



Inferring horizontal asymmetry of the bulk ice crystal fabric from phase-sensitive radar measurements

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Abstract. The bulk crystal orientation in ice influences the flow of glaciers and ice streams. The ice *c*-axes fabric is most reliably derived from ice cores. Because these are sparse, the spatial and vertical distribution of the fabric in the Greenland and Antarctic ice sheets is largely unknown. In recent years, methods have been developed to determine fabric characteristics from radar measurements. The aim of this paper is to present an improved method to infer the horizontal fabric asymmetry by precisely determining the travel-time difference using co-polarised phase-sensitive radar data. We applied this method to six radar measurements from the EastGRIP drill site on Greenland's largest ice stream to give a proof-of-concept by comparing the results with the horizontal asymmetry of the bulk crystal anisotropy derived from the ice core. This comparison shows an almost perfect agreement, which is a large improvement compared to previously used methods. Our approach is particularly useful for determining the vertical profile of the fabric asymmetry in higher resolution and over larger depths than was achievable with previous methods, especially in regions with strong asymmetry.

1 Introduction

The distribution of the crystallographic-axis (*c*-axis) orientation fabric (henceforth *fabric*) in glaciers and ice sheets is a result of ice deformation history and influences present-day ice-flow dynamics (Alley, 1988; Faria et al., 2014). Due to the mechanical anisotropy of ice crystals, the bulk viscosity is a directional quantity, spanning several orders of magnitude, depending on the orientation of stresses with respect to the fabric type and orientation (Cuffey and Paterson, 2010).

While some ice-flow models already account for fabric evolution and/or its effect on ice flow (e.g. Thorsteinsson, 2002; Gillet-Chaulet et al., 2006; Seddik et al., 2008; Martín et al., 2009), the validation of these models is obstructed by the lack of in-situ data.

Most reliably, the crystal fabric of ice can be determined from the analysis of ice core thin sections (e.g. Thorsteinsson et al., 1997; Azuma et al., 1999; Wang et al., 2002; Montagnat et al., 2014; Weikusat et al., 2017). Since deep ice cores are sparse in Greenland and Antarctica and often restricted to domes with rather undisturbed stratigraphy, little is known about the spatial



distribution of crystal fabric anisotropy of ice sheets. It is therefore of great importance to use other methods in order to infer the spatial and vertical distribution of the fabric asymmetry.

Ice crystals are uniaxially birefringent (Hargreaves, 1978). This means that crystals also have a dielectric anisotropy in addition to the crystal anisotropy and thus, the horizontal fabric asymmetry can also be surveyed using radar measurements (e.g. Fujita et al., 2006; Drews et al., 2012; Leinss et al., 2016; Brisbourne et al., 2019; Jordan et al., 2019, 2020; Young et al., 2021a, b; Ershadi et al., 2022; Jordan et al., 2022; Gerber et al., 2022), with certain limitations (Rathmann et al., 2022). Since radar measurements are easier to achieve than ice core analyses, they enable a good spatial coverage and thus offer the opportunity to examine the distribution of fabric asymmetry.

For vertically propagating radio waves, the relevant dielectric anisotropy is the difference between the bulk horizontal permittivities (Fujita et al., 2000). One way of inferring the horizontal fabric asymmetry is based on a polarimetric coherence method (Dall, 2010), which has recently been applied to radar data and compared with the fabric asymmetry from the NEEM ice core in Greenland (Jordan et al., 2019), WAIS divide ice core in West Antarctica (Young et al., 2021a) or the EDML and EDC ice cores in East Antarctica, respectively (Ershadi et al., 2022).

However, this method has some limitations (Leinss et al., 2016). Most importantly, the method can either only be used where the asymmetry of the fabric is weak or otherwise its application is severely limited to shallow depth (Jordan et al., 2022), which we discuss later in detail.

In this study, we infer the horizontal asymmetry of the bulk crystal fabric at the East Greenland Ice Core Project (EastGRIP) drill site from co-polarised phase-sensitive radar measurements by using a new, improved coherence method. This method differs from previously used analysis schemes and has the advantage that the asymmetry can be determined with much higher vertical resolution and, regardless of its strength, up to the onset of the noise level. We present a proof-of-concept by comparing the derived horizontal fabric asymmetry with those from the ice core analysis. A glaciological interpretation of the detected fabric asymmetry regarding the flow dynamics in the region of the EastGRIP drill site is part of a larger study by Gerber et al. (2022) and we refer to their study for ice-dynamical interpretations.

2 Data

In order to investigate ice flow dynamics of Greenland's largest ice stream, the Northeast Greenland Ice Stream (NEGIS), an ice core is being drilled as part of EastGRIP. In the vicinity of the EastGRIP drill site (75.63 °N, 35.99 °W in 2019), we performed polarimetric measurements with a phase-sensitive radio echo sounder (pRES; Brennan et al. (2014); Nicholls et al. (2015)) in 2019: within the drill trench next to the core location (*CL*; ~ 10m apart) and at five sites (called *GRID*) within 20 × 20m² approximately 360m from the drill site (Fig. 1). These five sites are labelled depending on their cardinal direction (N, E, S and W) compared to the centre point (C). The pRES is a ground-based nadir looking Frequency-Modulated Continuous-Wave (FMCW) radar, which allows to determine vertical displacements of reflections within firn and ice from repeated measurements with a high accuracy of ~ 1 mm. While the pRES is mainly operated to derive basal melt rates (e.g. Marsh et al., 2016; Stewart



et al., 2019; Zeising et al., 2022), it also can be used to estimate the ice fabric by applying polarimetric pRES measurements
 55 (Brisbourne et al., 2019; Jordan et al., 2020; Ershadi et al., 2022).

A polarimetric pRES measurement consists of several measurements with different antenna orientations. The pRES transmits
 linearly polarised electromagnetic waves via the transmitting skeleton slot antenna and records the reflected signals in one
 direction with another antenna. This allows co-polarised measurements to be made in which the two antennas are oriented in
 the same direction. While in a hh measurement the direction of polarization points to each other, in a vv measurement it is
 60 perpendicular to the hh measurement (Fig. 1).

We aligned the antennas at an arbitrary azimuthal angle of roughly 258° (at CL) and 168° (at GRID) clockwise to magnetic
 North (283° and 193° true North), respectively. The ice flow direction at EastGRIP is roughly 58° magnetic North. We per-
 formed multi-polarised measurements by rotating each antenna separately horizontally clockwise in 22.5° steps up to 157.5° .
 Here, we only considered the co-polarised measurements taken roughly in ice flow direction (hh ; 55° to magnetic North)
 65 and perpendicular to ice flow (vv ; 145° to magnetic North). During each measurement, the pRES transmitted a sequence of
 chirps by linearly increasing the transmitted frequency from 200 to 400 MHz within 1 s for each chirp. In order to achieve
 a higher signal-to-noise ratio, the measurement at CL contained 250 chirps and those of the GRID contained 100 chirps per
 measurement.

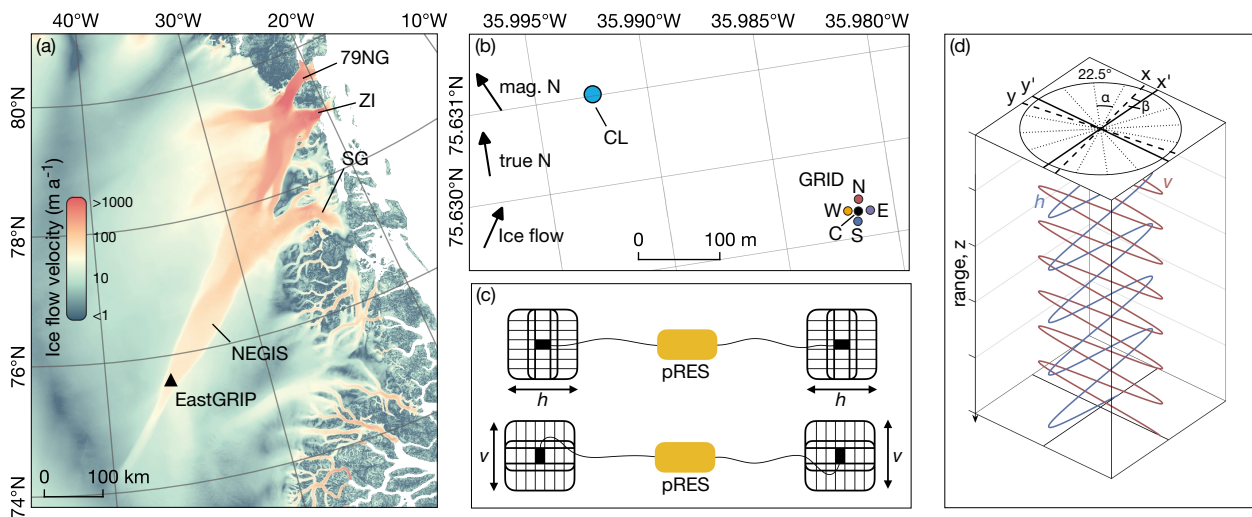


Figure 1. Location and orientation of polarimetric pRES measurements. (a) Surface ice flow velocity of the Greenland Ice Sheet (Joughin et al., 2016, 2018), showing the three major outlet glaciers of the Northeast Greenland Ice Stream (NEGIS): Nioghalvfjærdsbrae (79 N Glacier, 79NG), Zachariæ Isstrøm (ZI) and Storstrømmen Glacier (SG). The location of the EastGRIP drill site is denoted by the black triangle. (b) Location of polarimetric pRES measurements at CL and at GRID. Arrows shown direction of magnetic North, true North and ice flow direction. (c) Sketch of a polarimetric pRES measurement with hh and a vv antenna orientation. (d) Sketch of propagating waves with polarisations in x' (h) and y' (v) direction (solid line) in the x - y coordinate system (dashed line). Dotted lines show the (unused) multi-polarised measurements, separated by $\alpha = 22.5^\circ$. Ice flow is in x direction with an angular offset of β to the hh -measurement in x' direction.



3 Methods

70 3.1 Fabric anisotropy from ice core analysis

The fabric data was obtained from continuous 55 cm ice core sections which are basically vertical taken every 5–15 m of depth. Details of the sample preparation, data acquisition and processing are given in Stoll et al. (2021). The grain size weighted orientation of the c -axis can be represented by a second order orientation tensor. Its normalised eigenvalues

$$\lambda_1 + \lambda_2 + \lambda_3 = 1 \quad \text{and} \quad \lambda_1 \leq \lambda_2 \leq \lambda_3 \quad (1)$$

75 correspond to the length of the three principal axes. In order to derive the fabric asymmetry, we averaged those eigenvalues from all samples of each section and calculated the difference between the eigenvalues ($\lambda_2 - \lambda_1$ and $\lambda_3 - \lambda_1$).

3.2 Horizontal fabric asymmetry from radar measurements

If two electromagnetic waves, whose polarisation in x' and y' are perpendicular to each other, propagate through an anisotropic medium, their propagation velocities differ due to the horizontal dielectric anisotropy (Hargreaves, 1978). The resulting differ-
 80 ence in two-way travel-time $\Delta t_{x'y'}$ of a backscatter from a reflector at depth z is

$$\Delta t_{x'y'}(z) = t_{y'}(z) - t_{x'}(z) = \frac{2 \left(\sqrt{\bar{\varepsilon}_{y'}(z)} - \sqrt{\bar{\varepsilon}_{x'}(z)} \right)}{c_0} z, \quad (2)$$

where c_0 is the speed of light in vacuum, $t_{x'}$, $t_{y'}$ are the two-way travel times and

$$\bar{\varepsilon}_{x'}(z) = \frac{c_0^2}{4} \left(\frac{t_{x'}(z)}{z} \right)^2, \quad (3)$$

$$\bar{\varepsilon}_{y'}(z) = \frac{c_0^2}{4} \left(\frac{t_{x'}(z) + \Delta t_{x'y'}(z)}{z} \right)^2 \quad (4)$$

85 are the depth-averaged permittivities of the corresponding polarization directions x' , y' . In order to obtain the vertical profile of the horizontal dielectric anisotropy $\Delta \varepsilon_{x'y'} = \varepsilon_{y'} - \varepsilon_{x'}$ instead of the depth-average value, the vertical gradient of the travel-times has to be taken into account:

$$\varepsilon_{x'}(z) = \frac{c_0^2}{4} \left(\frac{dt_{x'}(z)}{dz} \right)^2, \quad (5)$$

$$\varepsilon_{y'}(z) = \frac{c_0^2}{4} \left(\frac{d(t_{x'}(z) + \Delta t_{x'y'}(z))}{dz} \right)^2. \quad (6)$$

90 According to Fujita et al. (2006) and Jordan et al. (2019), the horizontal dielectric anisotropy for the two polarisations in x' and y' direction, $\Delta \varepsilon_{x'y'}$, is a function of the horizontal fabric asymmetry $\Delta \lambda_{x'y'} = \lambda_{y'} - \lambda_{x'}$ and of the dielectric anisotropy of an ice crystal $\Delta \varepsilon'$:

$$\Delta \varepsilon_{x'y'}(z) = \varepsilon_{y'}(z) - \varepsilon_{x'}(z) = \Delta \varepsilon' (\lambda_{y'}(z) - \lambda_{x'}(z)) = \Delta \varepsilon' \Delta \lambda_{x'y'}(z). \quad (7)$$



Matsuoka et al. (1997) found $\Delta\varepsilon' = 0.034$ for ice-penetrating radar frequencies. Finally, the horizontal fabric asymmetry
 95 $\Delta\lambda_{x'y'}$ at depth z is given by

$$\Delta\lambda_{x'y'}(z) = \lambda_{y'}(z) - \lambda_{x'}(z) = \frac{\Delta\varepsilon_{x'y'}(z)}{\Delta\varepsilon'}. \quad (8)$$

Thus, the bulk dielectric anisotropy $\Delta\varepsilon_{x'y'}$, and based on this, the horizontal fabric asymmetry $\Delta\lambda_{x'y'}$ can be determined from
 the difference in the two-way travel time $\Delta t_{x'y'}$. The vertical resolution of $\Delta\lambda_{x'y'}$ depends on the precise determination of
 $\Delta t_{x'y'}$, which used to be a problem for previous radar systems that did not have the required resolution in the time domain. This
 100 is the main difference to and advantage of the in-depth analysis of the phase which is why polarimetric pRES measurements
 offer the chance to investigate the horizontal fabric asymmetry in the ice.

3.3 Phase-sensitive radar data analysis

For data processing, we followed Brennan et al. (2014) and Stewart et al. (2019) in order to get the complex valued signals
 s_{hh} and s_{vv} as a function of two-way travel time with the amplitude $|s_{hh}|$ and its phase. We convert s from time t to depth z
 105 domain by using dielectric permittivities derived from dielectric profiling (DEP) of the EastGRIP ice core by Mojtabavi et al.
 (2020).

The method we apply to compute the travel-time difference $\Delta t_{x'y'}$ is based on a cross-correlation of the co-polarised mea-
 surements. The same method is widely used to estimate vertical displacements for strain analysis from repeated pRES mea-
 surements as shown by e.g. Jenkins et al. (2006), Gillet-Chaulet et al. (2011) Stewart et al. (2019) and Zeising and Humbert
 110 (2021). We divided s_{hh} into segments of 12m depth overlapping by 9m and calculated for each the complex coherence

$$c_{hhvv}(z, l) = \frac{\sum_{j=z_n}^{z_n+N} s_{hh}(j) s_{vv}^*(j+l)}{\sqrt{\sum_{j=z_n}^{z_n+N} |s_{hh}(j)|^2} \sqrt{\sum_{j=z_n}^{z_n+N} |s_{vv}(j+l)|^2}}, \quad (9)$$

where z_n is the lower range bin index of the segment, N the number of bins in the segment, l the range-bin offset (lag) and *
 indicates the complex conjugate (Stewart et al., 2019). While the magnitude of the complex coherence $|c_{hhvv}|$ is the correlation
 between s_{hh} and s_{vv} , the argument is the coherence phase $\phi_{hhvv} = \arg(c_{hhvv})$ (Jordan et al., 2019).

115 Our *polarimetric cross-correlation* approach differs from the *coherence* method from Dall (2010) that was used by Jordan
 et al. (2019, 2020, 2022), Young et al. (2021a) and Ershadi et al. (2022). In their applications, the range-bin offset was set
 to zero ($l = 0$). Thus, these studies interpreted the *hhvv* coherence phase gradient of the same two-way travel-time. In our
 cross-correlation method, we shift s_{vv} by a number of integer bin offsets l (see eq. 9). This enables the possibility to analyse
 the travel-time difference of the same reflector. We identified the correct l of each segment by following the minimum phase
 120 difference from the surface downwards, indicated by high correlation values (Fig. 2b,c).

Next, we derived the travel-time difference $\Delta t_{x'y'}$ (Fig. 2d) from the sum of the range-bin offset and the corresponding
 phase shift

$$\Delta t_{x'y'}(z) = \frac{l(z)}{Bp} + \frac{\phi_{hhvv}(z)}{2\pi f_c} \quad (10)$$

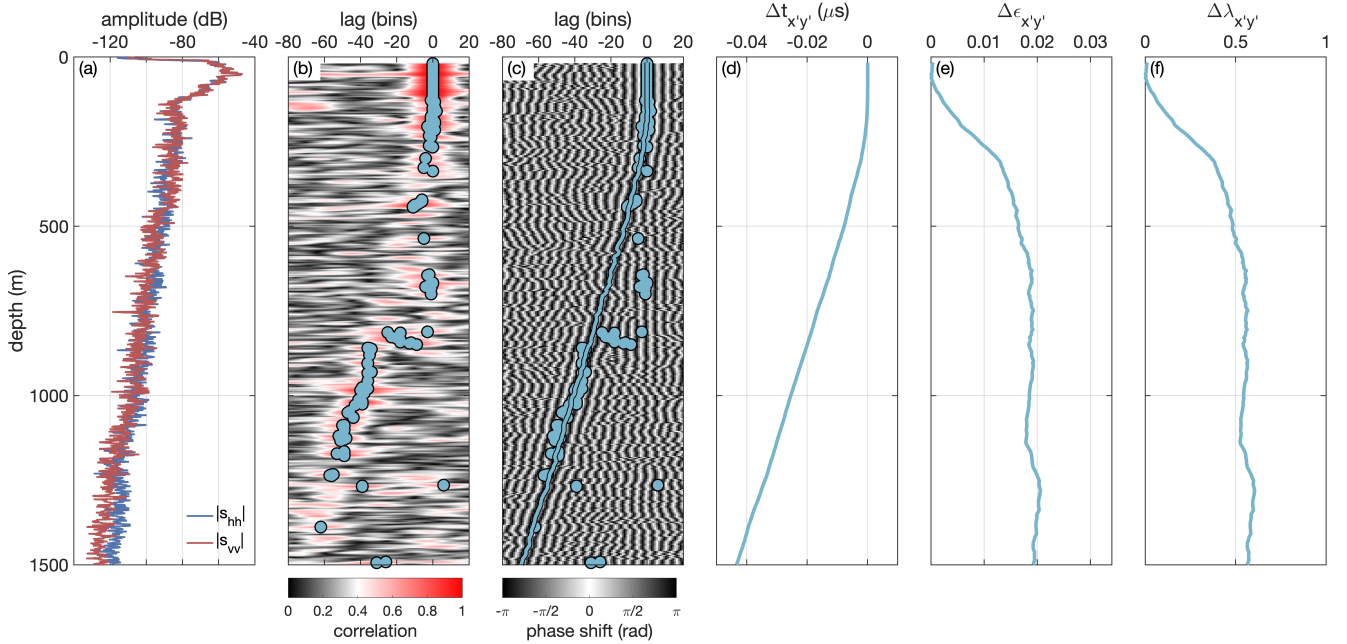


Figure 2. Analysis of the horizontal fabric asymmetry from polarimetric pRES measurements at the location CL (Fig. 1) next to the EastGRIP ice core. (a) Magnitude profiles of s_{hh} (blue line) and s_{vv} (red line) as a function of depth. (b) Cross-correlation $|c_{hhvv}|$ of s_{hh} and s_{vv} as a function of lag and depth. Blue dots mark the lag of best correlation for each segment exceeding a correlation of 0.65. (c) Coherence phase shift ϕ_{hhvv} as a function of lag and depth. The blue dots are the same as in (b). The blue line marks the tracked minimum phase shift. (d) Difference in two-way travel time between both measurements at the same depth after smoothing with a 100 m moving average filter. (e) Difference in horizontal dielectric anisotropy $\Delta\varepsilon_{x'y'}$. (f) Difference in horizontal eigenvalues $\Delta\lambda_{x'y'}$.

where $p = 8$ is a 'padding factor' that reduces the range-bin spacing, $B = 200$ MHz is the bandwidth and $f_c = 300$ MHz the centre frequency of the pRES measurement (see Brennan et al., 2014).

Since the travel-time difference is cumulative, we calculated the mean vertical gradient of the two-way travel times, $dt_{x'}/dz$ and $d(t_{x'} + \Delta t_{x'y'})/dz$, from a 200m moving window after smoothing $\Delta t_{x'y'}$ with a 100m moving average filter. Below 100m depth, we adopted the method to use a smaller, adaptive moving window that increases with depth. Finally, we compute the dielectric anisotropy $\Delta\varepsilon_{x'y'}$ from Eqs. 6 and 7 (Fig. 2e) and the horizontal fabric asymmetry $\Delta\lambda_{x'y'}$ from Eq. 8 (Fig. 2f).

130 4 Results

The horizontal fabric asymmetry from the polarimetric cross-correlation analysis at all measurement locations show the same vertical distribution with only minor differences (Fig. 3a). They indicate a rapid increase of $\Delta\lambda_{x'y'}$ from 0.06 to 0.4 between 125 and 320m depth, followed by a minor increase to 0.55, reached at a depth of 550m. Between this depth and 1400m, the



horizontal anisotropy remains at high level and varies between 0.52 and 0.62. Below the depth of 1400 m, a low signal-to-noise
 135 ratio prevented the analysis of the horizontal fabric asymmetry.

The pRES-derived vertical distribution matches the distribution of the difference of the weighted horizontal eigenvalues
 from the EastGRIP ice core analysis nearly perfectly (Fig. 3b). While the differences of the first two eigenvalues ($\lambda_2 - \lambda_1$)
 show the same rapid increase between 125 and 250 m depth, below, it is $\lambda_3 - \lambda_1$ that is of the same size than the pRES-derived
 values. This indicates that one of the horizontal eigenvalues becomes the largest value (λ_3 by definition) at a depth of 250 m
 140 and thus λ_3 switches from the vertical to one horizontal axis. However, since $\Delta\lambda_{x'y'}$ exceeds 0.5, it is obvious that a horizontal
 eigenvalue is the largest since the $\lambda_2 - \lambda_1$ is always ≤ 0.5 .

The mean deviation of the horizontal fabric asymmetry derived from polarimetric pRES measurements and the difference
 of the weighted horizontal eigenvalues from the ice core analysis ($\lambda_2 - \lambda_1$ between 120 and 250 m and $\lambda_3 - \lambda_1$ between 250
 and 1400 m) is only 0.03 and thus of the same size than the uncertainty of the ice core analysis. However, the mean deviation
 145 to the difference of the unweighted horizontal eigenvalue is higher, which is a result compatible to analyses of seismic waves
 by Kerch et al. (2018).

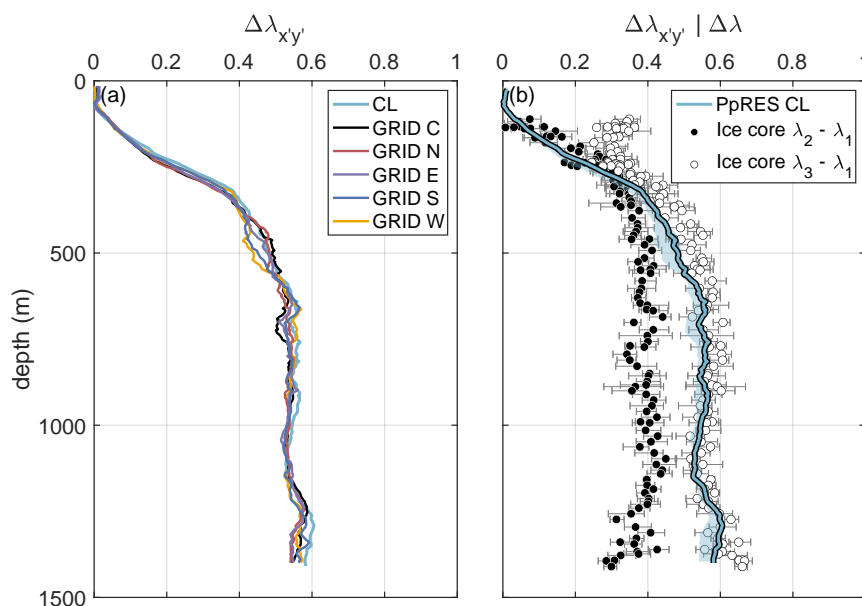


Figure 3. Horizontal fabric asymmetry $\Delta\lambda$ derived from polarimetric pRES measurements at CL (light blue line), compared with (a) analysed
 polarimetric pRES measurements at the 20×20 m GRID outside drill site (all analysed with our new approach) and (b) fabric asymmetry
 derived from the weighted horizontal eigenvalues from EastGRIP ice core (black and white dots). The blue shaded area in (b) marks the
 range of the polarimetric pRES-derived asymmetry from the measurements in the GRID and at CL.



5 Discussion

Our polarimetric cross-correlation method allows to resolve the travel-time difference of the co-polarised waves with sub-nanoseconds resolution. On this basis, the vertical distribution of the horizontal dielectric anisotropy as well as the bulk crystal fabric asymmetry can be determined. Despite the high range resolution, the uncertainty prevented a small-scale determination of the gradient of the travel-time difference. Thus, the derived horizontal anisotropy only represents a coarse distribution.

The determination of the horizontal asymmetry is not possible for every azimuthal angle. The azimuth angle of the antenna has to match the alignment of the orientation of the ice fabric principal axes sufficiently. If the direction of polarization is rotated 45° to the alignment of the principal axes, no anisotropy can be determined, as the propagation velocity is the same in x' and y' direction. The polarimetric pRES measurements at EastGRIP show that with an azimuthal rotation of the antennas with 22.5° increments up to 67.5° , a determination is possible in two of the four orientations and that the derived horizontal anisotropy is identical in both cases. In that sense, polarimetric pRES measurements allow to determine the fabric orientation within the rotation intervals (in this case 22.5°). However, Ershadi et al. (2022) presented a method which allows to perform measurements from only one azimuthal angle and to calculate the maximum horizontal asymmetry based on synthesizing the other angles.

Due to the ambiguities caused by phase wrapping, the previous methods which are based on the coherence phase gradient were limited to the derivation of phase shifts of a maximum of half a wavelength. This limits the depth to which fabric asymmetry can be determined: The more pronounced the asymmetry, the greater the difference in propagation velocity and thus the half-wavelength limit is reached earlier. At ice divides or domes with very little asymmetry, such as at NEEM (Jordan et al., 2019), WAIS divide (Young et al., 2021a) or EDC (Ershadi et al., 2022), the fabric asymmetry could successfully be determined using the previous coherence method up to the onset of noise. However, in fast moving areas like the Rutford Ice Stream, Antarctica (Jordan et al., 2022) or NEGIS, Greenland (this study) the previous method is limited to a few hundred meters below the surface. At EastGRIP, the phase shift at $l = 0$ exceeds half a wavelength at a depth of 250m, which significantly limits the application of the previous coherence method in contrast to our approach. Below 250m, the correlation drops at $l = 0$ (Fig. 2b), since the correlated reflections occurred from different reflectors that are more than half a wavelength apart. This may also be the reason for the low coherence magnitude in the analysis results from Jordan et al. (2022), which is why the evaluation of the pRES measurements at the Rutford Ice Stream, Antarctica was only possible down to a maximum depth of 300m. With our polarimetric cross-correlation method, which is based on phase unwrapping, this limitation can be overcome and phase shifts greater than half a wavelength can be evaluated. This allows the determination of even strong horizontal fabric asymmetries to a much greater depth.

Ershadi et al. (2022) presented a method to estimate horizontal ice fabric anisotropy based on a non-linear inverse approach and synthesized polarimetric radar data by using the coherence phase gradient and power anomaly. Here we tried to use this method on our data to compare the two methods directly. However, the ice fabric orientation in this area rotates several times at different depths of the ice column, which is the main limitation of the previous method using the inverse approach. Therefore,



180 the attempt for direct comparison has failed which is another proof that our method has improved the previous method and goes beyond its limits.

6 Conclusions

We presented a new method to infer the vertical profile of the horizontal fabric asymmetry from polarimetric phase-sensitive radar measurements. Our approach is based on a cross-correlation of co-polarised measurements to derive precisely travel-time
185 differences caused by dielectric anisotropy. In contrast to previous methods, this polarimetric cross-correlation approach allows to analyse even strong horizontal fabric asymmetries to a much greater depth.

The remarkable agreement between the vertical profile of the horizontal fabric asymmetry obtained by our analyses of multiple polarimetric pRES measurements and the fabric measured in the EastGRIP ice core demonstrates the robustness and precision of our method.

190 In the future, the applicability of our polarimetric cross-correlation method to other radar systems should be tested, in particular to polarimetric airborne radar measurements. If successful, this would increase the spatial coverage of mapped crystal fabric and its variability than would be possible with pointwise polarimetric pRES measurements. Such an application would significantly improve the understanding of the link between the stress state and crystal fabric evolution and allow to decrease uncertainties for response times of dynamically active glacial systems, like ice streams, to external perturbations.

195 *Code and data availability.* The MATLAB Code will be available before publication. Raw data of the polarimetric pRES measurements and the ice core-derived eigenvalues are submitted to the World Data Center PANGAEA.

Author contributions. OZ and AH designed the study and performed the radar measurements. OZ processed the data together with MRE and prepared the manuscript with contributions from all co-authors. OZ and TAG developed the method with support from OE. NS and IW prepared the ice core samples used for fabric analyses, performed the measurements, and processed and analysed the fabric data.

200 *Competing interests.* OE is an editor for TC but has not competing interests.

Acknowledgements. Data has been acquired at the EastGRIP camp that kindly hosted this activity as an associate project. EastGRIP is directed and organized by the Centre for Ice and Climate at the Niels Bohr Institute, University of Copenhagen. It is supported by funding agencies and institutions in Denmark (A. P. Møller Foundation, University of Copenhagen), USA (US National Science Foundation, Office of Polar Programs), Germany (Alfred Wegener Institute, Helmholtz Centre for Polar and Marine Research), Japan (National Institute of Polar
205 Research and Arctic Challenge for Sustainability), Norway (University of Bergen and Trond Mohn Foundation), Switzerland (Swiss Na-



tional Science Foundation), France (French Polar Institute Paul-Emile Victor, Institute for Geosciences and Environmental research), Canada (University of Manitoba) and China (Chinese Academy of Sciences and Beijing Normal University). Nicolas Stoll gratefully acknowledges funding from the Helmholtz Junior Research group “The effect of deformation mechanisms on ice sheet dynamics” (VH-NG-802).



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