



# 1 Simulated retreat of Jakobshavn Isbræ during the 21<sup>st</sup> 2 century

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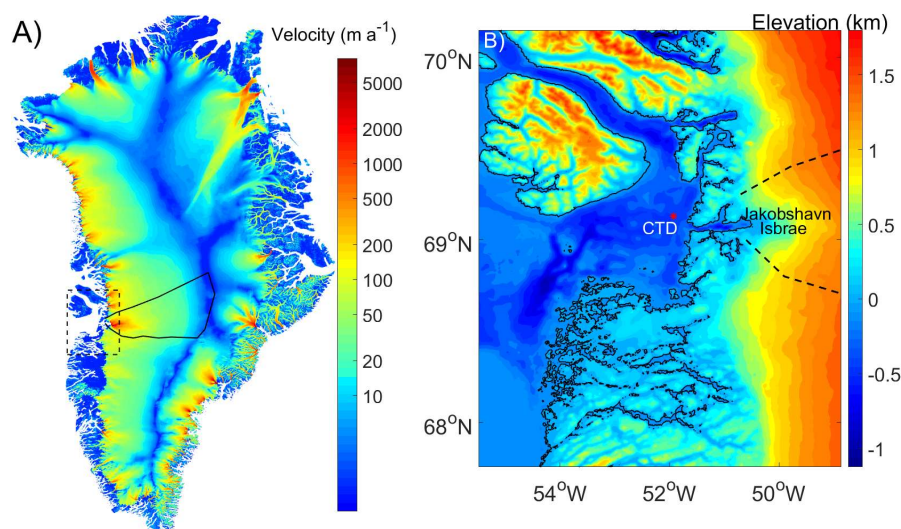
## 11 12 Abstract

13 The early in the 21<sup>st</sup> century retreat of Jakobshavn Isbræ, one of Greenland's largest outlet glaciers,  
14 into its over-deepened bedrock trough was accompanied by acceleration to unprecedented ice-  
15 stream speeds. Such dramatic changes suggested the possibility of substantial mass loss over the  
16 rest of this century. Using a three-dimensional ice-sheet model with parameterizations to represent  
17 the effects of ice mélange buttressing, crevasse-depth-based calving and submarine melting, we can  
18 reproduce its recent evolution. The model can accurately replicate its inter-annual variations in  
19 grounding line and terminus position, including new modes of seasonal fluctuations that emerged  
20 after arriving at the over-deepened basin and the disappearance of a persistent floating ice shelf. The  
21 shear margin induced decreases in ice viscosity we simulate are particularly important in  
22 reproducing the large observed inter-annual changes in terminus velocity. We use this model to  
23 project Jakobshavn's evolution over this century when forced by the IPCC RCP4.5 climate scenario  
24 and simulated by ocean temperatures from 7 Earth System Models along with surface runoff derived  
25 from RACMO. In our simulations, Jakobshavn's grounding line continues to retreat ~ 18.5 km by



the end of this century with total mass loss of  $\sim 2030$  Gt (5.6 mm sea-level-rise equivalent). Despite the relative success of the model in simulating the recent behavior of the glacier, the model does not simulate winter calving events that have become relatively more important.

## 1 Introduction



**Figure 1.** A) Greenland ice sheet flow speeds from Joughin et al. (2018), with the Jakobshavn drainage basin outlined by the solid black line and the area shown in panel B by the dashed box. B) Ilulissat Fjord and Disko Bay bathymetry from Jakobsson et al. (2012), with the CTD (Conductivity Temperature Depth) site used for ocean temperature here marked by the red star.

Jakobshavn Isbræ (Fig. 1) is Greenland's largest and fastest outlet glacier, with speeds of up to 17  $\text{km a}^{-1}$  (Joughin et al., 2014). Jakobshavn Isbræ drains  $\sim 6.5$  % of the Greenland Ice sheet (Krabbill et al., 2000), and it alone contributed  $\sim 1$  mm to global sea-level rise between 2000 and 2011 (Howat et al., 2011). Since 1997, measurements indicate that the water entering Ilulissat Fjord where Jakobshavn Isbræ terminates, is about  $1.1$  °C warmer than it was during 1987-1991 (Holland et al., 2008). This rise in water temperature coincided with the onset of dramatic thinning, speedup and



36 retreat of Jakobshavn Isbræ. By 2003 its velocity near the grounding line had reached  $\sim 12.6 \text{ km a}^{-1}$ ,  
37 <sup>1</sup>, more than double that of 1992, and the floating ice mélange in the fjord had disintegrated (Joughin  
38 et al., 2004). From 2005 to 2007, as it retreated inland, seasonal fluctuations in velocity 4 km inland  
39 from the calving front amounted to  $\pm 1 \text{ km a}^{-1}$ . The winter slowdowns and summer accelerations  
40 occurred in tandem with the calving front winter advance and summer retreat. By 2012 the seasonal  
41 velocity fluctuations 4 km upstream from the calving front was nearly  $\pm 8 \text{ km a}^{-1}$  and the grounding  
42 line of Jakobshavn Isbræ had reached the bottom of a sub-glacial bedrock trough after years of  
43 down-slope migration (Joughin et al., 2014).

44 Before 1997, Jakobshavn possessed a  $\sim 15 \text{ km}$  long floating ice mélange in front of its terminus ice  
45 cliff and experienced submarine melting on its ice-ocean interface (Amundson et al., 2010). After  
46 1998 the terminus became more crevassed, coinciding with acceleration of the glacier, implying that  
47 weakened buttressing had triggered its dramatic speed-up. A thinning rate of  $230 \pm 50 \text{ m a}^{-1}$  between  
48 the summers of 1984 and 1985 was deduced from photogrammetric surveys, mostly due to  
49 submarine melting (Motyka et al., 2011). The floating tongue thickened during the mid-1990s  
50 followed by progressive thinning after 1997 (Motyka et al., 2011). From 1997 to 2008, the average  
51 ocean temperature was  $1.1^\circ\text{C}$  higher than during the period 1980 – 1991, which raised its thinning  
52 rate substantially, affecting the whole ice mélange, and the ice shelf eventually collapsed in 2003.  
53 Many lines of evidence suggest that warm water was responsible for the submarine melting beneath  
54 the ice mélange and ice-shelf, brought by a buoyancy-driven, overturning circulation in Ilulissat  
55 fjord (Gladish et al., 2015).

56 Jakobshavn, in common with most outlet glaciers in Greenland, flows through a narrow, deeply  
57 incised bedrock trough at a far faster rate than the ice surrounding it (Joughin et al., 2010). Gravity  
58 surveys suggest a deep layer of soft till underlies much of the Jakobshavn trough (Block and Bell,



2011). This soft bed provides almost no resistance to ice flow and basal shear stress maps show that most of the gravitational driving force on the glacier is balanced by lateral drag (Shapiro et al., 2016).

Basal drag decreased from 1995 to 2006 (Habermann et al., 2013), possibly due to fast thinning that reduced the effective pressure, that is the ice overburden minus water pressure, at the bed. The effective pressure distribution under the glacier is important to basal drag and must be zero at the grounding line as it begins to float. Several sliding parameterizations (also termed sliding relations or sliding laws) have been used in the literature that assume basal drag depends on sliding speed (so-called Weertman sliding), or on effective pressure (Schoof, 2010; Gagliardini et al., 2014). Tsai et al. (2015) introduced a combined Weertman and Coulomb sliding law based on effective pressures with a boundary layer at the grounding line; this has a higher scaling of ice flux with grounding-line thickness compared with the Weertman. However, in the Jakobshavn case, both Weertman and Coulomb sliding produce very similar fluxes because the basal shear stresses along the main trough are typically only 2 % of the driving force.

Simulations using a flow-band model with a crevasse-depth-based calving parameterization (Viellet et al., 2011) demonstrated that loss of buttressing from the disintegration of its floating ice mélange or enhanced submarine melting could have triggered the dramatic changes seen in Jakobshavn Isbræ at the end of the 20th century. Later work (Muresan et al., 2016), using a simple calving model with dependence on the strain field at the terminus was able to reproduce the inter-annual retreat of Jakobshavn Isbræ until 2009, when the terminus arrived at the beginning of the reverse sloping bed. But retreat after 2010 was not captured by their model, and neither was the seasonal fluctuation in terminus position.

In this paper we use a three-dimensional ice-flow model with a treatment of calving that successfully



82 tracks the seasonal terminus position and its retreat into the over-deepened basin. We use historic  
83 observations of ocean temperature as forcing and ice tongue thinning rate to scale submarine melting  
84 rates for our model and thence make future projections. Our aim is to track the evolution of  
85 Jakobshavn Isbræ through the 21st century under a specific climate forcing scenario. In Section 2  
86 we describe the approach and calibration of our model, Section 3 shows the simulations for the  
87 period to 2100 under the IPCC RCP4.5 scenario (Moss et al., 2010), Section 4 is a discussion of our  
88 method with reference to other studies and suggestions for improvements, and we conclude in  
89 Section 5.

## 90 2 Methods and data

### 91 2.1 Ice sheet model

92 We model Jakobshavn Isbræ using the BISICLES continuum ice sheet dynamics model that is based  
93 on the vertically integrated stress balance formulation of Schoof and Hindmarsh (2010), which treats  
94 longitudinal and lateral stresses as depth-independent, but allows for vertical shear in the nonlinear  
95 rheology (Cornford et al., 2013). BISICLES is particularly useful for Jakobshavn Isbræ as it uses  
96 block-structured finite volume discretization with adaptive mesh refinement (Cornford et al., 2013)  
97 allowing for high resolution modeling of critical sections of the glacier. Jakobshavn Isbræ is fed by  
98 a ~ 400 km long and extensive drainage basin (Fig. 1), but the fast flow area is only around 10 km  
99 in width.

100 We assume Jakobshavn Isbræ to be in hydrostatic equilibrium, thus the upper surface elevation  $s$  is

$$101 \quad s = \max \left[ h + b, \left( 1 - \frac{\rho_i}{\rho_w} \right) h \right], \quad (1)$$



102 where  $\rho_i$  and  $\rho_w$  are the densities of ice and ocean water,  $h$  is ice thickness and  $b$  is bedrock  
 103 elevation relative to sea level. The ice thickness evolves in time as

$$104 \quad \frac{\partial h}{\partial t} + \nabla \cdot [\mathbf{u}h] = M_s - M_b, \quad (2)$$

105 where  $M_s$ ,  $M_b$  are surface mass balance (SMB) and submarine melt rate respectively and  $\mathbf{u}$  is the  
 106 depth-independent horizontal velocity. No basal melting over the grounded area is allowed. The  
 107 velocity  $\mathbf{u}$  satisfies an approximate stress balance equation

$$108 \quad \nabla \cdot [\phi h \bar{\mu} (2\dot{\epsilon} + 2\text{tr}(\dot{\epsilon})\mathbf{I})] - \boldsymbol{\tau}^b = \rho_i g h \nabla s, \quad (3)$$

109 where  $\mathbf{I}$  is the identity tensor,  $s$  is the ice surface elevation,  $g$  is the acceleration due to gravity,  $\dot{\epsilon}$  is  
 110 the horizontal strain-rate tensor defined by

$$111 \quad \dot{\epsilon} = \frac{1}{2} [\nabla \mathbf{u} + (\nabla \mathbf{u})^T], \quad (4)$$

112 and  $\boldsymbol{\tau}^b$  is the basal shear stress. The vertically integrated effective viscosity  $\phi h \bar{\mu}$  is given by

$$113 \quad \phi h \bar{\mu}(x, y) = \phi \int_{s-h}^s \mu(x, y, z) dz, \quad (5)$$

114 where the vertically varying effective viscosity  $\mu$  includes a contribution from vertical shear and  
 115 satisfies

$$116 \quad 2\mu A(T) (4\mu^2 \dot{\epsilon}^2 + |\rho_i g (s - z) \nabla s|^2)^{(n-1)/2} = 1, \quad (6)$$

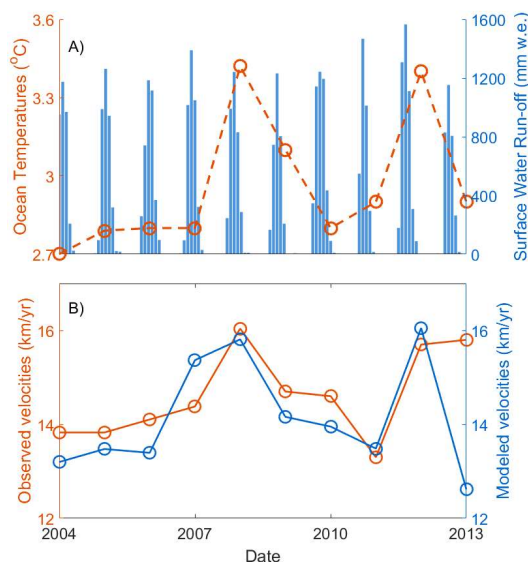
117 where  $n$  is the flow rate exponent, set to 3 in the current study, and  $A(T)$  is the rate factor, dependent  
 118 on the ice temperature  $T$  through an Arrhenius law (Cuffey and Paterson, 2010).  $\phi$  is a stiffening  
 119 factor estimated by solving an inverse problem (Cornford et al., 2015) using measured surface  
 120 velocities.



We use a viscous Weertman sliding relation to define the basal friction:

$$\tau^b = \begin{cases} -C|u|^{m-1}u & \text{if } \frac{\rho_l}{\rho_w}h > -b, \\ 0 & \text{otherwise} \end{cases}, \quad (7)$$

and here we assume a linear relation taking  $m=1$ . The basal traction coefficient  $C(x, y)$  is estimated simultaneously with the stiffening factor  $\phi$  by solving the inverse problem.  $C$  and  $\phi$  are adjusted iteratively to reduce the misfit with a set of 2010 surface velocity observations (Joughin et al. 2010). We hold the fields  $C$  and  $\phi$  constant over time throughout our simulations, although they must actually change as the glacier retreats.



128

Figure 2. A) Time series of ~300 m deep ocean temperature (red) from  
 http://ocean.ices.dk/HydChem/ near the mouth of Ilulissat fjord. Blue bars are simulated monthly  
 surface water run-off from the MAR regional surface mass and energy balance model (Alexander  
 et al. 2016). B) Measured ice front annual mean ice flow speeds (red) from Joughin et al. (2010),  
 compared with our modeled speeds (blue).



## 134 2.2 Climate Forcing

135 We use ocean temperatures collected from a CTD site close to the mouth of Ilulissat fjord (Fig. 1)  
 136 as an approximation of ocean temperatures near the glacier grounding line. A comprehensive study  
 137 focusing on ocean circulation within Ilulissat fjord validated this approximation (Gladish et al.  
 138 2015). A positive correlation ( $r=0.74$ ,  $p<0.05$ ) exists between deep ocean temperatures and flow  
 139 speed near the terminus of Jakobshavn Isbrae (Fig. 2) from 2004 onwards. There is no significant  
 140 correlation prior to 2004, the floating ice tongue period. As a working hypothesis we assume that  
 141 the correlation since 2004 reflects the effects of the sea ice and iceberg mélange in the fjord on the  
 142 flow speed near the terminus: a warmer ocean reduces mélange thickness and therefore buttressing.  
 143 There appears to be no lag between the glacier acceleration and change in deep ocean temperature,  
 144 suggesting mélange response times are faster than 1 year. When the floating ice tongue was present  
 145 lags in the system were likely longer, accounting for the lack of correlation between ocean  
 146 temperatures and glacier flow speed prior to 2004. It is also possible that ocean temperatures reflect  
 147 changes in surface runoff and basal lubrication for sliding, but we consider that the runoff more  
 148 strongly affects calving mechanisms as discussed later. We therefore modify the driving force (Eq.  
 149 3) on the grid cells next to the calving front by multiplying by a factor  $\alpha$  (tuned value shown by  
 150 Eq. 9 that is linearly related to ocean temperature ( $T$ ) as a means of representing the buttressing  
 151 effects of the ice mélange in the fjord.

$$152 \nabla \cdot [\phi h \bar{\mu} (2\dot{\epsilon} + 2\text{tr}(\dot{\epsilon})\mathbf{I})] + \tau^b = \alpha \cdot \rho_i g h \nabla s, \quad (8)$$

$$153 \alpha = 0.82 + 0.111 \cdot T, \quad (9)$$

154 We use a crevasse based calving parameterization (Benn et al., 2007; Nick et al., 2011) that calves  
 155 ice where the crevasse penetration depth ( $D_s$ ) is greater than upper surface elevation.  $D_s$  is defined





156 as

$$157 \quad D_s = \frac{S}{g \cdot \rho_i} + \frac{\rho_w}{\rho_i} \cdot \bullet \cdot \beta, \quad (10)$$

158 where  $S$  is the magnitude of extensional stress,  $\bullet$  is surface water run-off, and  $\beta$  is a tuning scalar.

159 We estimate runoff from the 25 km resolution regional climate model, MAR, (Alexander et al. 2016),

160 driven by the ERA-Interim reanalysis (Dec et al., 2011).

161 We characterize submarine melting as a linear function of ocean forcing

$$162 \quad M_b = \gamma T_f, \quad (11)$$

163 where  $T_f$  is the far field ocean forcing temperature, taken in Disko Bay (CTD in Fig. 1), relative to

164 pressure melting temperature under the ice shelf. We derive  $\gamma$  from the 1985 observed submarine

165 melt rate of  $1 \pm 0.2$  m day<sup>-1</sup> beneath the floating ice mélange of Jakobshavn Isbræ, when Disko Bay

166 ocean temperatures were 4.2°C warmer than the pressure melting point at the bottom of the floating

167 ice shelf (Motyka et al. 2011). We test the sensitivity of the modeled glacier to uncertainty in

168 submarine melt rate in section 2.4.

169 We force Jakobshavn Isbræ in the 21st century using SMB and run-off from the 11 km resolution

170 RACMO model (Van Angelen et al., 2013) driven by the RCP4.5 scenario (Moss et al. 2010). The

171 run-off values are averaged over the nine grid points nearest to the terminus of Jakobshavn (69.1°N,

172 50.0°W). The RACMO simulation was forced by the HadGEM2-ES Earth system model (Collins

173 et al., 2011), as this climate model was found to be the most realistic for present-day simulations of

174 the Greenland ice sheet (Van Angelen et al., 2013). Ocean forcing should relate to temperatures off

175 the continental shelf close to the fjord mouth. Cowton et al. (2018) achieved success in simulating

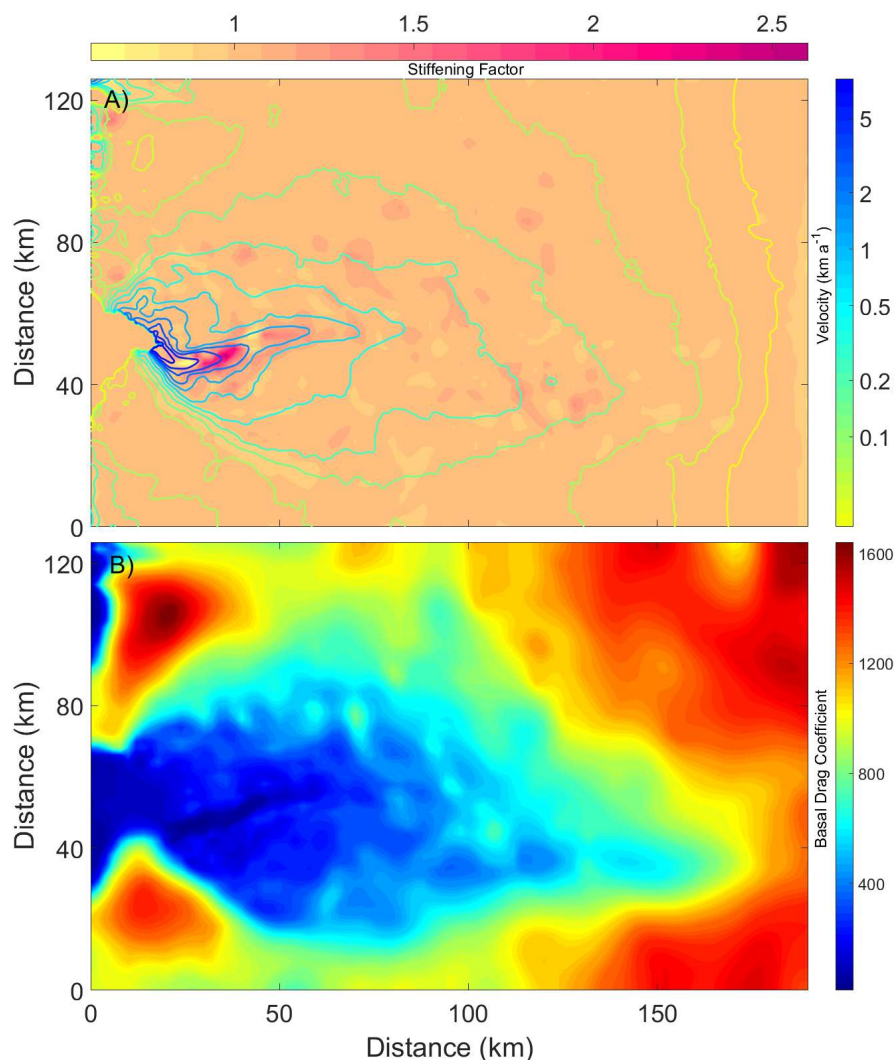
176 the terminus position and yearly variability of 10 glaciers along the east coast of Greenland using



mean 200–400 m depth temperatures from reanalysis data. For consistency with the RACMO results, we use deep ocean temperatures at  $\sim 300$  m depth from the  $0.83^\circ \times 1^\circ$  resolution HadGEM2-ES driven by the RCP 4.5 climate scenario from 2005 to 2100 at the 3 closest grids point to Disko Bay. We also compare this with results from 7 other climate model simulations of RCP4.5: HadGEM2-ES (Collin et al., 2011), BNU-ESM (Ji et al., 2014), MIROC-ESM (Watanabe et al., 2011), IPSL-CM5A-LR (Dufresne et al., 2013), CSIRO-Mk3L-1-2 (Gordon et al., 2002), NorESM1-M (Bentsen et al., 2012) and MPI-ESM-LR (Giorgetta et al., 2013).

## 2.3 Initialization Procedure

As we are interested in high resolution simulations and validating our model parameterizations with observations over the last 2 decades, we take care to initialize the model as accurately as possible. Detailed bedrock topography and ice thickness data in the year 2009 come from Gogineni et al. (2012). In 2004 the floating ice shelf disintegrated, making it a convenient starting point for simulations since we might expect the system to respond differently to forcing when there was a floating ice shelf compared with the situation of ocean forcing along a near-vertical ice cliff. This is consistent with the observed good correlation between ocean temperature and flow speed after 2004 but not before. The aim of this initialization was provide a state rather similar to 2004, that is barely retreating on inter-annual scales (Joughin et al., 2010) and small changes of annual mean velocity in the following 3 years. Therefore



195

196 **Figure 3. (A) Stiffening factor  $\Phi$  (Eq. 3) and (B) basal traction coefficient  $C$  (Eq. 7) over the**  
 197 **computational domain from solving the inverse problem. Contour lines in panel A show the**  
 198 **modeled velocity (logarithmic scale).**

199 1) We solved the inverse problem for basal conditions (Eq. 7) and stiffening factor using 2010  
 200 velocities (Joughin et al., 2010) and 2009 geometry (Gogineni et al., 2012). Our friction



201 coefficient field is shown in Fig. 3.

202 2) We made an initial guess for  $\beta$  (Eq. 10) based on short (less than one decade) transient  
 203 simulations incorporating calving but not SMB or ice shelf melt forcing, starting from our  
 204 inverted model state from step 1 above. This was forced by seasonal repetitions of runoff  
 205 from MAR for the year 2010.

206 3) Starting from the inversion of step 1 and using  $\beta$  (Eq. 10) from step 2, we let the model  
 207 glacier evolve freely without calving and with zero SMB and with sub-shelf melting forced  
 208 by repeating the observed 2004 ocean temperature for 11 years (that means the coefficient  
 209  $\gamma$  in Eq. 11 was set to be 1) until its surface elevation profile was similar to the known  
 210 profile of 1998 (Bamber et al. 2001).

211 4) We carried out several 10-year simulations with  $\beta$  gradually decreased from its initial value  
 212 given by step 2, to 45% of this value. These simulations were forced by repeatedly applying  
 213 the 2004 seasonal climate forcing. From these, we selected the  $\beta$  that best allowed  
 214 Jakobshavn Isbræ to reach a stable state by the 8th year (changes in the 9<sup>th</sup> and 10<sup>th</sup> years  
 215 of simulations were negligible) such that its calving front position closest to that observed  
 216 in 2004. The best  $\beta$  here is 53% of its value from step 2. This is our best guess for the 2004  
 217 state, but there are no thickness data available for 2004. Notice that the front positions and  
 218 May front velocities from 2004-2006 are stable (Figs 2 and 3), suggesting that the glacier  
 219 was reasonably close to steady state. This also makes 2004 a good time from which to start  
 220 transient simulations.

221 Basal friction coefficient values downstream of the 2010 grounding line were set equal to that  
 222 in the nearest 2010 grounded location. This was necessary because steps 2, 3 and 4 involved



223 grounding line advance beyond the region for which basal friction coefficients had been inferred.  
 224 The geometry after this spin up procedure, and the friction coefficient and stiffening factor  
 225 distribution from the inversion in step 1 were used as the initial condition for model calibration.

226

## 227 2.4 Model calibration

228 There are three essential parameters in the model,  $\alpha$ ,  $\beta$  and  $\gamma$  representing mélange buttressing,  
 229 crevasse depth sensitivity to surface runoff, and shelf melt sensitivity to ocean temperatures. In the  
 230 initialization, we fix  $\gamma$  to be 1. Therefore, we performed a suite of about 50 simulations to tune the  
 231 parameters  $\alpha$  and  $\beta$  with fixed  $\gamma=1$ . Our target was to best reproduce Jakobshavn Isbræ's calving  
 232 front position and surface velocity evolution for the 10 year period 2004-2013. Reproducing the  
 233 total distance of retreat and the temporary stable state after 2012 were secondary desirable features  
 234 to match.

235 Because only within a small range of  $\beta$  will modeled retreats make sense, firstly we estimate  
 236 reasonable range of  $\beta$  when hold  $\alpha=1.0$ ,  $\gamma=1.0$ . Secondly, we explore the parameter space centered  
 237 by ( $\alpha=1.0$ ,  $\beta=0.06$ ), which come from estimations above, to match observed retreats and general  
 238 velocities neglecting the inter-annual variations. The parameter space tested here is  $(1.0 \pm 0.25$ ,  
 239  $0.06 \pm 0.01)$ . As the discussion above, we further assume  $\alpha = \alpha_1 + \alpha_2 T$ , i.e., linearly related to deep  
 240 ocean temperature. With velocity depending on ocean temperature, degree of freedom of our  
 241 parameter space grow to three, which are  $\alpha_1$ ,  $\alpha_2$ ,  $\beta$ . However, we find  $\beta$  and  $\alpha_2$  behave quite  
 242 independently within the small  $\beta$  range estimated so far, which allow us finally reach a set of  
 243 parameters that can accurately reproduce both the total retreats and the velocity variations including  
 244 inter-annual fluctuation. The tuning was implemented manually. The best set of parameters are

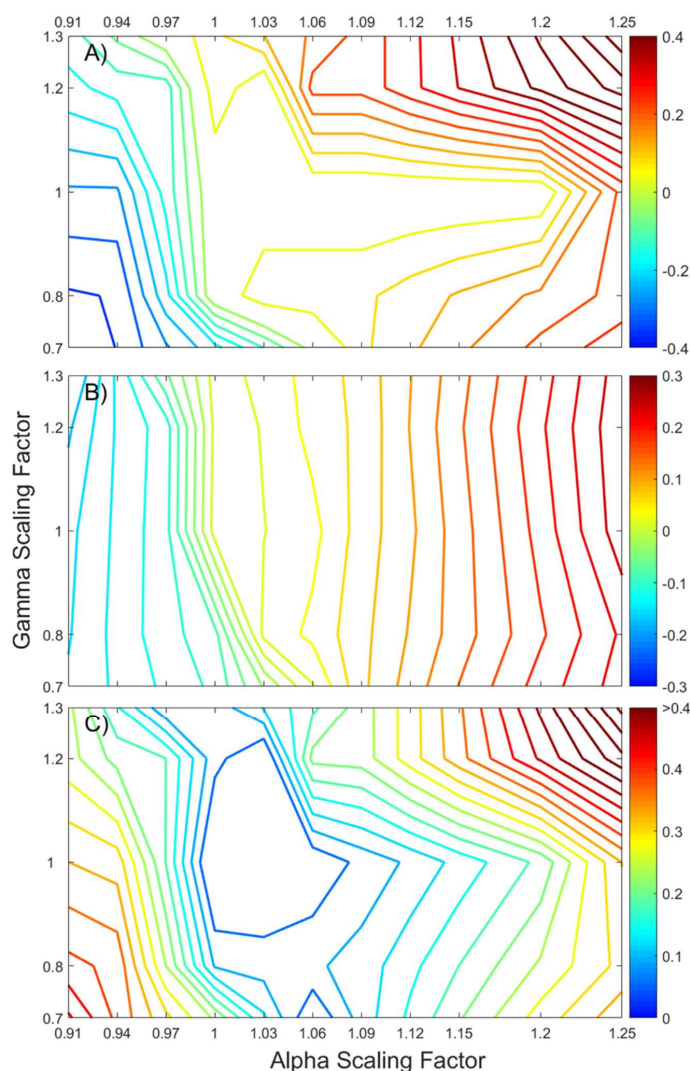


245  $\alpha_1=0.82$ ,  $\alpha_2=0.111$  and  $\beta=0.0638$ , as shown in Eq. 9.

246 We explore the glacier's sensitivity to two types of boundary perturbations. They are ice mélange  
 247 buttressing effect (defined by  $\alpha$ ) and submarine melting (defined by  $\gamma$ ). We scaled submarine melt  
 248 rates by multiplying  $\gamma$  by values from 0.8-1.2, based on the range of the observation uncertainty in  
 249 melt of  $\sim 20\%$  (Motyka et al. 2011). Also we varied  $\alpha$  by multiplying by factors from 0.91 to 1.25  
 250 to represent different buttressing strengths (Eq. 8). These multiplication factors were varied  
 251 systematically with typical intervals of 0.1 and 0.03 respectively for the  $\gamma$  and  $\alpha$  factors. We  
 252 calculated the following relative mismatches defined as (model-observations)/observations for each  
 253 simulation (shown in Fig 4):

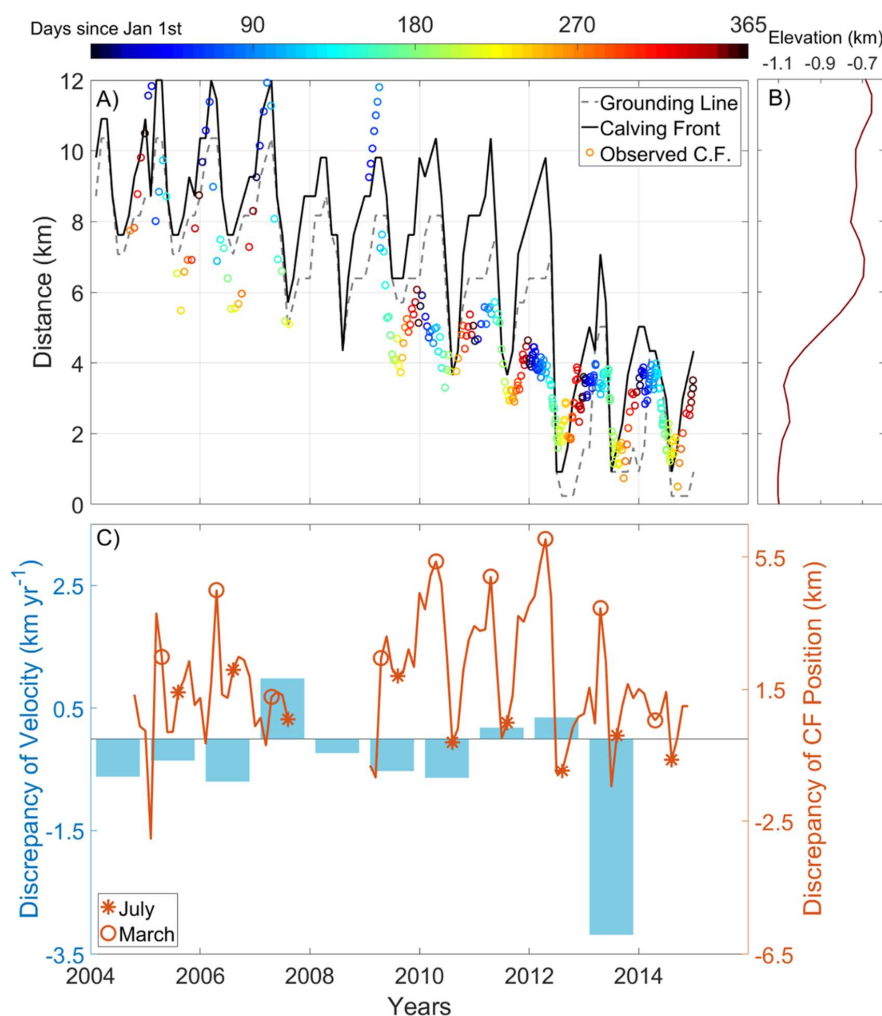
- 254 1. Total calving front retreat from 2004-2013
- 255 2. Annual mean front velocities
- 256 3. Vector sum of 1) and 2)

257 We used  $\beta$  (Eq. 10) from our optimal set of parameters. Our optimal value for  $\alpha$  is such that a  
 258 20% rise of its value does not affect modeled retreat when  $\beta$  and  $\gamma$  are kept to be their optimal  
 259 values (Fig. 4 A).



260

261 **Figure 4. Relative mismatches defined as (model-observed)/observed for A) total calving front**  
 262 **retreat, B) average of annual mean front velocity during 2004-2013, C) the vector sum of**  
 263 **mismatches in panels A and B,  $\sqrt{A^2+B^2}$  in our 2-D parameter space ( $\alpha, \gamma$ ) centered by the best**  
 264 **set ( $\alpha * 1.0, \gamma * 1.0$ ). X- and y-axis are multiplier of  $\alpha$  and  $\gamma$  used respectively in different**  
 265 **runs.**



266

267 **Figure 5. (A) Modeled retreat of the calving front (black solid line), grounding line (gray**  
 268 **dashed line), and observed calving front positions (color-coded circles and scale bar) from**  
 269 **Joughin et al. (2014). (B) Bedrock elevations. (C) Residuals (modeled minus observed) of**  
 270 **annual mean front velocity (blue bars, left axis) and of calving front position (red lines, right**  
 271 **axis) with typical timings of annual maximum (March) and minimum (July) extent marked.**  
 272 **The modeled front velocities and calving positions explain about 49% and 76% of the variance**





273 **in corresponding observations.**

274 The two biggest mismatches occur with the 2007 and especially 2013 velocities (Fig. 5). 2013 has  
275 the lowest simulated surface water run-off (Fig. 2) of all the years since 2004. The Benn calving  
276 model we use is sensitive to runoff, with reduced run-off leading to lower crevasse-penetration-  
277 depth and reduced terminus fracturing thus increasing its buttressing force. In the 2012/13 winter,  
278 the modeled glacier had an unprecedentedly high calving front and had been flowing fast the  
279 previous summer. This led to growth of an unusually long seasonal ice-shelf in the winter which  
280 caused low velocities near the end of front advancing season, and so accordingly low annual mean  
281 velocity. The low modeled velocity in 2009 can also be interpreted by the same over-growth effect,  
282 even though the ocean temperatures were high in 2009. Jakobshavn Isbræ did not in fact slow down  
283 very much in 2009 and 2013 because there were calving events that are unrepresented in our model.  
284 The relevant mechanisms are discussed later.

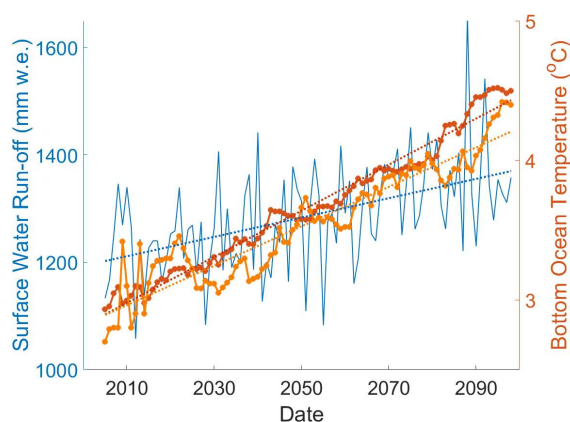
285 In 2007 high run-off caused more simulated calving and retreat than in reality. These retreat phases  
286 reduced the buttressing and lateral drag due to shear-margin-weakening, all of which lead to  
287 excessive speed-up near the terminus.

288 Modeled calving front retreat is ~ 7 km in total from 2004-2014 (Fig. 5), which is consistent with  
289 observations (Joughin et al. 2014). In 2009 a dramatic retreat brought the grounding line to the  
290 bottom of the bedrock slope, and since then it has gradually retreated with smaller seasonal  
291 fluctuations. The run-off forcing we applied triggered major retreats in the summers of 2007 and  
292 2012, due to large summer peak run-off (Fig. 2), demonstrating the sensitivity of our calving  
293 parameterization to run-off forcing. Modeled timings of maximum extent and minimum extent each  
294 year are in good agreement with observations, also demonstrating that summer, in particular, May  
295 to July, run-off determines much of the behavior of Jakobshavn Isbræ.

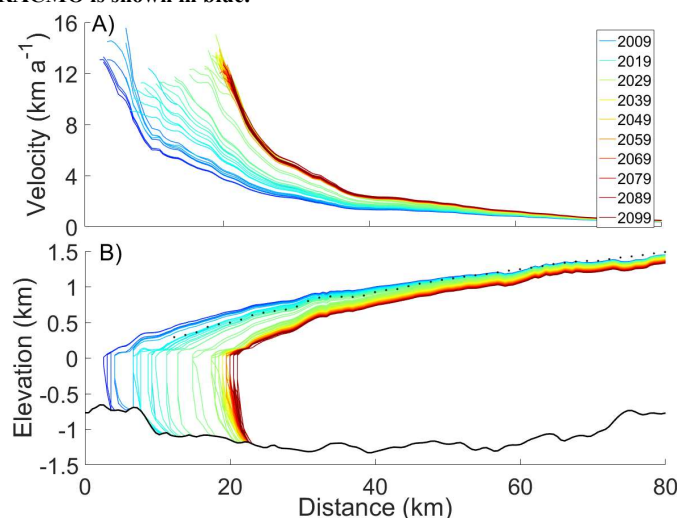


296 The modeled range of seasonal fluctuation in front position is  $\sim 5$  km, which is similar to  
 297 observations in the period before 2008. From January 2009 to December 2011, there was an abrupt  
 298 decrease in seasonal front fluctuation, with many winter calving events occurring, in contrast with  
 299 previous years (Cassotto et al. 2015). These winter calving events may explain the small observed  
 300 seasonal fluctuations because they limit the winter advance. Our model is unable to stimulate these  
 301 winter calving because there is no winter run-off, and hence calving is zero then. The largest  
 302 discrepancy of front position occurs during these winter calving periods (Fig. 5). Observations also  
 303 showed that from 2013 to 2017, Jakobshavn Isbræ barely retreated (Joughin et al. 2010). The decline  
 304 of run-off (Fig. 2) in 2014 suggests the reason. But since no RACMO run-off simulations are yet  
 305 available for 2015 and later, our parameterizations cannot be tested against this lack of retreat.

### 306 3 Future evolution



307 **Figure 6. Climate forcing for future projection under the RCP4.5 scenario taken as 300 m**  
 308 **depth ocean temperatures from HadGEM2-ES (orange) compared with the ensemble mean**  
 309 **(red) of 7 Earth System Models (HadGEM2-ES, BNU-ESM, MIROC-ESM, IPSL-CM5A-LR,**  
 310 **CSIRO-Mk3L-1-2, NorESM1-M and MPI-ESM-LR), (right axis), with their linear trends.**



**Figure 7. Modeled profiles of (A) January velocity and (B) January surface elevation along the center-flow-line of Jakobshavn Isbrae from 2004 to 2099 for the RCP4.5 scenario. Bedrock elevation is shown in black. Black dotted line is the surface elevation profile extracted from radar data measured around 2010 (Gogineni et al., 2012). Profiles are shown at intervals of 1 years. Profiles are color-coded in the legend and range from blue to green and red.**

Under the RCP4.5 scenario (Fig. 6) surface runoff slowly rises over the 21st century, with RACMO simulating slightly greater runoff during the second half than for the first 50 years. Runoff increases by 14% over the century. Bottom ocean temperature in the grid cell closest to Jakobshavn increases by 52%, and, as may be expected, has less variability than runoff.

Under this forcing, Jakobshavn Isbrae continues its retreat (Fig. 7) for 18 years after 2013, producing a total grounding line retreat of ~18 km upstream. As calving produces a steepening surface profile, terminus velocities increase, to reach a 21st century peak of ~19 km a<sup>-1</sup> in 2031 summer. Eventually the calving front becomes higher than the crevasse penetration depth in the calving parameterization.



326 This leads to a stable period with little inter-annual retreat and which lasts until the end of this  
327 century. During this period, nearly all of the seasonal retreats are offset by the following winter re-  
328 advances. Mass transport continually flattens and thins the ice geometry, leading to reduced flow  
329 speeds that eventually become half those of 2031, the 21<sup>st</sup> century peak.

330 The surprisingly high run-off anomaly in 2088 (Fig. 6) does not affect the stable state indicating  
331 run-off fluctuation alone cannot break this retreat pattern immediately. Once the inter-annual retreats  
332 cease in 2031, the dynamic thinning rate is greatly reduced because calving front height stops  
333 increasing.

334 **Table 1 Estimates of glacier mass loss and grounding line retreat from different sources.**

Source	Climate scenario	Mass loss 2004-2013 (10 years) (Gt)	Mass loss by 2100 (Gt)	Grounding line retreat 2004-2013 (km)	Grounding line retreat by 2100 (km)
This paper	RCP4.5	234	2068 (2044-2723)	7.0	18.5 (17.5-23.0)
Muresan et al. (2016)		220			
Nick et al. (2011)	A1B		1870 - 2281		14.0 - 26.0
Observations		225 ± 15		7.0	

335 Table 1 shows estimates of glacier mass loss and retreat. Under RCP4.5, total cumulative mass  
336 change of Jakobshavn Isbræ is 2029 Gt by 2100, using best set of  $\alpha$ ,  $\beta$  and  $\gamma$  with ocean temperature  
337 inputs from ensemble mean of 7 ESMs (Fig. 6). To estimate an upper bound for mass loss over this  
338 century, we scale the  $\alpha$  parameter by 1.2 giving 2680 Gt for the same forcing. Using the HadGEM2-  
339 ES forcing, which is the same model used to force RACMO with  $\alpha$  and  $\gamma$  set to their best estimates  
340 (Fig. 4) gives 2000 Gt. We suggest that this may be the lower reasonable bound of mass loss since  
341 the HadGEM-ES ocean temperatures rise notably slower than the ensemble mean.

342 Exploring the ( $\alpha$ ,  $\gamma$ ) parameter space we notice that values of (1.0, 0.8) produces a mass loss over  
343 this century of 2021 Gt with the HadGEM-ES ocean forcing, almost the same value as for the best



344 set of parameters. This implies that less submarine melting (determined by  $\gamma$ ) leads to larger ice loss  
345 by dynamic processes. The reason is that lesser submarine melt allows a larger ice thickness at the  
346 grounding line with stronger dynamic thinning in advancing season. Notice in our stress balance  
347 equation (Eq. 3), thickness contributes to driving force term, thus ice flux across the grounding line  
348 is highly nonlinear in ice thickness. This highly nonlinear relationship is also shown in our  
349 sensitivity tests (Fig. 4). Over the mismatch field measured by front velocity (Fig. 4, Panel B), the  
350 velocity is partly dominated by low values of  $\gamma$  around the line  $\alpha = 1.06$ , while  $\alpha$  is almost the only  
351 control on velocity over the region where  $\alpha < 1.09$ . Within our sample space, the non-linear and non-  
352 monotonic relationship between submarine melting and retreats is clear (Fig. 4, Panel A). Around  
353 the point ( $\alpha = 1.12$ ,  $\gamma = 1.0$ ), total retreat will increase no matter if  $\gamma$  is decreasing or increasing  
354 within the range  $0.8 < \gamma < 1.2$ . The area  $\alpha > 1.0$  in sample space is the very likely future condition  
355 for Jakobshavn Isbræ because increasing terminal ice cliff height caused by retreating into deep  
356 water will act as an amplifier to frontal driving force.

## 357 4 Discussion

### 358 4.1 Parameterization of Buttressing effect

359 The sudden  $1.1^{\circ}\text{C}$  rise in temperature of water entering Ilulissat fjord in 1997 (Holland et al., 2008)  
360 initiated rapid melting and disintegration of the floating ice mélange in 2003. This disintegration  
361 coincided with a near doubling of ice velocities. Modeling (Vieli et al., 2011) suggested that this  
362 was due to the reduction in buttressing from the floating ice-mélange. We can realistically reproduce  
363 the velocity variation of Jakobshavn Isbræ on seasonal and inter-annual scales using our  
364 parameterization of the buttressing effect from the ice mélange in the fjord.



365 Gladish et al. (2015) analyzed glacial flow speeds from 1998 to 2014, finding no correlation with  
366 Ilulissat fjord temperatures. This is because at the beginning of 2004, Jakobshavn's evolution entered  
367 a new phase with the disintegration of the ice mélange and floating ice shelf. We find good  
368 correlations between Disko Bay temperatures and ice velocities from 2004 to 2014. The  
369 improvement in correlation with temperatures may be explained by a faster response between the  
370 grounded glacier and the fjord water temperatures after loss of the floating ice shelf. Thus only  
371 freshly calved icebergs played roles in providing terminus resistance, and these could be reasonably  
372 supposed to react to seasonal fjord temperatures very quickly.

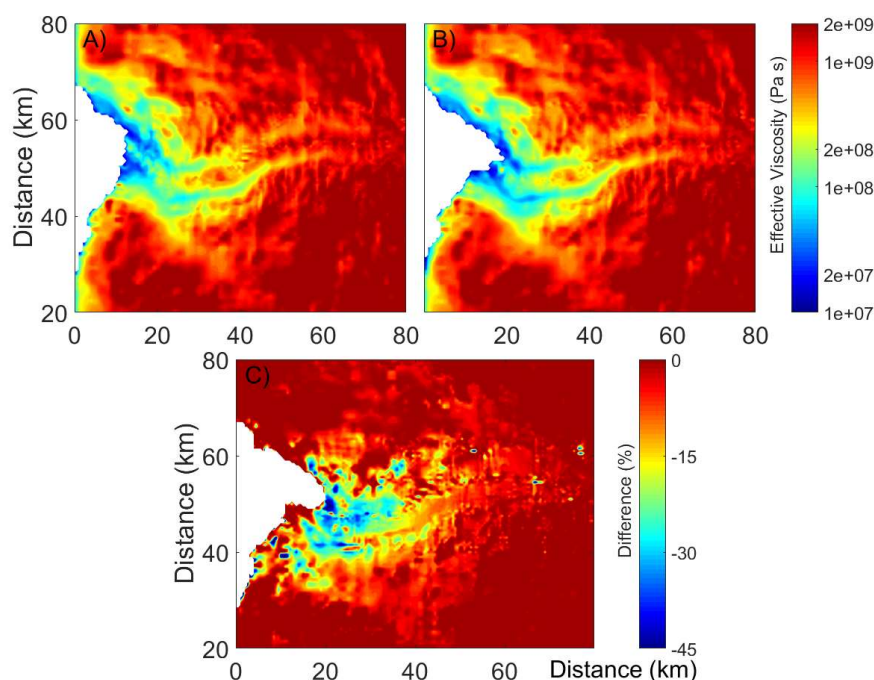
373 Buttrressing would affect the calving process by altering the longitudinal resistive stress in the glacier.  
374 Temperatures in Ilulissat Fjord will be warmer during the 21<sup>st</sup> century under essentially all climate  
375 scenarios, even those with modest emissions, due to the thermal inertia of the oceans. Thus a new  
376 floating ice shelf is unlikely to form. Prior to 2004, there were large changes in Jakobshavn: loss of  
377 ~15 km long stiff ice mélange and the sudden rise in fjord temperatures in 1998. There are fewer  
378 mechanisms to effect such dramatic changes in the future now that almost the entirety of the glacier  
379 is grounded. We therefore propose that our representation of the mélange buttressing mechanism,  
380 tuned for 2004-2013, is likely to maintain its validity during the 21<sup>st</sup> century.

## 381 **4.2 Horizontal shearing and viscosity**

382 Van Der Veen et al. (2011) estimated a maximum horizontal shear stress of ~800 kPa across the  
383 shear margin of Jakobshavn Isbræ where the horizontal velocity shear reaches the peak, while the  
384 bed stress is only 10-40 kPa in fast flowing regions (Shapiro et al., 2016). Given that the width of  
385 the Jakobshavn Isbræ fast flow region is typically under 5 km and its thickness is typically between  
386 1-2 km, these numbers indicate that the shear margins provide at least an order of magnitude greater  
387 total resistance than the bed. Thus, the shear margin, rather than the bed of Jakobshavn Isbræ



388 provides most of the resistance balancing the driving force. The main trunk of Jakobshavn Isbræ  
 389 exhibits considerable seasonal velocity changes, while the slow moving ice outside the shear margin  
 390 has little or no seasonal cycle. This flow structure implies speed gradients perpendicular to the flow  
 391 direction with large seasonal variation. These velocity shears would in turn generate large seasonal  
 392 variations in effective ice viscosity (Eq. 6). This mechanism implies a positive feedback on velocity  
 393 in the fast flow region: increases in the speed of fast flowing ice cause increases in horizontal shear  
 394 stress across the margins, reduced viscosity, and further increased horizontal velocity shear,  
 395 allowing further increase to speeds in the fast flow region. Observations show that, as the terminus  
 396 retreated into deeper water, seasonal fluctuations in terminus velocity increased (Joughin et al. 2008).  
 397 By 2012, the summer time peak terminus velocity was  $\sim 17 \text{ km a}^{-1}$ , more than twice the wintertime  
 398 minimum velocity (Joughin et al. 2014). This amplified seasonal velocity cycle was likely enhanced  
 399 by the shear-margin weakening mechanism.





401 **Figure 8. Modeled annual mean of vertically averaged effective viscosity  $\Phi\mu$  (Eq. 5) in 2004**  
402 **(A) and 2013 (B) and the percentage decreases from 2004 to 2013 (C).**

403 Our modeled shear margin weakening on decadal scales is consistent with other estimates from a  
404 thermomechanical ice flow model of Jakobshavn Isbræ forced by calving front positions (Bondzio  
405 et al., 2017). Their modeled viscosity drops between 2003 to 2015 reach  $\sim 40\%$  which is close to  
406 our maximum viscosity decrease of  $\sim 45\%$  between 2004 to 2013 (Fig. 8). The extreme calving  
407 season we simulated in summer 2012 was accompanied by  $\sim 12 \text{ km a}^{-1}$  variations in speed at the  
408 calving front, which were facilitated by the accompanying shear margin-induced ice viscosity  
409 reductions of 60% at the time of maximum terminus advance. Simpler models of Jakobshavn Isbræ,  
410 using a flowband model (Nick et al., 2011) or simple calving parameterizations with no seasonal  
411 cycle (Muresan et al., 2016) cannot produce these seasonal variations in shearing. However, our  
412 model accommodates both the seasonal forcing from calving and the three-dimensional seasonal  
413 velocity shear impacts on effective viscosity. Without this physical process, speedups during intense  
414 calving events would be under-estimated, and this would lead to under-estimated mass  
415 transportation during the retreat.

### 416 **4.3 Comparison with previous estimates**

417 The cumulative mass change of Jakobshavn Isbræ estimated from airborne and satellite laser  
418 altimetry for 1997–2014 was tabulated Muresan et al. (2016). The mass loss over the 10-year period  
419 2004–2013 modeled by Muresan et al. (2016) is closer to observations than ours (Table 1). This is  
420 partly due to different tuning targets: matching observed mass change was a stated target in their  
421 study, whereas our study targets ice front position and velocity. Their close match to observed mass  
422 loss may be partly due to cancelling errors: 1) their modeled calving front barely moves after 2006,  
423 which leads to under-estimation of mass change; and 2) the modeled fast flow widths are larger than





424 observations, which amplifies the mass flux across the calving front. These two biases will not  
425 always offset each other perfectly in the future.

426 Muresan et al. (2016) failed to simulate the retreat of Jakobshavn Isbræ after 2010. This may be due  
427 to the thickness threshold employed in their calving parameterization. Once Jakobshavn Isbræ  
428 terminus has retreated into the deeper part of the bedrock trough, the terminus height might never  
429 drop below their calving threshold of 375 m. In this case their calving rate will be solely due to the  
430 eigen parameterization of strain rates. Moreover, absence of seasonality in their calving front leads  
431 to under-estimated dynamic thinning, which is a key prerequisite for further calving. In contrast,  
432 our crevasse-depth calving model depends on stresses and surface water run-off with strong seasonal  
433 variation. As the terminus retreats and the surface slope steepens the enhanced surface stretching  
434 enhances the opening of crevasses in both calving parameterizations.

435 Nick et al. (2011) used a flow-band model to estimate a mass loss of 2280 Gt for Jakobshavn Isbræ  
436 by 2100 under the A1B climate scenario (Table 1). In our model we use RCP4.5 climate forcing,  
437 which has lower temperature rises than A1B, especially after 2050. Nick et al. (2011) prescribed a  
438 flow-band that has a near uniform width of 5 km near the terminus. Later modeling work using a  
439 similar model suggested that stability of the glacier is fundamentally controlled by geometry, and in  
440 reality the width varies along the ice-stream (Steiger et al. 2017). Nick et al. (2011) chose sets of  
441 parameters that produced small inter-annual retreats of Jakobshavn from 2000-2010, which may  
442 limit mass loss and retreat. The absence of the shear margin weakening feedback in their model also  
443 likely causes underestimation of mass loss. This could account for the comparable projected mass  
444 loss to our results, and less terminus retreat (Table 1), even though their climate forcing scenario  
445 was warmer.



## 4.4 Model improvements

We overestimate mass loss relative to observations over Jakobshavn Isbræ drainage basin for 2004-2013 (Table 1). The main reason is excessive dynamic thinning over the lowest  $\sim 20$  km of the main trunk due to over-estimated summer speed. For example, modeled front velocity soared to a peak of  $\sim 20$  km  $a^{-1}$  in summer 2012, while the observed maximum speed is only 18 km  $a^{-1}$  (Joughin et al., 2014). In this summer, we simulated a series of full-thickness calving that eventually left an unprecedented tall ice cliff. In reality, calving events do not always occur to full thickness, thus the glacier tends to form a shorter ice cliff that caters for lower velocity and less dynamic thinning.

Since the grounding line of Jakobshavn retreated to the bottom of a reverse bed slope in 2009, the height of the calving front has generally increased, causing larger mass flux downstream across the calving front. Instead of enhancing the seasonal fluctuation of calving front position, substantial winter calving events have occurred instead. Given the fact that these calving events have reduced the typical winter advance from  $\sim 6$  km to  $\sim 3$  km since 2010, winter calving is now likely as important as summer run-off-driven calving. During this period of low magnitude seasonal fluctuations, a series of retreats gradually moved the calving front position on inter-annual scale. In contrast, the inter-annual retreats before 2009 were mostly driven by single calving seasons, e.g., May to July 2009. Our model using the Benn calving model is better able to stimulate this earlier retreat pattern, which is largely determined by each year's peak surface water run-off.

The grounding line of Jakobshavn Isbræ is unlikely to return to shallow water in the remainder of the 21<sup>st</sup> century because bedrock elevations  $< -1000$  m beneath the main trunk further extend a further  $\sim 60$  km inland. Accordingly, the latest retreat pattern including winter calving, is likely closer to the pattern of future evolution of Jakobshavn Isbræ. A short floating part due to winter calving is always accompanied by weaker lateral drag and steeper surface slope near the grounding



469 line, all of which are conducive for faster ice-flow. So, winter calving would enhance the  
470 downstream mass transportation, a missing process in our model.

471 The process of winter calving must take place without any surface water. That calving must be  
472 generated by processes affecting ice cliff stability, and that is likely due to changes at the base rather  
473 than the surface. Evidence of calving by opening of basal crevasses and splitting comes from  
474 terrestrial radar showing the terminus lifting several days prior to a large calving (Xie et al., 2016;  
475 James et al., 2014). These observations suggest that the glacier is not in hydrostatic equilibrium  
476 during calving. Our simulation specifies the glacier is in hydrostatic equilibrium on timescales of  
477 the simulation. Our model cannot simulate the process of up-lifting. Instead we assume the upper  
478 and lower surface would instantly lift to the state of floating (Eq. 1). However, there is some  
479 evidence that Jakobshavn must behave super-buoyantly in winter. We observe that the simulated  
480 grounding line of Jakobshavn retreats even after cessation of calving front retreat (Fig. 3). These  
481 retreats can be explained by rapid dynamic thinning near the grounding line leading to its buoyancy  
482 exceeding gravity and, consequently, floating.

483 A combination of discrete element model and continuum ice-dynamic model (solving the 3-  
484 Dimensional full-stokes equation) is able to reliably replicate observed calving styles in the case of  
485 a super-buoyant terminus (Benn et al. 2017). The discrete element model allows investigation of  
486 calving processes in unprecedented detail by analyzing the stress pattern dominated by glacier  
487 geometry and boundary conditions. However, these calving processes are beyond the capability of  
488 calving parameterization based on surface crevasse depth assuming depth-independent flow. Better  
489 understanding of this buoyancy-driven calving and further model development to represent more  
490 details such as fracture propagation are needed to accurately simulate glacier's future evolution.



## 491 **5 Conclusion**

492 We use a three-dimensional dynamic ice-sheet model with a physically-based calving  
493 parameterization to model the evolution of Jakobshavn Isbræ. After tuning the parameters, our  
494 model can accurately reproduce Jakobshavn Isbræ's retreats and velocity changes from 2004-2013  
495 on both seasonal and inter-annual scale. We project Jakobshavn Isbræ's future dynamic changes  
496 with climate forcing data from RACMO (2014-2099) and an ensemble of 7 Earth System Models  
497 for the RCP4.5 scenario.

498 We successfully model two-dimensional ice-flow patterns and their seasonal variations for  
499 Jakobshavn Isbræ, which are missing from several previous modeling studies. Moreover, capturing  
500 these two-dimensional patterns allows us to handle the influence of horizontal velocity shear on  
501 effective ice viscosity, which impacts on speedup processes of Jakobshavn Isbræ.

502 Over most of the 21st century, Jakobshavn Isbræ's grounding line will, we predict, retreat along the  
503 deep parts of a basal trough where bedrock elevation is significantly lower than at the present  
504 grounding line. Using the current generation of calving parameterizations, which are essentially  
505 thickness threshold models, is challenging because of the increasing height of the calving front as  
506 Jakobshavn Isbræ retreats, meaning that crevasse penetration depths become too small to initiate  
507 calving. Our model successfully reproduced Jakobshavn Isbræ's retreat down a reverse bed slope  
508 with an elevation drop of  $\sim 400$  m and the subsequent temporarily stable calving front position in  
509 2013 and 2014.

510 Our results suggest that rapid dynamic thinning and calving caused by deep crevasse penetration  
511 are responsible for most of its recent mass loss, and will be a decisive process in future mass loss.  
512 Further exploration of the physics of calving and basal sliding of Greenland outlet glaciers are



513 required to improve future projections.

## 514 Acknowledgements

515 This study is supported by National Key Science Program for Global Change Research  
516 (2015CB953601), National Key Research and Development Program of China (2018YFC1406104)  
517 and National Natural Science Foundation of China (No. 41506212). We thank Stephen Cornford for  
518 his help in implementing some parameterizations used in our model.

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## 520 References

- 521 Alexander, P. M., and Luthcke, S. B.: Greenland Ice Sheet seasonal and spatial mass variability from  
522 model simulations and GRACE (2003-2012), *The Cryosphere*, 10(3), 1259, 2016.
- 523 Amundson, J. M., Fahnestock, M., Truffer, M., Brown, J., Lüthi, M. P., and Motyka, R. J.: Ice  
524 mélange dynamics and implications for terminus stability, Jakobshavn Isbræ, Greenland, *J. Geophys.*  
525 *Res.-Earth Surf.*, 115(F1), 2010.
- 526 Bamber, J. L., Layberry, R. L., and Gogineni, S. P.: A new ice thickness and bed data set for the  
527 Greenland ice sheet: 1. Measurement, data reduction, and errors, *J. Geophys Res.-Atmos*, 106(D24),  
528 33773-33780, 2001.
- 529 Benn, D. I., Åström, J., Zwinger, T., Todd, J., Nick, F. M., Cook, S., Hulton, N. R.J., and Luckman,  
530 A.: Melt-under-cutting and buoyancy-driven calving from tidewater glaciers: new insights from  
531 discrete element and continuum model simulations, *J. Glaciol.*, 63(240), 691-702, 2017.
- 532 Benn, D. I., Warren, C. R., and Mottram, R. H.: Calving processes and the dynamics of calving  
533 glaciers, *Earth-Sci Rev.*, 82(3), 143-179, 2007.
- 534 Bentsen, M., Bethke, I., Debernard, J. B., Iversen, T., Kirkevåg, A., Seland, Ø., Drange, H., Roelandt,  
535 C., Seierstad, I. A., Hoose, C., and Kristjánsson, J. E.: The Norwegian earth system model,  
536 NorESM1-M-Part 1: Description and basic evaluation, *Geosci. Model Dev.*, 5, 2843-2931, 2012.
- 537 Block, A. E., and Bell, R. E.: Geophysical evidence for soft bed sliding at Jakobshavn Isbrae, West  
538 Greenland, *The Cryosphere Discussions*, 5, 339-366, 2011.
- 539 Bondzio, J. H., Morlighem, M., Seroussi, H., Kleiner, T., Rückamp, M., Mouginot, J., Moon, T.,  
540 Larour, E. Y., and Humbert, A.: The mechanisms behind Jakobshavn Isbræ's acceleration and mass  
541 loss: A 3-D thermomechanical model study, *Geophys. Res. Lett.*, 44(12), 6252-6260, 2017.
- 542 Cassotto, R., Fahnestock, M., Amundson, J. M., Truffer, M., and Joughin, I.: Seasonal and  
543 interannual variations in ice mélange and its impact on terminus stability, Jakobshavn Isbrae,  
544 Greenland, *J. Glaciol.*, 61(225), 76-88, 2015.
- 545 Collins, W. J., Bellouin, N., Doutriaux-Boucher, M., Gedney, N., Halloran, P., Hinton, T., Hughes,



- 546 J., Jones, C. D., Joshi, M., Liddicoat, S., Martin, G., O'Connor, F., Rae, J., Senior, C., Sitch, S.,  
547 Totterdell, I., Wiltshire, A., and Martin, G.: Development and evaluation of an Earth-System model  
548 HadGEM2, *Geosci. Model Dev.*, 4(4), 1051-1075, 2011.
- 549 Cornford, S. L., Martin, D. F., Graves, D. T., Ranken, D. F., Le Brocq, A. M., Gladstone, R. M.,  
550 Payne, A., J., Ng, E. G., and Lipscomb, W. H.: Adaptive mesh, finite volume modeling of marine  
551 ice sheets, *J. Comput. Phys.*, 232(1), 529-549, 2013.
- 552 Cornford, S. L., Martin, D. F., Payne, A. J., Ng, E. G., Le Brocq, A. M., Gladstone, R. M., Edwards,  
553 T. L., Shannon, S. R., Agosta, C., Van Den Broeke, M. R., Hellmer, H. H., Krinner, G., Ligtenberg,  
554 S. R. M., Timmermann, R., and Hellmer, H. H.: Century-scale simulations of the response of the  
555 West Antarctic Ice Sheet to a warming climate, *The Cryosphere*, 9, 1-22, 2015.
- 556 Cowton, T. R., Sole, A. J., Nienow, P. W., Slater, D. A., and Christoffersen, P.: Linear response of  
557 east Greenland's tidewater glaciers to ocean/atmosphere warming, *P. Natl. Acad. Sci. USA*, 115(31),  
558 7907-7912, 2018.
- 559 Csatho, B., Schenk, T., Van Der Veen, C. J., and Krabill, W. B.: Intermittent thinning of Jakobshavn  
560 Isbrae, West Greenland, since the little ice age, *J. Glaciol.*, 54(184), 131-144, 2008.
- 561 Cuffey, K. M., and Paterson, W. S. B.: *The physics of glaciers*. Academic Press, 2010.
- 562 Dee, D., Uppala, S., Simmons, A., Berrisford, P., Poli, P., Kobayashi, S., Andrae, U., Balmaseda,  
563 M., Balsamo, G., Bauer, P., Bechtold, P., Beljaars, A. C. M., van de Berg, L., Bidlot, J., Bormann,  
564 N., Delsol, C., Dragani, R., Fuentes, M., Geer, A. J., Haimberger, L., Healy, S. B., Hersbach, H.,  
565 Hólm, E. V., Isaksen, I., Kållberg, P., Köhler, M., Matricardi, M., McNally, A. P., Monge-Sanz, B.  
566 M., Morcrette, J. J., Park, B. K., Peubey, C., de Rosnay, P., Tavolato, C., Thépaut, J. N., and Vitart,  
567 F.: The ERA-Interim reanalysis: Configuration and performance of the data assimilation system, *Q.*  
568 *J. Roy. Meteor. Soc.*, 137, 553–597, 2011.
- 569 Dufresne, J. L., Foujols, M. A., Denvil, S., Caubel, A., Marti, O., Aumont, O., Balkanski, Y., Bekki,  
570 S., Bellenger, H., Benshila, R., Bony, S., Bopp, L., Braconnot, P., Brockmann, P., Cadule, P., Cheruy,  
571 F., Codron, F., Cozic, A., Cugnet, D., De Noblet, N., Duvel, J. P., Eth'e, C., Fairhead, L., Fichefet,  
572 T., Flavoni, S., Friedlingstein, P., Grandpeix, J. Y., Guez, L., Guilyardi, E., Hauglustaine, D.,  
573 Hourdin, F., Idelkadi, A., Ghattas, J., Joussaume, S., Kageyama, M., Krinner, G., Labetoulle, S.,  
574 Lahellec, A., Lefebvre, M. P., Lefevre, F., Levy, C., Li, Z. X., Lloyd, J., Lott, F., Madec, G., Mancip,  
575 M., Marchand, M., Masson, S., Meurdesoif, Y., Mignot, J., Musat, I., Parouty, S., Polcher, J., Rio,  
576 C., Schulz, M., Swingedouw, D., Szopa, S., Talandier, C., Terray, P., Viovy, N., and Bony, S.:  
577 Climate change projections using the IPSL-CM5 Earth System Model: from CMIP3 to CMIP5,  
578 *Clim. Dynam.*, 40(9-10), 2123-2165, 2013.
- 579 Gagliardini, O., Passalacqua, O., and Werder, M. A.: Retroactions between Basal Hydrology and  
580 Basal Sliding from Numerical Experiments. In *AGU Fall Meeting Abstracts*, 2014.
- 581 Giorgetta, M. A., Jungclaus, J., Reick, C. H., Legutke, S., Bader, J., Böttinger, M., Brovkin, V.,  
582 Crueger, T., Esch, M., Fieg, K., Glushak, K., Gayler, V., Haak, H., Hollweg, H., Ilyina, T., Kinne,  
583 S., Kornbluh, L., Matei, D., Mauritsen, T., Mikolajewicz, U., Mueller, W., Notz, D., Pithan, F.,  
584 Raddatz, T., Rast, S., Redler, R., Roeckner, E., Schmidt, H., Schnur, R., Segsneider, J., Six,  
585 Katharina D., Stockhause, M., Timmreck, C., Wegner, J., Widmann, H., Wieners, K., Claussen, M.,  
586 Marotzke, J., Stevens, B., and Glushak, K.: Climate and carbon cycle changes from 1850 to 2100  
587 in MPI-ESM simulations for the Coupled Model Intercomparison Project phase 5, *J. Adv. Model*  
588 *Earth Sy.*, 5(3), 572-597, 2013.



- 589 Gladish, C. V., Holland, D. M., Rosing-Asvid, A., Behrens, J. W., and Boje, J.: Oceanic boundary  
590 conditions for Jakobshavn Glacier. Part I: Variability and renewal of Ilulissat Icefjord waters, 2001-  
591 14, *J. Phys. Oceanogr.*, 45(1), 3-32, 2015.
- 592 Gogineni, P.: CReSIS radar depth sounder data, Center for Remote Sensing of Ice Sheets, Lawrence,  
593 KS <https://data.cresis.ku.edu>, 2012.
- 594 Gordon, H. B., Rotstain, L. D., McGregor, J. L., Dix, M. R., Kowalczyk, E. A., O'Farrell, S. P.,  
595 Waterman, L. J., Hirst, A. C., Wilson, S. G., Collier, M. A., and Watterson, I. G.: The CSIRO Mk3  
596 climate system model, 2002.
- 597 Habermann, M., Truffer, M., and Maxwell, D.: Changing basal conditions during the speed-up of  
598 Jakobshavn Isbræ, Greenland, *The Cryosphere*, 7(6), 1679-1692, 2013.
- 599 Holland, D. M., Thomas, R. H., De Young, B., Ribergaard, M. H., and Lyberth, B.: Acceleration of  
600 Jakobshavn Isbrae triggered by warm subsurface ocean waters, *Nat. Geosci.*, 1(10), 659, 2008.
- 601 Howat, I. M., Ahn, Y., Joughin, I., van den Broeke, M. R., Lenaerts, J. T., and Smith, B.: Mass  
602 balance of Greenland's three largest outlet glaciers, 2000-2010, *Geophys. Res. Lett.*, 38(12), 2011.
- 603 Jakobsson, M., Mayer, L., Coakley, B., Dowdeswell, J. A., Forbes, S., Fridman, B., Hodnesdal, H.,  
604 Noormets, R., Pedersen, R., Rebesco, M., Schenke, H. W., Zarayskaya, Y., Accettella, D.,  
605 Armstrong, A., Anderson, R. M., Bienhoff, P., Camerlenghi, A., Church, I., Edwards, M., Gardner,  
606 J. V., Hall, J. K., Hell, B., Hestvik, O., Kristoffersen, Y., Marcussen, C., Mohammad, R., Mosher,  
607 D., Nghiem, S. V., Pedrosa, M. T., Travaglini, P. G., and Schenke, H. W.: The international  
608 bathymetric chart of the Arctic Ocean (IBCAO) version 3.0, *Geophys. Res. Lett.*, 39(12), 2012.
- 609 James, T. D., Murray, T., Selmes, N., Scharrer, K., and O'Leary, M.: Buoyant flexure and basal  
610 crevassing in dynamic mass loss at Helheim Glacier, *Nat. Geosci.*, 7(8), 593-596, 2014.
- 611 Ji, D., Wang, L., Feng, J., Wu, Q., Cheng, H., Zhang, Q., Yang, J., Dong, W., Dai, Y., Gong, D.,  
612 Zhang, R., Wang, X., Liu, J., Moore, J. C., Chen, D., and Zhang, R. H.: Description and basic  
613 evaluation of Beijing Normal University Earth system model (BNU-ESM) version 1, *Geosci. Model*  
614 *Dev.*, 7(5), 2039-2064, 2014.
- 615 Joughin, I., Abdalati, W., and Fahnestock, M.: Large fluctuations in speed on Greenland's  
616 Jakobshavn Isbrae glacier, *Nature*, 432(7017), 608, 2004.
- 617 Joughin, I., Howat, I. M., Fahnestock, M., Smith, B., Krabill, W., Alley, R. B., Stern, H., and Truffer,  
618 M.: Continued evolution of Jakobshavn Isbrae following its rapid speedup, *J. Geophys. Res.-Earth*  
619 *Surf.*, 113(F4), 2008.
- 620 Joughin, I., Smith, B. E., and Howat, I. M.: A complete map of Greenland ice velocity derived from  
621 satellite data collected over 20 years, *J. Glaciol.*, 64(243), 1-11, 2018.
- 622 Joughin, I., Smith, B. E., Howat, I. M., Scambos, T., and Moon, T.: Greenland flow variability from  
623 ice-sheet-wide velocity mapping, *J. Glaciol.*, 56(197), 415-430, 2010.
- 624 Joughin, I., Smith, B., Shean, D., and Floricioiu, D.: Brief communication: Further summer speedup  
625 of Jakobshavn Isbræ, *The Cryosphere*, 8, 209-214, 2014.
- 626 Krabill, W., Abdalati, W., Frederick, E., Manizade, S., Martin, C., Sonntag, J., Swift, R., Thomas,  
627 R., Wright, W., and Yungel, J.: Greenland ice sheet: High-elevation balance and peripheral thinning.



- 628 Science, 289(5478), 428-430, 2000.
- 629 Luckman, A., and Murray, T.: Seasonal variation in velocity before retreat of Jakobshavn Isbræ,  
630 Greenland, *Geophys. Res. Lett.*, 32(8), 2005.
- 631 Moss, R. H., Edmonds, J. A., Hibbard, K. A., Manning, M. R., Rose, S. K., Van Vuuren, D. P., Carter,  
632 T. R., Emori, S., Kainuma, M., Kram, T., Meehl, G. A., Mitchell, J. F.B., Nakicenovic, N., Riahi,  
633 K., Smith, S. J., Stouffer, R. J., Thomson, A. M., Weyant, J. P., and Wilbanks, T. J.: The next  
634 generation of scenarios for climate change research and assessment. *Nature*, 463(7282), 747-756,  
635 2010.
- 636 Motyka, R. J., Truffer, M., Fahnestock, M., Mortensen, J., Rysgaard, S., and Howat, I.: Submarine  
637 melting of the 1985 Jakobshavn Isbræ floating mélange and the triggering of the current retreat, *J.*  
638 *Geophys. Res.-Earth Surf.*, 116(F1), 2011.
- 639 Muresan, I. S., Khan, S. A., Aschwanden, A., Khroulev, C., Van Dam, T., Bamber, J., Van Den  
640 Broeke, M. R., Wouters, B., Munneke, P. K., and Kjær, K. H.: Modelled glacier dynamics over the  
641 last quarter of a century at Jakobshavn Isbræ, *The Cryosphere*, 10(2), 597-611, 2016.
- 642 Nick, F. M., Vieli, A., Andersen, M. L., Joughin, I., Payne, A., Edwards, T. L., Pattyn, F. and van de  
643 Wal, R. S.: Future sea-level rise from Greenland's main outlet glaciers in a warming climate. *Nature*,  
644 497(7448), 235, 2013.
- 645 Pollard, D., and DeConto, R. M.: Description of a hybrid ice sheet-shelf model, and application to  
646 Antarctica, *Geosci. Model Dev.*, 5(5), 1273, 2012.
- 647 Schoof, C., and Hindmarsh, R. C.: Thin-film flows with wall slip: an asymptotic analysis of higher  
648 order glacier flow models, *Q. J. Mech. Appl. Math.*, 63(1), 73-114, 2010.
- 649 Shapero, D. R., Joughin, I. R., Poinar, K., Morlighem, M., and Gillet-Chaulet, F.: Basal resistance  
650 for three of the largest Greenland outlet glaciers, *J. Geophys. Res.-Earth Surf.*, 121(1), 168-180,  
651 2016.
- 652 Sohn, H. G., Jezek, K. C., and van der Veen, C. J.: Jakobshavn Glacier, West Greenland: 30 years  
653 of spaceborne observations, *Geophys. Res. Lett.*, 25(14), 2699-2702, 1998.
- 654 Steiger, N., Nisancioglu, K. H., Åkesson, H., Fleurian, B. D., and Nick, F. M.: Simulated retreat of  
655 Jakobshavn Isbræ since the Little Ice Age controlled by geometry, *The Cryosphere*, 12(7), 2249-  
656 2266, 2018.
- 657 Tsai, V. C., Stewart, A. L., and Thompson, A. F.: Marine ice-sheet profiles and stability under  
658 Coulomb basal conditions, *J. Glaciol.*, 61(226), 205-215, 2015.
- 659 Van Angelen, J. H., M Lenaerts, J. T., Van den Broeke, M. R., Fettweis, X., and Meijgaard, E.: Rapid  
660 loss of firn pore space accelerates 21st century Greenland mass loss, *Geophys. Res. Lett.*, 40(10),  
661 2109-2113, 2013.
- 662 Van Der Veen, C. J., Plummer, J. C., and Stearns, L. A.: Controls on the recent speed-up of  
663 Jakobshavn Isbræ, West Greenland, *J. Glaciol.*, 57(204), 770-782, 2011.
- 664 Vieli, A., and Nick, F. M.: Understanding and modelling rapid dynamic changes of tidewater outlet  
665 glaciers: issues and implications, *Surv. Geophys.*, 32(4-5), 437-458, 2011.





666 Watanabe, S., Hajima, T., Sudo, K., Nagashima, T., Takemura, T., Okajima, H., Nozawa, T., Kawase,  
667 H., Abe, M., Yokohata, T., Ise, T., Sato, H., Kato, E., Takata, K., Emori, S., and Kawamiya, M.:  
668 MIROC-ESM 2010: Model description and basic results of CMIP5-20c3m experiments, Geosci.  
669 Model Dev., 4(4), 845, 2011.

670 Xie, S., Dixon, T. H., Voytenko, D., Holland, D. M., Holland, D., and Zheng, T.: Precursor motion  
671 to iceberg calving at Jakobshavn Isbræ, Greenland, observed with terrestrial radar interferometry, J.  
672 Glaciol., 62(236), 1134-1142, 2016.

673