., 2023). The amplitudes of SMB variations, as well as just as the SMB itself, vary strongly over space. They are influenced by ice sheet topography and atmospheric and oceanic conditions (Lenaerts et al., 2019). Antarctic SMB variability is associated with large-scale atmospheric circulation, such as the Amundsen Sea Low, the Southern Annular Mode and the El Niño Southern Oscillation (e.g. Cullather et al., 1996; Lenaerts et al., 2019; Noble et al., 2 . The strong interannual Antarctic SMB variability hampers the detection of statistically significant trends in the Antarctic (surface) mass balance. To separate long-term trends from short-term variability, the time period considered is essential (Wouters et al., 2013). Ice cores indicate an increase in SMB, in particular in West Antarctica, over the twentieth century (Thomas et al., 2017; Wang et al., 2019; Medley and Thomas, 2019). Over the shorter by oceanic and atmospheric conditions and circulations (Lenaerts et al., 2019; Noble et al., 2020; Kaitheri et al., 2021). Over the satellite period, on a decadal and multidecadal scale, possible climate trends are masked by the large short-term-interannual Antarctic SMB variability (Mottram et al., 2021; Gutiérrez et al., 2021). An improved quantification of interannual SMB variations in space and time is required in order to robustly resolve long-term SMB trends in the Antarctic SMB and overall mass balance (King and Watson,

2020). This is currently lacking (e.g. Mottram et al., 2021).

85

## 1.2 Modelling and observing SMB and firn thickness changes

To date, the SMB for the entire ice sheet is commonly simulated using regional climate models (RCMs) that are thoroughly evaluated against hundreds of in situ observations of SMB. Earth system models have recently caught up in this regard (Lenaerts et al., 2019). RCMs specialise in the physics of polar ice sheets (van Wessem et al., 2018; Agosta et al., 2019). They are commonly used to simulate the SMB for the entire ice sheet (Lenaerts et al., 2019). When the main goal of RCMs is to realistically simulate the ice sheet weather, as is the case here, they are forced by atmospheric reanalysis products which typically provide data from 1979 onwards (Gossart et al., 2019). Mottram et al. (2021) compared Antarctic SMB simulations from an ensemble of five different RCMs all forced by ERA-Interim (Dee et al., 2011). Model differences comprise e.g. the topography model, horizontal resolution and complexity in (sub)surface, snow and firn schemes. Mottram et al. (2021) find and thoroughly evaluated against hundreds of in situ observations of SMB (van Wessem et al., 2018; Agosta et al., 2019). Mottram et al. (2021) demonstrated that different RCMs provide similar outputs for annual to decadal SMB variations on a continental scale (Antarctica), as long as they are driven by the same reanalysis product. However, spatial variations in SMB show a poorer agreement. On a basin scale, the largest deviations are found for the Antarctic Peninsula and the basin that includes the Transantarctic Mountains and part of the interior of the East Antarctic Ice Sheet (EAIS) (basin 8 in this study; Fig. 3). For this basin, the ensemble standard deviation (std) of 40 Gt yr<sup>-1</sup> amounts to 37 % of the ensemble mean. Moreover, even when models provide similar basin-wide SMB estimates, their spatial patterns the spatial patterns of the different SMB 75 estimates differ substantially on a regional and local scale. The largest deviations between the models are mainly at the coastal margin of the entire grounded AIS.

Results results from RCMs are used to force firn models, that which simulate the temporal evolution of the Antarctic firn due to SMB and firn processes such as densification (Ligtenberg et al., 2011; Lundin et al., 2017). Firn elevation changes, or firn thickness changes, are an output of firn models. Verjans et al. (2021) examined differences in linear trends of There is a large spread between firn thickness changes between a range of 54 from different firn model setups for the EAIS. On a basin scale, the ensemble stds range from 0.2 to 1.0 cm yr<sup>-1</sup> and amount to 15 to 300% of the ensemble mean trends of their respective basins. Over the entire EAIS, the choice of climate forcing (RCM), firn compaction and surface snow density contribute to the ensemble spread by 72%, 20% and 4%, respectively, which highlights the importance and need for more precise RCMs, mainly because the uncertainty in the modelled SMB directly influences the modelled firn thickness (Verjans et al., 2021).

Besides modelling tools, satellite measurements are the only possibility to infer ice-sheet-wide changes in SMB and firm thickness. Observations from the satellite gravimetry missions GRACE and GRACE-FO are widely used to estimate Antarctic ice mass changes (e.g. Horwath and Dietrich, 2009; Velicogna and Wahr, 2013; Barletta et al., 2013; Groh et al., 2019). Comparisons between gravimetric ice-mass balance estimates and modelled SMB results (with additional consideration of ice-dynamics changes) were made for the entire AIS, its main regions, drainage basins, or glacier catchments (Mohajerani et al., 2018; Velicogna et al., 2018; However, gravimetric mass-balance estimates have to be corrected for superimposed signals such as glacial isostatic adjustment, involving large uncertainties (Shepherd et al., 2018; Whitehouse et al., 2019; Willen et al., 2020; Groh and Horwath, 2021). Moreover,

GRACE/GRACE-FO cannot resolve mass changes on smaller spatial seales and their observations are restricted to the period after 2002.

By contrast, observations from satellite satellite altimetry provide a higher high spatial resolution of several kilometres and go back to the year 1992 for covering most of the AIS (Wingham et al., 1998). They These measurements allow the derivation of temporal changes of the ice sheet's surface elevation and are therefore sensitive ice-sheet surface elevation changes due to volume changes of the AIS and to the deformation of the solid Earth, with the latter negligible compared to the former (Willen et al., 2021). Most of the altimetry missions utilise (d)-radar waves (e.g. Envisat, CryoSat-2). Since 2003 laser altimeters are also used (e.g. ICESat-2ICESat). While laser altimeters rely on good atmospheric conditions (no thick clouds or blowing snow) radar altimetry is independent of weather conditions (Otosaka et al., 2023a). On the other hand, laser signals are reflected at or near the ice-sheet surface, independently of its properties, while radar signals penetrate into the upper snow/firm layers. Radar altimetry results can thus be biased by firm layer. This may cause biases and artificial variations in radar altimetry results depending on the time-variable dielectric properties of the ice-sheet surface (Davis and Ferguson, 2004; Rémy et al., 2012). If elevation changes due to changing ice flow can either be neglected or subtracted, altimetric elevation changes can be compared to modelled elevation changes due to SMB and firm processes provided by firm models (Kuipers Munneke et al., 2015; Medley et al., 2022a) -firm and the data processing choices to account for them (Davis and Ferguson, 2004; Rémy et al., 2012).

While the rates of modelled and observed elevation changes agree well when averaged over large drainage basins and over 25 years, the correlation diminishes significantly on a grid scale and over 5 years (supplement to Shepherd et al., 2019). Recently, Veldhuijsen et al. (2023) reported that agreement between altimetry and firn modelling results improved when an updated firn model was employed. Nevertheless, discrepancies still remain. (See Section 2.3 for further details on comparisons between altimetry and firn models.) Inconsistencies between models and altimetry also affect the derivation of altimetric ice mass changes, as this depends on modelling results. Using the models in a rigorous, deterministic manner (Kuipers Munneke et al., 2015) resulted in altimetric mass changes characterised by widespread signals of dynamic imbalance that are not deemed fully realistic (supplement to McMillan et al., 2016; Shepherd et al., 2019; Kappelsberger et al., 2021). The reason likely lies in errors in the involved altimetry and modelling results. Therefore, a simplified approach using a (steady-state) density model is commonly applied (e.g. Sørensen et al., 2011; McMillan et al., 2016; Schröder et al., 2019a; Shepherd et al., 2019; Kappelsberger et al., 2

#### 1.3 Previous work

100

105

110

125

Using SMB and firn modelling outputs alone to quantify interannual variations in SMB and firn thickness introduces large uncertainties: the inter-model spread is large, and the model outputs also differ from observational data (Section 1.2). satellite observations (Veldhuijsen et al., 2023). The latter is particularly true at local spatial scales (supplement to Shepherd et al., 2019). Likewise, interannual variations analysed using only data from gravimetry and altimetry satellite observations are strongly affected by their errors (Horwath et al., 2012; Mémin et al., 2015; Su et al., 2018; Shi et al., 2022). Moreover, it is difficult to relate the variations derived from observations alone observed variations to their physical causes. Therefore, the studies of Sasgen et al. (2010), Bodart and Bingham (2019), Kim et al. (2020), Kaitheri et al. (2021) and Zhang et al. (2021) compared or

combined observational and modelling/meteorological dataspace-based geodetic observations with meteorological fields from atmospheric reanalysis data or RCMs. However, their derived interannual variations are spatially coarsely resolved coarsely resolved in space (at about 400 km) and mainly limited to the GRACE/period of the satellite gravimetry missions GRACE and GRACE-FOperiod.

#### 1.4 Purpose

130

140

145

155

This study focuses on the interannual variations in firn thickness on a regional to local scale. Knowledge of interannual variations is required to isolate long-term trends in ice volume or mass changes(Section 1.1). To identify the underlying glaciological processes and separate SMB and firn signals from ice dynamics, the spatial patterns of interannual variations and long-term trends need to be resolved. As the analysis of basin integrals is not sufficient for this purpose, we work at 10 km grid-scale level. We characterise and quantify firn thickness variations in space and time by combining results from satellite altimetry and firn modelling. By combining both data sets, we expect to reduce uncertainties and errors compared with compared to the variations derived from altimetry or models alone. For the first time, the entire spatial full spatial and temporal information present in both the altimetry products and modelling outputs, is exploited together with the high (monthly) temporal resolution of gridded altimetry products, is jointly exploited modelling results. Apart from determining firn thickness variations empirically, our analysis provides information on the error characteristics of both the altimetry products and the model outputs.

## 2 Data

# 2.1 Altimetry

We use the altimetry product from Technische Universität Dresden (TUD) (Schröder et al., 2019a), referred to as TUD altimetry. As an alternative data set, we use the product from Jet Propulsion Laboratory (JPL) (Nilsson et al., 2022), referred to as JPL altimetry. products from Schröder et al. (2019a) and Nilsson et al. (2022)derived. Both studies provide monthly resolved elevation changes of the grounded AIS from a multi-mission satellite altimetry analysis. The elevation changesrepresent By elevation changes, or elevation anomalies, as they we refer to the difference between the elevation at time t and the elevation at a reference epoch to Chosen reference epoch. We use elevation changes over the time period May 1992 to December 2017 containing data from pulse-limited radar altimetry ERS-1, ERS-2, Envisat and CryoSat-2 low resolution mode (LRM), from radar altimetry CryoSat-2 in synthetic aperture radar interferometric (SARIn) mode and from laser altimetry ICESat. As each altimetry mission differs in its orbit configuration, its maximum southern latitude differs. Thus, the lower time limit May 1992 is set to ensure spatial data coverage While the orbit configurations of the missions entail different limits of coverage close to the poles, all mentioned missions cover at least up to 81.5° S. Grid We exclude grid cells with large gaps in the altimetry time series, such as the area south of 81.5° S and the Antarctic Peninsulaare excluded. The upper time limit December 2017 is set to ensure an overlapping period of TUD and JPL altimetry. In the following, the main altimetry processing steps are summarised and differences between TUD and JPL pointed out, coverage by both products.

Schröder et al. (2019a) and Nilsson et al. (2022) corrected the measurements from pulse-limited radar altimetry for sloping terrain with the relocation method (Roemer et al., 2007; Nilsson et al., 2016) using different digital elevation models (Helm et al., (2014)) versus Fretwell et al. (2013)). Both studies applied a threshold retracker for the offset center of gravityamplitude (Wingham et al., 1986) 160 to the radar-applied their own retracking and slope correction to the return signal (waveform). While the TUD product adopts a very low threshold at 10 % to reduce the sensitivity to variations in firn pack properties (Schröder et al., 2019a), the JPL product is based on a 30 % threshold for ERS-1, ERS-2 and Envisat data. CryoSat-2 LRM data were treated similarly for both products (using a 10 % threshold). of the pulse-limited radar altimeters to derive elevation measurements. Data from the CryoSat-2 SARIn mode was processed by Helm et al. (2014)and Nilsson et al. (2016) for TUD and JPL, respectively. The height. 165 Nilsson et al. (2022) used the pre-processed elevation measurements of the 'Geophysical Data Record' (Brockley et al., 2017) for ERS-1, ERS-2 and Envisat. They applied their own processing to the CryoSat-2 data (Nilsson et al., 2016). The elevation measurements were analysed using repeat-track altimetry on a polar-stereographic grid to derive elevation time series. For this analysis, Schröder et al. (2019a) and Nilsson et al. (2022) used different grid spacing and different search radii(constant versus varying/mission-dependent)search radii. Further differences refer to the removal of time-invariant topography (bilinear 170 surface versus varying models/mean, bilinear or biquadratic surface) and the correction for time-variable radar signal penetration and scattering effects (backscatter correction versus backscatter, leading edge width and trailing edge slope correction). While Schröder et al. (2019a) performed these two-steps in one least-squares fit, Nilsson et al. (2022) implemented two separate fits for this purpose fitted them separately.

To derive a continuous time series of elevation changes, intermission fand intermode calibration offsets must be solved. This 175 is a major difference between both altimetry products: TUD is based on While Schröder et al. (2019a) used overlapping epochs or subtracting subtracted a technique-specific reference elevation and JPL, Nilsson et al. (2022) used a least-squares adjustment based on all altimetric measurements. In general, also the weighting between measurements from different missions, in particular the weighting ratio between Envisat and ICESat (Table 1 Schröder et al. (2019a) versus Table 1 Nilsson et al. (2021) differs and then selected overlapping epochs with special treatment of the less than four months Envisat-CryoSat-2 overlap. Moreover, Nilsson et al. (2022) scaled the seasonal amplitudes of the time series of ERS-1, ERS-2 and Envisat to the seasonal 180 amplitudes derived from CryoSat-2 to mitigate artificial seasonal variations caused by time-variable signal penetration. Finally, Schröder et al. (2019a) smoothed the processed data by a three-month moving average and a  $10 \,\mathrm{km}$  one- $\sigma$  Gaussian weighting function. This reduced the spatial grid resolution to 10 km x 10 km. Nilsson et al. (2022) interpolated the processed data with collocation (max, search radius of 50 km, correlation length of 20 km) on a spatial grid with a formal resolution of 1920 m x 1920 m. We interpolate the JPL product to the spatial grid of TUD by averaging the data data from Nilsson et al. (2022) to 185 conform to the product from Schröder et al. (2019a). Therefore, we average the data spatially over 10 km x 10 km. We smooth the JPL time series by a three-month moving average in order to conform to the TUD product. The use of the JPL altimetry data is further restricted to and temporally over three months. We only use those points in time and space where TUD altimetry data is available data are available from both products.

In addition to TUD and JPL, Shepherd et al. (2019) published a long-term, multi-mission altimetry product. However, we do not use their product because it is not resolved on a monthly basis (consecutive 5-year intervals are provided) Since we focus

190

on the interannual to decadal time scales, we fit and remove the offset, linear, quadratic and seasonal signals from the monthly elevation changes for every  $10 \,\mathrm{km} \times 10 \,\mathrm{km}$  grid cell. Seasonal signals are modelled by annual and semi-annual cosine and sine functions. Thereby, we fit different seasonal amplitudes for the time periods before and after 2003. In this way we account for the inconsistency in the seasonal amplitudes between the older pulse-limited radar altimetry missions (ERS-1, ERS-2) and the newer missions (Envisat, ICESat, CryoSat-2) (Nilsson et al., 2022), as the corrections for time-variable penetration effects on the radar return signal are imperfect in reducing unrealistic seasonal amplitudes in particular for the older missions (Ligtenberg et al., 2012). The fitted parameters are presented in Fig. S1–S4. After subtracting the offset, linear, quadratic and seasonal signals, we are left with the interannual elevation changes, which we refer to as altimetric variations,  $hv^A$ .

## 200 2.2 Firn models

195

205

220

We use the firn model-thickness changes from the firn models IMAU-FDM v1.2A of Veldhuijsen et al. (2023), referred to as IMAU (Institute for Marine and Atmospheric Research Utrecht) firn model, which is an update of v1.1 (Ligtenberg et al., 2011). As an alternative data set we involve the GSFC-FDM v1-Ligtenberg et al. (2011), and GSFC-FDMv1.2.1 of Medley et al. (2022a), referred to as GSFC (Goddard Space Flight Center) firn model. It which uses the Community Firn Model framework of Stevens et al. (2020, 2021). Here, one output of the models is used, the firn thickness changes. Firn thickness changes represent firn thickness anomalies, as they refer to the difference between firn thickness at time t and the mean firn thickness over a certain reference period (see below). The IMAU model outputs Outputs from Veldhuijsen et al. (2023) are given every ten days and on a regular grid with a spacing of 27 km from 1979 to 2020. The GSFC model outputs Outputs from Medley et al. (2022a) are given every five days and on a regular grid with a spacing of 12.5 km from 1980 to 2021. In accordance with the altimetry data, we involve use firn thickness changes from May 1992 to December 2017 and from the grounded AIS excluding the Antarctic Peninsulaand the period May 1992 to December 2017. We adapt the temporal resolution to that of the altimetry product by calculating monthly means and applying a three-month moving average smoothing. In the following, the main firn model set ups are summarised and differences between IMAU and GSFC pointed out.

The IMAU firn model firn model from Veldhuijsen et al. (2023) is forced with 3-hourly three-hourly fields of surface temperature, 10 m wind speed and SMB components (snowfall, rainfall, sublimation, snowdrift erosion, snowmelt) from the ERA5 atmospheric reanalysis data (Hersbach et al., 2020) dynamically downscaled with RACMO2.3p2 (van Wessem et al., 2018)to-. RACMO2.3p2 uses a spatial resolution of 27 km x 27 km. The GSFC firn model and is forced by the ERA5 atmospheric reanalysis data (Hersbach et al., 2020). The firn model from Medley et al. (2022a) is forced with hourly fields of snowfall, total precipitation, evaporation, 2 m air temperature , skin temperature and runoff from and skin temperature from a downscaled version (12.5 km x 12.5 km) of the MERRA-2 atmospheric reanalysis data (Gelaro et al., 2017) downscaled to a spatial resolution of 12.5 km x 12.5 km (Gelaro et al., 2017; Tian et al., 2017). The firn layer was initialised by looping over the forcing data of the reference period 1979–2020 (for the IMAU model) (for Veldhuijsen et al., 2023) and 1980–2019 (for the GSFC model) (for Medley et al., 2022a) until the firn column was refreshed at least once. This implies the assumption that the reference period represents stable climatic conditions and the current firn layer is in equilibrium. However,

Veldhuijsen et al. (2023) noted that the assumption of a steady-state firm layer in regions where precipitation has increased over the last centuries, such as the Antarctic Peninsula and Ellsworth Land.

Both firn models use the same semi-empirical equation of Arthern et al. (2010) to model dry-snow densification but their procedure procedures for deriving the empirical correction terms differs. IMAU derives differ. Veldhuijsen et al. (2023) derived this empirical correction from observations in Antarctica, while GSFC employs Medley et al. (2022a) employed observations from both Antarctica and Greenland. Furthermore, the two firn models use a different parameterisation for surface snow density. Veldhuijsen et al. (2023) use the formulation of Lenaerts et al. (2012), which depends on instantaneous surface temperature and 10 m wind speed, but with updated constants derived from their own calibration. Medley et al. (2022a) built a new model parameterisation depending on snow accumulation, air temperature, total wind speed, and specific humidity. In general Overall, they follow the approach from Helsen et al. (2008), which incorporates mean annual parameters. Both firn models take into account include the processes of meltwater percolation and refreezing.

## 2.3 Previous comparisons between altimetry and firn models

225

230

235

240

245

250

255

The data sets used in this study have been compared mainly on the basis of multi-year to decadal rates and seasonal amplitudes. The results of these comparisons are briefly summarised below.

Nilsson et al. (2022) reported that elevation change rates of TUD, JPL and Shepherd et al. (2019) are generally in good agreement and within their uncertainties (over 1992–2016). By excluding regions of dynamic imbalance, they also compared rates between altimetric elevation changes and We subtract the offset, linear, quadratic and seasonal signals from the modelled firn thickness changes derived from the IMAU-FDM v1.1 (Ligtenberg et al., 2011) forced by ERA-Interim reanalysis data. In Dronning Maud Land and Enderby Land the thickening patterns of modelled and observed rates are in good overall agreement the same way as we do for the altimetric time series, except that the observed rates show stronger magnitudes than the modelled rates of FDM v1.1. However, in the region of Wilhelm II Land and Wilkes Land the differences between observed and modelled rates are larger. Rates are of opposite sign (altimetry: positive rates; FDM v1.1: negative rates). The three altimetry products agree in magnitude and sign (Nilsson et al., 2022). Recently, the update from FDM v1.1 to v1.2A (forced by ERA5 reanalysis data; Section 2.2) lead to an improved agreement with the observed rates (evaluated by the TUD product over 2003–2015). With the update, the modelled rates were found to be more positive in Dronning Maud Land and Enderby Land and less negative in Wilhelm II Land and Wilkes Land. However, discrepancies between altimetric and modelled rates remain, in particular for the Antarctic Peninsula and Ellsworth Land (Veldhuijsen et al., 2023).

Estimates of the average ice sheet seasonal amplitude in firn thickness give discrepant results for different altimetry products, firn models and time periods (with the latter also involving different spatial coverage): 5.1, 2.7 and 2.9 cm for TUD altimetry, JPL altimetry and the IMAU-FDM v1.1, respectively, over 1992–2016 (Nilsson et al., 2022), and, 5.2, 3.1 and 3.0 cm for TUD altimetry, the IMAU firn model (v1.2A) and the GSFC firn model, respectively, over 2003–2015 (Veldhuijsen et al., 2023) we assume constant seasonal amplitudes for the entire period. The subtracted parameters are presented in Fig. S1–S4. This leaves us with firn thickness variations on interannual time scales, which we refer to as modelled firn thickness variations,  $fv^{\rm M}$ .

## 2.4 Illustration of data sets

260

265

275

280

285

In this section, we illustrate and compare the original data sets through basin-mean time series and root mean square (rms) maps. This recalls typical ways of previous comparisons between altimetry and firn modelling results and serves as a reference for our subsequent exploration of a wider range of spatio-temporal scales.

Fig. 2 (dash-dotted lines) shows the basin-mean time series for the original elevation changes from TUD altimetry,  $h^{\rm A1}$ , and the IMAU firn model,  $f^{\rm Ma}$ . (For the JPL altimetry,  $h^{\rm A2}$ , and the GSFC firn model,  $f^{\rm Mb}$ , similar time series are shown in Fig. S1.) Agreement between  $h^{\rm A1}$  and  $f^{\rm Ma}$  is generally good on interannual scales. Differences appear in the long-term trends. The trend differences are greatest for basin 10 (Amundsen Sea Embayment region). This is due to the effect of changing ice flow (Mouginot et al., 2014; Gardner et al., 2018; Diener et al., 2021), reflected in  $h^{\rm A1}$ , while this effect is purposely not considered by  $f^{\rm Ma}$ . A further difference is that, prior to 2003 the seasonal amplitudes of altimetry exceed those of the firn model (Ligtenberg et al., 2012; Nilsson et al., 2022). Fig. 2 shows maps of rms values for  $h^{\rm A1}$  and  $f^{\rm Ma}$  over the entire period 1993-2017. The rms values include the effect of a linear component which dominates e.g.  $h^{\rm A1}$  in the Amundsen Sea Embayment region. Besides long-term influences, the overall spatial patterns of  $h^{\rm A1}$  and  $f^{\rm Ma}$  are related to the spatial variability of the SMB, with values increasing from the AIS interior to the margin (van Wessem et al., 2014; Lenaerts et al., 2019). The rms values of  $h^{\rm A1}$  are generally larger than those of  $f^{\rm Ma}$ . This is mainly due to the higher noise level in altimetry measurements before 2003 (Schröder et al., 2019a; Nilsson et al., 2022). For the period after 2003, the rms values of altimetry and the firn model are in better agreement (Fig. S2). (Fig. S2 also shows the rms for all data sets used (A1, A2, Ma, Mb) separately for the periods before and after 2003.)

## 3 Methods

## 3.1 Regression approach Basic approach

We jointly analyse satellite altimetry and firn modelling results while focusing on interannual to decadal time scales. The the interannual elevation changes from satellite altimetry and firn modelling results. Fig. 1 gives an overview of the workflow. The new combination approach is a regression of altimetric elevation changes against several signals explained in the following. For each  $10 \, \mathrm{km} \times 10 \, \mathrm{km}$  grid cell, we describe the time series of monthly elevation changes from altimetry  $h^{\mathrm{A}}$  by

$$\frac{h^{A}(t) = \underbrace{a + bt + c(0.5t^{2})}_{+H_{1}(t) [d_{1}\cos(\omega t) + d_{2}\sin(\omega t) + d_{3}\cos(2\omega t) + d_{4}\sin(2\omega t)]}_{+H_{2}(t) [d_{5}\cos(\omega t) + d_{6}\sin(\omega t) + d_{7}\cos(2\omega t) + d_{8}\sin(2\omega t)]}$$

$$+ \sum_{n=1}^{N} e_{n}^{A} P C_{n}^{M}(t)$$

$$+ r^{A}(t)$$

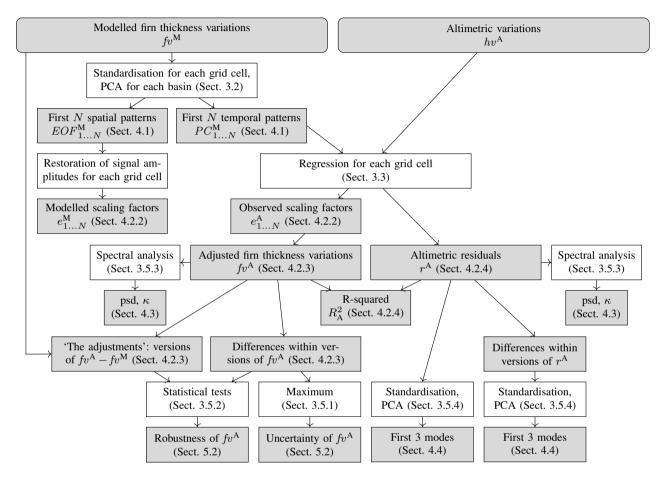


Figure 1. Basin-mean time series Workflow of the original elevation changes from analysis. Grey boxes: the IMAU firn model (Ma) results,  $f^{\text{Ma}}$ , (dash-dotted, black line) their notation and from TUD altimetry (A1),  $h^{\text{A1}}$ , (dash-dotted, eyan line) the section where they are first presented. Basin-mean time series of modelled firn thickness variations from Ma,  $fv^{\text{Ma}}$ , (solid, black line) White boxes: the main methodological steps to derive these results and of adjusted firn thickness variations based on A1a,  $fv^{\text{A1a}}$ , (solid, eyan line) the sections where they are explained.

with 
$$H_1(t) = \begin{cases} 1, & \text{if } t < 2003 \\ 0, & \text{if } t > 2003 \end{cases}$$

290 and 
$$H_2(t) = \begin{cases} 0, & \text{if } t < 2003 \\ 1, & \text{if } t > 2003. \end{cases}$$

295

300

The regression parameters a (offset), b (linear trend), c (acceleration),  $d_{1,\dots,8}$  (amplitudes of annual and semi-annual harmonic signals, with  $\omega=2\pi/1\,\mathrm{yr}$ ) and  $e_{1\dots N}^A$  (scaling factors for dominant temporal patterns in modelled firm thickness variations) are estimated by least squares adjustment. They are adjusted w.r.t. the reference epoch,  $t_0$ , September 2010 no matter of data coverage. The definition of N dominant temporal patterns in modelled the altimetric variations,  $hv^A$ , against dominant signals in the firm thickness variations  $PC_1^M,\dots,PC_N^M$  depends on the drainage basin to which the considered location belongs. It is explained in Section 3.2. The residuals  $r^A$  are the difference between the elevation changes  $h^A$  and the fitted model.

Seasonal signals are modelled by annual and semi-annual cosine and sine functions. By applying the masks H<sub>1</sub> and H<sub>2</sub>, we fit different seasonal amplitudes for the time periods before and after 2003. In this way we account for the inconsistency in the seasonal amplitudes between the older pulse-limited radar altimetry missions (ERS-1, ERS-2) and the newer missions of different techniques (Envisat, ICESat, CryoSat-2) (Nilsson et al., 2022). The corrections for the influence of the ice sheet surface dielectric properties on the radar return signal (Section 2.1) are only partly able to reduce artificially large seasonal amplitudes in particular for the older missions (Ligtenberg et al., 2012).

We subtract the adjusted offset, linear, quadratic and seasonal signals from  $h^{A}$  to derive elevation changes on interannual time scales from altimetry according to

The interannual elevation changes are termed altimetric variations,  $hv^{A}$ .

We perform a weighted regression. Observations h<sup>A</sup> after 2003 are weighted by 1, while observations prior to 2003 are given a different (usually lower) weight, according to the finding of a generally higher noise level of the results from the older altimetry missions (Schröder et al., 2019a; Nilsson et al., 2022). The weight prior to 2003 is defined (individually for every grid point) by the ratio of the noise variance of h<sup>A</sup> prior to 2003 and after 2003. The noise variance ratio is assessed empirically from the variance of the high-pass filtered time series (cf. Groh et al., 2019). The high-pass filtering (performed separately for the period prior to 2003 and after 2003) consists in removing linear and seasonal signals and subsequently removing a low-pass filtered version of the time series, where the low-pass filter is a Gaussian filter with a 6σ = 12 months filter width.

## 3.1.1 Adjusted firn thickness variations

fv<sup>M</sup>. Our regression approach relies on the ability of firn models to capture the timing of dominant variations in SMB and firn processes across basins. However, the amplitudes and spatial patterns of the variations are adjusted to satellite altimetry results. We trust the temporal more than the spatial give more trust to the temporal patterns of the firn model than to their spatial patterns for the following reasons. Mottram et al. (2021) as well as Lenaerts et al. (2019) and Gutiérrez et al. (2021) have pointed out that the spatial patterns of RCMs, which force firn models, show a large spread between models but not their temporal patterns (Section 1.2)while there is less spread between the temporal patterns. While spatially resolved differences (between models, between observations and between models and observations) are substantial, the differences have been shown to be are reduced when basin averages are used (Agosta et al., 2019; Shepherd et al., 2019; Willen et al., 2021). Also, the The overall good agreement of basin-mean time series on interannual scales (between the data sets used here) has also been noted in Section 2.4, of  $fv^{M}$  and  $hv^{A}$  is supported in Fig. S1. 2.

For each grid cell, the adjusted firn thickness variations  $fv^{A}$  are determined by the linear combination in Eq.1:

$$\mathit{fv}^{\mathrm{A}}(t) = \sum_{n=1}^{N} e_n^{\mathrm{A}} P C_n^{\mathrm{M}}(t).$$

330 The

335

340

320

325

## 3.2 Principal component analysis of modelled firn thickness variations

We identify dominant temporal patterns in firn thickness variations  $\frac{PC_n^M}{n}$ , are identified by principal component analysis (PCA) of the firn modelling data. PCA, also called empirical orthogonal function (EOF) analysis, is applied to identify dominant modes of variability, represented by pairs of a principal component (PC) and an EOF, where EOFs and the corresponding, uncorrelated PCs represent the spatial and temporal which represent the temporal and spatial patterns, respectively. Comprehensive and general references for PCA are Preisendorfer (1988) and Jolliffe (2002), while e.g. Forootan and Kusche (2012) or Boergens et al. (2014) apply PCA and extensions of PCA to geodetic data. (Preisendorfer, 1988; Jolliffe, 2002; Forootan and Kusche, 20

Prior to applying PCA to the firn modelling data, we remove offset, linear, quadratic and seasonal signals from the modelled firn thickness changes  $f^{M}$  according to

$$\begin{split} \underline{f^{\mathrm{M}}(t)} &= & \underline{a+bt+c\left(0.5\,t^2\right)} \\ &+ d_1\cos(\omega t) + d_2\sin(\omega t) + d_3\cos(2\omega t) + d_4\sin(2\omega t) \\ &+ fv^{\mathrm{M}}(t), \end{split}$$

where  $a, b, c, d_{1,...,4}$  are estimated by an ordinary least-squares adjustment. The residuals,  $fv^{M}$ , are referred to as firn thickness variations. The The PCA is performed on these the modelled firn thickness variations,  $fv^{M}$ , after their standardisation. We standardise the time series of  $fv^{M}$  for each grid cell, i.e. we shift and scale it such that it has zero mean and a std-standard

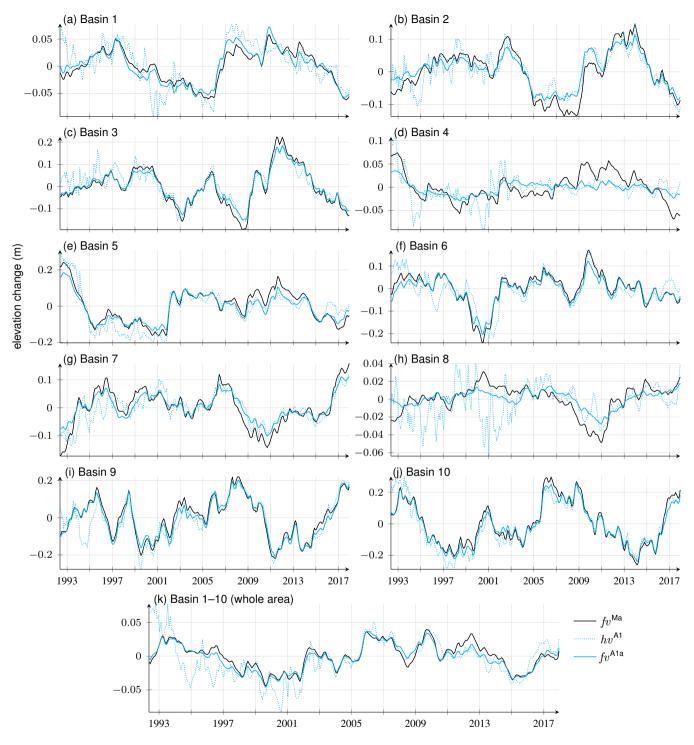


Figure 2. Root mean square (rms) of the original time series of elevation change over the entire period for (a) TUD altimetry,  $h^{\rm Al}$ , and (b) the IMAU firm model,  $f^{\rm Ma}$ . The color scale is logarithmic. Basin-mean time series of modelled firm thickness variations from Veldhuijsen et al. (2023),  $fv^{\rm Ma}$ , (solid, black), of altimetric variations from Schröder et al. (2019a),  $hv^{\rm Al}$ , (dotted, cyan), and of adjusted firm thickness variations based on Ala,  $fv^{\rm Ala}$ , (solid, cyan). Basin definitions are shown in Fig. 3.

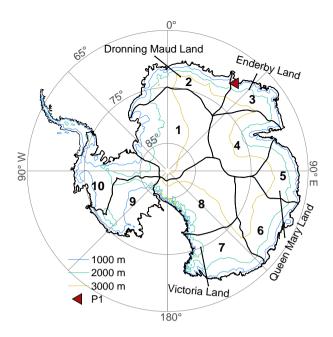


Figure 3. Drainage basins of the EAIS and WAIS used in this study (thick black lines) following, slightly modified from the definition of Rignot et al. (2011a, b). The outline of Antarctic Peninsula is indicated by a thin black line. Contour lines of the ice sheet surface are shown at 1000 m, 2000 m and 3000 m. Highlighted in red are the circle at constant latitude of 72° S (profile 1) and the line at constant longitude at 115° E (profile 2). Grid point P1 is located at lon = 37.7° E, lat = 70.2° S. We use the polar stereographic projection EPSG:3031 (WGS84, latitude of true scale: 71° S, central meridian: 0°). All further maps are displayed in the same projection and with the same spacing of longitude grid lines (every 45°) and latitude grid lines (every 10°).

deviation (std) of one, because we aim to equally represent the patterns of temporal evolution regardless of location or absolute amplitudes. Otherwise, PCA results would mainly reflect patterns that are dominant at the margins as where the amplitudes of SMB and firn thickness variations exhibit much larger amplitudes at the margins are much larger than in the interior (van Wessem et al., 2018; Lenaerts et al., 2019). To regain interpretable magnitudes of the EOFs, the EOFs are multiplied by the std of the time series of  $fv^M$  for each grid cell, which was previously used for standardisation. After this restoration of the signal amplitudes, we no longer speak of EOFs but of modelled scaling factors,  $e^M$ .

PCA is applied individually. We separately apply the PCA to  $fv^M$  for 10 selected regions (basins that together cover the East Antarctic Ice Sheet (EAIS) and the WAIS (Fig. 3). To define the regions basins, we make use of the drainage basin definition by Rignot et al. (2011a, b). We aggregate and aggregate neighbouring basins smaller than  $\sim 600,000 \, \mathrm{km^2}$  with those neighbouring basins where we find strongest correlation between their. The decision which of the original 15 basins are aggregated is guided by the correlations between the first three PCs. This step reduces the original number of 15 drainage basinsfor the EAIS and WAIS to 10 of a preliminary PCA per original basins. For each (aggregated) basin of the 10 basins, we choose the first N modes that contain explain at least 90% of the total variance of the (standardised) data. In addition, North's rule of thumb (North et al., 1982) is applied to test whether the eigenvalues of these N patterns are well separated with respect to their errors. The

355

360

first N dominant temporal patterns  $\frac{PC_n^M}{n}$  enter Eq. 1,  $PC_{1,N}^M$ , enter the regression approach in normalised form.

## 3.2.1 Goodness of fit

To examine how well a regression fits the observations,

## 3.3 Regression approach

For each  $10 \,\mathrm{km} \times 10 \,\mathrm{km}$  grid cell, we describe the time series of monthly altimetric variations,  $hv^{\mathrm{A}}$ , by

$$hv^{A}(t) = a + \sum_{n=1}^{N} e_{n}^{A} PC_{n}^{M}(t) + r^{A}(t).$$
 (1)

The scaling factors  $e_1^A$  and the offset a are estimated by least squares adjustment. The dominant temporal patterns in modelled firn thickness variations,  $PC_n^M(t)$ , refer to the basin to which the grid cell belongs. The residuals of the fit are  $r^A$ .

We define a combined product by the linear combination of Eq. 1, evaluated per grid cell and time:

370 
$$fv^{A}(t) = \sum_{n=1}^{N} e_n^{A} P C_n^{M}(t).$$
 (2)

We refer to  $fv^{A}(t)$  as the 'adjusted firn thickness variations'.

The stochastic model of our regression in Eq. 1 prescribes a different weighting of observations from two time periods. As results from the older altimetry missions generally have a higher noise level (Schröder et al., 2019a; Nilsson et al., 2022),  $hv^A$  after 2003 are weighted by 1, while  $hv^A$  before 2003 are given a different (usually lower) weight, which is defined, individually for every grid point, by the ratio of the noise variance of  $hv^A$  before and after 2003. We assess the noise by the high-pass filtered version of  $hv^A$  separately for both periods (cf. Groh et al., 2019). The high-pass filtering consists of removing a low-pass filtered version of  $hv^A$ , where the low-pass filter is a Gaussian filter with a  $6\sigma = 12$  months filter width.

To assess the goodness of fit, we calculate the coefficient of determination values of R-squared,  $R^2$  (R squared), as, as

$$\underline{\underline{R^2}}_{\mathsf{A}}^2 = 1 - \underbrace{\frac{SS(r)}{SS(h_{\mathsf{tot}})}}_{\mathsf{SS}(h_{\mathsf{tot}})} \cdot \underbrace{\frac{SS(r^{\mathsf{A}})}{SS(hv^{\mathsf{A}})}}_{\mathsf{SS}(hv^{\mathsf{A}})} = 1 - \underbrace{\frac{SS(r^{\mathsf{A}})}{SS(fv^{\mathsf{A}} + r^{\mathsf{A}})}}_{\mathsf{SS}(fv^{\mathsf{A}} + r^{\mathsf{A}})}, \tag{3}$$

SS(r) and  $SS(h_{tot})$  where  $SS(r^A)$  and  $SS(hv^A)$  are the residual and total sum of squares, respectively.  $SS(r)/SS(h_{tot})$   $SS(r^A)/SS(hv^A)$  describes the proportion of unexplained variance. Here, we We calculate  $R^2$  for every grid cell individually and exclude the adjusted linear, seasonal and quadratic signals in  $h_{tot}$ . Thus, Eq. 5 specifies to

$$\underline{R_{\rm A}^2} = 1 - \frac{SS(r^{\rm A})}{SS(hv^{\rm A})} = 1 - \frac{SS(r^{\rm A})}{SS(fv^{\rm A} + r^{\rm A})}.$$

÷

**Table 1.** Names of the four different versions of adjusted firm thickness variations,  $fv^A$ , regression results derived by applying the regression approach Eq. 1 with different data sets. Differences to A1a are indicated in bold.

Name	$h^{A} h v^{A}$ from	$PC^{M}$ from $\overset{*}{\sim}$
A1a	TUD altimetry (A1 ) (Schröder et al., 2019a)	IMAU firn model (Ma )* Ma (Veldhuijsen et al., 2023)
<mark>A2a</mark> -A2a	JPL altimetry (A2) A2 (Nilsson et al., 2022)	IMAU firn model (Ma )* Ma (Veldhuijsen et al., 2023)
Alb Alb	TUD altimetry (A1 ) (Schröder et al., 2019a)	GSFC firn model (Mb)* Mb (Medley et al., 2022a)
<b>A2b</b> <u>A2b</u>		GSFC firn model (Mb)*-Mb (Medley et al., 2022a)

<sup>\*</sup> standardised fv<sup>M</sup> (Sect. 3.2)

385

390

395

400

405

## 3.4 Different versions of adjusted firn thickness variations

We derive two different sets of  $PC^{\rm M}$  depending on the firn model incorporated. In our annotation, we distinguish Our annotation distinguishes the firn models by superscripts 'Ma' and 'Mb' for the IMAU and GSFC model model by Veldhuijsen et al. (2023) and Medley et al. (2022a), respectively. The regression approach (Eq. 1) is applied with each set of  $PC^{\rm M}$  and equally to each of the two altimetry products  $h^{\rm A}$  from TUD and JPL products of  $hv^{\rm A}$  from Schröder et al. (2019a) and Nilsson et al. (2022), which we distinguish by superscripts 'A1' and 'A2', respectively. Thus, depending on the combination All combinations of data sets used result in four applications of the regression approach (Table 1). Thus, we obtain four versions of adjusted firn thickness variations ( $fv^{\rm A1a}$ ,  $fv^{\rm A2a}$ ,  $fv^{\rm A1b}$  and  $fv^{\rm A2b}$ ). This also results in four versions of altimetric residuals ( $r^{\rm A1a}$ ,  $r^{\rm A2a}$ ,  $r^{\rm A1b}$  and  $fv^{\rm A2b}$ ) and of associated coefficients of determination R-squared ( $R^2_{\rm A1a}$ ,  $R^2_{\rm A2a}$ ,  $R^2_{\rm A1b}$  and  $fv^{\rm A2b}$ . Table 1 gives an overview of the applications of the regression approach.

We additionally fit a regression similar to Eq. 1 to the firn model data,  $f^{\rm M}$ , after their interpolation to the altimetric grid of  $10\,{\rm km}$  spacing. The same deterministic model Eq. 1 is used, but no weighting is applied. In this way, the regression parameters  $a,b,c,d_{1,\ldots,8}$  and in particular the scaling factors adjusted to altimetry,  $e^{\rm A}_n$ , can be directly compared to the scaling factors derived from the firn models,  $e^{\rm M}_n$ . Replacing  $e^{\rm A}_n$  by  $e^{\rm M}_n$  in Eq. 2 would then lead to a variant of modelled firn thickness variations, restricted to the dominant temporal modes found in the PCA. We refer to this variant as truncated modelled firn thickness variations,  $fv^{\rm M}_{90}$ . The suffix 90 indicates that the dominant patterns were chosen such that they cover at least  $90\,\%$  of the variance of standardised time series within the specific basin.

In Appendix A1, we additionally assess three alternative ways of defining 'adjusted' firn thickness variations. These alternatives are: (E1) Accept the modelled firn thickness variations,  $fv^{\rm M}$ , without any adjustment to altimetry. (E2) Instead of using PCA-based dominant temporal patterns use the modelled time series of firn thickness variations at every grid cell and scale it to fit the altimetry. These alternative variations are called We refer to the results as scaled firn thickness variations. We refer to them by  $fv^{\rm E2}$ . (E3) Identify the dominant temporal patterns of modelled firn thickness variations by a PCA without prior standardisation of the time series. These alternative variations are called Omit the standardisation step prior to the PCA and proceed according to Eq. 1 and 2. We refer to the result as modified adjusted firn thickness variations. We refer to them by

 $fv^{E3}$ . See Table B1 for an overview of the defined symbols and their terminology. Note, that we do not introduce  $fv^{E1}$  as this 0 would correspond to  $fv^{M}$ .

## 3.5 Assessment methods

420

425

# 3.5.1 Uncertainty of adjusted firn thickness variations

To We assess the impact of the choice of data sets and thus the influence of different errors on the adjusted firn thickness variations,  $fv^A$ , differences between by using differences between the time series of the various versions of firn thickness variations, fv, (Section 3.4) are used. In general, from . For each time series of differences we can calculate the temporal root mean square (rms). This procedure is applied to time series differences evaluated for every is done for time series per grid cell and also for differences in basin-mean time series.

time series of basin averages. To assess the uncertainty of the adjusted firn thickness variations,  $fv^A$ , we consider the maximum deviation within the different versions of  $fv^A$  (Table 1). For this purpose, we form all possible combinations of differences between the four versions of  $fv^A$ . It results in six combinations of time series differences and thus, six (temporal) rms values, where we choose the one that is maximum the six possible differences from  $fv^{A1a}$ ,  $fv^{A2a}$ ,  $fv^{A1b}$ , and  $fv^{A2b}$ , and take the maximum of the rms differences.

## 3.5.2 Robustness of adjusted firn thickness variations

The adjusted-We refer to the differences between adjusted and modelled firn thickness variations  $fv^A$ , can be considered an improved representation of firn thickness variations compared with the modelled variations,  $fv^M$ , if we can statistically demonstrate that as 'the adjustments'  $(fv^A - fv^M)$ . We consider these 'adjustments' to be 'improvements' over the firn models, if the differences within different versions of  $fv^A$  are significantly smaller than the differences to  $fv^M$ . To investigate this, we perform statistical tests comparing distributions of temporal rms of differences within  $fv^A$  to differences  $fv^A - fv^M$ .

We work with adjustments. We test for significance by comparing the distributions of their temporal rms. We use a two-sample, one-sided Kolmogorov-Smirnov test which is a non-parametric hypothesis test as the differences in fv do not follow a normal distribution. The Kolmogorov-Smirnov test uses the empirical cumulative distribution function (cdf), which is the integral of the probability density function (pdf), to compare the distributions of two samples (Massey, 1951; Miller, 1956; Marsaglia et al., . The null hypothesis (H0) reads: both samples, the data of both differences to be compared, are from the same continuous distribution. Thus, the alternative hypothesis (H1) reads: the empirical cdf of sample one (the differences within  $fv^A$ ), is larger than the empirical cdf of sample two (the differences between  $fv^A$  and  $fv^M$  adjustments), that is the differences within  $fv^A$  tend to be smaller than the differences between  $fv^A$  and  $fv^M$ .

## 3.5.3 Spectral analysis of regression results

The We analyse the time series of altimetric residuals,  $r^A$ , and the adjusted firn thickness variations,  $fv^A$ , are analysed in the spectral domain to characterise their stochastic properties. We calculate the through their power spectral density (psd) and the

their spectral indices,  $\kappa$ , of the underlying time series of  $r^A$  and  $fv^A$ . We use the software HECTOR v1.7.2 (Bos et al., 2012) to estimate  $\kappa$  (Bos et al., 2012). As  $r^A$  and  $fv^A$  do not yield a white noise behaviour we use the formulation of power-law noise to approximate their stochastic properties. (For example, power-law with  $\kappa = -1$  and  $\kappa = -2$  represents flicker and random walk noise, respectively.)

## 3.5.4 Principal component analysis of altimetric residuals

- The four versions of altimetric residuals (Section 3.4) altimetric residuals,  $r^A$ , are further analysed in the spatio-temporal domain. First, we perform PCA on the altimetric residuals themselves to further identify dominant signals related to ice sheet processes not considered or incorrectly represented by the firn models. (Note that the residuals may additionally contain signals related to variations in ice flow dynamics or subglacial hydrology.)—Second, we perform PCA on the residual differences to further detect and investigate prevailing uncertainties in the altimetry analysis. Only data after 2003 is used because of the higher noise level in the altimetry measurements of the older satellite missions. Inclusion—Test experiments showed that errors of pre-2003 data would result in more noisy dominant patterns and therefore could distort detected dominant modes bias the dominant modes and hardly helps to distinguish between signal and error. We standardise the time series of residuals and residual differences, as we did previously when identifying dominant patterns in modelled firn thickness variations (Section Sect. 3.2).
- The first PCA is applied to four versions of standardised residuals ( $r^{A1a}$ ,  $r^{A1b}$ ,  $r^{A2a}$  and  $r^{A2b}$ ). The second PCA is applied to two versions of standardised residual differences ( $r^{A1a} r^{A2a}$  and  $r^{A1b} r^{A2b}$ ). For each PCA, we set up one aggregated 'super data matrix' in which we arrange the time series of residuals/residual differences for all pixels and for the different versions for each grid cell and each version into a single set of time series. PCA Specifically, our data sets comprises m = 90638 points in space (entire area under investigation) and p = 108 points in time (2003–2017). Thus, for the first and second PCA, the super data matrix has the size of  $4m \times p$  and  $2m \times p$ , respectively. The PCA is conducted to identify the dominant temporal patterns (PCs), which are shared by all versions, together with their space-dependent and version-dependent amplitudes, i.e. their spatial patterns . spatial patterns (EOFs). Each identified mode thus consists of one joint PC (1 x p) and four, or two, EOFs (4m x 1 or 2m x 1) in the case of the first, or second, PCA, respectively.

#### 4 Results

## 465 4.1 Dominant patterns in modelled firn thickness variations

Fig. 4 shows the PCA results for the example of basin 3 and the IMAU firn model input data,  $fv^{Ma}$ . The figure shows the dominant spatial patterns (EOFs) and temporal patterns (PCs) together with their share of the total variance. We recall that PCA is performed individually for each basin and that  $fv^{Ma}$  are standardised prior to PCA. A comprehensive presentation of results for all basins and for the two alternative input firn models IMAU and GSFC is given by Fig. S5—S9. Depending on the basin, different numbers of modes (i.e. PC-EOF pairs) are required to We can explain at least 90% of the total variance : two

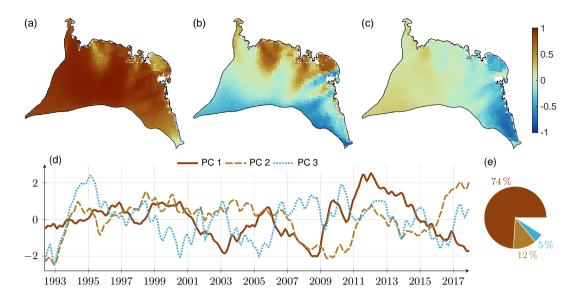


Figure 4. PCA results of basin 3: dominant patterns in firn thickness variations identified from standardised firn modelling data (Ma). (a, b, ea-c) First, second and third three spatial pattern patterns (EOFEOFs). (d) First three temporal patterns (PCs). (e) Associated percentages of the basin's total data variance. We define the PCs as standardised time series (mean of zeroResults for all basins and for both firn models, std of 1) Ma and without a unit while the EOFs have the unit of metreMb, are given by Fig. \$5–89.

modes for of the modelled firn thickness variations,  $fv^{Ma}$ , with two modes (basin 5), three modes for (basins 1, 3 and 6), four modes for (basins 2, 4 and 8) and five modes for (basins 7, 9 and 10(based on Ma). The first, second, third, fourth and fifth modes describe 58 to 74 %, 11 to 21 %, 4 to 12 %, 3 to 5 % and 3 to 4 % of the data variance, respectively (based on Ma).

The PCs and EOFs modes (i.e. the PC-EOF pairs) reveal a typical hierarchy of modes of an autocorrelated geophysical signal. The first temporal patterns,  $PC_1^M$ , show a longer wavelength signal than the following PCs, as shown in Fig. 4 for the region of Dronning Maud Land (basin 3). The first EOFs show an approximately uniform distribution, while the following EOFs are more complex and change sign. For basin 3, the first three EOFs exhibit auniform behaviour, a north-south gradient and EOF is almost uniform over the entire basin (Fig. 4a). The spatial features of the second EOF follow the topography from north to south (Fig. 4b) and the third EOF exhibits an east-west gradient , respectively (Fig. 4c). The first mode of basins 2 and 3 (the region of Dronning Maud Land and Enderby Land) capture the PC shows a longer wavelength signal than the following PCs. All three PCs fluctuate over time similar to an integrated random walk process (Fig. 4d). In the case of basin 3, 74 % of the variance is explained by the first mode, which captures the accumulation events in 2009 and 2011 (Boening et al., 2012; Lenaerts et al., 2013) . Their temporal patterns,  $PC_1^M$ , show a characteristic increase as shown by the characteristic increase in the PC during these years (Fig. 4d). All subsequent modes are more difficult to interpret as a geophysical signal because of the fact that their determination is governed by the mathematical orthogonality property of PCs.

485

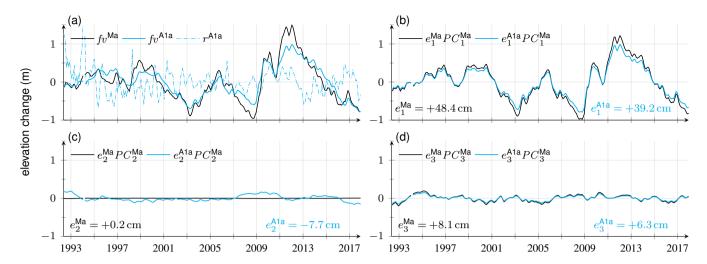


Figure 5. Illustration of the generation of adjusted firn thickness variations  $fv^A$  Regression results for the grid point P1 (Fig. 3). Cyan and black curves show regression results from the adjustment to TUD altimetry (A1a) and, for adirect comparison, to the IMAU firn model (Ma), respectively. (a) Original time series,  $h^{A1}$  and  $f^{Ma}$ . (b) Modelled firn thickness variations,  $fv^{Ma}$  (dashedsolid, black), truncated modelled adjusted firn thickness variations,  $fv^{Ma}$  (solid, blackcyan), and adjusted firn thickness variationsaltimetric residuals,  $fv^{A1a}$  (soliddash-dotted, cyan). (eb-d), (d), (e) Scaled first, second, and third dominant temporal patterns in  $fv^{Ma}$ . Hence,  $PC^M$  from the solid regression version A1a (cyan) and the model Ma (black/). The solid cyan curve in (ba) is the sum of the black/cyan curves in (e-eb-d). (f) Time series of the regression residuals. The black curve a larger subset of selected grid points ( $r^{Ma}$ Fig. S10) equals  $fv^M$   $fv^{Ma}$  are shown in Fig. S11 and S12.

## 4.2 Regression results

490

495

#### 4.2.1 Time series for a selected grid point

Fig. 5 exemplifies the derivation of adjusted firn thickness variations for a selected grid point, P1, and based on the regression A1a (Table 1). P1 (37.7° E, 70.2° S) is located in basin 3 close to the ice sheet margin at  $\sim 1080\,\mathrm{m}$  height (Fig. 3). There, the adjusted and modelled firn thickness variations,  $fv^{\mathrm{A1a}}$  and  $fv^{\mathrm{Ma}}$ , have a standard deviation (std.) of 41.0 std of 40.5 and 51.5 cm, respectively (Fig. 5b). In addition to  $fv^{\mathrm{Ma}}$ , we illustrate the time series of truncated modelled firn thickness variations,  $fv^{\mathrm{Ma}}$  (Section 3.4), which has a std of 49.1 cm. The difference between  $fv^{\mathrm{Ma}}$  and  $fv^{\mathrm{Ma}}_{90}$  equals  $r^{\mathrm{Ma}}$  and is shown in Fig. 5f. a). By construction, the scaling factors  $e_{1,\ldots,3}e_{1\ldots3}$  equal the std of the respective associated scaled dominant temporal patterns. (In the case In the presence of data gaps in the altimetry time series, this equality holds approximately. ). Both  $fv^{\mathrm{A1a}}$  and  $fv^{\mathrm{Ma}}_{90}$   $fv^{\mathrm{Ma}}_{90}$  are dominated by  $PC_1^{\mathrm{Ma}}$  of basin  $3PC_1^{\mathrm{Ma}}$  (Fig. 5b), as this pattern is scaled by  $e_1^{\mathrm{A1a}} = 39.6\,\mathrm{cm}$  39.2 cm (altimetry) and  $e_1^{\mathrm{Ma}} = 48.4\,\mathrm{cm}$  (firn model). For  $e_2$ , altimetry and The second pattern,  $PC_2^{\mathrm{Ma}}$  (Fig. 5c) is very small in the firn model have opposite signs, yet small values, so that they  $(e_2^{\mathrm{Ma}} = 0.2\,\mathrm{cm})$ , while somewhat larger and of opposite sign for altimetry  $(e_2^{\mathrm{A1a}} = -7.7\,\mathrm{cm})$ , but still small enough to contribute little to  $fv^{\mathrm{A1a}}$ .

The std of altimetric residuals,  $r^{\rm A1a}$  is 31.0 cm, is 31.7 cm, less than the std of  $fv^{\rm A1a}$ . The coefficient of determination, (Fig. 5a). The R-squared value  $R^2_{\rm A1a}$  (Eq. 3) is 0.610.601. When calculated separately for the time before and after 2003,  $R^2_{\rm A1a}$  equals -0.06 and 0.84-0.004 and 0.831, respectively. Thus, the adjusted firm thickness variations,  $fv^{\rm A1a}$ , describe less of the variance of altimetric variations altimetry variance before 2003 while after 2003 they explain 84%. Distinguishing the time before and after 2003 is reasonable as we include different weights for the altimetry observations 82%. Because of the different weighting of  $hv^{\rm A}$  before and after 2003 (Section 3.3). Sect. 3.3),  $R^2_{\rm A}$  can indeed be negative and distinguishing the two periods is reasonable.

For a larger subset of selected grid points (Fig. S10), time series of the original elevation changes h and the regression results are shown in Fig. S9–S12. While the following Section 4.2.2 focuses on the adjusted scaling factors  $e_{1,...,N}$ , maps of the other regression parameters (adjusted linear, quadratic and seasonal terms) are presented in Fig. S1–S3.

# 4.2.2 Scaling factors e

515

520

525

530

Figure 6. Adjusted scaling factors along profile 1 (left) and profile 2 (right). (a, b)  $e_1$ , (c, d)  $e_2$  and (c, f)  $e_3$ . Cyan and black curves show the scaling factors adjusted to TUD altimetry (A1a) and to the IMAU firn model (Ma), respectively. Note the different scaling of the y-axes of profile 2.

Fig. 6 shows the spatial variation of the scaling factors  $e_{1...3}$  along two selected profiles marked in Fig. 3. Profile 1 is along the circle of latitude at  $72^{\circ}$  S. Profile 2 is along the meridian at  $115^{\circ}$  E. The absolute magnitude of both scaling factors (from A1a and Ma) is largest at the ice sheet margin. This applies for profile 1 across basin 2, in the middle part of basin 4 and at the end part of basin 7 as well as for profile 2 at the end part of basin 6. Observed factors,  $e_{1...3}^{A1a}$ , reveal stronger variations along both profiles than modelled factors,  $e_{1...3}^{Ma}$ . Discontinuities across basin borders arise because the scaling factors refer to basin-specific patterns.

The scaling factors  $e_{1...3}^{A1a}$  and  $e_{1...3}^{Ma}$  per grid cell are mapped for the example of basin 3 in Fig. 6. The patterns of the factors, like the EOFs (Fig. 4), follow a typical hierarchy already discussed in Section discussed in Sect. 4.1. Overall, the patterns of  $e_{1...3}^{Ma}$  are in a good agreement agree for a large part with  $e_{1...3}^{A1a}$ . However, the first spatial pattern from the model extends further towards into the ice sheet interior than the pattern from altimetry -(Fig. 6d versus Fig. 6a). In general, scaling factors from the model show a smoother and more blurred pattern than the ones adjusted to altimetry. Patterns from altimetry reveal a higher level of detail and a more localised spatial distribution. At certain regions the spatial distributions also differs. In the area at and around differ, e.g. for the second pattern in the vicinity of P1 (marked as a triangle), altimetry observes the second temporal pattern with a negative amplitude  $e_2$ , while the firm model suggests an amplitude near zero. (Fig. 6b versus Fig. 6e). A comprehensive presentation of The spatial variation of the scaling factors along two selected profiles is given by Fig. S13 and a comprehensive representation of the scaling factors for all basins and for with the different choices of input data are is given by Fig. S5 and S14.

## 4.2.3 Firn thickness variations and their sensitivity to the choice of data sets

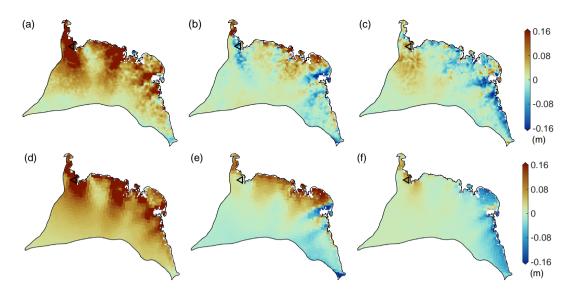


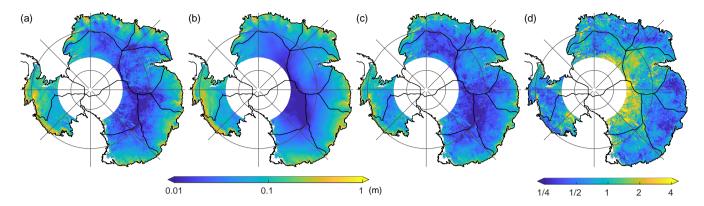
Figure 6. Adjusted scaling Scaling factors for basin 3. (a–c)  $e_{1...3}^{A1a}$ , first three observed factors adjusted to TUD altimetry from the regression A1a. (d–f)  $e_{1...3}^{Ma}$ , first three modelled factors adjusted to from Ma. (d–f) is the HMAU firm modelsame as Fig. 4a–c but with restored signal amplitudes for each grid cell. The location of P1 is shown by the black triangle.

We calculate the rms of the time series of firn thickness variations, fv, for each grid cell. Fig. 7a and 7b show the rms of Ingeneral, the spatial patterns of the rms of the adjusted firn thickness variations based on A1a,  $fv^{A1a}$  (Table 1),  $fv^A$ , and the rms of modelled firn thickness variations based on Ma,  $fv^{Ma}$ , respectively. (The rms of all versions of  $fv^A$  and  $fv^M$  is illustrated in  $fv^M$ , are similar (Fig. S15a –d and Fig. S17a , b, respectively.) In general, the spatial patterns of  $fv^A$  and  $fv^M$  are similar 7a and 7b). Rms values are largest at the ice sheet margin and smallest over the plateau of the EAIS. For grid cells in the elevation ranges (1) below  $1000 \,\mathrm{m}$ , (2)  $1000 \,\mathrm{to} \,2000 \,\mathrm{m}$ , (3)  $2000 \,\mathrm{to} \,3000 \,\mathrm{m}$  and (4) above  $3000 \,\mathrm{m}$ , median rms values are in the range of (1)  $13.2 \,\mathrm{to} \,16.4 \,12.2 \,\mathrm{to} \,16.4$ , (2)  $8.7 \,\mathrm{to} \,10.9 \,8.3 \,\mathrm{to} \,10.9$ , (3)  $3.7 \,\mathrm{to} \,5.1 \,3.5 \,\mathrm{to} \,5.1$  and (4)  $2.2 \,\mathrm{to} \,2.4 \,\mathrm{cm} \,2.1 \,\mathrm{to} \,2.3 \,\mathrm{cm}$ , respectively. Fig. 7e and 7d show the rms of the differences  $fv^{A1a} - fv^{Ma}$  in an absolute and relative way, respectively. Differences between adjusted and modelled variations reveal highest absolute rms values at lower elevations, near the AIS margins (with median rms differences in the range of  $11.5 \,\mathrm{to} \,12.7 \,\mathrm{cm}$  below  $1000 \,\mathrm{m} \,13.4 \,\mathrm{to} \,14.7 \,\mathrm{cm}$  below  $1000 \,\mathrm{m} \,(\mathrm{Fig.} \,7c)$ . In a relative sense, largest mismatch is found in the interior of the EAIS but also at some locations at the ice sheet margin  $-(\mathrm{Fig.} \,7d)$ .

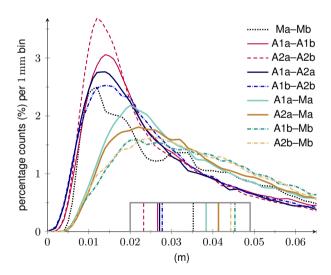
535

540

To evaluate the sensitivity of fv to the choice of data sets, we calculate the difference between various versions of fv (Section Sect. 3.5.1). Fig. 8 shows, and compare the distributions of the rms of differences between various versions of fv. (Corresponding rms maps of differences are displayed in these differences (Fig. S17–S19). In addition to the distributions their median values are presented in Fig. 8 and listed in Table 2. 8). In total, differences within  $fv^A$  are smallest, followed by differences within  $fv^A$  while differences between  $fv^A$  and  $fv^M$  are largest –(Fig. 8, Table 2). Differences within  $fv^A$  indicate a smaller influence by different firn model data than by different altimetry data. Differences between  $fv^A$  and  $fv^M$  are smallest for A1a (adjustment over the IMAU firn model through TUD altimetry) and firn model Ma through altimetry A1), largest for



**Figure 7.** Root mean square (rms) of the times series of (a) adjusted firn thickness variations based on A1a,  $fv^{A1a}$ , and (b) modelled firn thickness variations based on Ma,  $fv^{Ma}$ . (c) Rms of the time series of the differences  $fv^{A1a} - fv^{Ma}$ . (d) Rms of the time series of the differences  $fv^{A1a} - fv^{Ma}$  divided by the rms of  $fv^{Ma}$ . All versions of  $fv^{A}$  and  $fv^{M}$  are illustrated in Fig. S15 and S16.



**Figure 8.** Histograms of the temporal rms, assessed at each grid cell, of differences between various versions of firn thickness variations. (a) Histograms. Vertical lines in the box indicate median values. (b) Cumulative histogramsCorresponding rms maps of differences are displayed in Fig. S15–S17.

**Table 2.** Overview of the comparison between various versions of firn thickness variations, as detailed in Fig. 8. Column 1 indicates the addressed comparison: between versions of adjusted firn thickness variations  $fv^A$  (row 1–4), between modelled firn thickness variations  $fv^M$  (row 5), and between  $fv^A$  and  $fv^M$  (row 6–9). For each comparison, column 2 gives the median (over all grid cells) of the rms (over time) of differences between the two time series evaluated at each grid cell. The table is ordered by the median values (from small to large). Column 3 also gives the median of the rms of differences but as a relative measure. For each grid cell, the rms of differences are divided by the rms of  $fv^{Ma}$ . Then, the median over all grid cells is calculated. Column 4 gives a short description or possible causes.

Difference	Median		Description/Cause	
	absolute	relative		
A2a-A2b	$2.3\mathrm{cm}$	0.46 0.47	influence of different firn model setups based on A2	
Ala-Alb	<del>2.6 cm </del> 2.7 cm	$\textcolor{red}{\textbf{0.51}} \textcolor{red}{0.52} \textcolor{gray}{\sim}$	influence of different firn model setups based on A1	
A1bA1a-A2b-A2a	2.8 cm - 2.7 cm	<del>0.55</del> - <u>0.54</u>	different altimetry analysis based on Mb	
$\underbrace{\text{A1a}\underbrace{\text{A1b}}}_{\text{A2a}}\underbrace{\text{A2b}}_{\text{A2b}}$	<del>2.9 cm -</del> 2.8 cm	$\underline{0.58} \underline{0.54}_{\scriptscriptstyle{\sim}}$	different altimetry analysis based on Ma	
Ma-Mb	$3.5\mathrm{cm}$	0.65	different firn model setups	
Ala-Ma	$3.8\mathrm{cm}$	0.73	Adjustment adjustment over Ma through A1*	
A2a-Ma	4.2 cm 4.1 cm	$0.82 \ 0.80$	Adjustment adjustment over Ma through A2*	
<del>A1b</del> A2b-Mb	$4.4\mathrm{cm}$	$0.83 \ 0.87$	Adjustment adjustment over Mb through A1*	
A2bA1b-Mb	$4.5\mathrm{cm}$	0.87-0.85	Adjustment adjustment over Mb through A2*	

<sup>\*</sup> due to firn signals not correctly represented by the models (firn model errors) and/or due to errors in the altimetry products

A2b (adjustment over the GSFC firn model through JPL altimetry )firn model Mb through altimetry A2) in a relative sense and largest for A1b (adjustment over the firn model Mb through altimetry A1) in an absolute sense (Table 2). The differences between the various versions of fv reflect errors in the firn models and in the altimetry products. These are further discussed in Sections Sect. 5.3 and 5.4.

## 4.2.4 Goodness of fit

The altimetric residuals are used to calculate the goodness of fit or R-squared (Eq. 3). The rms of the altimetric residual time series is and the values of R-squared based on the regression A1a,  $R_{A1a}^2$ , are presented in Fig. 9a (estimated per grid cell over the full period ). The altimetric residuals are used to calculate the goodness of fit (Section 3.1.2). Here, we distinguish between the periods before and after 2003. As mentioned in Section 4.2.1, this is useful and b, respectively, for the period after 2003. The rms of the residuals after 2003 is generally smaller than before 2003 (Fig. S18) due to the different noise levels and weighting of altimetry observations during the altimetry observations in the two periods (Section 3.3). The rms of the residuals Sect. 3.3), so that  $R^2$  is generally higher after 2003 than before 2003 (Fig. 9b)are generally smaller than from the ones over the full period. S19).

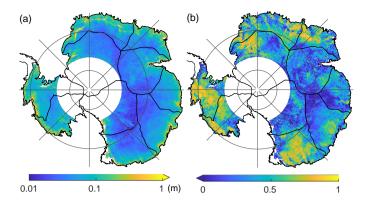


Figure 9. (a, b) Rms of the residual altimetric time series,  $r^{A1a}$ , for (a) the full period and (b) the period after 2003. (eb) Coefficients Values of determination R-squared for the regression A1a,  $R^2_{A1a}$ , considering the period after 2003. Colour bar arrows indicate that the value range exceeds the limits of the colour scale.

The spatial distribution of the coefficients of determination based on the regression A1a,  $R_{\rm A1a}^2$ , and for the period after 2003 is displayed by Fig. 9c. After the individual calculation of  $R_s$   $R_s^2$  for each grid cell, basin-mean values are derived and listed in Table 3 for all versions of regression. (Fig. S18 and Fig. S19 further shows maps of the residuals rms and of  $R^2$  for different versions of regression and both time periods. Table S1 lists basin averages of  $R^2$  for the period before 2003.) Averaged over the entire area,  $R_{\rm A1a}^2$  is 0.46-0.40 after 2003 (Table 3). This means that on average 46%40% of the variance of altimetric variations is captured by the regression model, i.e. by  $fv^{\rm A1a}$ . Depending on the basin,  $fv^{\rm A1a}$  eapture 30% (basin captures 26% (basins 4 and 8) to 62%58% (basin 10) of the data variance. In general, the goodness of fit decreases slightly when using JPL altimetry instead of TUD altimetry (column A1a versus A2a and column A1b versus A2b of Table 3) or we find less agreement with altimetry when incorporating the GSFC Mb firn model instead of the IMAU Ma firn model (Table 3, column A1a versus A1b and column A2a versus A2b).

565

570

580

The impact of methodological changes to the regression approach (E1, E2 and E3 as summarised in SectionSect. 3.4) is presented elaborated in Appendix A2. There, Fig. A1 compares  $R^2$  values of the modified approaches to  $R^2_{A1a}$  for each grid cell and Table A1 lists basin-averaged  $R^2$  values of the modified approaches. The methodological changes result in smaller average  $R^2$  values , (Fig. A1, Table A1), so less of the data variance could can be explained. For this reason, the modified approaches are not preferable to the chosen regression approach presented in SectionSect. 3.3.

By now, the presented  $R^2$  values are based on calculations per grid cell in accordance with the regression approach Eq. 1. For basin average time series,  $R^2$  become becomes larger. Fig. 2 shows the basin-averages basin averages of adjusted firm thickness variations, which we may compare to the basin-averages of the altimetric variations through the coefficients of determination values of R-squared given in Table 3, last column. Indeed,  $fv^{A1a}$  could capture 51% (basin 4) to 97% captures up to 96% (basins 9 and 10; West AIS) of the variance of basin average altimetry variations. (Basin-mean time series of all regression results and versions are presented in Fig. S4–S21.) 20 and S21.

Table 3. Explained variance or eoefficients of determination,  $R^2$ , for each basin and each version of regression (Table 1) over the period after 2003. Apart from the last column  $\overline{A1a}$ ,  $R^2$  is first calculated for each grid cell according to Eq. 3 and then averaged over each basin. Values of  $\overline{A1a}$  are calculated by first averaging the regression results over each basin and then applying Eq. 3. Basin averages of  $R^2$  for the period before 2003 are listed by Table S1.

Basin	Ala	A2a	A1b	A2b	Ala
01	0.46 0.42	0.43 0.40	0.41 0.29	0.36-0.32	0.79
02	0.53 0.50	0.48-0.45	0.49-0.39	0.42 $0.40$	<del>0.94</del> <u>0.92</u>
03	0.48 0.45	0.48 0.46	0.48 0.43	0.46 0.44	<del>0.94</del> <u>0.93</u>
04	0.26	0.36	0.41-0.13	0.27-0.26	$\underbrace{0.33\ 0.51}_{\textstyle \bullet}\underbrace{0.11}_{\textstyle \bullet}$
05	0.40 0.27	0.39 0.33	0.36 0.24	0.39-0.35	<del>0.79</del> <u>0.54</u>
06	0.42 0.32	0.29 0.27	0.36 0.21	0.30 0.25	0.82 - 0.74
07	0.57-0.52	0.47-0.43	0.51-0.40	0.41-0.34	0.94-0.92
08	0.30-0.26	0.37	0.30-0.11	$\underbrace{0.37\text{-}0.31}_{}$	0.66-0.58
09	0.57-0.54	0.50-0.46	0.53-0.45	0.47-0.44	<del>0.97</del> -0.96
10	0.62 0.58	0.56-0.53	0.56-0.44	0.48 0.44	0.97-0.96
01–10*	0.46-0.40	0.43-0.40	0.42 0.29	0.39 0.34	0.83-0.79

<sup>\*</sup> refers to the entire area (considered as a single basin)

However, on the level of individual grid cells the altimetric residuals,  $r^{A}$ , still contain a large proportion of the variance of altimetric variations. For example, for A1a and the period after 2003, an average ratio of 54%60% of the altimetric variations are unexplained. Therefore, the residuals  $r^{A}$  are further investigated in the following Sections Sect. 4.3 and 4.4.

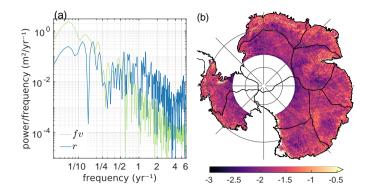
## 4.3 Spectral analysis of regression results

585

590

595

Fig. 10a shows the power spectral density (psd) of the altimetric residuals,  $r^{A1a}$ , and the adjusted firn thickness variations, We find a stronger autocorrelation for the time series of  $fv^{A1a}$ , for the selected grid point P1. The underlying time series are displayed by Fig. 5. (For the larger subset of selected grid points, Fig. S22 and S23 display the psd of the regression results from A1a and A2a, respectively.) The psd of both  $fv^{A1a}$  and  $r^{A1}$  generally decreases from low to high frequencies. The slope of the psd is steeper for  $fv^{A1a}$  than for  $r^{A1}$ . This means that the underlying time series of  $fv^{A1a}$  have stronger autocorrelation than that of than for that of  $r^{A1a}$ , or in other words, the underlying time series of i.e.  $r^{A1a}$  are is closer to white noise behaviour than  $fv^{A1a}$ , since the power spectral density (psd) for  $fv^{A1a}$  shows a steeper decrease with frequency than for  $r^{A1a}$  (Fig. 10a). At low frequencies the psd of  $fv^{A1a}$  generally exceeds the psd of  $r^{A1a}$ , while above a certain frequency ( $\sim 0.5 \text{ yr}^{-1}$  for P1) the psd of  $r^{A1a}$  exceeds that of  $fv^{A1a}$  (Fig. 10a). For P1 that means on time scales shorter than  $\sim 2 \text{ yr } r^{A1a}$  includes more power then  $fv^{A1a}$ . At P1, the The spectral indices  $\kappa$  adjusted to determined for  $r^{A1a}$  and  $fv^{A1a}$  (Section 3.5.3) are -1.75 and  $\leq -3$ , respectively. For each grid cell, -3, respectively, at P1. Over the entire area, the mean value of  $\kappa$  of  $fv^{A1a}$  is calculated to be



**Figure 10.** (a) Lomb-Scargle power spectral density (psd) of altimetric residuals,  $r^{A1a}$ , (blue) and adjusted firn thickness variations,  $fv^{A1a}$ , (green) for grid point P1. See Fig. S22 and S23 for the larger subset of selected grid points. (b) Spectral index  $\kappa$  for power-law noise adjusted to the residuals  $r^{A1a}$  of every grid cell. Colour bar arrows indicate that the value range exceeds the limits of the colour scale.

-3 or more negative. (HECTOR only yields numerical stable results for  $\kappa \geq -3$ .) For each grid cell,  $\kappa$  of for  $r^{A1a}$  are shown in is -1.72 (Fig. 10b. The mean value over the entire area amounts to -1.72. It), which indicates temporally correlated residuals with characteristics close to random-walk noise. For  $fv^{A1a}$ , in contrast, the value of  $\kappa$  is -3 at each grid cell. The employed software to estimate  $\kappa$  (Bos et al., 2012) has -3 as its minimum output value. Hence,  $fv^{A1a}$  has  $\kappa \leq -3$  and therefore a stronger autocorrelation than  $r^{A1a}$ .

# 4.4 Dominant patterns in altimetric residuals

Fig. 11 and 12 show results of the PCA performed on the altimetric residuals and residual differences, respectively (Section 3.5.4).

The first three modes explain together 22 % of the residual variance and 20 % dominant modes explain 23 % of the variance of altimetric residuals (Fig. 11e) and 19 % of the variance of residual differences — (Fig. 12e). The first mode of the residual differences captures 10 % and its temporal pattern reveals a prominent drop between July 2010 and January 2011. Due to 2011 (Fig. 12d). Due to the data standardisation prior to PCAthe spatial patterns, the EOFs cannot be directly interpreted as amplitudes in elevation changeof the respective temporal patterns. For this reason, we rescale the spatial patterns by multiplying them with. For their presentation (Fig. 11a—c and 12a—c), we restored the signal amplitudes for each grid cell by multiplying the std of each time series (the time series, which was used beforehand to normalise the time series). Thereby, we regain interpretable magnitudes of the spatial patterns. Fig. 11a—f and 12a—f illustrate the version-dependent original and rescaled spatial patterns for r<sup>A1a</sup> and r<sup>A1a</sup>, respectively. (For all versions and both PCA, the original and rescaled patterns are illustrated in Fig. S24—S26).

## 615 5 Discussion

600

#### 5.1 Interannual firm thickness variations

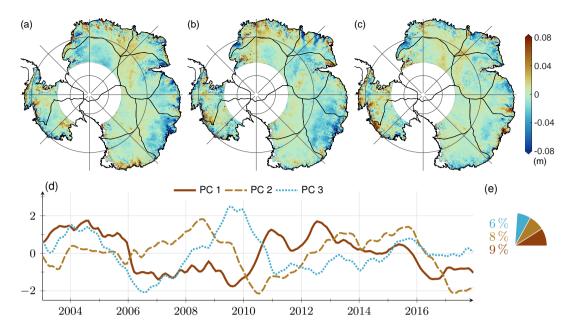


Figure 11. PCA results of standardised altimetric residuals for the period after 2003. (a–c) First three spatial patterns (EOFs) – version-dependent, shown here for  $r^{A1a}$  and with restored signal amplitudes for each grid cell. (d–f) Rescaled first three The EOFs for  $r^{A1a}$  of all versions are illustrated in Fig. S24 and S25. (gd) First three temporal patterns (PCs) determined from the aggregated data sets of  $r^{A1a}$ ,  $r^{A1b}$ ,  $r^{A2a}$  and  $r^{A2b}$ . (he) Associated percentages of the total residual variance considering the respective PC–EOF pairs. We define the PCs as standardised time series (mean of zero, std of 1) and without a unit while the EOFs have the unit of metre.

Adjusted In general, adjusted firn thickness variations  $fv^A$  (e.g. Fig. 7a for version A1a) and modelled firn thickness variations  $fv^M$  (e.g. Fig. 7b for Ma) share the same general spatial patterns. The largest magnitudes are found at lower elevations near the ice sheet margins with median rms values in the range of decimetres. The smallest magnitudes are found over the plateau of the EAIS with median rms values in the range of centimetres (Section Sect. 4.2.3). This general spatial pattern was to be expected, as it is related to the spatial variability of SMB. Snowfall, the main driver of Antarctic SMB variability, increases from the dry, relatively flat and homogeneous interior to the steep and complex topography of the wetter coastcoastal conditions. High snowfall at the ice sheet margins occurs due to orographic precipitation, influenced by the winds and topography of the AIS (Lenaerts et al., 2019)(van Wessem et al., 2014; Lenaerts et al., 2019).

620

625

630

The power spectral density (psd) of adjusted firm thickness variations  $fv^A$  decreases from low to high frequencies with reveal a strong temporal autocorrelation through the strong decrease of their psd with frequency, with spectral indices  $\kappa \le -3$  for a power-law noise (Section 4.3, Fig. 10a). The strong temporal autocorrelation observed in the interannual firm signals go model (Sect. 4.3). This is in line with the findings of King and Watson (2020). They estimated the power-law noise parameter  $\kappa$  in the range of -2.3 to -2.2 and -3.0 to -2.6 based on SMB estimates from RACMO2.3p2 and ice core composites, respectively. (Unlike our analysis, they did not co-estimate a quadratic or seasonal term.) only co-estimated a linear trend.

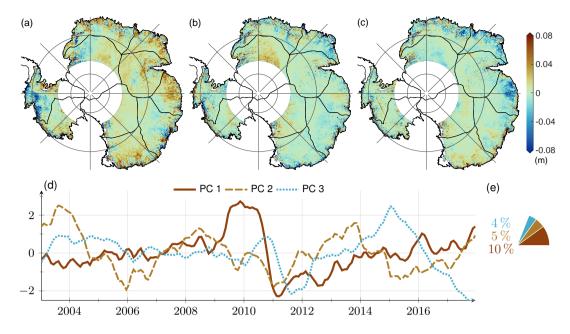


Figure 12. PCA results of standardised altimetric residual differences for the period after 2003. (a–c) First three spatial patterns (EOFs) – version-dependent, shown here for  $r^{A1a} - r^{A2a}$  and with restored signal amplitudes for each grid cell. (d–f) Rescaled first three The EOFs for  $r^{A1a} - r^{A2a}$  of all versions are illustrated in Fig. S26 and S27. (gd) First three temporal patterns (PCs) – the joint basis of  $r^{A1a} - r^{A2a}$  and  $r^{A1b} - r^{A2b}$ . (he) Associated percentages of the total variance of residual differences considering the respective PC–EOF pairs.

In the following, we compare how much variance of altimetric variations (for the period after 2003) can be explained according to the applied approach and the two different spatial considerations used previously: First, namely, first, the percentages assessed from grid cell time series and then averaged over the entire area. Second, and second, the percentages from time series averaged over the entire area ('mean Antarctic' time series, Fig. 13). The modelled firn thickness variations,  $fv^{Ma}$ , explain 11% and  $\frac{63\%64\%}{64\%}$  for the two spatial considerations, respectively (Table A1, columns E1 and  $\overline{E1}$ ). The scaled firm thickness variations,  $fv^{E2}$ , explain  $\frac{35\%}{40\%}$  and  $\frac{73\%31\%}{40\%}$  and  $\frac{71\%}{40\%}$  (Table A1, columns E2 and  $\overline{E2}$ ), respectively. The modified adjusted firm thickness variations,  $fv^{E3}$ , explain  $\frac{42\%}{40\%}$  and  $\frac{82\%}{40\%}$  and  $\frac{79\%}{40\%}$  for the two spatial considerations (Table 3, columns A1a and  $\overline{A1a}$ ). Our regression approach (Eq. 1), which generates  $fv^{A1a}$ , explains the greatest a larger part of the variance of altimetric variations compared with than the other approaches. This applies not only for the estimates considering each grid cell equally, but also for the estimates based on time series averaged over basins or the entire area. Furthermore, the The spatial scale investigated is crucial for the results, as the estimates from the basin-mean time series explain more of the altimetry variance than the estimates considering each grid cell equally. However, the latter are needed to enable the investigation and further interpretation of regression results based on their spatial patterns understand the spatial patterns of firm variations.

635

640

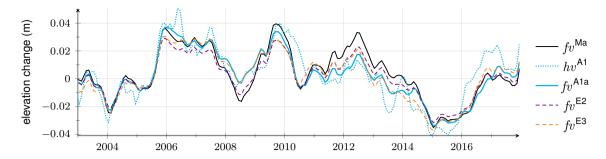
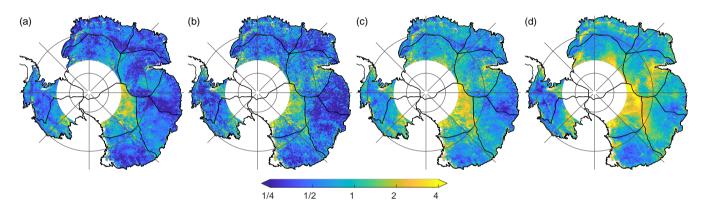


Figure 13. Mean Antarctic interannual elevation changes depending on the applied approach. Altimetric Modelled firm thickness variations ( $fv^{\text{Ma}} fv^{\text{Ma}}$ ), modelled firm thickness altimetric variations ( $fv^{\text{Ma}} hv^{\text{Al}}$ ), adjusted firm thickness variations ( $fv^{\text{E1}}$ ) and modified adjusted firm thickness variations ( $fv^{\text{E3}}$ ).

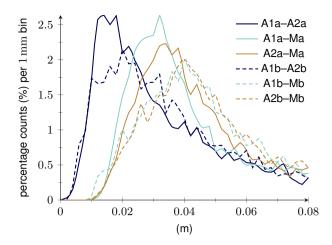


**Figure 14.** (a) Rms of (the time series of) the differences  $fv^{A1a} - fv^{A1b}$ . (b) Rms of the differences  $fv^{A1a} - fv^{A2a}$ . (c) Uncertainty estimate of  $fv^A$ : Maximum maximum rms of any combination of differences within versions of  $fv^A$ . (d) Rms of the residual differences  $r^{A1a} - r^{A2a}$  considering only the period after 2003. All rms maps (a–d) are normalised by the rms of  $fv^{Ma}$ .

## 5.2 Uncertainty and robustness of adjusted firn thickness variations

645

The adjusted firn thickness variations,  $fv^A$ , include the effects of firn model errors and altimetry errors. The differences  $fv^{A1a} - fv^{A1b}$  (Fig. 14a) and  $fv^{A2a} - fv^{A2b}$ , evaluated at every grid cell, are used to assess the influence of different firn model setups on  $fv^A$ . The median values (over all grid cells) of absolute and relative differences (A1a-A1b) are ~2.6 cm and ~51 % are in the range of 2.3 to 2.7 cm and 47 to 52 %, respectively (Table 2, Fig. 8). The differences  $fv^{A1a} - fv^{A2a}$  (Fig. 14b) and  $fv^{A1b} - fv^{A2b}$ , evaluated at every grid cell, are used to assess the influence of different altimetry analysis on  $fv^A$ . The median values (over all grid cells) of absolute and relative differences (A1a-A2a) are ~2.9 cm and ~58 % are in the range of 2.7 to 2.8 cm and ~54 %, respectively (Table 2, Fig. 8). Both the firn model and altimetry errors are discussed in Sections Sect. 5.3 and 5.4 separately. To assess the combined influence of firn model and altimetry errors on  $fv^A$ , the maximum deviation within the different versions of  $fv^A$  is used (Section Sect. 3.5.1). Fig. 14c shows the map of the maximum rms values. The median values (over



**Figure 15.** Histograms of the temporal rms of differences between various versions of firn thickness variations assessed at each grid cell of basin 3.(a) Histograms. (b) Cumulative histograms.

all grid cells) of absolute and relative (maximum) differences are ~4.3 cm and ~82 %~4.2 cm and ~80 %, respectively. In addition, median values are calculated for every basin ; i. e. over all grid cells within the respective basins. separately. The absolute and relative uncertainties range from 2.3 cm 2.2 cm (basin 8) to 10.9 cm 10.6 cm (basin 10) and from 59 % 54 % (basin 5) to 189 % 186 % (basin 8), respectively. We consider these estimates to be rough, but rather conservative uncertainty assessments for the adjusted firn thickness variations. In addition to the evaluation at grid cell level, the uncertainty of  $fv^A$  is assessed by time series differences of the basin means. (See Fig. S20 for the basin-mean time series of the four versions of  $fv^A$ ). The associated uncertainties per basin range from 1.0 cm 0.9 cm (basin 4) to 6.7 cm 6.4 cm (basin 10). The relative uncertainties are in the range of 21 % 20 % (basin 2) to 111 % 108 % (basin 8). For mean Antarctic  $fv^A$  an absolute and relative uncertainty of ~1.4 cm and ~71 %~1.3 cm and ~66 %, respectively, are estimated.

665

670

To-We assess the robustness of  $fv^A$ , statistical tests were carried out (Section 3.5.2). In particular through statistical tests according to Sect. 3.5.2. For each basin, four tests per basinare conducted, each comparing the temporal rms of the following pair of differences in firn thickness variations are conducted: Test (1) compares A1a-A2a to A1a-Ma, test (2) compares A1a-A2a to A2a-Ma, test (3) compares A1b-A2b to A1b-Mb and test (4) compares A1b-A2b to A2b-Mb. For all 40 tests, H0 is rejected (at the 5% significance level) and thus, H1 is accepted. This means that the differences within  $fv^A$  are significantly smaller than the adjustments, i.e. the differences between  $fv^A$  and  $fv^M$ , and that  $fv^A$  can be described as an improvement over  $fv^M$ . Fig. 15 a exemplifies the distributions of the differences for basin 3. (The histograms and cumulative histograms for all basins are shown in Fig. S28 and S29, respectively.) The results of the statistical tests demonstrate that  $fv^A$  is relatively robust to the choice of data sets, firn models and altimetry products. The choice of data sets does not significantly influence  $fv^A$ . Consequently, the assumption that  $fv^A$  represents a significant improvement over the modelled variations is reasonable. Limitations are discussed below.

## 675 5.3 Firn model errors

680

685

690

695

700

Firn model errors arise due to firn signals not (correctly) represented by either from firn signals that are not simulated or not correctly represented by the firn model or its input from RCMs or even and reanalysis data. Differences between  $fv^{Ma}$  and  $fv^{Mb}$  They are partly reflected in the differences  $fv^{Ma} - fv^{Mb}$  (Fig. S16) as well as differences between and the adjustments over the firn models, i.e. any version of  $fv^{A}$  and  $fv^{M}$   $fv^{A} - fv^{M}$  (Fig. S17)reflect firn model uncertainties and errors. Partly,  $fv^{A} - fv^{M}$  the adjustments also include errors related to the altimetry measurements and analysis of the altimetry products, as discussed in SectionSect. 5.4. (See also Table 2 for an overview of the various differences in firn thickness variations fv and their description.) Firn models generally show a smoother, more blurred spatial pattern than altimetry. This can be seen in the different fv (Fig. 7b versus 7a) and also in the different scaling factors e (cf. Fig. 6d–f versus 6a–b). One reason for this may be the lack of a–b and also Fig. 7b versus 7a). Reasons may be small-scale, mainly wind-driven processes that are missing in the model physics or not resolved in the firn modelling outputs (Lenaerts et al., 2012, 2019), leading to a blurred spatial distribution of modelled firn thickness variations ame detail due to the coarser spatial resolution of the models (Lenaerts et al., 2012, 2019)

The spatial patterns of absolute differences within  $fv^M$  and between any version of  $fv^A$  and  $fv^M$  of the adjustments (the adjustments, e.g. Fig. 7c), follow the spatial pattern of the signal itself. The greatest differences occur at the margins, where the climate is wetter and temperatures and accumulation are higher than inland. Especially in these coastal regions of high-relief topography, the horizontal resolution of the models, probably together with its physics, play an important role (Mottram et al., 2021). There, the differences between altimetry and firn models may be influenced by an incorrect or inaccurate spatial distribution of the modelled firn thickness variations (Fig. 6).

The modelled SMB components and their uncertainties have a direct impact on the modelled firn thickness. By assessing the spread of an ensemble of modelled firn thickness changes, Verjans et al. (2021) identified the RCMs as the largest contributor to the ensemble uncertainty. A precise parameterisation of firn compaction and surface snow density gains in importance in regions with high snowfall and large spatial variability of climatic conditions, such as Dronning Maud Land and Enderby Land (Verjans et al., 2021). However, the firn compaction rate in the IMAU and GSFC firn model both firn models used here is determined by constant mean annual accumulation and not by instantaneous overburden pressure. This lessens the actual firn compaction variability potentially across all the areas of large accumulation variability (Kuipers Munneke et al., 2015).

In a relative sense, the adjustments over the firn models (that is, any version of  $fv^A - fv^M$ , (e.g. Fig. 7d) generally increase from the coast to the EAIS interior as the magnitude of the signal, the firn thickness variation, is very small in the interior due to the cold and dry climate. In these areas of low snowfall, the relative uncertainties in the firn models are virtually unaffected by the formulation of firn densification and surface snow density, but the input of RCM components is essential (Verjans et al., 2021). Scambos et al. (2012) argue that RCMs might overestimate SMB in wind-glazed areas. These areas feature wind-polished glazed surfaces at the top of a coarsely recrystallised firn layer and are formed by constant katabatic winds. They have near-zero SMB and occur on leeward faces of ice-sheet undulations and megadunes (Scambos et al., 2012). Large wind glazed areas are located across basin 4 and 8, where all four versions of adjustments reveal highest relative values (Fig. S17e–h).

In basin 4, towards the border to boundary with basins 1 and 3, the large relative adjustments (Fig. S17e-h) indicate dis-710 agreement between the models and altimetry. Neither the uncertainties due to different altimetry analysis nor the influence of the different firn model setups have a strong impact on  $fv^A$  in this region, whereas the four versions of altimetry agree (Fig. S15i-l). Furthermore, within  $fv^{\rm M}$  there are no large differences in this region of basin 4 and the two models agree (Fig. S16d). The two models agree and the four versions of altimetry agree, but the models and altimetry do not agree. The reasons why discrepancies occur particularly in this region reasons for this are not yet clear. Basin 8 is characterised by large megadune fields (Fahnestock et al., 2000; Dadic et al., 2013). Megadune fields cover more than 500,000 km<sup>2</sup> of the East AIS 715 plateau. The megadunes Megadunes typically have an amplitude of 2 to 4 m and wavelengths of 2 to 5 km and are formed by a complex interaction of surface topography, snow accumulation and redistribution due to highly persistent katabatic winds. While leeward slopes are wind glazed, windward slopes accumulate and are characterised by streamlined bumps or grooves (sastrugi) sastrugi up to 1.5 m in height (Fahnestock et al., 2000; Frezzotti et al., 2002). The discrepancy between altimetry and the firn models across basin 8 can partly be explained by the lacking modelling of the formation of the complex spatial 720 pattern of megadunes and their migration over time in the firn models. For In case of basin 8, not only do the models and altimetry not match, but the relative differences between fv Ma and fv Mb disagree (Fig. S17e-h), as well as the different versions of  $tv^{M}$  (Fig. S16d) and between the different versions of  $tv^{A}$  (discussed in Section 5.4) are also high Fig. S15i–I). The latter is discussed in Sect. 5.4.

Discrepancies within the four versions of adjustments adjustments (i.e. within versions of  $fv^A - fv^M$ ) can further indicate which firn model(, or which dominant patterns of one firn model), fits the altimetry better. Overall, the adjustments are smaller when involving Ma, the IMAU firn model (Fig. 8, Table 2). Amongst the different basins, this applies in particular for basins 4-6 ((see Fig. S28 d-f and 29Sd-f) and S29 solid green/brown versus dash-dotted green/brown), this applies in particular for basins 4-6. Across basin 4-2 the adjustments tend to be slightly smaller when involving Mb, the GSFC firn model (Fig. S28a and S29a).

725

730

735

Altimetric residuals,  $r^A$ , still include a non-negligible part ( $\frac{53\%}{60\%}$  for A1a) of the variance of altimetric variations (Fig. 9c, Table 3). It Since the dominant patterns were chosen such that they cover at least 90% of the variance of  $fv^M$ ,  $r^A$  could partially contain real firn signals captured by firn models in the remaining  $\sim 10\%$  of the data variance. However, it is likely that a larger part of  $r^A$  still include includes real firn signals not captured by the dominant temporal patterns of the firn models. The psd of the underlying time series of  $r^{A1a}$  yield a spectral index of -1.7 (Section Sect. 4.3, Fig. 10b). The remaining autocorrelation in the residuals suggests that temporally correlated signals such as real firn signals are still present. Also, the spatial patterns of the most dominant modes of  $r^A$  reveal topography-dependent magnitudes and patterns, as one would expect from SMB and its variations (Section Sect. 4.4, Fig. 11d-fa-c). Besides other firn signals, the altimetric residuals additionally include altimetry errors (discussed in Section Sect. 5.4) and probably also further signals related to variations in ice flow dynamics or subglacial hydrology (not further discussed discussed further).

## 5.4 Altimetry errors

750

755

760

The differences between any version of  $fv^A$  and  $fv^M$  (the adjustments adjustments, e.g. Fig. 7c) may include effects of altimetry errors, in addition to firn model errors. Measurement noise in altimetry Noise in the altimetry measurements might explain another part of the fact that firn models show a smoother spatial pattern of variations than altimetry. Noise in altimetry can be a problem, especially in the interior of the EAIS where the signal-to-noise ratio is low(Section 5.5). Over megadune areas (widely located in the interior across basin 8), conventional radar altimetry with pulse-limited footprints of 1.5 to 2.5 km in diameter may not be capable of adequately observing the time-varying spatial patterns of megadunes.

A further limitation in radar altimetry is that measurements refer to the local topographic maxima within their footprints. Especially at the margins over complex topography, this can lead to sampling issues, as the elevation changes acquired there cannot capture the larger changes often found in the valleys. Laser altimeters are not affected by this since their footprints are much smaller (in the range of decimetres) sampling issue. However, ICESat had to be since ICESat operated in campaign mode (Abshire et al., 2005). Thus, the sampling in areas with steep slopes can vary strongly during the period 2003–2009 as some of the months rely only on radar altimetry measurements while other months include measurements from radar and laser, with some months including laser altimetry and some months relying exclusively on radar altimetry. Moreover, radar altimetry results are affected by the time-varying radar waveform shape due to time-varying signal penetration (Davis and Ferguson, 2004; Rémy et al., 2012). Even though errors related to these effects are accounted for in the altimetry processing, they are not fully eliminated and related residual errors may have an impact on the adjustments. In addition, these time-variable errors also likely to be These errors, which tend to be correlated in time, are likely included in the altimetric residuals  $\tau^A$ , because  $\tau^A$  are temporally correlated just as the errors (Sectionwhich may explain, to some part, the temporal correlation of  $\tau^A$  (Sect. 4.3, Fig. 10b).

Discrepancies within the adjustments (any version i.e. within versions of  $fv^A - fv^M$ ) can indicate which altimetry solution is closer to the firn models. Overall, the adjustments are smaller when involving A1, TUD altimetry However, results are equivocal (Fig. 8, Table 2). Amongst the different basins this applies in particular for basins 1, 5 and 6 (When involving the Ma firn model, the adjustments through A1 are smaller than those through A2 for most basins (see Fig. S28 b, e, f and S29 b, e, f). Across basin 8 the adjustments tend to be smaller when involving green solid versus brown solid). When involving the Mb firn model instead, the adjustments are in the same order of magnitude for A1 and A2, JPL altimetry and it depends on the basin whether the adjustments are smaller with A1 or A2.

Uncertainties due to a different analysis of the altimetry measurements are reflected by the differences in  $fv^A$  (Fig. 829h).

The differences in  $fv^A$  and the altimetric residuals, 14b) and  $r^A$ , (Fig. 14d) between solutions based on the same firn model (A1a-A2a or A1b-A2b)are displayed in Fig. 14b and d, respectively. They mirror the altimetry uncertainties due to a different analysis of the altimetry measurements. The median values (over all grid cells) of the absolute and relative residual differences  $r^{A1} - r^{A2}$  are  $\sim 4.7$  cm rms differences  $r^{A1a} - r^{A2a}$  in the time period after 2003 are  $\sim 4.9$  cm and  $\sim 96\%$ , respectively. The residual differences are evaluated for every grid cell and only the time period after 2003 is considered, in an absolute and

relative sense, respectively. If the entire period was considered, the median values would increase considerably ( $\sim$ 7.2 cm and  $\sim$ 162% $\sim$ 7.3 cm and  $\sim$ 163%). For both periods, the residual differences are greater than the differences  $fv^{A1} - fv^{A2} = fv^{A1a} - fv^{A2a}$  (Table 2, Fig. 14b) and also greater than the uncertainty estimate of  $fv^{A}$  (Section Sect. 5.2, Fig. 14c). Thus, the altimetry uncertainties in the residuals are greater than the combined uncertainties of firn modelling and altimetry affecting the adjusted firn thickness variations.

The differences between  $fv^{A1}$  and  $fv^{A2}$  as well as between  $r^{A1}$  and  $r^{A2}$  mostly result from the combined effect of the various differences between the altimetry analysis of TUD and JPL (Section Schröder et al. (2019a) and Nilsson et al. (2022) (Sect. 2.1). The rms of  $fv^{A1a} - fv^{A2a}$  is shown in Fig. 14b in a relative sense. The largest relative differences occur in regions of complex topography, such as in Victoria Land (at the margin of basin 7) and next to the Amery Ice Shelf (at the margin of basin 4) and over almost the entire basin 8, for which we already discussed the possible influence of megadunes. In addition, stripes related to the satellite ground tracks are visible in the region of basins 1 to 2 (Fig. 14b). They seem to appear predominantly in  $fv^{A2}$  (Fig. S15b and d).

The following features may likely be quite clearly attributed to a difference in intermission /and intermode calibration between TUD and JPL altimetry the two altimetry products. The mode change of CryoSat-2 (LRM/SARIn mode; see e.g. Fig. 5 in Slater et al. (2018) for the mode boundaries) is reflected in the difference of the residuals (Fig. 14d). Here, the main influence seems to come from JPL A2 altimetry, as the areas at the mode boundary in basins 5–7 and 9–10, characterised by a higher rms value, are mainly visible in  $r^{A2}$  (Fig. S18f and h). In addition, the mode transition also appears to be reflected in  $fv^{A2}$  particularly at basins 5 and 6 (Fig. S15b and d). The PCA carried out on  $r^{A1a} - r^{A2a}$  and  $r^{A1b} - r^{A2b}$  reveal a prominent drop between July 2010 and January 2011 together with overall linear trends before and after this drop in the first PC (Fig. 12gd). The corresponding spatial pattern (Fig. 12a S26a and b) is most pronounced and coherent over the EAIS. The pattern of the first mode is an indicator for uncertainties/differences differences and uncertainties in deriving intermission offsets, as CryoSat-2 measurements begin in July 2010. The errors in the altimetry are not only seen in the first modes of the PCA of the residual differences. It is likely that the first modes of the PCA of the residuals themselves are also affected by also contain altimetry errors. A comparison of the dominant modes of the residuals (Fig. 11) with those of the residual differences (Fig. 12) indicates partly similar features, which suggests similar causes. For example, there are also remarkably large fluctuations in the first temporal patterns of the residuals between July 2009 and January 2011 (Fig. 11gd).

## 5.5 Limitations of the approach

780

795

800

805

In regions of low signal-to-noise ratio the regression approach has a limited capability to distinguish between signal and error. This applies in particular for the interior of the EAIS (basin 8 and parts of basin 1 and 4). In these areas, the regression of the altimetry data to  $PC^{M}$  (the dominant temporal patterns in modelled firn thickness variations) may be dominated by noise in the altimetry data. In this study, we work with a constant spatial grid resolution of  $10 \,\mathrm{km} \times 10 \,\mathrm{km}$  regardless of the signal magnitude each grid cell. To improve the signal-to-noise ratio, further work may geographically vary and adjust the spatial resolution choose a geographically varying spatial resolution adapted to the spatial variability of the glaciological processes, that is in general a higher resolution on the coast and a which would probably imply a coarser resolution in the interior.

We included altimetry measurements only over the period May 1992 to December 2017 as this represents the common period of TUD and JPL altimetry (Section both altimetry products (Sect. 2). JPL A2 altimetry data, however, are available until December 2020. Further investigations may include the JPL data after December 2017, hence extend the period to the more recent past. These may incorporate accurate laser measurements from ICESat-2 characterised by low noise level and near-zero signal penetration (Nilsson et al., 2022; Otosaka et al., 2023a).

The stochastic model in the regression approach does not include co-variances temporal error covariances in altimetry (SectionSect. 3.3), although errors in the altimetry time series exhibit temporal correlations, as shown by Ferguson et al. (2004) and also in this study (SectionSect. 4.3). The consideration of temporal correlations is essential for a proper, realistic uncertainty estimation of assessing more realistic uncertainties. In particular, this is the case for long-term trends in particular (Williams et al., 2014). Thus, for inferring potentially statistically significant long-term signals in satellite altimetry future work may extend the stochastic model. This requires a comprehensive error characterisation for altimetry products, which is not given-provided up to now. Nevertheless, An empirical error characterisation could apply different noise models (e.g., power-law, Generalized Generalised Gauss Markov, auto-regressive) could be considered to empirically identify and apply the best fitting noise model to the to the regression approach (Bos et al., 2012; King and Watson, 2020). Another possibility for characterising errors could be the consideration Alternatively, the spread of an ensemble of altimetry solutions and their spread as demonstrated by Willen et al. (2022) could be considered (Willen et al., 2022).

Our study does not include independent observations to validate the benefits of  $fv^A$ . Most of the

## 5.6 Outlook

810

815

820

825

830

835

840

We do not aim here to compare our results with in situ data, as the ground-based SMB observations are mostly single point measurements and have a very sparse spatial and temporal coverage (Eisen et al., 2008). Thus, a validation However, future investigations may assess the benefits of  $fv^A$  could only be performed for selected, distinctly local regions and/or certain time intervals. A conceivable comparison could make use of stakes observations, as in the studies of Mottram et al. (2021) across Antarctica and Richter et al. (2021) in the Lake Vostok region certain regions with in situ data, e.g. by making use of stake observations (Mottram et al., 2021; Richter et al., 2021).

## 5.6 Outlook

To improve firn model outputs, we underline the importance of refining the it is important to refine the horizontal spatial resolution of RCMs and to simulate surface processes at a higher spatial distribution resolution (Lenaerts et al., 2019). For Greenland, Noël et al. (2016) statistically downscaled outputs from RACMO2.3 at 5.5 and 11 km to a high-resolution product resolution of 1 km, leading which led to, e.g.to, increased melt over certain areas. Similar work is in progress for Antarctica, downscaling RACMO2.3p2 at 27 km to 2 km (Noël et al., 2023). Furthermore, a more detailed physical parameterisation of the processes already considered and the inclusion of processes not yet simulated can improve the models (Agosta et al., 2019; Gutiérrez et al., 2021). An update of RACMO2.3p2 to RACMO2.4 with enhanced physics may soon be available. This includes

several new and updated parameterisations, such as a cloud, aerosol and radiation scheme or a new spectral albedo and radiative transfer scheme in the snow scheme (van Dalum and van de Berg, 2023).

To improve altimetry products, measurement noise noise in the altimetry measurements and correlated altimetry errors related in particular to time-variable radar signal penetration and scattering effects could be reduced by improving the methods of analysisneed to be reduced. Helm et al. (2023) developed a new retracker-processing scheme (retracker) based on a deep convolutional neural network architecture, resulting in presumably strongly reduced time-variable signal penetration. The new retracker effects, which could significantly improve the accuracy of elevation change products from the entire sequence of radar altimetry missions. Furthermore, improving the methods for intermission calibration would reduce uncertainties in altimetry estimates at various time scales Moreover, the intermission calibration needs further investigation. The patterns of estimated intermission offsets are spatially variant and are related to the waveform parameters(, possibly associated to topography and surface properties play a role here). However, this relation is not fully understood, so that no functional relationship has yet been found and intermission offsets are determined empirically (Zwally et al., 2005; Khvorostovsky, 2012; Schröder et al., 2019a; Nilsson et al., 2022). Therefore, intermission calibration still remains one of the most challenging processing steps for inferring a long-term, multi-mission satellite altimetry estimate.

Future developments in firn modelling, satellite altimetry analysis and altimetry mission sensors will allow to identify interannual firn signals to be identified and quantified with higher accuracy and, thereby, to better isolate and quantify long-term trends. This will further impact improve long-term estimates and reduce their uncertainties. The regression approach presented in this study may set the stage for isolating long-term signals in satellite altimetry from the large interannual variations. For this reason To this end, future studies should extend the approach with an appropriate stochastic model that accounts for covariances in altimetry to derive statistically significant long-term trends over 25 to 30 years. Longer (altimetry) time series will then further reduce trend uncertainties. With longer time series, trend uncertainties will be further reduced (Wouters et al., 2013). In this way, large uncertainties in inferring mass balance estimates of the EAIS (Otosaka et al., 2023b) may be reduced and the question whether the EAIS is currently thickening or thinning (Nilsson et al., 2021) may be answered in the future.

# 6 Conclusions

845

850

855

860

We developed a new approach for combining that combines satellite altimetry and firn modelling estimates results to resolve Antarctic firn thickness variations at a high temporal (monthly) and spatial (grid scale of 10 km) resolutionand spatial resolution, namely by monthly 10 km grids. On the one hand, our approach incorporates the strengths of the firn model models, above all the capability to capture the timing of firn thickness variations. On the other hand, our approach compensates for shortcomings of the firn model, foremost the accurate models, foremost in the simulation of the location-dependent amplitudes of the variations. To do so, we fitted dominant temporal patterns of interannual to decadal variations in Antarctic firn thickness inferred from the firn models IMAU (Veldhuijsen et al., 2023) and GSFC (Medley et al., 2022a) from Veldhuijsen et al., 2023) and Medley et al., (2022a) to satellite altimetry observations from TUD (Schröder et al., 2019a) and JPL (Nilsson et al., 2022)

Schröder et al. (2019a) and Nilsson et al. (2022). In this way, we generated a new, combined data set product, which we named the adjusted firm thickness variations,  $fv^A$ .

Our guiding question was: How well can satellite altimetry and firn models resolve Antarctic firn thickness variations? Well it depends. This study shows that firn models and altimetry products provide complementary information on firn thickness variations. The combined data set,  $fv^A$ , characterises spatially resolved variations better than either (1) firn models alone or (2) altimetry alone. (1) The adjusted firn thickness variations,  $fv^A$ , outperform the modelled firn thickness variations,  $fv^M$ , because Compared with  $fv^M$  for improves the amplitudes of the variations compared with  $fv^M$ . The amplitudes represent an improvement because they are observed by the altimeter satellites and their patterns actually indicate more spatial and thereby meaningful information. However, one caveat should be noted. The the improved observed amplitudes may also include effects of altimetry errors due to firn penetration. This is because the temporal variations of these errors correlate with the temporal variations of the signal, as both the time-variable signal and the these errors are influenced by the SMB and firn processes and are thus temporally correlated. (2) The adjusted firn thickness variations,  $fv^A$ , outperform the altimetric variations,  $hv^A$ , because  $fv^A$  eliminates a large part of the altimetry errors. If one were to take  $hv^A$  alone, this would also incorporate all the errors of  $hv^A$ . Over Antarctica, or rather the entire area studied, this would introduce median absolute and relative uncertainties of  $\sim 7.2$  cm and  $\sim 162\% \sim 7.3$  cm and  $\sim 163\%$ , respectively (evaluated on at grid cell level). However, one caveat should be noted. By by choosing  $fv^A$  instead of  $hv^A$ , part of the observed firn signal is ignored.

How well the  $fv^A$  resolve resolves real Antarctic firn thickness variations depends strongly on the region under investigation. Over all grid cells of Antarctica, median absolute and relative uncertainties of  $fv^A$  are ~4.3 cm and ~82%, respectively(evaluated on grid cell level). Over the basin areas ~4.2 cm and ~80%, respectively. Over all grid cells of individual basins, the median relative uncertainties range from 59% (are lowest for basin 5) to 189% (basin 8). Across basin 8, we also spatially resolved disagreements between  $fv^A$  and  $fv^M$ . (the region of Queen Mary Land), and highest for basin 8. The large uncertainty and the disagreement are in basin 8 is likely due to the presence of megadune fields. Overall, the differences between We find the smallest adjustment that  $fv^A$  and requires over  $fv^M$  are smallest when using the TUD altimetry and the IMAU firn model. Amongst the different basins, this is especially true altimetry data from Schröder et al. (2019a) and the firn model Veldhuijsen et al. (2023) and this is most prominent for basins 5 and 6. From the spectral analysis of the altimetry residuals,  $r^A$ , we find still autocorrelated signals that we could not attribute to firn thickness variations using the firn models. We attribute this to a combination of altimetry errors(, in particular time-variable signal penetration , and errors in intermission offsets) and, and to firn model errors(incorrectly simulated /missing processes).

We identified regions of discrepancy between the firn models and the altimetry products and within the models or altimetry, and discussed the underlying errors in both the models and the altimetry. These results shall help modellers and altimetry data processors to improve their simulations and processing methods (Sect. 5.6), and help users to better understand the nature of the modelling and altimetry data and to apply and interpret them knowing their strengths and limitations.

# Appendix A: Impact of methodological changes

#### A1 Methods

925

To investigate the impact of methodological changes on determining adjusted firm thickness variations,  $fv^A$ , three modifications of to the original regression approach are tested.

In the first experiment E1, we simply subtract the altimetric variations,  $hv^{\rm El}$  modelled firn thickness variations,  $fv^{\rm M}$ , from the modelled firn thickness variations,  $fv^{\rm M}$  altimetric variations,  $hv^{\rm A}$ , according to

$$r^{\mathrm{EI}}(t) = \frac{hv^{\mathrm{EI}}hv^{\mathrm{A}}(t) - fv^{\mathrm{M}}(t). \tag{A1}$$

fv<sup>M</sup> is derived by least squares fit according to Eq. 4. hv<sup>E1</sup> is derived by least squares fit according to

915 
$$\underline{h(t)^{\rm A}} = \underbrace{\frac{a + bt + c(0.5t^2)}{+ H_1(t) \left[ d_1 \cos(\omega t) + d_2 \sin(\omega t) + d_3 \cos(2\omega t) + d_4 \sin(2\omega t) \right]}_{+ H_2(t) \left[ d_5 \cos(\omega t) + d_6 \sin(\omega t) + d_7 \cos(2\omega t) + d_8 \sin(2\omega t) \right]}_{+ hv^{\rm E1}(t)}$$

with the parameters  $a, b, c, d_{1,...,8}$  and the masks  $H_1, H_2$  as in Eq. 1.

In the second experiment E2,  $fv^{M}$  at any grid cell is simply scaled to fit the altimetric variations. The regression reads

$$\underline{hhv}^{A}(t)_{-}^{A} = \underline{a + bt + c(0.5t^{2}) + H_{1}(t) d_{1}\cos(\omega t) + d_{2}\sin(\omega t) + d_{3}\cos(2\omega t) + d_{4}\sin(2\omega t) + H_{2}(t) d_{5}\cos(\omega t) + d_{6}\sin(\omega t) + d_{7}\cos(2\omega t)}$$
(A2)

where e is the scaling factor. We refer to  $e f v^{M} = f v^{E2}$  as scaled firn thickness variations.

In the third experiment E3, we do not change the principle of the deterministic model Eq. 1 but we modify the dominant temporal patterns  $PC^{M}$ . Originally,  $PC^{M}$  are derived from standardised  $fv^{M}$  by PCA. In E3,  $fv^{M}$  are not standardised prior to the PCA. The resulting modified adjusted firn thickness variations are referred to by  $fv^{E3}$ . See also Table B1 for an overview of the defined symbols and their terminology.

We consider that regression method as best whose eoefficient of determination R-squared value,  $R^2$ , is maximum, i.e. which is able to describe most of the data variance. For the three experiments, the general from of Eq. 5 specifies Eq. 3 modifies to

$$R_{\rm E1}^2 = 1 - \frac{SS(r^{\rm E1})}{SS(fv^{\rm Ma} + r^{\rm E1})} = 1 - \frac{SS(r^{\rm E1})}{SS(hv^{\rm E1})} \approx 1 - \frac{SS(r^{\rm E1})}{SS(hv^{\rm A})},\tag{A3a}$$

930 
$$R_{\rm E2}^2 = 1 - \frac{SS(r^{\rm E2})}{SS(efv^{\rm Ma} + r^{\rm E2})} = 1 - \frac{SS(r^{\rm E2})}{SS(fv^{\rm E2} + r^{\rm E2})} \approx 1 - \frac{SS(r^{\rm E2})}{SS(hv^{\rm E1})} \approx 1 - \frac{SS(r^{\rm E2})}{SS(hv^{\rm A})},$$
 (A3b)

$$R_{\rm E3}^2 = 1 - \frac{SS(r^{\rm E3})}{SS(fv^{\rm E3} + r^{\rm E3})} \approx 1 - \frac{SS(r^{\rm E3})}{SS(hv^{\rm E1})} \approx 1 - \frac{SS(r^{\rm E3})}{SS(hv^{\rm A})}.$$
 (A3c)

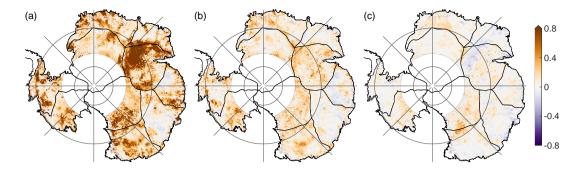


Figure A1. Differences between the coefficients of determination R-squared values from A1a and the experiments E1, E2 and E3. (a) A1a-E1, (b) A1a-E2 and (c) A1a-E3. Colour bar arrows indicate that the value range exceeds the limits of the colour scale.

Assuming that the changes in the adjusted parameters  $b, e, d_{1,...,8}$  due to different versions of regression are negligible, the following approximations are reasonable:  $(efv^{\text{Ma}} + r^{\text{E2}}) \approx (fv^{\text{E3}} + r^{\text{E3}}) \approx hv^{\text{E1}} \approx hv^{\text{A}}$ .

# A2 Results

The impact of methodological choices on the goodness of fit is tested based on the three modifications/experiments E1–E3 (Section Sect. A1). The results are given for using the IMAU\_Ma firn model and TUD\_A1 altimetry and should, therefore, be compared to the results from the regression approach A1a.

For every grid cell, Fig. A1 compares the coefficients of determination R-squared values from the regression approach A1a,  $R_{A1a}^2$ , to the coefficients of determination R-squared values  $R_{E1}^2$ ,  $R_{E2}^2$  and  $R_{E3}^2$ .  $R_{A1a}^2$  is larger than  $R_{E1}^2$ ,  $R_{E2}^2$  and  $R_{E3}^2$  over 96, 81 and 69 % 88, 78 and 66 % of the total area, respectively. After calculating  $R_{E1}^2$ ,  $R_{E2}^2$  and  $R_{E3}^2$  for each grid cell, (basin) mean values are derived and listed by Table A1, columns 2–4. Averaged over the entire area, E1, E2 and E3 have mean  $R^2$  values of 0.09, 0.35 and 0.430.11, 0.31 and 0.37. For all three modifications,  $R^2$  is smaller than  $R_{A1a}^2$  (Table 3, column A1a) and thus, their regression approaches describe less of the data variance than the original regression approach A1a. E2 and E3 describe describes slightly more of the data variance than A1a for one out of 10 basins (E1, basin 5: 44 versus 43 %; E2, basin basin 3: 49 versus 47 %47 versus 45 %). Moreover, Table A1 (columns 6–7) lists values of  $R^2$  derived from basin averages time series (E1, E2 and E3). Values derived from basin averages time series are larger than values based on the calculations per grid cell, similar to the regression approach A1a (Table 3, column  $\overline{A1a}$  versus A1a).

The simple scaling factor e adjusted during the regression approach after experiment E2 is displayed in Fig. S30.

# **Appendix B: List of symbols**

950 List of symbols and their terminology (columns 1–2). Sections and equations where the symbols are explained and defined (column 3). Different versions (see Table 1) of the respective symbols (column 4). Symbol Terminology References Versions  $h^{A}$  Altimetric elevation changes\* Section 3.1, Eq. 1 A1, A2  $hv^{A}$  Altimetric variations Section 3.1,2, Eq. 2 A1a, A2a, A1b,

**Table A1.** Explained variance or eoefficients of determination,  $R^2$ , for each basin and each experiment E1, E2, E3 of methodological changes to the regression approach over the period after 2003.  $R^2$  is first calculated for each grid cell according to Eq. A3a–A3c and after averaged over each basin. Values of  $\overline{E1}$ ,  $\overline{E2}$  and  $\overline{E3}$  are calculated by first averaging the results from the experiments over each basin and then applying Eq. A3a–A3c.

Basin	E1	E2	E3	E1	E2	<del>E3</del>
01	0.20	0.35-0.34	0.43 0.38	0.71	<del>0.75</del> <u>0.73</u>	<del>0.79</del> <u>0.78</u>
02	0.21	0.41-0.38	0.48-0.45	<del>0.76</del> <u>0.77</u>	0.92_0.88	0.91-0.88
03	0.21	0.44-0.42	0.50-0.47	0.88 0.89	0.94	0.95
04	-0.29	0.15-0.12	0.27-0.24	<del>-5.63</del> -5.52	<del>-0.08</del> 0.22	<del>0.10</del> -0.06
05	$0.02 \cdot 0.06$	0.42 0.26	0.34 0.26	<del>-0.50</del> - <u>0.19</u>	0.66 0.36	<del>0.47</del> -0.31
06	$\underbrace{0.20}_{}\underbrace{0.21}_{}$	0.38-0.29	$\underbrace{0.40}_{\textstyle \bullet}\underbrace{0.32}_{\textstyle \bullet}$	$\underbrace{0.70}_{0.73}\underbrace{0.73}_{0.00}$	0.86 0.76	0.90-0.80
07	0.22-0.23	0.44 - 0.40	0.54-0.49	0.68-0.72	0.91 $0.87$	0.94-0.92
08	-0.08	$\underbrace{0.13}_{}\underbrace{0.11}_{}$	0.17	0.23	<del>0.19</del> 0.48	0.54 0.56 0.50
09	0.32	0.42 0.39	<del>0.50</del> 0.47	0.94	$0.94 \underbrace{0.93}_{\bullet}$	<del>0.97</del> -0.96
10	0.27	0.49-0.46	0.59-0.56	<del>0.92</del> <u>0.93</u>	<del>0.98</del> 0.97	0.96
01-10*	0.11	0.35-0.31	0.42-0.37	0.63_0.64	0.73_0.71	0.82-0.79

<sup>\*</sup> refers to the entire area (considered as a single basin)

A2b  $f^{\rm M}$  Modelled firn thickness changes\* Section 3.1.1, Eq. 4 Ma, Mb  $fv^{\rm M}$  Modelled firn thickness variations Section 3.1.1, Eq. 4 Ma, Mb  $PC_{1...N}^{\rm M}$  Modelled firn thickness variations Section 3.1.1, Eq. 1, 3 Ma, Mb  $e^{\rm A}_{1...N}$  N observed scaling factors Section 3.1, Eq. 1, 3 A1a, A2a, A1b, A2b  $e^{\rm M}_{1...N}$  N modelled scaling factors Section 3.2 Ma, Mb  $fv^{\rm M}$  Truncated modelled firn thickness variations Section 3.2 Ma, Mb  $fv^{\rm A}$  Adjusted firn thickness variations Section 3.1.1, Eq. 3 A1a, A2a, A1b, A2b  $fv^{\rm E2}$  Scaled firn thickness variations Appendix A1  $fv^{\rm E3}$  Modified adjusted firn thickness variations Appendix A1

Data availability. The altimetry products from Schröder et al. (2019a) and Nilsson et al. (2022) are available at https://doi.pangaea.de/10. 1594/PANGAEA.897390 (Schröder et al., 2019b) and https://doi.org/10.5067/L3LSVDZS15ZV (Nilsson et al., 2021), respectively. The firn model data from Medley et al. (2022a) is available at https://doi.org/10.5281/zenodo.7054574 (Medley et al., 2022b). The code of the firn model from Veldhuijsen et al. (2023) is available at https://github.com/brils001/IMAU-FDM and https://zenodo.org/records/5172513 (Brils et al., 2021). The firn model data from Veldhuijsen et al. (2023) and the results of this study can be obtained from the authors without conditions.

Author contributions. CRediT Taxonomy: MK: conceptualization, data curation, formal analysis, investigation, methodology, software, visualization, writing - original draft. MH: conceptualization, funding acquisition, methodology, supervision, writing - review and editing. EB: formal analysis, writing - review and editing. MW: data curation, writing - review and editing. LS: funding acquisition, writing - review and editing. SV, PKM, MvdB: resources (provision of the firn model data), writing - review and editing.

Competing interests. MvdB is a member of the editorial board of the journal. All other authors declare that they have no conflict of interest.

Acknowledgements. MK and EB were funded by the Deutsche Forschungsgemeinschaft (DFG) as part of the Special Priority Program (SPP) 1158 'Antarctic Research with Comparative Investigations in Arctic Ice Areas' through grant HO 4232/10-1 (project number 442929109) and grant SCHE 1426/26-1 and 2 (project number 404719077), respectively. MW was funded through grant HO 4232/4-2 (project number 313917204) from the DFG as part of the SPP 1889 'Regional Sea Level Change and Society (SeaLevel)'. SV acknowledges funding by the Netherlands Organisation for Scientific Research (NWO) grant OCENW.GROOT.2019.091.

# 975 References

- Abshire, J., Sun, X., Riris, H., Sirota, J., MCGarry, J., Palm, S., Yi, D., and Liiva, P.: Geoscience Laser Altimeter System (GLAS) on the ICESat Mission: On–orbit measurement performance, Geophys. Res. Lett., 32, L21S02, https://doi.org/10.1029/2005GL024028, 2005.
- Agosta, C., Amory, C., Kittel, C., Orsi, A., Favier, V., Gallée, H., van den Broeke, M., Lenaerts, J., van Wessem, J., van de Berg, W., and Fettweis, X.: Estimation of the Antarctic surface mass balance using the regional climate model MAR (1979–2015) and identification of dominant processes. The Cryosphere, 13, 281–296, https://doi.org/10.5194/tc-13-281-2019, 2019.
- Arthern, R., Vaughan, D., Rankin, A., Mulvaney, R., and Thomas, E.: In situ measurements of Antarctic snow compaction compared with predictions of models, J. Geophys. Res., 115, https://doi.org/10.1029/2009JF001306, 2010.
- Barletta, V., Sørensen, L., and Forsberg, R.: Scatter of mass changes estimates at basin scale for Greenland and Antarctica, The Cryosphere, 7, 1411–1432, https://doi.org/10.5194/tc-7-1411-2013, 2013.
- 985 Bodart, J. and Bingham, R.: The Impact of the Extreme 2015-16 El Niño on the Mass Balance of the Antarctic Ice Sheet, Geophys. Res. Lett., 46, 13 862–13 871, https://doi.org/10.1029/2019GL084466, 2019.
  - Boening, C., Lebsock, M., Landerer, F., and Stephens, G.: Snowfall-driven mass change on the East Antarctic ice sheet, Geophys. Res. Lett., 39, L21 501, https://doi.org/10.1029/2012GL053316, 2012.
- Boergens, E., Rangelova, E., Sideris, M., and Kusche, J.: Assessment of the capabilities of the temporal and spatiotemporal ICA method for geophysical signal separation in GRACE data, J. Geophys. Res. Solid Earth, 119, 4429–4447, https://doi.org/10.1002/2013JB010452, 2014.
  - Bos, M., Fernandes, R., Williams, S., and Bastos, L.: Fast error analysis of continuous GNSS observations with missing data, J. Geod., 87, 351–360, https://doi.org/10.1007/s00190-012-0605-0, 2012.
- Brils, M., Kuipers Munneke, P., Van de Berg, W., and Van den Broeke, M.: IMAU-FDM v1.2G GDM release (v1.2G), https://doi.org/10.5281/zenodo.5172513, 2021.
  - Brockley, D., Baker, S., Femenias, P., Martinez, B., Massmann, F.-H., Otten, M., Paul, F., Picard, B., Prandi, P., Roca, M., Rudenko, S., Scharroo, R., and Visser, P.: REAPER: Reprocessing 12 Years of ERS-1 and ERS-2 Altimeters and Microwave Radiometer Data, IEEE Trans. Geosci. Remote Sens., 55, 5506–5514, https://doi.org/10.1109/TGRS.2017.2709343, 2017.
- Cullather, R., Bromwich, D., and van Woert, M.: Interannual variations in Antarctic precipitation related to El Niño-Southern Oscillation, J. Geophys. Res., 101, 19–109, 1996.
  - Dadic, R., Mott, R., Horgan, H., and Lehning, M.: Observations, theory, and modeling of the differential accumulation of Antarctic megadunes, J. Geophys. Res. Earth Surf., 118, 2343–2353, https://doi.org/10.1002/2013JF002844, 2013.
  - Davis, C. and Ferguson, A.: Elevation change of the Antarctic ice sheet, 1995-2000, from ERS-2 satellite radar altimetry, IEEE Trans. Geosci. Remote Sens., 42, 2437–2445, https://doi.org/10.1109/TGRS.2004.836789, 2004.
- Davison, B., Hogg, A., Rigby, R., Veldhuijsen, S., van Wessem, J., van den Broeke, M., Holland, P., Selley, H., and Dutrieux, P.: Sea level rise from West Antarctic mass loss significantly modified by large snowfall anomalies, Nature Communications, 14, https://doi.org/10.1038/s41467-023-36990-3, 2023.
  - Dee, D., Uppala, S., Simmons, A., Berrisford, P., Poli, P., Kobayashi, S., Andrae, U., Balmaseda, M., Balsamo, G., Bauer, P., Bechtold, P., Beljaars, A., van de Berg, L., Bidlot, J., Bormann, N., Delsol, C., Dragani, R., Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hersbach,
- 1010 H., Hólm, E., Isaksen, L., Kållberg, P., Köhler, M., Matricardi, M., McNally, A., Monge-Sanz, B., Morcrette, J.-J., Park, B.-K., Peubey,

- C., de Rosnay, P., Tavolato, C., Thépaut, J.-N., and Vitart, F.: The ERA-Interim reanalysis: configuration and performance of the data assimilation system, Quart. J. R. Met. Soc., 137, 553–597, https://doi.org/10.1002/qj.828, 2011.
- Diener, T., Sasgen, I., Agosta, C., Fürst, J., Braun, M., Konrad, H., and Fettweis, X.: Acceleration of Dynamic Ice Loss in Antarctica From Satellite Gravimetry, Front. Earth Sci., 9, https://doi.org/10.3389/feart.2021.741789, 2021.
- Eisen, O., Frezzotti, M., Genthon, C., Isaksson, E., Magand, O., van den Broeke, M., Dixon, D., Ekaykin, A., Holmlund, P., Kameda, T., Karlöf, L., Kaspari, S., Lipenkov, V., Oerter, H., Takahashi, S., and Vaughan, D.: Ground-based measurements of spatial and temporal variability of snow accumulation in East Antarctica, Rev. Geophys., 46, https://doi.org/10.1029/2006RG000218, 2008.
  - Fahnestock, M., Scambos, T., Shuman, C., Arthern, R., Winebrenner, D., and Kwok, R.: Snow megadune fields on the East Antarctic Plateau: Extreme atmosphere-ice interaction, Geophys. Res. Lett., 27, 3719–3722, https://doi.org/10.1029/1999GL011248, 2000.
- Ferguson, A., Davis, C., and Cavanaugh, J.: An Autoregressive Model for Analysis of Ice Sheet Elevation Change Time Series, IEEE Trans. Geosci. Remote Sens., 42, 2426–2436, https://doi.org/10.1109/TGRS.2004.836788, 2004.
  - Forootan, E. and Kusche, J.: Separation of global time-variable gravity signals into maximally independent components, J. Geod., 86, 477–497, https://doi.org/10.1007/s00190-011-0532-5, 2012.
- Fox-Kemper, B., Hewitt, H., Xiao, C., Drijfhout, S., Edwards, T., Golledge, N., Hemer, M., Kopp, R., Krinner, G., Mix, A., Notz, D., Nowicki, S., Nurhati, I., Ruiz, L., Sallée, J.-B., Slangen, A., and Yu, Y.: Climate Change 2021: The Physical Science Basis. Contribution of Working Group I to the Sixth Assessment Report of the Intergovernmental Panel on Climate Change, chap. Ocean, Cryosphere and Sea Level Change, pp. 1211–1362, Cambridge University Press, Cambridge, United Kingdom and New York, NY, USA, https://doi.org/10.1017/9781009157896.011, 2021.
- Fretwell, P., Pritchard, H., Vaughan, D., Bamber, J., Barrand, N., Bell, R., Bianchi, C., Bingham, R., Blankenship, D., Casassa, G., Catania,
  G., Callens, D., Conway, H., Cook, A., Corr, H., Damaske, D., Damm, V., Ferraccioli, F., Forsberg, R., Fujita, S., Gim, Y., Gogineni,
  P., Griggs, J., Hindmarsh, R., Holmlund, P., Holt, J., Jacobel, R., Jenkins, A., Jokat, W., Jordan, T., King, E., Kohler, J., Krabill, W.,
  Riger-Kusk, M., Langley, K., Leitchenkov, G., Leuschen, C., Luyendyk, B., Matsuoka, K., Mouginot, J., Nitsche, F., Nogi, Y., Nost, O.,
  Popov, S., Rignot, E., Rippin, D., Rivera, A., Roberts, J., Ross, N., Siegert, M., Smith, A., Steinhage, D., Studinger, M., Sun, B., Tinto,
  B., Welch, B., Wilson, D., Young, D., Xiangbin, C., and Zirizzotti, A.: Bedmap2: improved ice bed, surface and thickness datasets for
  Antarctica, The Cryosphere, 7, 375–393, https://doi.org/10.5194/tc-7-375-2013, 2013.
  - Frezzotti, M., Gandolfi, S., and Urbini, S.: Snow megadunes in Antarctica: Sedimentary structure and genesis, J. Geophys. Res., 107, https://doi.org/10.1029/2001JD000673, 2002.
  - Gardner, A., Moholdt, G., Scambos, T., Fahnstock, M., Ligtenberg, S., van den Broeke, M., and Nilsson, J.: Increased West Antarctic and unchanged East Antarctic ice discharge over the last 7 years, The Cryosphere, 12, 521–547, https://doi.org/10.5194/tc-12-521-2018, 2018.
- Gelaro, R., McCarty, W., Suárez, M., Todling, R., Molod, A., Takacs, L., Randles, C., Darmenov, A., Bosilovich, M., Reichle, R., Wargan, K., Coy, L., Cullather, R., Draper, C., Akella, S., Buchard, V., Conaty, A., da Silva, A., Gu, W., Kim, G.-K., Koster, R., Lucchesi, R., Merkova, D., Nielsen, J., Partyka, G., Pawson, S., Putman, W., Rienecker, M., Schubert, S., Sienkiewicz, M., and Zhao, B.: The Modern-Era Retrospective Analysis for Research and Applications, Version 2 (MERRA-2), Journal of Climate, 30, 5419–5454, https://doi.org/10.1175/JCLI-D-16-0758.1, 2017.
- 1045 Gossart, A., Helsen, S., Lenaerts, J., Broucke, S., van Lipzig, N., and Souverijns, N.: An Evaluation of Surface Climatology in State-of-the-Art Reanalyses over the Antarctic Ice Sheet, Journal of Climate, 32, 6899–6915, https://doi.org/10.1175/JCLI-D-19-0030.1, 2019.
  - Groh, A. and Horwath, M.: Antarctic Ice Mass Change Products from GRACE/GRACE-FO Using Tailored Sensitivity Kernels, Remote Sens., 13, 1736, https://doi.org/10.3390/rs13091736, 2021.

- Groh, A., Horwath, M., Horvath, A., Meister, R., Sørensen, L., Barletta, V., Forsberg, R., Wouters, B., Ditmar, P., Ran, J., Klees, R., Su, X., Shang, K., Guo, J., Shum, C., Schrama, E., and Shepherd, A.: Evaluating GRACE Mass Change Time Series for the Antarctic and Greenland Ice Sheet—Methods and Results, Geosciences, 9, 415, https://doi.org/10.3390/geosciences9100415, 2019.
  - Gutiérrez, J., Jones, R., Narisma, G., Alves, L., Amjad, M., Gorodetskaya, I., Grose, M., Klutse, N., Krakovska, S., Li, J., Martínez-Castro, D., Mearns, L., Mernild, S., Ngo-Duc, T., van den Hurk, B., and Yoon, J.-H.: Climate Change 2021: The Physical Science Basis. Contribution of Working Group I to the Sixth Assessment Report of the Intergovernmental Panel on Climate Change, chap. Atlas, pp. 1927–2058,

Cambridge University Press, Cambridge, United Kingdom and New York, NY, USA, https://doi.org/10.1017/9781009157896.021, 2021.

Helm, V., Humbert, A., and Miller, H.: Elevation and elevation change of Greenland and Antarctica derived from CryoSat-2, The Cryosphere, 8, 1539–1559, https://doi.org/10.5194/tc-8-1539-2014, 2014.

- Helm, V., Dehghanpour, A., Hänsch, R., Loebel, E., Horwath, M., and Humbert, A.: AWI-ICENet1: A convolutional neural network retracker for ice altimetry, The Cryosphere Discuss., 2023, 1–61, https://doi.org/10.5194/tc-2023-80, 2023.
- Helsen, M., van den Broeke, M., van de Wal, R., van de Berg, W., van Meijgaard, E., Davis, C., Li, Y., and Goodwin, I.: Elevation Changes in Antarctica Mainly Determined by Accumulation Variability, Science, 320, 1626–1629, https://doi.org/10.1126/science.1153894, 2008.
  - Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas, J., Peubey, C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan, X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara, G., Dahlgren, P., Dee, D., Diamantakis, M., Dragani, R., Flemming, J., Forbes, R., Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hogan, R., Hólm, E., Janisková,
- M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay, P., Rozum, I., Vamborg, F., Villaume, S., and Thépaut, J.-N.: The ERA5 global reanalysis, Quart. J. R. Met. Soc., 146, 1999–2049, https://doi.org/10.1002/qj.3803, 2020.
  - Horwath, M. and Dietrich, R.: Signal and error in mass change inferences from GRACE: the case of Antarctica, Geophys. J. Int., 177, 849–864, https://doi.org/10.1111/j.1365-246X.2009.04139.x, 2009.
- Horwath, M., Legrésy, B., Rémy, F., Blarel, F., and Lemoine, J.-M.: Consistent patterns of Antarctic ice sheet interannual variations from ENVISAT radar altimetry and GRACE satellite gravimetry, Geophys. J. Int., 189, 863–876, https://doi.org/10.1111/j.1365-246X.2012.05401.x, 2012.
  - Horwath, M., Gutknecht, B., Cazenave, A., Palanisamy, H., Marti, F., Marzeion, B., Paul, F., Le Bris, R., Hogg, A., Otosaka, I., Shepherd, A., Döll, P., Cáceres, D., Müller Schmied, H., Johannessen, J., Nilsen, J., Raj, R., Forsberg, R., Sandberg Sørensen, L., Barletta, V., Simonsen, S., Knudsen, P., Andersen, O., Ranndal, H., Rose, S., Merchant, C., Macintosh, C., von Schuckmann, K., Novotny, K., Groh, A.,
- Restano, M., and Benveniste, J.: Global sea-level budget and ocean-mass budget, with a focus on advanced data products and uncertainty characterisation, Earth System Science Data, 14, 411–447, https://doi.org/10.5194/essd-14-411-2022, 2022.
  - IPCC: Climate Change 2021: The Physical Science Basis. Contribution of Working Group I to the Sixth Assessment Report of the Intergovernmental Panel on Climate Change, chap. Summary for Policymakers, pp. 3–32, Cambridge University Press, Cambridge, United Kingdom and New York, NY, USA, https://doi.org/10.1017/9781009157896.001, 2021.
- Jolliffe, I.: Principal Component Analysis, Springer Series in Statistics, Springer New York, NY, 2 edn., https://doi.org/10.1007/b98835, 2002.
  - Kaitheri, A., Mémin, A., and Rémy, F.: Inter-Annual Variability in the Antarctic Ice Sheets Using Geodetic Observations and a Climate Model, Remote Sens., 13, 2199, https://doi.org/10.3390/rs13112199, 2021.
- Kappelsberger, M., Strößenreuther, U., Scheinert, M., Horwath, M., Groh, A., Knöfel, C., Lunz, S., and Khan, S.: Modeled and Observed
   Bedrock Displacements in North-East Greenland Using Refined Estimates of Present-Day Ice-Mass Changes and Densified GNSS Measurements, J. Geophys. Res. Earth Surf., 126, e2020JF005 860., https://doi.org/10.1029/2020JF005860, 2021.

- Khvorostovsky, K.: Merging and Analysis of Elevation Time Series Over Greenland Ice Sheet From Satellite Radar Altimetry, IEEE Trans. Geosci. Remote Sens., 50, 23–36, https://doi.org/10.1109/TGRS.2011.2160071, 2012.
- Kim, B.-H., Seo, K.-W., Eom, J., Chen, J., and Wilson, C.: Antarctic ice mass variations from 1979 to 2017 driven by anomalous precipitation accumulation, Scientific Reports, 10, 2045–2322, https://doi.org/10.1038/s41598-020-77403-5, 2020.
  - King, M. and Watson, C.: Antarctic Surface Mass Balance: Natural Variability, Noise, and Detecting New Trends, Geophys. Res. Lett., 47, https://doi.org/10.1029/2020GL087493, 2020.
  - Kuipers Munneke, P., Ligtenberg, S., Noël, B., Howat, I., Box, J., Mosley-Thompson, E., McConnell, J., Steffen, K., Harper, J., Das, S., and van den Broeke, M.: Elevation change of the Greenland Ice Sheet due to surface mass balance and firn processes, 1960–2014, The Cryosphere, 9, 2009–2025, https://doi.org/10.5194/tc-9-2009-2015, 2015.

1095

- Lenaerts, J., van den Broeke, M., Déry, S., van Meijgaard, E., van de Berg, W., Palm, S., and Sanz Rodrigo, J.: Modeling drifting snow in Antarctica with a regional climate model: 1. Methods and model evaluation, J. Geophys. Res., 117, https://doi.org/10.1029/2011JD016145, 2012.
- Lenaerts, J., van Meijgaard, E., van den Broeke, M., Ligtenberg, S., Horwath, M., and Isaksson, E.: Recent snowfall anomalies in Dronning Maud Land, East Antarctica, in a historical and future climate perspective, Geophys. Res. Lett., 40, 2684–2688, https://doi.org/10.1002/grl.50559, 2013.
  - Lenaerts, J., Medley, B., Broeke, M., and Wouters, B.: Observing and Modeling Ice Sheet Surface Mass Balance, Rev. Geophys., 57, https://doi.org/10.1029/2018RG000622, 2019.
- Ligtenberg, S., Helsen, M., and van den Broeke, M.: An improved semi-empirical model for the densification of Antarctic firn, The Cryosphere, 5, 809–819, https://doi.org/10.5194/tc-5-809-2011, 2011.
  - Ligtenberg, S., Horwath, M., van den Broeke, M., and Legrésy, B.: Quantifying the seasonal "breathing" of the Antarctic ice sheet, Geophys. Res. Lett., 39, L23 501, https://doi.org/10.1029/2012GL053628, 2012.
  - Lundin, J., Stevens, C., Arthern, R., Buizert, C., Orsi, A., Ligtenberg, S., Simonsen, S., Cummings, E., Essery, R., Leahy, W., Harris, P., Helsen, M., and Waddington, E.: Firn Model Intercomparison Experiment (FirnMICE), J. Glac., 63, 401–422, https://doi.org/10.1017/jog.2016.114, 2017.
  - Marsaglia, G., Tsang, W., and Wang, J.: Evaluating Kolmogorov's Distribution, Journal of Statistical Software, 8, 1–4, https://doi.org/10.18637/jss.v008.i18, 2003.
  - Massey, F.: The Kolmogorov-Smirnov Test for Goodness of Fit, Journal of the American Statistical Association, 46, 68–78, 1951.
- McMillan, M., Leeson, A., Shepherd, A., Briggs, K., Armitage, T., Hogg, A., Kuipers Munneke, P., van den Broeke, M., Noël, B., van de Berg, W., Ligtenberg, S., Horwath, M., Groh, A., Muir, A., and Gilbert, L.: A high-resolution record of Greenland mass balance, Geophys. Res. Lett., 43, 7002–7010, https://doi.org/10.1002/2016GL069666, 2016.
  - Medley, B. and Thomas, E.: Increased snowfall over the Antarctic Ice Sheet mitigated twentieth-century sea-level rise, Nature Clim. Change, 9, 34–39, https://doi.org/10.1038/s41558-018-0356-x, 2019.
- Medley, B., Neumann, T., Zwally, H., Smith, B., and Stevens, C.: Simulations of firn processes over the Greenland and Antarctic ice sheets: 1980-2021, The Cryosphere, 16, 3971–4011, https://doi.org/10.5194/tc-16-3971-2022, 2022a.
  - Medley, B., Neumann, T., Zwally, H., Smith, B., and Stevens, C.: NASA GSFC Firn Densification Model version 1.2.1 (GSFC-FDMv1.2.1) for the Greenland and Antarctic Ice Sheets: 1980-2021, https://doi.org/10.5281/zenodo.7054574, 2022b.
  - Mémin, A., Flament, T., Alizier, B., Watson, C., and Rémy, F.: Interannual variation of the Antarctic Ice Sheet from a combined analysis of satellite gravimetry and altimetry data, Earth Planet. Sci. Lett., 422, 150–156, https://doi.org/10.1016/j.epsl.2015.03.045, 2015.

1125 Miller, L.: Table of Percentage Points of Kolmogorov Statistics, Journal of the American Statistical Association, 51, 111–121, 1956.

1130

- Mohajerani, Y., Velicogna, I., and Rignot, E.: Mass Loss of Totten and Moscow University Glaciers, East Antarctica, Using Regionally Optimized GRACE Mascons, Geophys. Res. Lett., 45, 7010–7018, https://doi.org/10.1029/2018GL078173, 2018.
- Mottram, R., Hansen, N., Kittel, C., van Wessem, J., Agosta, C., Amory, C., Boberg, F., van de Berg, W., Fettweis, X., Gossart, A., van Lipzig, N., van Meijgaard, E., Orr, A., Phillips, T., Webster, S., Simonsen, S., and Souverijns, N.: What is the surface mass balance of Antarctica? An intercomparison of regional climate model estimates, The Cryosphere, 15, 3751–3784, https://doi.org/10.5194/tc-15-3751-2021, 2021.
- Mouginot, J., Rignot, E., and Scheuchl, B.: Sustained increase in ice discharge from the Amundsen Sea Embayment, West Antarctica, from 1973 to 2013, Geophys, Res. Lett., 41, 1576–1584, https://doi.org/10.1002/2013GL059069, 2014.
  - Nilsson, J., Gardner, A., Sandberg Sørensen, L., and Forsberg, R.: Improved retrieval of land ice topography from CryoSat-2 data and its impact for volume-change estimation of the Greenland Ice Sheet, The Cryosphere, 10, 2953–2969, https://doi.org/10.5194/tc-10-2953-2016, 2016.
  - Nilsson, J., Gardner, A., and Paolo, F.: MEaSUREs ITS\_LIVE Antarctic Grounded Ice Sheet Elevation Change, Version 1, https://doi.org/10.5067/L3LSVDZS15ZV, 2021.
  - Nilsson, J., Gardner, A., and Paolo, F.: Elevation change of the Antarctic Ice Sheet: 1985 to 2020, Earth System Science Data, 14, 3573–3598, https://doi.org/10.5194/essd-14-3573-2022, 2022.
- Noble, T., Rohling, E., Aitken, A., Bostock, H., Chase, Z., Gomez, N., Jong, L., King, M., Mackintosh, A., McCormack, F., McKay, R., Menviel, L., Phipps, S., Weber, M., Fogwill, C., Gayen, B., Golledge, N., Gwyther, D., Hogg, A., Martos, Y., Pena-Molino, B., Roberts, J., van de Flierdt, T., and Williams, T.: The sensitivity of the Antarctic Ice Sheet to a changing climate: Past, present and future, Rev. Geophys., 58, https://doi.org/10.1029/2019RG000663, e2019RG000663 2019RG000663, 2020.
- Noël, B., van de Berg, W., Machguth, H., Lhermitte, S., Howat, I., Fettweis, X., and van den Broeke, M.: A daily, 1 km resolution data set of downscaled Greenland ice sheet surface mass balance (1958–2015), The Cryosphere, 10, 2361–2377, https://doi.org/10.5194/tc-10-2361-2016, 2016.
  - Noël, B., van Wessem, J., Wouters, B., Trusel, L., Lhermitte, S., and van den Broeke, M.: Higher Antarctic ice sheet accumulation and surface melt rates revealed at 2 km resolution, Nature Communications, 14, https://doi.org/10.1038/s41467-023-43584-6, 2023.
- North, G., Bell, T., Cahalan, R., and Moeng, F.: Sampling errors in the estimation of empirical orthogonal functions, Month. Weath. Rev., 1150 110, 699–706, 1982.
  - Otosaka, I., Horwath, M., Mottram, R., and Nowicki, S.: Mass Balances of the Antarctic and Greenland Ice Sheets Monitored from Space, Surv. Geophys., https://doi.org/10.1007/s10712-023-09795-8, 2023a.
  - Otosaka, I., Shepherd, A., Ivins, E., Schlegel, N.-J., Amory, C., van den Broeke, M., Horwath, M., Joughin, I., King, M., Krinner, G., Nowicki, S., Payne, A., Rignot, E., Scambos, T., Simon, K., Smith, B., Sørensen, L., Velicogna, I., Whitehouse, P., A, G., Agosta, C., Ahlstrøm, A.,
- Blazquez, A., Colgan, W., Engdahl, M., Fettweis, X., Forsberg, R., Gallée, H., Gardner, A., Gilbert, L., Gourmelen, N., Groh, A., Gunter, B., Harig, C., Helm, V., Khan, S., Kittel, C., Konrad, H., Langen, P., Lecavalier, B., Liang, C.-C., Loomis, B., McMillan, M., Melini, D., Mernild, S., Mottram, R., Mouginot, J., Nilsson, J., Noël, B., Pattle, M., Peltier, W., Pie, N., Roca, M., Sasgen, I., Save, H., Seo, K.-W., Scheuchl, B., Schrama, E., Schröder, L., Simonsen, S., Slater, T., Spada, G., Sutterley, T., Vishwakarma, B., van Wessem, J., Wiese, D., van der Wal, W., and Wouters, B.: Mass balance of the Greenland and Antarctic ice sheets from 1992 to 2020, Earth System Science Data,
  15, 1597–1616, https://doi.org/10.5194/essd-15-1597-2023, 2023b.
  - Preisendorfer, R.: Principal Component Analysis in Meteorology and Oceanography, Elsevier Science Publishers B.V., Amsterdam, 1988.

- Rémy, F., Flament, T., Blarel, F., and Benveniste, J.: Radar altimetry measurements over antarctic ice sheet: A focus on antenna polarization and change in backscatter problems, Adv. Space Res., 50, 998–1006, https://doi.org/10.1016/j.asr.2012.04.003, 2012.
- Richter, A., Ekaykin, A., Willen, M., Lipenkov, V., Groh, A., Popov, S., Scheinert, M., Horwath, M., and Dietrich, R.: Surface Mass Balance Models Vs. Stake Observations: A Comparison in the Lake Vostok Region, Central East Antarctica, Front. Earth Sci., 9, 388, https://doi.org/10.3389/feart.2021.669977, 2021.
  - Rignot, E., Mouginot, J., and Scheuchl, B.: Ice Flow of the Antarctic Ice Sheet, Science, 333, 1427–1430, https://doi.org/10.1126/science.1208336, 2011a.
- Rignot, E., Mouginot, J., and Scheuchl, B.: Antarctic grounding line mapping from differential satellite radar interferometry, Geophys. Res. Lett., 38, https://doi.org/10.1029/2011GL047109, 2011b.
  - Rignot, E., Mouginot, J., Scheuchl, B., van den Broeke, M., Van Wessem, J., and Morlighem, M.: Four decades of Antarctic Ice Sheet mass balance from 1979–2017, Proc. Natl. Acad. Sci. USA, 116, 1095–1103, https://doi.org/10.1073/pnas.1812883116, 2019.
  - Roemer, S., Legrésy, B., Horwath, M., and Dietrich, R.: Refined analysis of radar altimetry data applied to the region of the subglacial Lake Vostok / Antarctica, Remote Sens. Environ., 106, 269–284, https://doi.org/10.1016/j.rse.2006.02.026, 2007.
- Sasgen, I., Dobslaw, H., Martinec, Z., and Thomas, M.: Satellite gravimetry observation of Antarctic snow accumulation related to ENSO, Earth Planet. Sci. Lett., 299, 352–358, https://doi.org/10.1016/j.epsl.2010.09.015, 2010.
  - Scambos, T., Frezzotti, M., Haran, T., Bohlander, J., Lenaerts, J., Van Den Broeke, M., Jezek, K., Long, D., Urbini, S., Farness, K., Neumann, T., Albert, M., and Winther, J.-G.: Extent of low-accumulation 'wind glaze' areas on the East Antarctic plateau: Implications for continental ice mass balance, J. Glac., 58, 633–647, https://doi.org/10.3189/2012JoG11J232, 2012.
- 1180 Schlegel, N.-J., Seroussi, H., Schodlok, M., Larour, E., Boening, C., Limonadi, D., Watkins, M., Morlighem, M., and van den Broeke, M.: Exploration of Antarctic Ice Sheet 100-year contribution to sea level rise and associated model uncertainties using the ISSM framework, The Cryosphere, 12, 3511–3534, https://doi.org/10.5194/tc-12-3511-2018, 2018.
  - Schröder, L., Horwath, M., Dietrich, R., Helm, V., van den Broeke, M., and Ligtenberg, S.: Four decades of Antarctic surface elevation changes from multi-mission satellite altimetry, The Cryosphere, 13, 427–449, https://doi.org/10.5194/tc-13-427-2019, 2019a.
- Schröder, L., Horwath, M., Dietrich, R., Helm, V., van den Broeke, M., and Ligtenberg, S.: Gridded surface elevation changes from multi-mission satellite altimetry 1978-2017, https://doi.org/10.1594/PANGAEA.897390, supplement to: Schröder, L. et al. (2019): Four decades of Antarctic surface elevation changes from multi-mission satellite altimetry. The Cryosphere, 13(2), 427-449, https://doi.org/10.5194/tc-13-427-2019, 2019b.
- Shepherd, A., Ivins, E., Rignot, E., Smith, B., van den Broeke, M., Velicogna, I., Whitehouse, P., Briggs, K., Joughin, I., Krinner, G., Nowicki,
  S., Payne, T., Scambos, T., Schlegel, N., A, G., Agosta, C., Ahlstrøm, A., Babonis, G., Barletta, V., Blazquez, A., Bonin, J., Csatho, B.,
  Cullather, R., Felikson, D., Fettweis, X., Forsberg, R., Gallee, H., Gardner, A., Gilbert, L., Groh, A., Gunter, B., Hanna, E., Harig, C.,
  Helm, V., Horvath, A., Horwath, M., Khan, S., Kjeldsen, K., Konrad, H., Langen, P., Lecavalier, B., Loomis, B., Luthcke, S., McMillan,
  M., Melini, D., Mernild, S., Mohajerani, Y., Moore, P., Mouginot, J., Moyano, G., Muir, A., Nagler, T., Nield, G., Nilsson, J., Noel, B.,
  Otosaka, I., Pattle, M., Peltier, W., Pie, N., Rietbroek, R., Rott, H., Sandberg-Sørensen, L., Sasgen, I., Save, H., Scheuchl, B., Schrama,
- E., Schröder, L., Seo, K.-W., Simonsen, S., Slater, T., Spada, G., Sutterley, T., Talpe, M., Tarasov, L., van de Berg, W., van der Wal, W., van Wessem, J., Vishwakarma, B., Wiese, D., and Wouters, B.: Mass balance of the Antarctic Ice Sheet from 1992 to 2017, Nature, 558, 219–222, https://doi.org/10.1038/s41586-018-0179-y, 2018.
  - Shepherd, A., Gilbert, L., Muir, A., Konrad, H., McMillan, M., Slater, T., Briggs, K., Sundal, A., Hogg, A., and Engdahl, M.: Trends in Antarctic Ice Sheet Elevation and Mass, Geophys. Res. Lett., 46, 8174–8183, https://doi.org/10.1029/2019GL082182, 2019.

- 1200 Shi, T., Fukuda, Y., Doi, K., and Okuno, J.: Extraction of GRACE/GRACE-FO observed mass change patterns across Antarctica via independent component analysis (ICA), Geophys. J. Int., 229, 1914–1926, https://doi.org/10.1093/gji/ggac033, 2022.
  - Slater, T., Shepherd, A., McMillan, M., Muir, A., Gilbert, L., Hogg, A., Konrad, H., and Parrinello, T.: A new digital elevation model of Antarctica derived from CryoSat-2 altimetry, The Cryosphere, 12, 1551–1562, https://doi.org/10.5194/tc-12-1551-2018, 2018.
- Sørensen, L., Simonsen, S., Nielsen, K., Lucas-Picher, P., Spada, G., Adalgeirsdottir, G., Forsberg, R., and Hvidberg, C.: Mass balance of the

  Greenland ice sheet (2003–2008) from ICESat data the impact of interpolation, sampling and firn density, The Cryosphere, 5, 173–186, https://doi.org/10.5194/tc-5-173-2011, 2011.
  - Stevens, C., Verjans, V., Lundin, J., Kahle, E., Horlings, A., Horlings, B., and Waddington, E.: The Community Firn Model (CFM) v1.0, Geoscientific Model Development, 13, 4355–4377, https://doi.org/10.5194/gmd-13-4355-2020, 2020.
- Stevens, M., Vo, H., Emmakahle, and Jboat: UWGlaciology/CommunityFirnModel: Version 1.1.6 (v1.1.6), https://doi.org/10.5281/zenodo.5719748, supplement to: https://github.com/UWGlaciology/CommunityFirnModel/tree/v1.1.6, 2021.
  - Su, X., Shum, C., Guo, J., Howat, I., Kuo, C., Jezek, K., Duan, J., and Yi, Y.: High-Resolution Interannual Mass Anomalies of the Antarctic Ice Sheet by Combining GRACE Gravimetry and ENVISAT Altimetry, IEEE Trans. Geosci. Remote Sens., 56, 539–546, https://doi.org/10.1109/TGRS.2017.2751070, 2018.
- Thomas, E., van Wessem, J., Roberts, J., Isaksson, E., Schlosser, E., Fudge, T., Vallelonga, P., Medley, B., Lenaerts, J., Bertler, N., van den

  Broeke, M., Dixon, D., Frezzotti, M., Stenni, B., Curran, M., and Ekaykin, A.: Regional Antarctic snow accumulation over the past 1000 years, Climate of the Past, 13, 1491–1513, https://doi.org/10.5194/cp-13-1491-2017, 2017.
  - Tian, B., Lee, H., Waliser, D. E., Ferraro, R., Kim, J., Case, J., Iguchi, T., Kemp, E., Wu, D., Putman, W., and Wang, W.: Development of a Model Performance Metric and Its Application to Assess Summer Precipitation over the U.S. Great Plains in Downscaled Climate Simulations, Journal of Hydrometeorology, 18, 2781–2799, https://doi.org/10.1175/JHM-D-17-0045.1, 2017.
- van Dalum, C. and van de Berg, W.: First results of RACMO2.4: A new model version with updated surface and atmospheric processes, in: EGU General Assembly 2023, Vienna, Austria, 24–28 Apr 2023, EGU23-13907, https://doi.org/10.5194/egusphere-egu23-13907, 2023. van den Broeke, M.: Depth and density of the Antarctic firn layer, Arctic, Antarctic, and Alpine Research, 40, 432–438, 2008.

- van den Broeke, M., Enderlin, E., Howat, I., Kuipers Munneke, P., Noël, B., van de Berg, W., van Meijgaard, E., and Wouters, B.: On the recent contribution of the Greenland ice sheet to sea level change, The Cryosphere, 10, 1933–1946, https://doi.org/10.5194/tc-10-1933-2016, 2016.
- van Wessem, J., Reijmer, C., Morlighem, M., Mouginot, J., Rignot, E., Medley, B., Joughin, I., Wouters, B., Depoorter, M., Bamber, J., Lenaerts, J., De Van Berg, W., Van Den Broeke, M., and Van Meijgaard, E.: Improved representation of East Antarctic surface mass balance in a regional atmospheric climate model, J. Glac., 60, 761–770, https://doi.org/10.3189/2014JoG14J051, 2014.
- van Wessem, J., van de Berg, W., Noël, B., van Meijgaard, E., Amory, C., Birnbaum, G., Jakobs, C., Krüger, K., Lenaerts, J., Lhermitte,

  S., Ligtenberg, S., Medley, B., Reijmer, C., van Tricht, K., Trusel, L., van Ulft, L., Wouters, B., Wuite, J., and van den Broeke, M.:

  Modelling the climate and surface mass balance of polar ice sheets using RACMO2 Part 2: Antarctica (1979–2016), The Cryosphere,

  12, 1479–1498, https://doi.org/10.5194/tc-12-1479-2018, 2018.
  - Veldhuijsen, S., van de Berg, W., Brils, M., Kuipers Munneke, P., and van den Broeke, M.: Characteristics of the 1979–2020 Antarctic firn layer simulated with IMAU-FDM v1.2A, The Cryosphere, 17, 1675–1696, https://doi.org/10.5194/tc-17-1675-2023, 2023.
- 1235 Velicogna, I. and Wahr, J.: Time-variable gravity observations of ice sheet mass balance: Precision and limitations of the GRACE satellite data, Geophys. Res. Lett., 40, 3055–3063, https://doi.org/10.1002/grl.50527, 2013.

- Velicogna, I., Mohajerani, Y., A, G., Landerer, F., Mouginot, J., Noel, B., Rignot, E., Sutterley, T., van den Broeke, M., van Wessem, J., and Wiese, D.: Continuity of ice sheet mass loss in Greenland and Antarctica from the GRACE and GRACE Follow-On missions, Geophys. Res. Lett., 47, https://doi.org/10.1029/2020GL087291, e2020GL087291 2020GL087291, 2020.
- 1240 Verjans, V., Leeson, A. M., Stevens, C., van Wessem, J. M., van de Berg, W., van den Broeke, M., Kittel, C., Amory, C., Fettweis, X., Hansen, N., Boberg, F., and Mottram, R.: Uncertainty in East Antarctic Firn Thickness Constrained Using a Model Ensemble Approach, Geophys. Res. Lett., 48, e2020GL092 060, https://doi.org/10.1029/2020GL092060, e2020GL092060 2020GL092060, 2021.

1245

1255

- Wang, Y., Huai, B., Thomas, E., van den Broeke, M., van Wessem, J., and Schlosser, E.: A New 200-Year Spatial Reconstruction of West Antarctic Surface Mass Balance, Journal of Geophysical Research: Atmospheres, 124, 5282–5295, https://doi.org/10.1029/2018JD029601, 2019.
- Whitehouse, P., Gomez, N., King, M., and Wiens, D.: Solid Earth change and the evolution of the Antarctic Ice Sheet, Nature Communications, 10, https://doi.org/10.1038/s41467-018-08068-y, 2019.
- Willen, M., Horwath, M., Schröder, L., Groh, A., Ligtenberg, S., Kuipers Munneke, P., and van den Broeke, M.: Sensitivity of inverse glacial isostatic adjustment estimates over Antarctica, The Cryosphere, 14, 349–366, https://doi.org/10.5194/tc-14-349-2020, 2020.
- Willen, M., Broerse, T., Groh, A., Wouters, B., Kuipers Munneke, P., Horwath, M., van den Broeke, M., and Schröder, L.: Separating Long-Term and Short-Term Mass Changes of Antarctic Ice Drainage Basins: A Coupled State Space Analysis of Satellite Observations and Model Products, J. Geophys. Res. Earth Surf., 126, e2020JF005 966, https://doi.org/10.1029/2020JF005966, 2021.
  - Willen, M., Horwath, M., Groh, A., Helm, V., Uebbing, B., and Kusche, J.: Feasibility of a global inversion for spatially resolved glacial isostatic adjustment and ice sheet mass changes proven in simulation experiments, J. Geod., 96, 1–21, https://doi.org/10.1007/s00190-022-01651-8, 2022.
  - Williams, S., Moore, P., King, M., and Whitehouse, P.: Revisiting GRACE Antarctic ice mass trends and accelerations considering autocorrelation, Earth Planet. Sci. Lett., 385, 12–21, https://doi.org/10.1016/j.epsl.2013.10.016, 2014.
  - Wingham, D., Rapley, C., and Griffiths, H.: New techniques in satellite altimeter tracking systems, in: ESA Proceedings of the 1986 International Geoscience and Remote Sensing Symposium(IGARSS'86) on Remote Sensing: Today's Solutions for Tomorrow's Information Needs., vol. 3, pp. 1339–1344, 1986.
  - Wingham, D., Ridout, A., Scharroo, R., Arthern, R., and Shum, C.: Antarctic Elevation Change from 1992 to 1996, Science, 282, 456–458, https://doi.org/10.1126/science.282.5388.456, 1998.
  - Wouters, B., Bamber, J., van den Broeke, M., Lenaerts, J., and Sasgen, I.: Limits in detecting acceleration of ice sheet mass loss due to climate variability, Nature Geosci., 6, 613–616, https://doi.org/10.1038/ngeo1874, 2013.
- Zhang, B., Yao, Y., Liu, L., and Yang, Y.: Interannual ice mass variations over the Antarctic ice sheet from 2003 to 2017 were linked to El Niño-Southern Oscillation, Earth Planet. Sci. Lett., 560, 116796, https://doi.org/10.1016/j.epsl.2021.116796, 2021.
  - Zwally, H., Giovinetto, M., Li, J., Cornejo, H., Beckley, M., Brenner, A., Saba, J., and Yi, D.: Mass changes of the Greenland and Antarctic ice sheets and shelves and contributions to sea-level rise: 1992-2002, J. Glac., 51, 509–527, https://doi.org/10.3189/172756505781829007, 2005.