

1 Simulating the Holocene deglaciation across a marine terminating portion of southwestern 2 Greenland in response to marine and atmospheric forcings

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14 **Abstract**

15
16 Numerical simulations of the Greenland Ice Sheet (GrIS) over geologic timescales can greatly
17 improve our knowledge of the critical factors driving GrIS demise during climatically warm
18 periods, which has clear relevance for better predicting GrIS behavior over the upcoming
19 centuries. To assess the fidelity of these modeling efforts, however, observational constraints of
20 past ice-sheet change are needed. Across southwestern Greenland, geologic records detail
21 Holocene ice retreat across both terrestrial-based and marine terminating environments, providing
22 an ideal opportunity to rigorously benchmark model simulations against geologic reconstructions
23 of ice-sheet change. Here, we present regional ice sheet modeling results using the Ice-sheet and
24 Sea-level System Model (ISSM) of Holocene ice sheet history across an extensive fjord region in
25 southwestern Greenland covering the landscape around the Kangiata Nunaata Sermia (KNS)
26 glacier and extending outward along the 200 km Nuup Kangerula (Godthåbsfjord). Our
27 simulations, forced by reconstructions of Holocene climate and recently implemented calving
28 laws, assess the sensitivity of ice retreat across the KNS region to atmospheric and oceanic
29 forcing. Our simulations reveal that the geologically reconstructed ice retreat across the terrestrial
30 landscape in the study area was likely driven by fluctuations in surface mass balance in response
31 to Early Holocene warming – and likely not influenced significantly by the response of adjacent
32 outlet glaciers to calving and ocean-induced melting. The impact of ice calving within fjords,
33 however, plays a significant role by enhancing ice discharge at the terminus, leading to interior
34 thinning up to the ice divide that is consistent with reconstructed magnitudes of Early Holocene
35 ice thinning. Our results, benchmarked against geologic constraints of past ice margin change,
36 suggest that while calving did not strongly influence Holocene ice margin migration across
37 terrestrial portions of the KNS forefield, it strongly impacted regional mass loss. While these
38 results imply that the implementation and resolution of ice calving in paleo ice flow models is
39 important towards making more robust estimations of past ice mass change, they also illustrate the
40 importance these processes have on contemporary and future long term ice mass change across
41 similar fjord-dominated regions of the GrIS.

42 43 **1. Introduction**

45 Over the past few decades, the Greenland Ice Sheet (GrIS) has experienced accelerating ice mass
46 loss driven by increases in surface melt, runoff, and dynamic ice loss at marine terminating margins
47 (IMBIE, 2019). While projected mass loss from the GrIS is expected to be driven increasingly by
48 its surface mass balance (SMB; Enderlin et al., 2014; Vizcaino et al., 2015; Goelzer et al., 2020)
49 and attendant meltwater runoff (Fettweis et al., 2008; Lenaerts et al., 2018), considerable
50 uncertainty exists regarding how oceanic forcing will influence GrIS mass loss, particularly
51 through ice calving processes (Goelzer et al., 2020; Choi et al., 2021). The satellite-based
52 observational record of GrIS change only spans a few decades making it difficult to identify and
53 disentangle the key drivers of GrIS mass change, and to understand over which timescales they
54 operate. Fortunately, geologic records detailing the retreat history of the GrIS provide an
55 important metric for evaluating numerical ice sheet models and help pinpoint the contributions of
56 various driving mechanisms to GrIS change. When combined, numerical ice sheet models and
57 geologic reconstructions can provide key insights into GrIS behavior in a warming climate across
58 centennial to millennial timescales.

59
60 The current interglacial, the Holocene (the last 11.7 ka), is characterized by prolonged warmth
61 with proxy records suggesting that temperatures during the early to Middle Holocene were 3 ± 1 °C
62 warmer than the pre-industrial period (Briner et al., 2016; Lecavalier et al., 2017), which drove
63 widespread retreat of the GrIS margin at a rate of ice mass loss exceeding 20th century values
64 (1900-2000 CE Young and Briner, 2015; Briner et al., 2020). Across southwestern Greenland, a
65 detailed geologic record of Holocene ice-margin retreat encompassing both terrestrial and marine
66 terminating environments exists, providing an ideal testbed for ice sheet models to test the
67 sensitivity of past ice margin migration to atmospheric and marine forcings (Larsen et al., 2014;
68 Lesnek et al., 2020; Young et al., 2020; Young et al., 2021). Where land-based ice existed, well-
69 dated moraine sequences constrain ~120 km of ice retreat from the present-day coastline to just
70 outboard of the present-day ice margin (Lesnek et al., 2020; Young et al., 2020), and have been
71 shown by ice sheet models to be driven by negative SMB in response to Early Holocene warming
72 (Cuzzone et al., 2019; Downs et al., 2020; Briner et al., 2020).

73
74 Unlike the land-based portions of Southwest Greenland however, across the marine based region
75 covering the forefield around Kangiata Nunaata Sermia (KNS; Figure 1), it remains unknown what
76 drove this rapid ice margin retreat during the Early Holocene (Young et al., 2021). While links
77 between atmospheric warming and runoff induced terminus retreat have been implicated as
78 reasons for the most recent historical retreat across the KNS region (Lea et al., 2014a,b), the longer
79 term triggers of rapid Holocene ice retreat are not constrained by the geologic data alone. Because
80 of the well dated chronology detailing Holocene ice retreat across this region however, ice sheet
81 models are well poised to address questions surrounding the scales of influence atmospheric and
82 oceanic forcings play on long term ice margin and mass change. However, as many paleo ice flow
83 models employ model grids that are relatively coarse (10 km or greater), ice margin migration and
84 ultimately ice discharge through fjord systems may be poorly simulated or not captured as many
85 models cannot resolve the complex and narrow fjord geometries found across the GrIS (Cuzzone
86 et al., 2019).

87
88 Building on recent advances in calving front dynamics in the Ice Sheet and Sea-level System
89 Model (ISSM; Larour et al., 2012), we use a high-resolution regional ice sheet model to investigate
90 the Holocene ice retreat across the KNS forefield. Our simulations build on prior ice modeling

91 efforts across southwestern Greenland that were driven by novel reconstructions of past climate
92 (Badgeley et al., 2020; Briner et al., 2020). Where our past ice flow modeling efforts excluded ice
93 ocean-interactions (Briner et al., 2020), our simulations presented here take advantage of recent
94 implementation of physically based calving schemes in ISSM to specifically address how
95 Holocene ice retreat across the KNS forefield was influenced by marine and atmospheric forcing's.
96 Moreover, this work provides a foundation for future experiments using ISSM to simulate the
97 influence of ice-ocean interactions on the Holocene variability of the broader GrIS.
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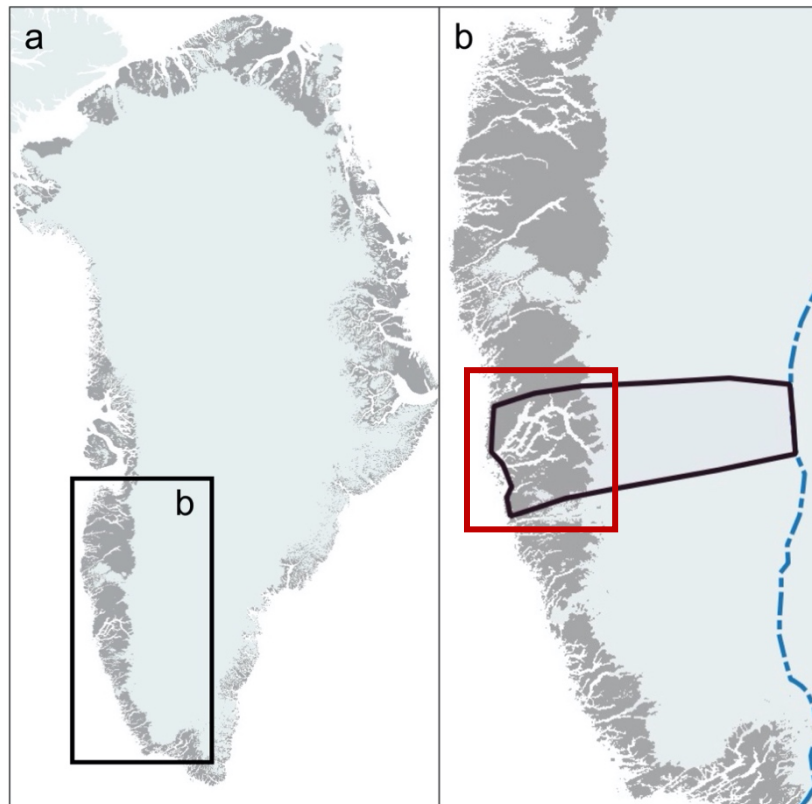


Figure 1. a.) The Greenland Ice Sheet. Highlighted is southwestern Greenland, where the ice model domain resides. b.) Southwestern Greenland. The ice model domain is outlined (bold black line), extending between the present-day coastline and ice divide (dashed blue line; Rignot and Mouginot, 2012). The red box corresponds to the area in figures 5, 6 and 8.

99

100 **2. Holocene Ice Retreat across the KNS Forefield**

101

102 Decades of radiocarbon dating and, more recently, cosmogenic ^{10}Be dating, track the retreat of the
103 GrIS in the KNS region through the Holocene (Weidick et al., 2012 and references therein; Larsen
104 et al., 2014; Young et al., 2021). Minimum-limiting radiocarbon ages from the outer coast near
105 Nuuk range from ca. 11.2 to 10.6 ka BP., which is mimicked by ^{10}Be ages of ca. 10.7 and 10.4 ka
106 BP (Figure 2). Between the outer coast and the modern GrIS margin at KNS are numerous
107 radiocarbon and ^{10}Be ages that are largely indistinguishable and require rapid deglaciation of the
108 region spanning about a millennium (Weidick et al., 2012; Larsen et al., 2014; Young et al., 2021).
109 Perhaps most relevant here are ^{10}Be ages in the immediate KNS region from just beyond the
110 historical ice limit that suggest KNS had retreated within or near its current position by ca. 10.3 ka

111 (Young et al., 2021). Radiocarbon ages from raised marine deposits, which require ice-free
 112 conditions, adjacent to the main KNS fjord appear slightly younger than regional ^{10}Be ages. These
 113 radiocarbon ages, however, are minimum-limiting ages and an upfjord radiocarbon age of ca. 10.2
 114 ka from a bivalve reworked by a KNS readvance requires that the main fjord deglaciated prior to
 115 ca. 10.2 ka (Figure 2). Collectively, the radiocarbon and ^{10}Be ages suggest rapid and synchronous
 116 deglaciation of both the landscape and fjord systems between the outer coast near Nuuk and the
 117 modern margin at KNS. Lastly, ^{10}Be ages from slightly beyond the historical limit to the north and
 118 south of KNS are slightly younger suggesting that these ice margins may have lagged behind ice
 119 retreat in the immediate KNS region (Figure 2).
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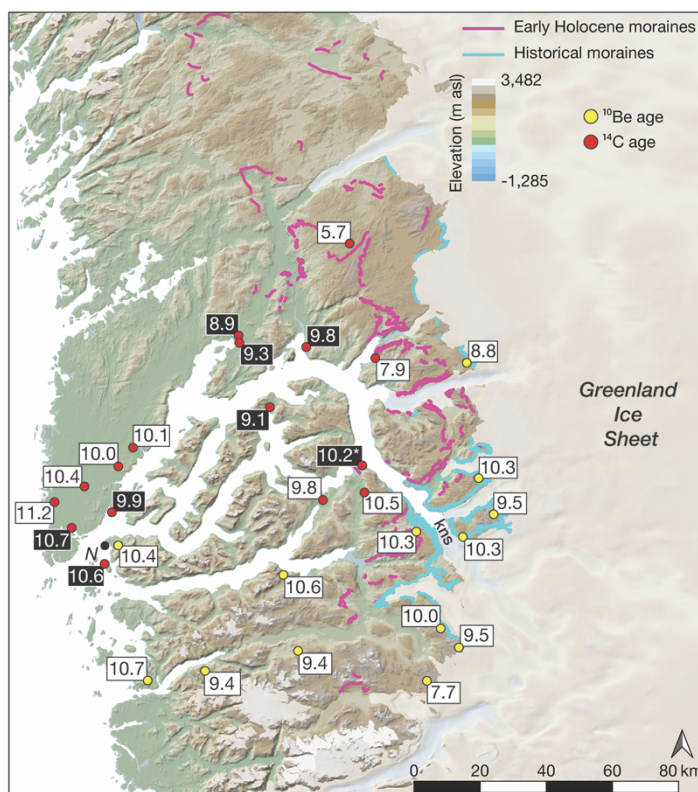


Figure 2. KNS region with geological constraints that track GrIS retreat in the Early Holocene. Radiocarbon ages (red circles) and ^{10}Be ages (yellow circles) are from Weidick et al. (2012), Larsen et al. (2014), and Young et al. (2021). For figure clarity, we only show the mean deglaciation age at each site (see Young et al., 2021 for full site descriptions). Radiocarbon and ^{10}Be across the immediate KNS region are similar and reveal that deglaciation of the coast occurred ca. 11.2-10.7 ka and KNS had retreated near or within its modern extent by ca. 10.3 ka. Radiocarbon ages in white text and black background are from marine deposits and constrain the timing of retreat within the main fjord. Figure has been modified from Young et al. (2021).

121
 122 **3. Model description and setup**
 123
 124 **3.1 Ice Sheet Model**
 125

126 We rely on ISSM, a thermomechanical finite-element ice sheet model, to simulate the Holocene
 127 ice history across the KNS forefield, and follow similar published model setups (Cuzzone et al.,
 128 2019; Briner et al., 2020). The higher-order approximation of Blatter (1995) and Pattyn (2003) is
 129 used to solve the momentum balance equations. Our model domain centers on the KNS and
 130 Godthåbsfjord forefield, extending from the present-day coastline, where geologic observations
 131 show ice resided at the end of the Younger Dryas (Larsen et al., 2014; Lesnek et al., 2020) to the
 132 present-day ice divide (Figure 1b; Rignot and Mouginot, 2012). The northern and southern
 133 boundaries of our model domain are chosen to represent regions of minimal north-south across
 134 boundary flow based on Holocene ice sheet simulations of southwestern Greenland (Briner et al.,
 135 2020). Anisotropic mesh adaptation is used to create a non-uniform model mesh that varies based
 136 upon gradients in bedrock topography from BedMachine v3 (Morlighem et al., 2017). Because
 137 fjord width across our domain is often <5 km and high-resolution grids are necessary for capturing
 138 grounding line dynamics (1 km; Seroussi and Morlighem, 2018), the horizontal mesh resolution
 139 varies from 1 km in fjords and areas of high bedrock relief to 15 km where the bedrock relief is
 140 low (Figure 3).

141
 142 To capture the thermal evolution of the ice, our model uses an enthalpy formulation (Aschwenden
 143 et al., 2012) that captures both temperate and cold ice. We impose transient air temperatures at
 144 the surface and a constant but spatially varying geothermal heat flux at the base (Shapiro and
 145 Ritzwoller, 2004) and our model contains only five vertical layers in order to reduce computational
 146 load (Cuzzone et al., 2018; Cuzzone et al., 2019). In order to capture sharp thermal gradients near
 147 the base and simulate the vertical distribution of temperature within the ice, we use quadratic finite
 148 elements (P1xP2) along the z-axis for the vertical interpolation following Cuzzone et al. (2018).
 149 This methodology has been successfully applied to simulate the transient behavior of the GrIS
 150 across geologic timescales and the contemporary period (Cuzzone et al., 2019; Briner et al., 2020;
 151 Smith-Johnson et al., 2020).

152
 153 We use a linear friction law and, similar to Briner et al. (2020), we construct a spatially varying
 154 basal friction coefficient (k) under areas covered by the present-day ice sheet using inverse
 155 methods (Morlighem et al., 2010; Larour et al., 2012) that satisfies the best match between
 156 modeled and satellite-derived surface velocities (Rignot and Mouginot, 2012):

$$157 \tau_b = -k^2 N v_b \tag{1}$$

158
 159 where τ_b represents the basal stress, N represents the effective pressure, and v_b is the magnitude
 160 of the basal velocity. For contemporary ice-free areas, a spatially varying basal friction coefficient
 161 is constructed to be proportional to bedrock elevation following Åkesson et al., 2018:

$$162 k = 100 \times \frac{\min[\max(0, z_b + 800), z_b]}{\max(z_b)} \tag{2}$$

163
 164 where z_b is the height of the bedrock with respect to sea level. For these parametrizations, the
 165 friction coefficient is low within fjords and is larger over areas of high topographic relief. This
 166 basal friction coefficient is allowed to vary through time based upon changes in the simulated basal
 167 temperature following Cuzzone et al. (2019). As simulated basal ice temperatures decrease with
 168 respect to present day, the friction coefficient will increase, and therefore sliding will decrease.
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171 The opposite occurs when simulated basal temperatures are warm relative to present day. Lastly,
 172 the ice rheology parameter B is temperature-dependent, following rate factors in Cuffey and
 173 Paterson (2010), and is initialized by solving for a present day thermal steady state and allowed to
 174 vary during transient simulations (Cuzzone et al., 2018; 2019).

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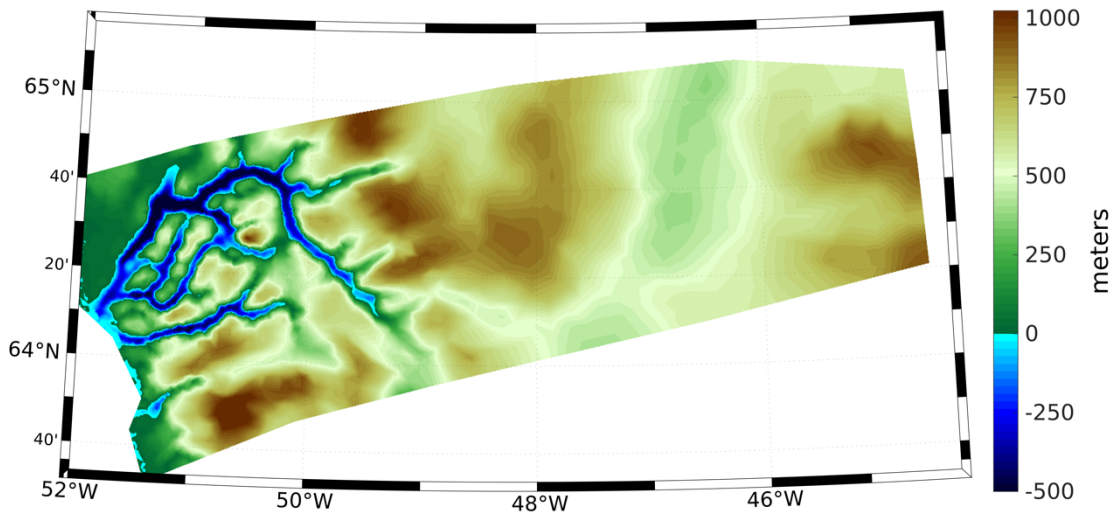


Figure 3. Bedrock topography for the model domain. Blue colors indicate areas that are below present-day sea level.

181
 182 **3.2 Ice Front Migration and Calving**

183
 184 We use the level-set method to track the motion of the ice front (Bondzio et al., 2016). The velocity
 185 of the moving ice front is calculated as:

186
 187
$$v_f = v - (c + \dot{M}) n \tag{3}$$

188
 189 where v_f is the ice velocity vector, v is the ice velocity vector at the ice front, c is the calving rate,
 190 \dot{M} is the melting rate of the calving front, and n is the unit normal vector pointing horizontally
 191 outward from the calving front. For these simulations, we assume that the melting rate at the
 192 calving front is negligible compared to the calving rate.

193
 194 To simulate calving, we rely on the physically-based Von Mises stress calving (Morlighem et al.,
 195 2016), whereby the calving rate is related to tensile stresses within the ice:

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 197
$$c = \|v\| \frac{\tilde{\sigma}}{\sigma_{max}} \tag{4}$$

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where $\tilde{\sigma}$ is the von Mises tensile strength, $\|v\|$ is the magnitude of the horizontal ice velocity, and σ_{max} is the maximum stress threshold, which has separate values for grounded and floating ice. Under this formulation, the ice front will remain stable when $\tilde{\sigma} = \sigma_{max}$, will retreat when $\tilde{\sigma} > \sigma_{max}$, and will advance when $\tilde{\sigma} < \sigma_{max}$. Tensile strength measurements of ice show a range of possible σ_{max} , ranging between 150 kPa to 3100 kPa (Petrovic 2003). For this study we choose $\sigma_{max} = 600$ kPa for grounded ice and 200 kPa for floating ice, which is within the ranges used by recent studies across Greenland (Bondzio et al., 2016; Morlighem et al., 2016; Choi et al., 2020).

3.3 Climate and Surface Mass Balance Reconstruction

We rely on a novel gridded paleoclimate reanalysis product that reconstructs the necessary climate variables of temperature and precipitation needed to calculate the surface mass balance history through the Holocene (Badgeley et al., 2020). Temperature was derived from oxygen-isotope records from eight ice cores, and five ice core accumulation records were used to reconstruct precipitation. This reanalysis relies on a data assimilation framework that combines the information from ice core proxies with climate-model simulations of the last deglaciation (Liu et al., 2009; He et al., 2013) to create a spatially complete (e.g., GrIS wide) and temporally consistent reconstruction of past temperature and precipitation. This reconstruction agrees well with independent proxies and previously published paleoclimate reconstructions (Badgeley et al. (2020)). For new simulations presented here, we chose two end members of reconstructed precipitation and temperature from Badgeley et al. (2020). The high temperature reconstruction was chosen, which has a greater magnitude of Early Holocene warming, and the low temperature scenario, which has a more muted Early Holocene warming (Figure 4a). Additionally, we choose the high and low precipitation scenarios (Figure 4b), which differ in the magnitude and timing of peak Holocene precipitation. These reconstructions span a plausible range of temperature and precipitation scenarios as discussed in Badgeley et al. (2020).

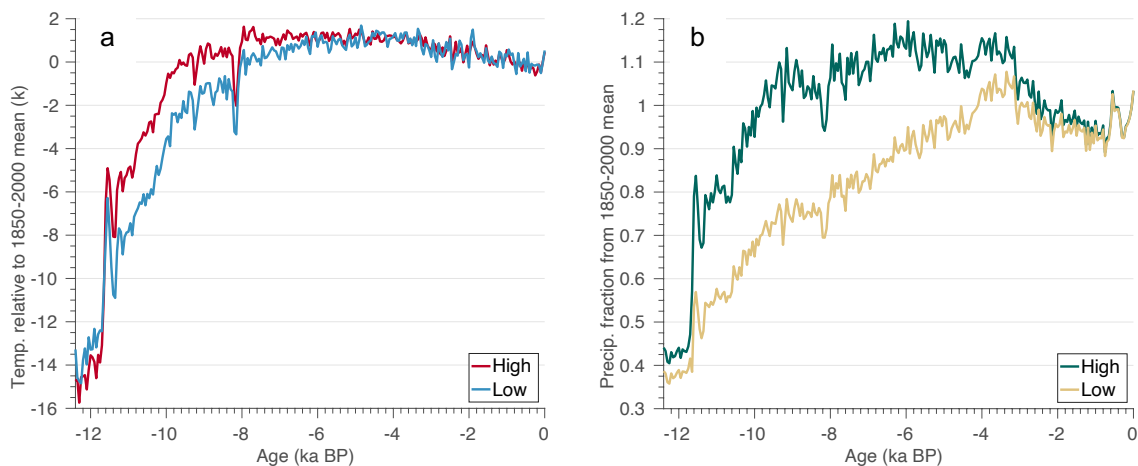


Figure 4. a) Area-averaged (over model domain) mean annual temperature anomaly (k) relative to the 1850-2000 mean for the High and Low temperature reconstructions from Badgeley et al. (2020). b) Area-averaged (over model domain) mean annual precipitation as a fraction from the 1850-2000 mean for the High and Low reconstructions from Badgeley et al. (2020).

226 The simulations discussed below use a combination of these forcings to address the possible role
 227 of varying climatic conditions. Following prior work (Cuzzone et al., 2019; Briner et al., 2020),
 228 we compute the surface mass balance over the Holocene using a positive degree day (PDD) method
 229 (Tarasov and Peltier, 1999; Le Morzadec et al., 2015). For this scheme, snow melts first at 4.3
 230 $\text{mm}^{-1}\text{C}^{-1}\text{day}^{-1}$ and the remaining positive degree days are used to melt bare ice at 8.3 $\text{mm}^{-1}\text{C}^{-1}\text{day}^{-1}$.
 231 A lapse rate of 6 $^{\circ}\text{C}/\text{km}$ is used to adjust the temperature of the climate forcings to ice-surface
 232 elevation, while an allowance for the formation of superimposed ice is permitted following
 233 Janssens and Huybrechts (2000).

234

235 **3.4 Experimental Setup**

236

237 For the reanalysis discussed in section 3.3, temperature is expressed as anomalies from the AD
 238 1850-2000 mean, and precipitation is expressed as a fraction of the AD 1850-2000 mean (Figure
 239 4). Following Briner et al. (2020), we apply these anomalies onto the 1850-2000 monthly mean
 240 climatology of temperature and precipitation from Box et al. (2013) to produce the necessary
 241 Holocene temperature and precipitation forcings:

242

$$243 T_t = \bar{T}_{(1850-2000)} + \Delta T_t \quad (5)$$

244

$$245 P_t = \bar{P}_{(1850-2000)} \times \Delta P_t \quad (6)$$

246

247 where $\bar{T}_{(1850-2000)}$ and $\bar{P}_{(1850-2000)}$ are the monthly mean temperature and precipitation over AD
 248 1850-2000 from Box et al. (2013) and ΔT_t and ΔP_t are the monthly anomalies from Badgeley et al.
 249 (2020). We perform four transient model simulations using four combinations of possible climate
 250 scenarios shown in Table 1. For each climate scenario, we run two simulations. First, simulations
 251 are performed where the calving parameterization is turned on (denoted as ‘Calving On’). Second,
 252 simulations are performed where the calving parameterization is turned off (denoted as ‘Calving
 253 Off’). For these simulations, we apply a temporally constant melting rate under floating ice of 40
 254 m/yr , which is consistent with contemporary melt rates derived near the grounding line of floating
 255 ice shelves across the GrIS (Wilson et al., 2017). We also perform additional simulations discussed
 256 further in section 4.4 to assess sensitivity to the calving maximum stress thresholds and ocean-
 257 induced melt-rates.

258

259 We initialize our regional ice-sheet model using present-day ice-surface elevation from the
 260 Greenland Ice Mapping Project digital elevation model (Howat et al., 2014). A constant climate
 261 from 12,400 years ago is then applied for each experiment, allowing our model to reach
 262 equilibrium in ice volume and basal temperature, which takes 20,000 years. Since our simulations
 263 are regional in scale, we use boundary conditions of temperature, ice velocity, and thickness from
 264 a recent ice sheet simulation of West-Southwest Greenland (Briner et al., 2020) and impose these
 265 as Dirichlet boundary conditions at the southern, northern, and ice-divide boundaries. The eastern
 266 boundary of our model domain extends outward to the present-day ice divide (Rignot and
 267 Mouginot, 2012), with the northern and southern boundary of our model domain extending to
 268 cover the KNS forefield. While the catchment for KNS may have changed during the Holocene
 269 and thus may have impacted ice flux into our domain, those changes are not constrained.
 270 Therefore, since we use consistent boundary conditions across our experiments, we consider that
 271 our results are primarily influenced by the surface climate and oceanic boundary conditions applied

272 and not influenced by model domain extent. These boundary conditions are forced transiently
 273 throughout the Holocene simulations and use similar model setups and climate forcings as
 274 discussed here. Each model is then run transiently through time from 12,400 years ago to AD
 275 1850 using the climatologies discussed above, and then from 1850 to 2013 we use monthly
 276 temperature and precipitation fields from Box et al. (2013). We use an adaptive timestep, which
 277 varies between 0.02 and 0.1 years, depending on the Courant–Friedrichs–Lewy criterion (Courant
 278 et al., 1928). Discussed further in Section 5.3, we do not include glacial isostatic adjustment (GIA)
 279 in these simulations. Although GIA can influence the underlying bedrock topography and
 280 ultimately surface mass balance gradients and grounding line stability, changes during the
 281 Holocene across our domain are likely small (i.e. on the order of 100 meters; Caron et al., 2018),
 282 and therefore we expect this to have a minimal impact on our simulated ice histories.

	Temperature Scenario	Precipitation Scenario	Calving Parameterization
Experiment I	High	High	On
			Off
Experiment II	High	Low	On
			Off
Experiment III	Low	High	On
			Off
Experiment IV	Low	Low	On
			Off

294 **Table 1.** Description of model experiments. See Figure 4 for a display of the temperature and precipitation
 295 forcings scenarios.

298 4. Results

300 We spin up each model as described above (section 3.4) without the ice calving parametrization
 301 turned on. Only when we begin the transient simulation through the Holocene do we turn on the
 302 ice calving parametrization for the ‘Calving On’ scenarios (Table 1). Our transient simulations
 303 begin 12,400 years ago with the ice margin residing along the present-day coastline for all
 304 experiments, which is approximately consistent with where geologic constraints place the ice
 305 margin at that time (Young et al., 2021 and references therein).

307 4.1 Simulated Deglaciation

309 First, we assess how our simulated deglaciation compares with geologic reconstructions of ice
 310 sheet change in the KNS region. Geological constraints outlined above reveal that ice retreated
 311 across the KNS forefield rapidly in the Early Holocene. While relatively little direct information
 312 exists detailing ice retreat within the fjords, the terrestrial portion of our domain (i.e., the inter-
 313 fjord bedrock landscape) became ice-free between ~ 11.2 ka and 9.5 ka as ice retreated from the
 314 modern coastline towards, and eventually surpassing, what is now the modern ice margin.

316 To compare against the geologic constraints, we determine when in time portions of our model
 317 domain become ice free (Figure 5). Since ice can readvance over areas that had been deglaciated
 318 during our simulations, we take the youngest age from which locations in our simulations became

319 ice free. Our simulations illustrate clear differences in the timing of deglaciation across terrestrial
 320 surfaces above sea-level and within the fjords. For the high and low temperature scenarios,
 321 terrestrial surfaces deglacialate up to a few millennia earlier than the adjacent fjords. This difference
 322 in timing between the fjords and terrestrial surfaces is perhaps unsurprising given how fjord
 323 systems act as conduits draining the ice interior. This persistence of ice extent within the fjords
 324 despite elevated warming experienced during early to Middle Holocene illustrates the role of ice
 325 dynamics, which is explored further in section 4.3.

326
 327 For the high and low temperature scenarios, there is little difference between the age of
 328 deglaciation on terrestrial surfaces for simulations that allow (Figures 5a and 5b; Figures 6a and
 329 6b) and do not allow calving (Figures 5d and 5e; Figures 6d and 6e). In contrast, deglaciation of
 330 terrestrial surfaces occurs later in Holocene for the simulations using the high precipitation

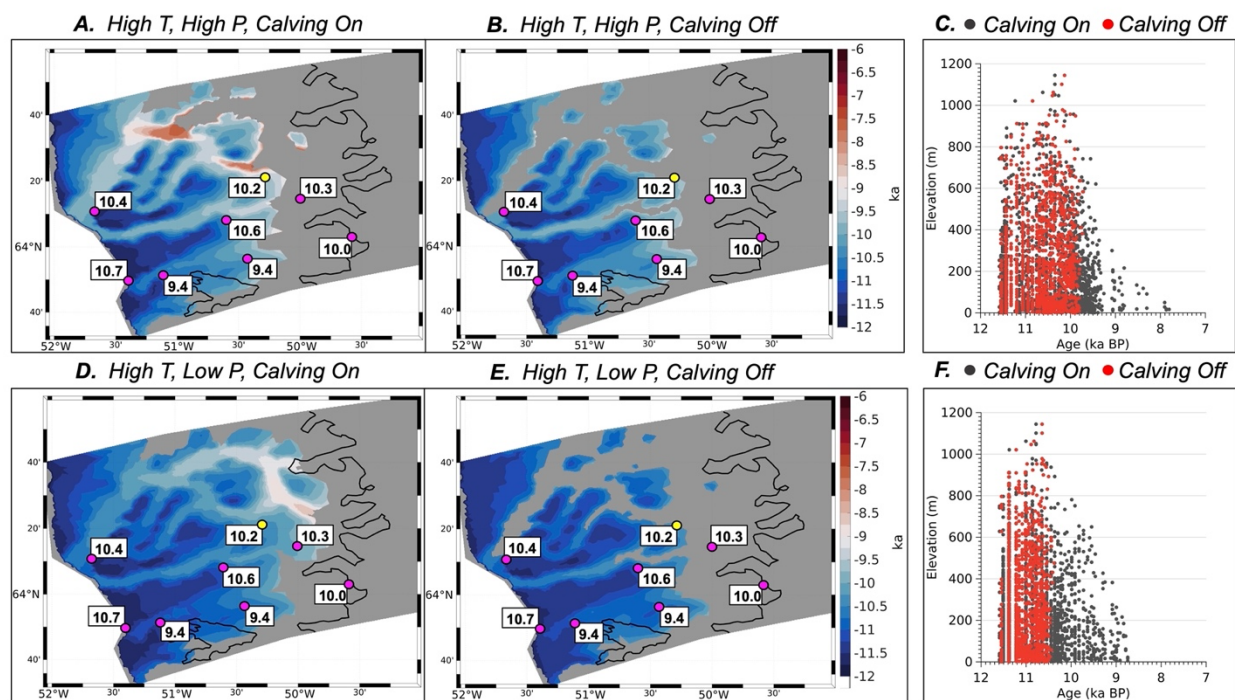


Figure 5. Map of simulated deglaciation ages for High temperature scenarios with A.) High precipitation, Calving On. B.) High precipitation, Calving Off. D.) Low precipitation, Calving On, and E.) Low precipitation – Calving off. Gray mask is the simulated ice extent at present day and the black line denotes the actual present day ice extent (Rignot and Mouginot, 2012). Magenta circles are the best estimate of the timing of deglaciation at that point based on ^{10}Be surface exposure ages in thousands of years ago and the yellow dot shows minimum limiting radiocarbon age (Young et al., 2021). Scatter plot of simulated deglaciation age (above sea level) versus bedrock elevation for C) High temperature, high precipitation, and F) High temperature, low precipitation. Red dots are from simulations without calving and black dots are for simulations with calving.

331 scenario than for those simulations using the low precipitation scenario. For simulations using the
 332 high temperature scenario, these differences are up to 500 years (Figure 5). For the low temperature
 333 scenarios, terrestrial surfaces deglacialate up to 1000 years later for simulations using the high
 334 precipitation forcing (Figure 6).

335

336 Lastly, it is important to note that simulations that allow calving have a more reduced ice extent
 337 (gray mask) at the end of each simulation, which may indicate that calving limits ice front
 338 readvance within the fjord.

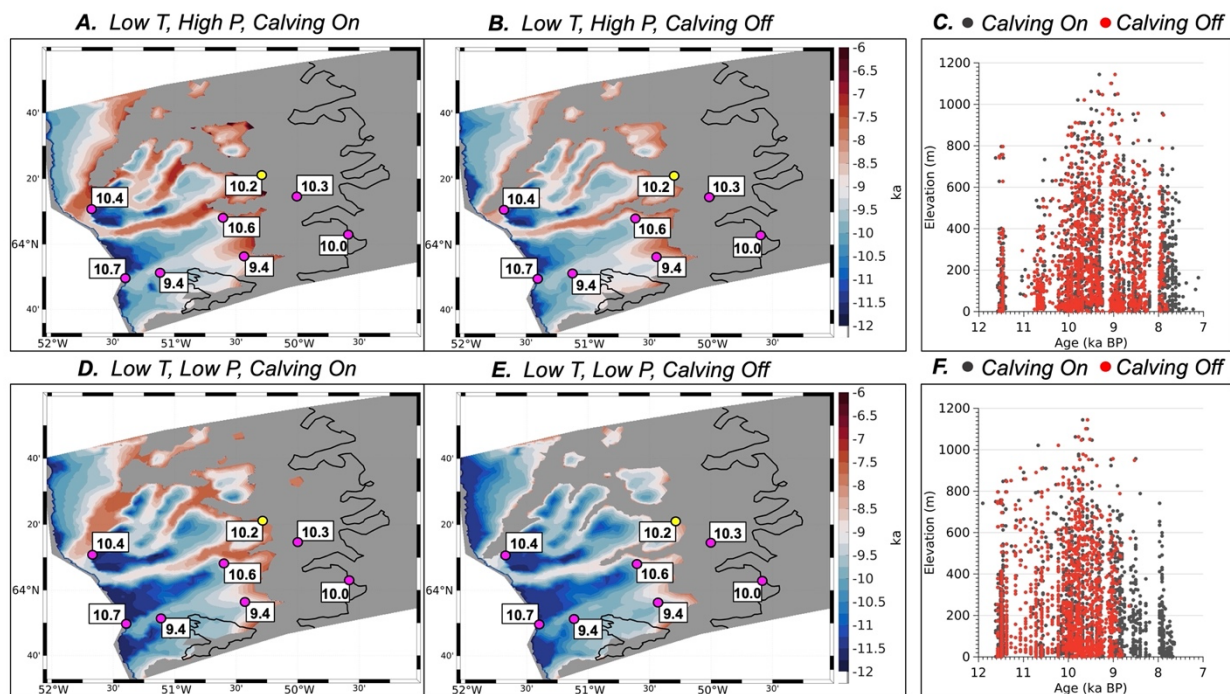


Figure 6. Map of simulated deglaciation ages for Low temperature scenarios with A.) High precipitation, Calving On. B.) High precipitation, Calving Off. D.) Low precipitation, Calving On, and E.) Low precipitation – Calving off. Gray mask is the simulated ice extent at present day and the black line denotes the actual present day ice extent (Rignot and Mouginot, 2012). Magenta circles are the best estimate of the timing of deglaciation at that point based on ^{10}Be surface exposure ages in thousands of years ago and the yellow dot shows a minimum limiting radiocarbon age that requires ice free conditions in the fjord at that time (Weidick et al., 2012; Young et al., 2021). Scatter plot of simulated deglaciation age (above sea level) versus bedrock elevation for C) Low temperature, high precipitation, and F) Low temperature, low precipitation. Red dots are from simulations without calving and black dots are for simulations with calving.

339
 340 The manner in which deglaciation occurs on terrestrial surfaces can be an important factor in
 341 determining the pace and magnitude of the ice margin response to warming. Geologic archives
 342 constraining ice retreat across the KNS forefield span an elevational range of 1300 m, yet, no
 343 elevational dependence on the age of deglaciation is evident (Larsen et al., 2014; Young et al.,
 344 2021). To compare our simulated deglaciation history as a function of elevation against the
 345 geologic data, we plot the simulated age of deglaciation against elevation, and restrict our
 346 datapoints to terrestrial surfaces above sea level (Figure 5c and 5f; Figure 6c and 6f). In general,
 347 our simulations agree with the geologic data indicating that there was no elevational dependence
 348 on the age of deglaciation; if there were any indication of an elevation dependence on the age of
 349 deglaciation, we would observe that high elevation sites would become ice free first, followed by
 350 low elevation sites. Instead, all of the plots show that deglaciation happens simultaneously at
 351 discrete time intervals across all elevation bands, indicating that ice surface lowering was rapid

352 and coincident with ice margin pullback. These elevation-time diagrams also highlight how the
353 higher precipitation scenarios have later mean deglaciation ages across terrestrial surfaces (Figure
354 5c and 6c) than corresponding simulations using the low precipitation scenario (Figure 5f and 6f).
355 We also note that for simulations where calving is turned off (red dots), ice retreat appears to stop
356 earlier than for those simulations with calving turned on (black dots). This occurs because the
357 simulations without calving experience a larger Late Holocene ice readvance than those
358 simulations where calving is turned on (black dots). As a consequence of this, model grid points
359 that would have otherwise deglaciated prior to the readvance are overrun with ice and therefore
360 are not marked as deglaciated in the simulation.

361
362 Lastly, each of our experiments end with a simulated present-day ice extent that is beyond
363 (westward of) the actual present-day ice extent (Figure 5 and 6). Yet, the simulated ice-margin
364 position in the fjords is less extensive for all experiments where calving is permitted. Those
365 experiments that allow calving and used the high temperature scenario (Figure 5a and 5d) simulate
366 a present-day ice extent that is closer to the observed present-day margin when compared to
367 simulations using the low temperature forcing (Figure 6a and 6d).

368

369 **4.2 Ice mass evolution and minimum ice extent**

370

371 Broadly, scenarios that allow calving undergo greater ice mass loss than those simulations where
372 calving is not allowed (Figure 7; black lines). The differences in simulated ice mass also vary
373 depending on the climate scenarios used. For example, during Early Holocene warming (12 ka -
374 8 ka), simulations that allow calving and use the high temperature scenarios (Figure 7a, b)
375 experience ice mass loss, while simulations that do not allow calving experience a period of ice
376 mass stability (Figure 7a, b; dashed red line), which is more prolonged in the simulation using the
377 high precipitation scenario (Figure 7a).

378

379 For the simulations using the low temperature scenario (Figure 7c, d), initial ice mass loss is
380 interrupted by brief increases in ice mass during the Early Holocene (between 11 ka-10 ka). This
381 increase in ice mass occurs for both scenarios with and without calving (Figure 7c, d; black and
382 dashed red line), although the simulations without calving experience larger increases in ice mass
383 during this period. Accordingly, the low temperature simulation with higher precipitation (Figure
384 7c) experiences larger ice mass gain than the simulation using the low precipitation scenario
385 (Figure 7d). During this interval, precipitation is approximately 20-30% more for the high
386 precipitation scenario during the Early Holocene than the low precipitation scenario. Much of this
387 mass gain is due to ice thickening over the interior of the model domain, where despite Early
388 Holocene warming, colder temperatures (at higher elevations on the ice sheet) support snowfall
389 (see section 4.3).

390

391 Throughout the remainder of the Holocene, the evolution of ice mass for experiments using the
392 high temperature scenario (Figure 7a, b) differ from those simulations using the low temperature
393 scenario (Figure 7c, d). Simulations using the high temperature scenario (Figure 7a, b) reach a
394 minimum ice volume between 7.6-7.2 ka. For the simulation using the high precipitation scenario,
395 ice mass increases slightly following this minimum, and remains generally stable throughout the
396 remainder of the Holocene (Figure 7a), whereas the simulation using the low precipitation scenario
397 experiences large ice mass gain following this minimum, with steady growth occurring throughout

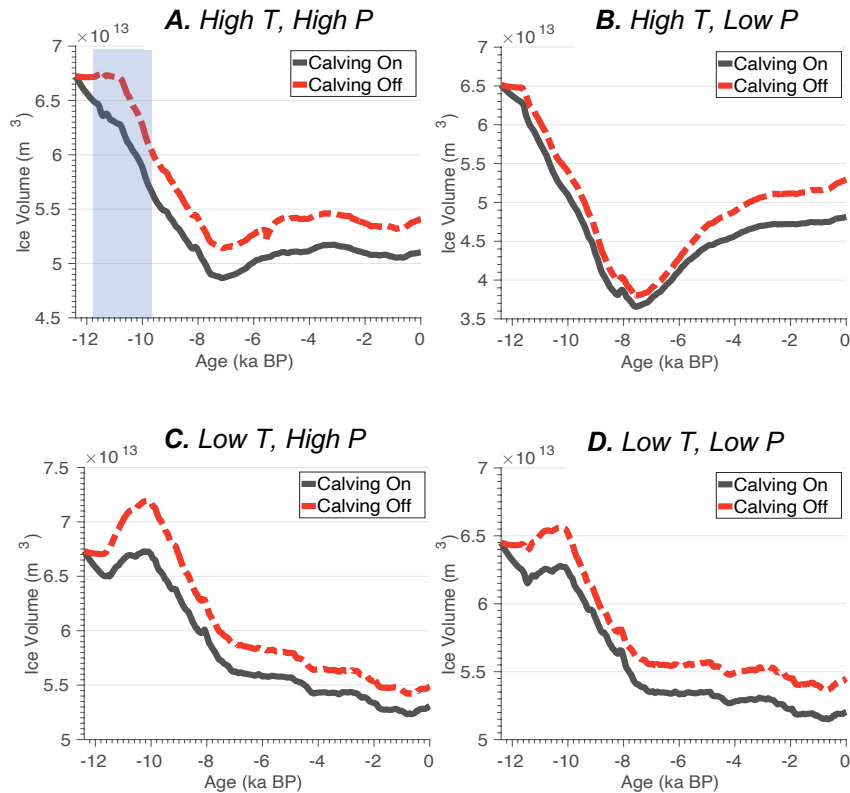


Figure 7. Holocene ice volume ($\times 10^{13} \text{ m}^3$) evolution for each model experiment. Refer to Table 1 for a summary of the climate forcings used in each experiment. Black lines denote those simulations with the calving parametrization turned on. Dashed red lines denote those simulations with the calving parametrization turned off. The vertical blue bar above marks a time period (12 ka – 10 ka) used for analysis presented in Figure 8 and 9.

398 the remainder of the simulation (Figure 7b). It is important to note, however, that for the high
 399 temperature scenarios, this ice mass gain is more muted for simulations that allow calving. In
 400 contrast, the simulations using the low temperature scenario (Figure 7c, d) lose the majority of ice
 401 mass by 8-7 ka, with ice mass loss either continuing through the Holocene (Figure 7c) or remaining
 402 relatively stable before reaching a minimum at 0.6-0.4 ka (Figure 7d).

403
 404 Regional relative sea-level records reveal that sea level fell below modern between 4-3 ka, before
 405 rising towards modern values (Long et al., 2011), interpreted to represent the re-loading of the
 406 Earth's crust as the GrIS readvanced during the Late Holocene following a mid-Holocene
 407 minimum. In addition, radiocarbon-dated lake sediments from southwestern Greenland suggest
 408 that this sector of the GrIS likely achieved its minimum extent after ca. 5 ka, and that eastwards
 409 retreat of the ice margin was likely minimal (Larsen et al., 2015; Young and Briner, 2015; Lesnek
 410 et al., 2020; Young et al., 2021). Although no direct geological constraints on the minimum GrIS
 411 ice extent during the Holocene exist, available constraints suggest that the magnitude of large-
 412 scale ice margin retreat inboard of the present-day extent as simulated by some ice sheet models
 413 in this sector (20-40 km; Tarasov and Peltier, 2002; Lecavalier et al., 2014) is likely too extreme.

414 Relying on these geologic constraints, we can crudely assess the temporal and spatial patterns of
415 the simulated ice mass and minimum extent.

416
417 None of our simulations accurately capture the exact timing of the GrIS minimum in the KNS
418 region, but some simulations are likely better representations than others. Simulations using the
419 high temperature scenario (Figure 7a, b) achieve an ice mass minimum prior to 5 ka followed by
420 ice regrowth. The high temperature-low precipitation scenario depicts an extreme GrIS minimum
421 followed by significant regrowth. While of the overall pattern of a GrIS minimum followed be

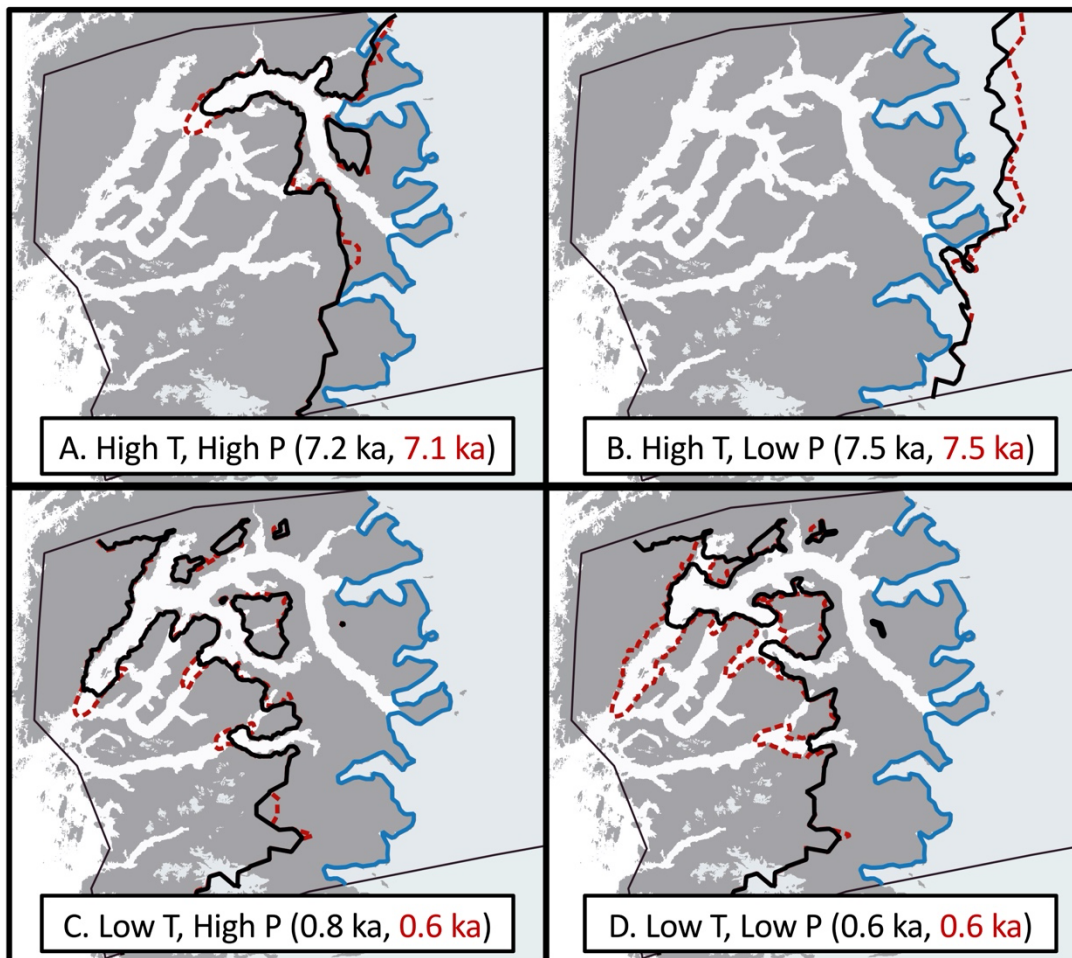


Figure 8. Age of minimum ice extent for each simulation (black text: simulations with calving, red text: simulations without calving). The black line denotes the minimum ice extent for simulations with calving. The dashed red line denotes the minimum ice extent for simulations without calving. The present-day ice extent is shown as the blue line.

422 regrowth is consistent with the geologic record, the magnitude of simulated change is likely
423 inconsistent with geological records, pointing to a rather modest GrIS minimum; although we do
424 acknowledge that minimal ice retreat as constrained by the geologic record does not necessarily
425 equate to muted mass loss. In contrast, the high temperature-high precipitation experiment depicts
426 an ice-mass minimum that is likely too early, but the magnitude of this minimum is less (Figure
427 7A). Moreover, ice regrowth following this minimum is restricted with only modest change

428 occurring over the last 6 kyr (Figure 7A). Although this simulated minimum is likely too early, a
429 simulated ice mass that undergoes minimal change over the last ~6 kyr is broadly consistent with
430 the geological record that depicts a minimum closer to ca. 4-3 ka, but where the GrIS margin likely
431 did not undergo significant change between ca. 7-3 ka (Young et al., 2021). Both low temperature
432 scenarios are inconsistent with the geological record as both show continued ice mass loss through
433 the Holocene. Although it is possible, but unlikely, that continued ice loss through the Holocene
434 could still be achieved if the ice margin retreated inland followed by a readvance toward its present
435 position, mass loss through the Holocene is inconsistent with relative sea-level records.

436
437 The minimum ice margin extent achieved in our simulations is shown in Figure 8. For the high
438 temperature scenarios (Figure 8a, b), the simulated minimum ice extent is either just outboard of
439 the present-day ice margin (Figure 8a; high precipitation) or inboard of the present-day ice margin
440 (Figure 8b; low precipitation). Because the geologic evidence supports that the Holocene ice
441 extent minimum was close to and perhaps slightly inboard of the present-day ice margin (Young
442 et al., 2021), both simulations are broadly consistent with the geological record. But, again, the
443 high temperature – high precipitation scenario depicts significant ice regrowth resulting in a
444 present-day ice margin significantly more extended than modern (Figure 5).

445 **4.3 Early Holocene Thinning**

446
447
448 Figures 9 and 10 show the simulated ice elevation changes for the time period between 12 ka to
449 10 ka for each experiment (highlighted in Figure 7a as the light blue vertical bar). During this time
450 period, widespread Early Holocene warming drove increased ice melt along the margin of the
451 model domain. This pervasive thinning along the margin is captured in all model experiments
452 (Figure 9 and 10), although the amplitude of ice thinning is greatest for the experiments using the
453 high temperature scenario (Figure 9). Across all experiments, inland thickening occurs, however,
454 the magnitude of interior thickening is not solely influenced by the SMB, but is also influenced by
455 calving. For our experiments that allow calving, interior thickening is reduced and ultimately
456 influences the trend and magnitude of changes in simulated ice volume; simulations that allow
457 calving either experience increased ice mass loss (Figure 7a, b) or more muted ice mass gain
458 between 12 ka and 10 ka (Figure 7c, d). Additionally, the spatial pattern of elevation changes
459 shows that marginal thinning propagates farther upstream and into the ice sheet interior for
460 simulations that allow ice calving. This relationship continues throughout the remainder of the
461 Holocene, as experiments with calving either result in more mass loss than simulations without
462 calving, or more muted ice mass gain (see Figure 7). These variations in simulated Holocene ice
463 mass and ice surface elevation change can be linked to the influence ice calving has on ice front
464 position and stability, and ultimately the rate at which ice can flux through the fjord system. During
465 the time period of 12 ka to 10 ka, ice velocity differences for simulations with and without calving
466 are in excess of 200 m/yr along many fjords within the KNS region (Figure 11). Calving at the
467 ice front leads to increases in ice velocity within outlets across the model domain, thereby
468 promoting increased mass flux and transport from the ice interior to the margin. Thus, even though
469 the large-scale ice margin migration across our model domain is relatively insensitive to calving,
470 the overall mass budget and surface profile of the ice is strongly influenced by calving.

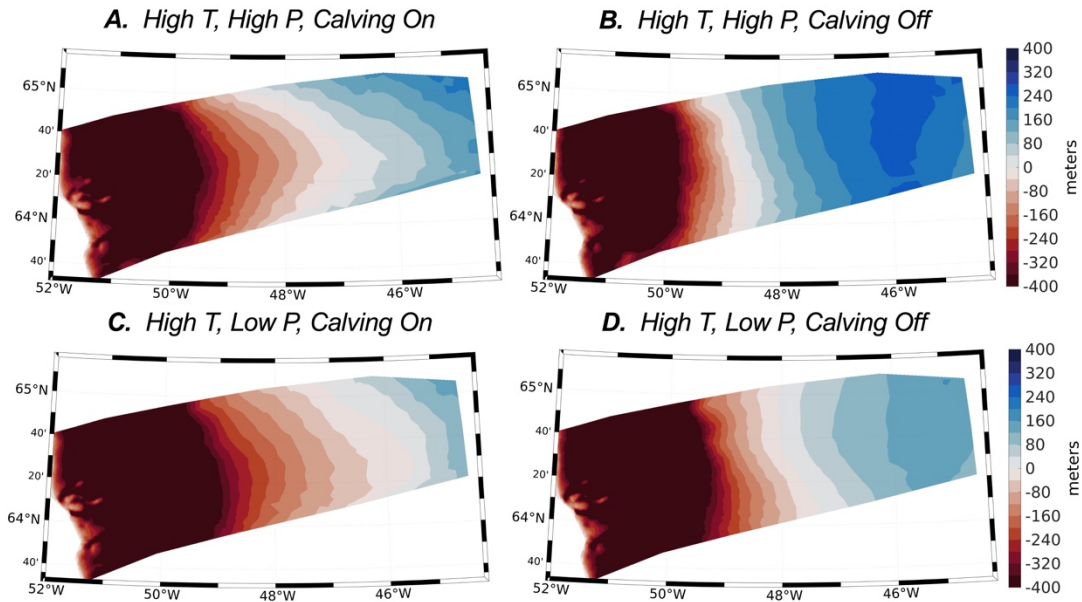


Figure 9. Simulated elevation changes (in meters) during period 12ka – 10ka shown for experiments using the high temperature forcing.

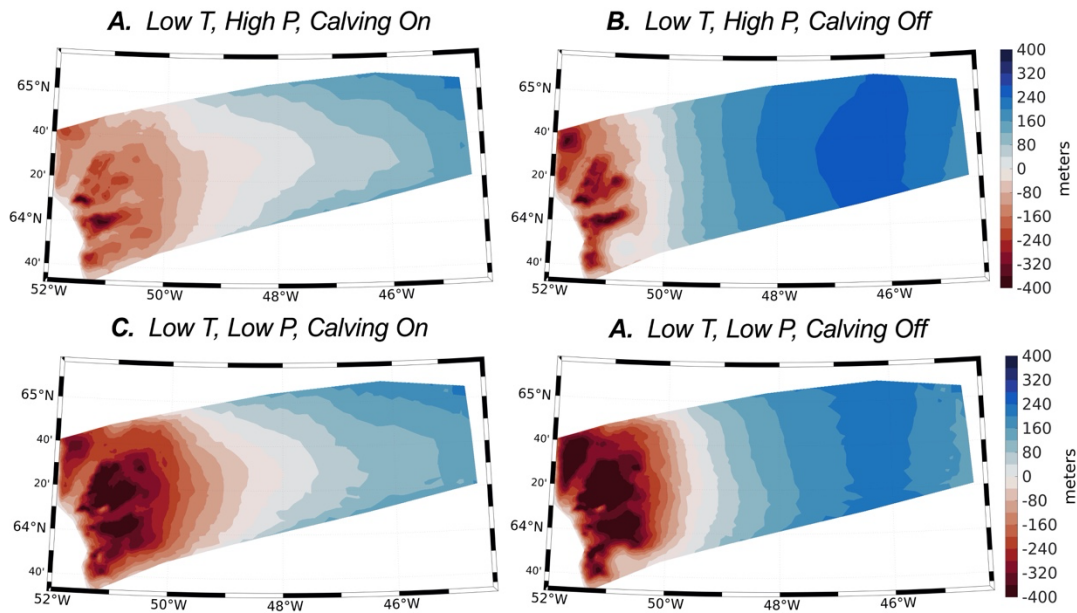


Figure 10. Simulated elevation changes during period 12ka-10ka shown for experiments using the low temperature forcing.

471
 472 Reconstructions of Holocene ice thickness across the GrIS are limited, but ice-core records provide
 473 a long-term perspective of dynamic changes in GrIS elevation at locations at or near the ice divide
 474 (Vinther et al., 2009; Lecavalier et al., 2017). For example, some locations experienced more rapid
 475 thinning in response to Holocene warming (i.e. Camp Century, Dye 3) while other locations

476 experienced more muted ice elevation changes (i.e. GRIP, NGRIP). A feature of many of these
 477 records, however, is the presence of Early Holocene thickening, potentially triggered by increased
 478 snowfall at higher elevation sites as the climate warmed or by elevation-mass balance feedbacks
 479 driven by isostatic uplift (Vinther et al., 2009). Across all model experiments, our simulated timing
 480 of inland thickening coincides with thickening experienced at high elevation ice core locations
 481 (Vinther et al., 2009). The magnitude of Early Holocene thickening from ice core records (Vinther
 482 et al., 2009; 11.7 ka-10 ka) is on the order of 30 – 70 meters. Therefore, our simulations that allow
 483 calving display inland thickening (<120 m) over the time interval 12 ka – 10 ka that is more
 484 consistent with thickening estimated from ice cores than simulations with no calving (>200m).

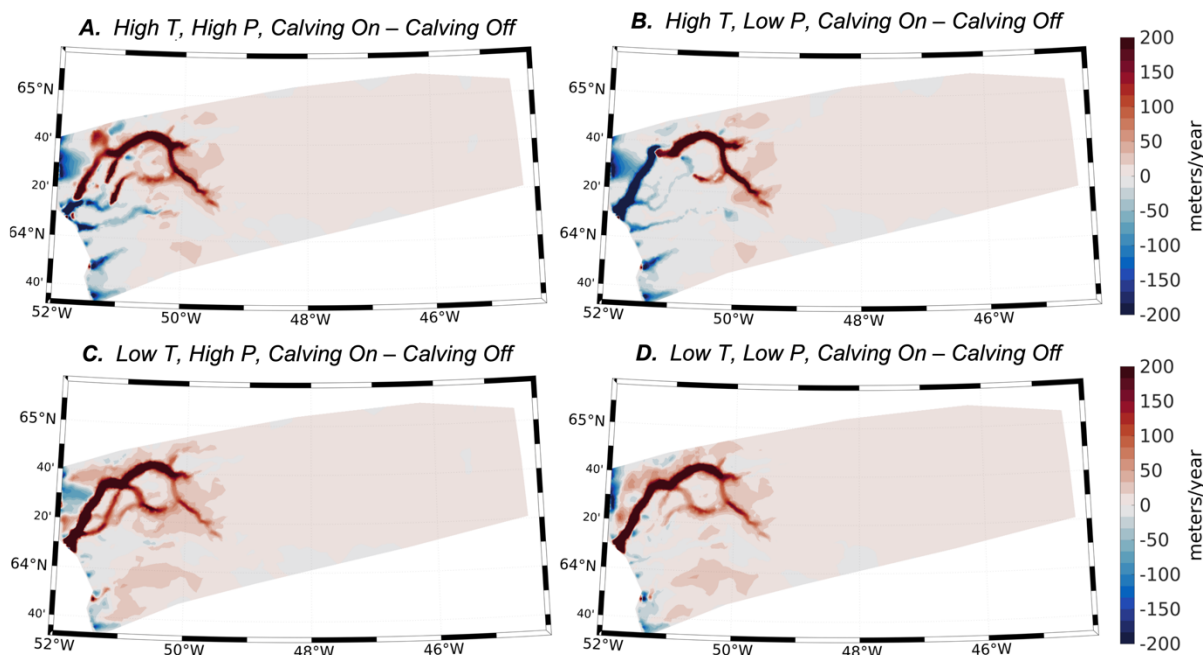


Figure 11. Simulated ice velocity differences between simulations with and without calving for each experiment over the time period 12 ka to 10 ka. Red colors denote an increase in ice velocity for simulations with calving relative to simulations without calving. Blue colors denote a decrease in ice velocity for simulations with calving relative to simulations without calving.

485 486 **4.4 Sensitivity to marine forcing** 487

488 Experiments on the tensile strength of ice show that stress thresholds can vary between 150 kPa
 489 and 3100 kPa (Petrovic, J. 2003), with modeling experiments on Jakobshavn glacier suggesting
 490 that the stress threshold for grounded ice can vary between 100 kPa to 4 MPa seasonally (Bondzio
 491 et al., 2017). Here, our grounded ice stress threshold is set to 600 kPa. Because our model setup
 492 incurs large computational expense, we did not perform a full uncertainty analysis on these
 493 parameterizations. Due to the nature of modeled variation in calibrated stress thresholds across
 494 Greenland (Choi et al., 2021), however, we ran a small set of experiments where we set the calving
 495 stress threshold on grounded ice to 1 MPa. We performed the transient simulations on the high
 496 and low temperature scenario cases using the high precipitation forcing (see Table 1).
 497 Additionally, we ran a set of experiments where the basal melt rate on floating ice was set to 120

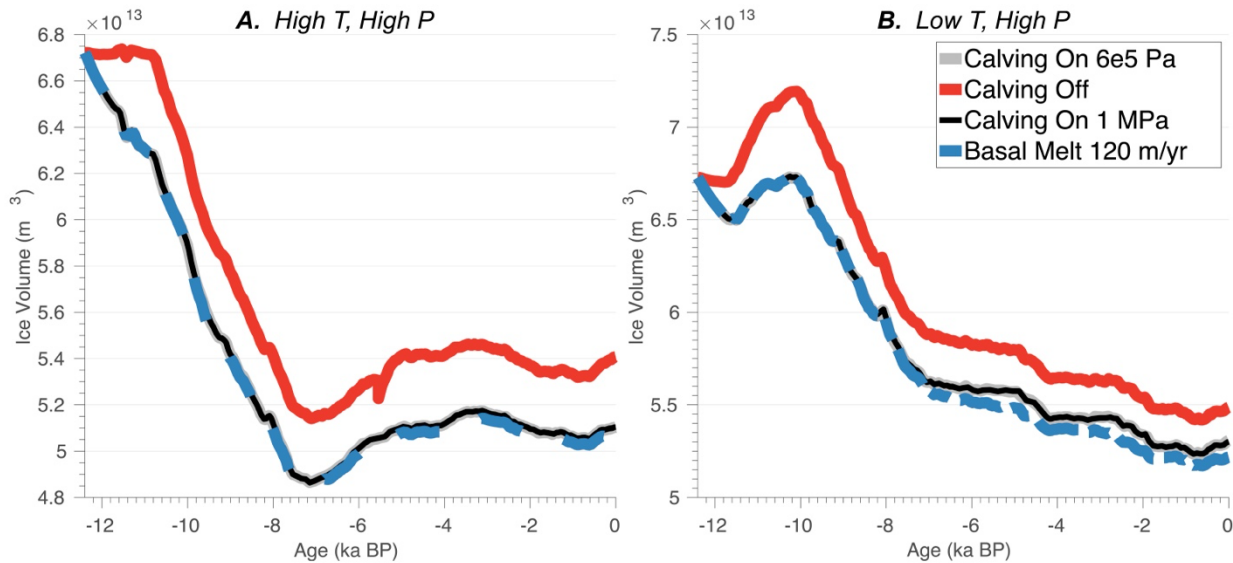


Figure 12. Sensitivity to the calving stress threshold for grounded ice and basal melt rates on floating ice. Red line: ice volume evolution for the simulations where the calving parameterization was turned off. Black line: ice volume evolution for the simulations where the calving stress threshold for grounded ice is 1 MPa. Gray line: ice volume evolution for the simulations where the calving stress threshold for grounded ice is 6 kPa. Dashed blue line: ice volume evolution for the simulations where the basal melt rate on floating ice was set to 120 m/yr (with calving stress threshold for grounded ice = 600 kPa).

498 m/yr. Figure 12 shows the simulated ice volumes for these experiments where the calving stress
 499 threshold of grounded ice and basal melt rate on floating ice were changed. These experiments
 500 reveal that adjusting the stress threshold from 600 kPa to 1 MPa has no effect on the evolution of
 501 the simulated ice volume. Accordingly, increasing the basal melt rate on floating ice has minimal
 502 effect on the simulated ice volume (Figure 12). Ice only begins to float in our experiments when
 503 the ice front retreats into the deeper fjord bathymetry within the KNS forefield (see Figure 3), and
 504 therefore submarine melting of floating ice seems to have limited influence on simulated ice mass
 505 changes.

506
 507 **5. Discussion**

508
 509 **5.1 Terrestrial vs. Marine ice retreat**

510
 511 Southwestern Greenland hosts a rich record of geologic constraints on past ice-sheet change
 512 (Lesnek et al., 2020). Whereas a series of well-defined moraines constrain Early Holocene ice
 513 retreat across portions of southwestern Greenland dominated by terrestrial ice-margin settings
 514 (Larsen et al., 2014; Lesnek et al., 2020; Young et al., 2020; Young et al., 2021), the Kapisigdlit
 515 moraine system (Figure 2: Early Holocene moraines) near the present-day ice margin is the only
 516 regionally traceable moraine within the marine-dominated KNS forefield. Instead, ice-margin
 517 retreat across the KNS forefield is constrained primarily by minimum limiting radiocarbon ages
 518 and ¹⁰Be surface exposure ages on deglaciated bedrock surfaces and glacial erratics (Larsen et al.,
 519 2014; Young et al., 2021). The lack of moraine systems between coast and ice is consistent with
 520 the relatively high rate of deglaciation estimated from the existing chronology. These
 521 chronological constraints detail widespread and rapid retreat of the ice margin across this domain

522 in the Early Holocene, with the ice margin retreating from the coastline around 12 ka to near the
523 present-day ice margin between 10-9.5 ka (Young et al., 2021). This relatively rapid retreat based
524 on geological observation is consistent with the lack of elevation-age relationship in our
525 simulations of ice margin change.

526
527 While the rapid retreat of the terrestrial ice margin is well constrained, how ice retreated up the
528 fjords is less certain. Our simulations depict a pattern of ice retreat across the landscape that was
529 largely independent of ice retreat within fjords, which lagged by 0.5 – 2 ka. For our simulations,
530 scenarios using the same climate forcing show little difference (<1 ka) in the simulated age of ice
531 retreat on terrestrial ice margins regardless of whether calving is allowed (Figures 5 and 6). The
532 timing and rate of Holocene ice retreat across terrestrial portions of the KNS forefield, however,
533 is strongly dependent on the climate forcing used, and ultimately the SMB. The earliest ice retreat
534 occurs in simulations that use the high temperature scenario. Ice retreat occurs later in simulations
535 that use the low temperature scenario, which has a delay in the timing and magnitude of Holocene
536 warming (Figure 4). The pace and magnitude of ice retreat is shown to be modulated depending
537 on precipitation similar to the findings of Briner et al. (2020) and Downs et al. (2020), with delayed
538 and less rapid ice retreat in scenarios with higher precipitation (Figures 5a and 6a). These results
539 point to the strong influence that climate and, in particular, precipitation can have on modulating
540 the temperature driven response of Holocene deglaciation. Indeed, select proxy records suggest
541 that southwestern Greenland may have experienced a prolonged period of anomalously high
542 snowfall in the Early Holocene, perhaps driven by increased moisture flux from Baffin Bay and
543 the Labrador Sea as sea-ice extent declined (Thomas et al., 2016). Ice flow modeling across
544 southwestern Greenland has also revealed that elevated precipitation may have accompanied Early
545 Holocene warming (Downs et al., 2020). And recent evidence from a shallow ice core in western
546 Greenland reveal that significant variations in precipitation occurred in the last two thousand years
547 across the margins of the GrIS, whereas this variability is not present in ice core data at the interior
548 of the GrIS (Osman et al., 2021). Because current climate reconstructions employed in
549 paleoclimate ice flow modeling use either simple scaling approaches to reconstruct past climate or
550 rely on information from interior ice cores, large hydroclimate shifts that occur at the ice sheet
551 margin may not be captured (Badgley et al., 2020). Continued progress in reconstructing past
552 climate will certainly improve our understanding of climatic controls on the long-term response of
553 the GrIS.

554
555 In general, simulations using the high temperature scenario experience terrestrial ice retreat that
556 occurs during 11.5 ka to 9 ka, a time window consistent with the geological record of ice-margin
557 change in our domain (Larsen et al., 2014; Young et al., 2021). Simulations using the low
558 temperature scenario reveal terrestrial ice retreat also beginning ca. 11.5 ka, but deglaciation of
559 our model domain continues until ~7.5 ka. In comparison, geological constraints suggest that by
560 ~10.3-9 ka BP the ice margin in the immediate KNS region had already retreated back to, and
561 likely behind, what is the present-day ice margin (Young et al., 2021). Ice surface lowering is
562 captured in all of our simulations, which indicate that on terrestrial surfaces ice retreat was
563 synchronous across low and high elevations. Therefore, the simulated ice retreat could indicate
564 large scale ice margin retreat in response to rapid ice surface lowering, but certainly precludes
565 scenarios where ice surface lowering occurred slowly exposing high elevation sites well before
566 low elevation sites. While ice calving does not seem to significantly influence the rate and timing
567 of ice retreat across terrestrial portions of our domain, Late Holocene ice readvance within fjords

568 is more restricted in those simulations that use the calving parametrization. Accordingly, flowband
569 modeling of KNS over the period historical period of 1761 to 2012 suggests that marine ice-front
570 retreat was primarily influenced by atmospheric warming and runoff, which helped to trigger ice
571 front retreat via a crevasse-depth calving criterion, with submarine melting only playing a minor
572 role on historical retreat (Lea et al., 2014; Lea et al., 2014). These results do suggest though that
573 climate anomalies were the main driver of historical ice terminus advance and retreat across KNS
574 (Lea et al, 2014), with our results suggesting that the longer-term Holocene ice terminus position
575 was also primarily driven by atmospheric warming and not through oceanic melting.

576

577 **5.2 Role of ice calving on mass transport**

578

579 Mass transport from the ice sheet interior to the margin plays an important role in ice sheet mass
580 change and ultimately its contribution to sea-level rise. Contemporary satellite-derived
581 measurements show inland thickening at high elevations across portions of the GrIS in response
582 to increased snowfall despite pervasive thinning at lower elevations (Smith et al., 2020). Although
583 the response of marine terminating portions of the GrIS and how it translates to interior ice mass
584 loss can be spatially varying (Williams et al., 2021), thinning at the ice margin due to dynamic- or
585 SMB-driven ice loss can elicit changes in driving stresses, which can propagate up glacier and into
586 the interior of the ice sheet (Price et al., 2008; Schlegel et al., 2013; Csatho et al., 2014; Felikson
587 et al., 2020; Williams et al., 2021).

588

589 While there is no apparent influence of ice calving on the Holocene ice retreat across the KNS
590 forefield over terrestrial surfaces, our simulations show that ice calving has a significant influence
591 on the evolution of the total ice volume. Ultimately, ice calving leads to an acceleration of ice
592 flow within outlet glaciers that promotes local ice thinning first, followed by propagation of this
593 thinning into the interior of the ice sheet, consistent with contemporary observations (Csatho et
594 al., 2014; Williams et al., 2021). Initially, interior ice surface elevation increases in our
595 simulations, with simulations that allow calving being more consistent with ice-core derived
596 surface height records (Vinther et al., 2009). Surface lowering near the ice margin driven by a
597 more negative SMB in response to Early Holocene warming causes the ice surface slopes to
598 steepen in our domain, increasing driving stresses and mass transport. This helps drive interior ice
599 thinning, as shown by elevation changes in simulations that allow ice calving (Figures 9 and 10),
600 leading to increased ice flux at the margin through the ice streams (Figure 11). This increased
601 mass transport helps limit thinning within outlet glaciers, and where terrestrial locations of our
602 domain become ice free early in the Holocene, ice front retreat within the fjords lag (Figure 5 and
603 6).

604

605 Our results suggest that, while calving did not play a significant role in the observed Holocene ice
606 retreat across the KNS forefield, it played an important role on the overall ice mass change across
607 our model domain. These results highlight that the inclusion of physically based ice calving
608 parameterizations is an important step towards modeling the fidelity of simulated ice mass change
609 across paleoclimate timescales. However, the choice of which ice calving parameterization is best
610 suited to Greenland over such timescales is still not well constrained (Goelzer et al., 2017). It
611 remains important though, that models maintain high enough spatial resolution in order to capture
612 fjord environments, associated bathymetry, and ultimately ice calving and grounding line

613 migrations over paleoclimate timescales (Cuzzone et al., 2019) as the model resolution can impact
614 simulated ice discharge significantly (Rückamp et al., 2020; Ashwanden et al., 2019).

615

616 **5.3 Model limitations**

617

618 Fjord systems in Greenland are typically <5 km in width, making it necessary to implement high-
619 resolution meshes to resolve these features. Our model setup relies on a high-resolution mesh that
620 is able to capture the fjord geometry within the KNS forefield, making it possible to simulate
621 grounding line migration and calving. The calving parameterization used does ignore frontal
622 melting at the grounded ice front. Frontal melt at the base of a calving face has been shown to
623 induce undercutting of the ice front, and greatly increases calving rates (O’leary and
624 Christofferson, 2013). For the present day, many of southwestern Greenland’s marine terminating
625 glaciers are not strongly influenced by undercutting (Wood et al., 2021), but this may have been
626 different as ice retreated up fjord to its present-day location through the Holocene. While proxy
627 records indicate changing sea surface temperatures during the Holocene proximal to our model
628 domain Axford et al. (2021), due to a lack of constraints on the long-term subsurface ocean thermal
629 forcing needed to implement undercutting in our simulations, we opted to disregard this. To
630 circumvent this shortcoming, we set our calving stress threshold on grounded ice to a number (600
631 kPa) that is on the lower end of measured tensile stresses of ice (Petrovic, 2003). Since there was
632 no discernable difference in our simulated ice mass change when a higher calving stress threshold
633 of grounded ice was used (1 MPa), we cautiously assume that implementation of undercutting
634 would have a negligible effect on the calving rates and overall Holocene mass change and ice
635 retreat across our domain. Future work will use a basal melt-rate parametrization (PICOP; Pelle
636 et al. 2019), employed in ISSM currently, to estimate oceanic melt rates from far field variations
637 in Holocene subsurface temperature and salinity in order to more robustly estimate the impact of
638 oceanic warming Holocene deglaciation across the GrIS.

639

640 At the time of this work, ISSM is undergoing improvements and new implementation of solid earth
641 and sea-level feedbacks. While we did not include time dependent forcings (e.g. Caron et al.,
642 2018) that account for relative sea-level change as we have in prior research (Cuzzone et al., 2019;
643 Briner et al., 2020), future simulations using ISSM will explore the influence of coupled solid
644 Earth-ice feedbacks on ice retreat. Recent ice sheet modeling (Kajanto et al., 2020) showed that
645 the Holocene retreat of Jakobshavn Isbræ was insensitive to relative sea-level (RSL) variations, as
646 RSL changes were small in comparison to fjord depth. RSL changes during the Holocene across
647 this domain were relatively small (~60-100 meters at 12.4 ka and decreasing through the Holocene;
648 Caron et al., 2018) compared to fjord depths. Given that ice calving did not seem to largely
649 influence terrestrial ice retreat, we only expect that inclusion of Holocene RSL changes may have
650 influenced ice front retreat that migrated into deeper waters where floating extensions of the ice
651 front could occur. However, in our sensitivity tests, basal melting on floating ice plays a trivial
652 role in total ice volume changes (Figure 11) as most of the ice within fjords is grounded during the
653 Holocene retreat.

654

655 **6. Conclusions**

656

657 Understanding how climate, calving, and marine processes contribute to ice sheet change across
658 paleoclimate timescales is challenging. Models with lower resolution meshes are typically favored

659 to ensure computational needs are satisfied. This ultimately leads to poor representation of
660 bedrock topography (Cuzzone et al., 2019; Jones et al., 2021) and grounding line migration
661 (Seroussi et al., 2018) that control ice flow (i.e., fjords), making the assessment of how ice calving
662 influences large scale ice margin change difficult. Moreover, while ice core records provide
663 snapshots of a changing climate at the ice-sheet interior, there remain a relative lack of
664 paleoclimate records from the ice sheet margin of sufficient resolution that can be easily
665 incorporated into an ice sheet model's climate forcing.

666
667 Here, we presented results from a high-resolution 3D thermomechanical regional ice sheet model
668 that evaluated controls on the behavior of the southwestern GrIS during the Holocene in the
669 vicinity of the KNS forefield, an area with extensive geologic constraints on past ice margin
670 change. Experiments were driven by novel reconstructions of Holocene climate (Badgeley et al.,
671 2020) and included a physically based ice calving parametrization (Morlighem et al., 2016).

672
673 Our modeling results shed light on the well constrained observations of Holocene ice retreat across
674 the KNS forefield. These simulations agree well with observations that ice retreat on terrestrial
675 bedrock surfaces occurred rapidly between 11.5 ka to 9.5 ka in response to Early Holocene
676 warming. The variations in the timing and magnitude of ice retreat on terrestrial bedrock surfaces
677 across this region are found to be insensitive to calving within the fjords that intersect this
678 landscape. Instead, the terrestrial ice retreat is more sensitive to the SMB, with warmer climate
679 reconstructions providing the best fit between the modeled and observed ice retreat. Calving
680 however does play a significant role in the simulated Holocene ice volume change across this
681 domain. Acting as conduits for mass transport and ice flux, ice velocity within the fjords in the
682 KNS forefield increases when the ice front is allowed to calve. Calving helps promote further ice
683 mass transport from the interior of the domain to the ice front which helps to thicken ice within
684 the fjords, allowing the ice front to persist longer than adjacent terrestrial margins similar to the
685 ice response simulated for the Holocene retreat of Jakobshavn Isbræ (Kajanto et al., 2020). These
686 results suggest that paleo ice flow models that do not sufficiently resolve fjord geometry may not
687 capture dynamic processes that are critical towards understanding long term ice mass change
688 across the GrIS. Recent ice flow modelling has suggested that despite increased ice mass loss due
689 to a more negative SMB, ice discharge from GrIS marine terminating glaciers will play a
690 significant role in overall GrIS mass change well into the future (Choi et al., 2021). These results
691 confirm that over paleoclimate timescales, while the SMB may dictate large scale ice margin
692 migration as captured in geologic observations, ice discharge has the ability to greatly influence
693 the rate and magnitude of ice mass change. However, as all simulations depict contemporary ice
694 extent that is too extensive, uncertainties in the reconstruction of past climate and model parametric
695 uncertainties ultimately contribute to misfits that are difficult to quantify given our
696 computationally expensive model setup. Future paleoclimate ice flow modelling with ISSM will
697 aim to take advantage of recent advances in statistical emulation (e.g., Edwards et al., 2021) to
698 better quantify the influence of model parametric uncertainty on simulated Holocene ice retreat.

699
700 Geologic archives serve an important role in our understanding of glacier and ice sheet response
701 to climate change. In turn, ice sheet modeling can help improve our understanding of the climatic
702 and ice dynamical factors that led to ice sheet changes preserved by the geologic record. Our
703 modeling results present an exploration of the factors that may have contributed to the observed
704 pattern of Holocene ice retreat across the KNS forefield, echoing that model–data comparisons

705 between ice sheet models and geologic reconstructions can help improve our understanding of
706 long-term ice sheet sensitivity to climatic and dynamic forcing mechanisms.

707

708 **Data Availability**

709

710 The simulations performed for this paper made use of the open-source Ice Sheet System Model
711 (ISSM) version 4.19 and are publicly available at <https://issm.jpl.nasa.gov/> (Larour et al., 2012).

712

713 **Acknowledgements**

714

715 Funding for this study was provided by the National Science Foundation Grant ARC no. 2105960
716 to JC and no. 1503959 to NEY. We thank the editor Caroline Clason, and James Lea and an
717 anonymous reviewer for their constructive feedback regarding this work.

718

719

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