Very old firn air linked to strong density layering at Styx Glacier, coastal Victoria Land, East Antarctica

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Abstract

Firn air provides plenty of old air from the near past, and can therefore be useful for understanding human impact on the recent history of the atmospheric composition. Most of the existing firn air records cover only the last several decades (typically 40 to 55 years) and are insufficient to understand the early part of anthropogenic impacts on atmosphere. In contrast, a few firn air records from inland sites, where temperatures and snow accumulation rates are very low, go back in time about a century. In this study, we report an unusually old firn air effective CO\textsubscript{2} age of 93 years from Styx Glacier, near the Ross Sea coast in Antarctica. This is the first report of such an old firn air age (> 55 years) from a warm coastal site. The lock-in zone thickness of 12.4 m is larger than at other sites where snow accumulation rates and air temperature are similar. High-resolution X-ray density measurements demonstrate a high variability of the vertical snow density at Styx Glacier. The CH\textsubscript{4} mole fraction and total air content of the closed pores also indicate large variations in cm-scale depth intervals, indicative of layering. We hypothesize that the large density variations in the firn increase the thickness of the lock-in zone and, consequently, increase the firn air ages because the age of firn air increases more rapidly with depth in the
lock-in zone than in the diffusive zone. Our study demonstrates that all else being equal, sites where weather
conditions are favourable for the formation of large density variations at the lock-in zone preserve older air
within their open porosity, making them ideal places for firn air sampling.

1 Introduction

Bubbles trapped in ice cores preserve ancient air and allow direct measurements of the atmospheric
composition in the past (e.g., Petit et al., 1999). However, it is difficult to obtain air samples over the past several
decades from ice cores since the more recent air has not yet been completely captured into bubbles closed off
from the atmosphere. In contrast, we can obtain the recent records from the interstitial air in the porous,
unconsolidated snow layer (firn) on top of glaciers and ice sheets (Schwander et al., 1989, 1993). In addition,
we can take advantage of the very large amount of firn air because it allows us to accurately analyze isotopic
ratios of greenhouse gases and many trace gases such as man-made CFCs, HCFCs and SF$_6$ (Buizert et al., 2012a;
Laube et al., 2012). However, reported firn air ages date back only several decades at the sites where snow
accumulation rates are relatively high (Table 1). Old firn air (> 55 years) was observed only at sites where
surface temperatures and snow accumulation rates are low such as South Pole (Battle et al., 1996) and inland
Antarctic Megadunes (Severinghaus et al., 2010) (Table 1); however, even under such circumstances very old
firn air is not guaranteed, as demonstrated by Dome C (Table 1).

In the firn layer, air moves through the open pores and is occluded into the adjacent ice at at total porosity of
~0.1 (Schaller et al., 2017). Firn air moves downward with the adjacent ice (advection), but is furthermore mixed
by diffusion, and affected by thermal and gravitational fractionation (Craig et al., 1988; Johnsen et al., 2000;
Severinghaus et al., 2001; Goujon et al., 2003). In addition, the gradual bubble trapping in the firn affects the
movement of the air. As a result, at each depth there is a gas age distribution (Schwander et al. 1993; Trudinger
et al., 1997), rather than a single gas age. Therefore, studying firn air is also important for interpreting the record
of ancient air trapped in ice cores.

The firn column is generally divided into three zones: convective, diffusive and lock-in, depending on the
mechanisms of firn air movement (Sowers et al., 1992). The gravitational enrichment in $^{15}$N of N$_2$ is traditionally
used to define the boundaries between these zones. The convective zone is the upper part of the firm where the
air can ventilate with the overlying atmosphere. With stronger wind pumping, there can be a deeper convective
zone (Kawamura et al., 2013). This zone has the same $\delta^{15}$N of $N_2$ value as that of the atmosphere. The diffusive
zone is located under the convective zone, where molecular diffusion is the dominant mechanism of trace gas
transport in interstitial air (Blunier and Schwander, 2000). The age of the firn air increases slowly with depth in
the diffusive zone because of continued gas exchange with atmospheric air via diffusion. Heavier isotopes are
enriched with depth due to the gravitational fractionation in the diffusive layer. Thus, $\delta^{15}$N of $N_2$ gradually
increases with depth in the diffusive zone. In the lock-in zone (LIZ) below the diffusive zone, gas diffusion is
strongly impeded although the bubbles are not entirely closed. The top of the lock-in zone is called lock-in depth
(LID), where the gravitational fractionation ceases, so that the $\delta^{15}$N of $N_2$ becomes constant. The bottom of
the LIZ is defined as the full close-off depth ($z_{COD}$), where all air bubbles are closed off and firn becomes sealed
ice. The $z_{COD}$ can be estimated in two different ways. First, we can calculate the $z_{COD}$ from firn densification
models. Typically, the close-off occurs when the density of ice reaches about 830 kg m$^{-3}$ (Blunier and Schwander,
2000), equivalent to a critical porosity of around 0.1 (Schaller et al., 2017). Also, if temperature is known, the
average density at close-off can be estimated from empirical relations (Martinerie et al., 1992). Second, the
deepest position where air can be sampled from the firn column is commonly considered as (just above) the
$z_{COD}$. In theory, the $z_{COD}$ is the depth at which all pores are closed, but it can be ambiguous to specify the $z_{COD}$
in the field because firn air can be sampled at a slightly deeper depth than that of the shallowest impermeable
snow layer due to the existence of permeable layers at deeper depths – this effect is due to density layering
(Mitchell et al., 2015).

The gas ages in the LIZ increase with depth faster than in the diffusive zone. In the LIZ, firn air moves
downward at nearly the same rate as the surrounding ice, and therefore the age of the air increases with depth
at nearly the same rate as the age of ice increases.

The age of the firn air is directly related to the movement of the firn air. Firn air models help calculate the
firn air age using some parameters such as temperature and accumulation rate. However, several studies found
that layering also affects the movement of firn air (e.g., Mitchell et al., 2015; Schaller et al., 2017). This implies
that physical properties of the ice may affect the age of the firn air as well.

With regard to the lock-in and close-off processes, recent studies have focused on snow layers and
microstructure of the firn (Hörhold et al., 2011; Gregory et al., 2014; Mitchell et al., 2015; Schaller et al., 2017).
Density variability on millimeter to tens of cm scales is observed in all polar sites. Hörhold et al. (2011)
demonstrate that density variability is caused by physical snow properties in the firn column. Several studies have dealt with how snow density variations affect the transport of firn air (Hörhold et al., 2011; Mitchell et al., 2015). Mitchell et al. (2015) showed that the firn layering can affect the closure of pores and the thickness of LIZ, but the relation between snow density variations and range of firn air ages was not quantitatively examined. In this study, we present firn air composition and δ\textsuperscript{15}N-N\textsubscript{2} from Styx Glacier, East Antarctica to better understand the role of snow density variations on the age of firn air. We also present X-ray density data with millimeter resolution and compare them with δ\textsuperscript{18}O\textsubscript{ice} and the closed-pore air composition in the LIZ.

We hypothesize that large snow density variations make the LIZ thicker and facilitate preservation of old firn air at the Styx Glacier. This study will help us better understand how the snow density layers of firn column affects movement and preservation of firn air, and provide guidance on selecting good sites for future firn air studies.

2 Materials and Methods

2.1 Firn air sampling and gas mole fractions analysis

The firn air and ice core were collected at the Styx Glacier, East Antarctica (73° 51.10’ S, 163° 41.22’ E, 1623 m asl) in December of 2014 (Fig. 1). This site is located 85 km north of the Korean Jang Bogo Station in the Southern Cross Mountains near the Ross Sea (Han et al., 2015). The snow accumulation rate is ~10 cm ice year\textsuperscript{1}, calculated from the Styx16b ice chronology based on methane correlation and tephra age tie-point and thinning functions (Yang et al., 2018). The mean annual surface temperature was measured as -31.7°C by borehole temperature logging at 15 m depth, two years after the ice core drilling (Yang et al., 2018). Table 1 lists the characteristics of the Styx Glacier and other firn air sampling sites. A total of 13 samples from the surface to 64.8 m depth were collected. The firn air sampling device was constructed, following the design of the University of Bern, Switzerland (Schwander et al., 1993). Three vacuum pumps (two diaphragm pumps and one metal bellows pump), several pressure gauges, stainless steel lines, and vacuum valves were housed in an aluminum case to transfer to the polar site. The pump system plays four major roles: (1) purging modern air from the bottom of a borehole, (2) inflating the bladder to block the deep firn layers from the atmosphere, (3) removing the contaminated air and extracting the firn air, (4) transporting firn air to a CO\textsubscript{2} analyzer for
measurements of gas mole fractions and store it in firn air containers. The bladder system is designed to be lowered into the borehole to seal the deep firn layer(s) being sampled from the atmosphere. The bladder consists of a 4 m-long rubber tube and metal caps on top and bottom of the rubber tube. The bladder’s external diameter is 119.5 mm and internal diameter is 114.5 mm. The material of the tube is butyl rubber (BIIR) which can endure low temperatures, providing no risk of sample contamination.

The firn air samples were collected in 3-liter glass flasks at all collection depths. However, to test preservation ability of the sample air containers, Silcocan canisters were also used at 4 depths (0, 35.36, 43.42, 53.95 m). Accurate mole fractions of CO$_2$, CH$_4$, and SF$_6$ were measured at US National Oceanic and Atmospheric Administration (NOAA; https://www.esrl.noaa.gov/). The results for the two types of containers show good agreements. δ$^{15}$N of N$_2$ was analyzed at Scripps Institution of Oceanography for correcting gravitational fractionation effect (Severinghaus et al., 2010).

### 2.2 Firn air transport model

We used the Center for Ice and Climate (CIC) firn air model which is a 1-dimensional advection-diffusion model to simulate how the air moves in Styx firn column. In this model, there are four types of air transport in the open porosity: (1) molecular diffusion, (2) vigorous mixing in the convective zone, (3) advection, and (4) dispersion in the deep firn (Buizert, 2012b, Buizert and Severinghaus, 2016). The model uses the stochastic bubble trapping formulation described by Mitchell et al. (2015).

### 2.3 CH$_4$ in closed bubbles and total air content measurements

CH$_4$ mole fraction in the (closed) air bubbles in the firn was measured at Seoul National University using a melt-refreeze air extraction method (Yang et al., 2017). 124 discrete firn samples (cross section of 8.5 cm × 3 cm, length of 3 cm, ~35 g) were prepared from 4 different depth intervals in the lock-in zone (54.59-55.34, 58.11-59.05, 59.86-60.55, 64.02-65.25 m). All ice samples were cut and trimmed by ~2.5 mm with a band saw to remove contaminants on the surface ice. Then, the ice samples were inserted into the glass flasks attached to the gas extraction line. The pump system evacuated air in the flask placed in a cooled ethanol bath at -70 °C for 20 min. The evacuation time was limited to 20 minutes to prevent from gas loss due to pore openings by
sublimation. After the pressure dropped below 0.2 mTorr, the ice samples in the glass flask were melted and air in the bubbles was extracted. After the melting was finished, we refroze the ice using a cooled ethanol bath to release the gas dissolved in the ice melt. Finally, the extracted air was injected into the sample loop of the gas chromatograph equipped with a flame ionization detector (FID). The calibration curve of the GC-FID was calculated by the standard air prepared at NOAA with a CH$_4$ mole fraction of 895 ppb on the NOAA04 scale (Dlugokencky et al., 2005).

Total air content of the firn ice samples was analyzed simultaneously with CH$_4$ mole fraction using the wet extraction system at SNU. The total air content was expressed as the volume of air trapped in the closed pores of unit mass of firn ice sample (in unit of ml per gram ice at STP conditions). The volume of air extracted from a firn ice sample was calculated by the ideal gas law with the internal pressure, volume and temperatures of the sample flasks and vacuum lines. The pressure of extracted air was measured by a pressure manometer connected to the sample loop of the GC-FID. As no direct measure of temperature was available, the temperature of extracted air was assumed to be identical to the surrounding temperatures; the ethanol temperature was used for the sample flasks, room temperature for vacuum lines, and valve box temperature (50°C) for the sample loop. In this study, the corrections for bubble-cut effect and thermal gradient within vacuum line were not considered. More detailed description of the protocols of total air content measurements is described in Yang (2019).

2.4 Analysis for stable isotopes of ice

After completing the measurements of the CH$_4$ mole fraction in air, the melt water was put into cleaned 125 ml bottles and analyzed for water stable isotope ratios at Korea Polar Research Institute (KOPRI) using a Cavity Ring-Down Spectroscopy (CRDS, L1102-i, Picarro, USA) system. We performed the same analysis for the snow pit samples, but without CH$_4$ analysis. The data are here presented as δ-notations:

$$\delta^{18}O = \left(\frac{^{18}O}{^{16}O}\right)_{\text{sample}}/\left(\frac{^{18}O}{^{16}O}\right)_{\text{VSMOW}} - 1$$  (1)

The firn ice melt was filled into a 400 μl insert in a 2 ml glass vial using a syringe filter. The auto sampler transported the ice melt samples in the insert to the vaporizer about 180 nl at a time. The samples with the liquid state were transferred to the cavity after being converted into the water vapor in a vaporizer at 110 °C. The measurement precision evaluated by measuring an in-house standard repeatedly (n=12) was 0.08‰ (1 sigma
standard deviation).

2.5 X-ray firn density measurement

We obtained high-resolution density data using the X-ray transmission method reported by Hori et al. (1999) for the firn ice at various depth intervals. This method is advantageous because it can measure continuously and non-destructively. The X-ray beam penetrates the ice samples and the detector on the opposite side analyzes the intensity of the beam. To make equal thickness for each core section, upper and side parts of the half circle-shape core were shaved by a microtome. After putting the precut ice core on a rack, we set the rate of measurement at 50 mm min$^{-1}$, and finally obtained 1mm-resolution density data.

3 Results

3.1 Layered stratigraphy

We examined a snow pit, located 10 m away from the main ice core borehole, 2 years after drilling to understand the physical properties such as layers, density, and ice grain size of the upper firn at Styx site. We scratched the snow wall by hand to remove soft layers and enhance the visibility of hard layers (Fig. 2a). The soft layers are presumed to be depth hoar, and the hard ones are wind crusts (Fig. 2b). The alternating layers repeat with intervals of few centimeters to 20 centimeters. The top boundaries of the hard layers are sharp and extend horizontally about a meter, but the bottom boundaries are not well defined due to gradual density changes. 10 cm-resolution density data were obtained by a density cutter (Proksch et al., 2016). The soft layers are coarse-grained, while the hard ones are fine-grained (Fig. 2b-d).

3.2 Firn gas sampling and the age of firn air

We calibrate the depth-diffusivity profile in the model using trace gases with a well-known atmospheric history (Buizert et al., 2012a; Trudinger et al., 2013; Witrant et al., 2012). The atmospheric time series from well-dated firn air (MacFarling Meure et al., 2006) and instrument measurement records (NOAA; https://www.esrl.noaa.gov/) were used for calibration. The simulated mole fraction profiles match well with the observations (Fig. 3). CO$_2$, CH$_4$, SF$_6$ and $\delta^{15}$N-N$_2$ distributions in firn air were modeled. The model does not
include thermal fractionation, and therefore provides a poor fit to the $\delta^{15}\text{N}-\text{N}_2$ data in the upper firn where seasonal temperature gradients fractionate the gases. Fitting the barometric equation to the $\delta^{15}\text{N}$ data of the upper diffusive zone suggests a convective zone thickness of approximately 3 m. This is within the typical range of observed convective zones, but perhaps lower than expected for a very windy site (Kawamura et al. 2006). The firn air age (black curves in Fig. 3) slowly increases with depth at the diffusive zone because it mixes with fresh atmospheric air on the surface mostly by molecular diffusion (Blunier and Schwander, 2000). In contrast, the firn air age rapidly increases within the LIZ at a rate similar to that of the ice age. The gas age distribution of Styx ice at $z_{\text{COD}}$ is narrower than the other sites where old firn air is reported (Fig. 4); we simulate a spectral width of 15.9, 22.8 and 45.5 years at Styx, South Pole, and Megadunes, respectively. This means that the past atmospheric history of trace gases can in principle be reconstructed with higher resolution at Styx than at the other old-air firn sites.

We estimate the age of samples in two ways. First, after calibrating the firn air model, we can derive the mean sample age from the simulated gas age distribution. At the deepest Styx sampling depth (64.8 m) we simulate a mean CO$_2$ age of 102 years, and a mean CH$_4$ age of 97 years; the CH$_4$ age is younger than the CO$_2$ age due to the higher diffusivity of CH$_4$. Second, we can estimate the sample ages by comparing the measured trace gas concentrations directly to the atmospheric histories of these gases – this age has been called the “effective age” (Trudinger et al. 2013). The lowest CO$_2$ mole fraction of 305.18 ppm at depth of 64.8 m (304 ppm after correcting for gravitational enrichment) corresponds to the year 1921 an effective age of 93 years (relative to sampling year 2014) on the Law Dome ice core record (MacFarling Meure et al., 2006; Rubino et al., 2019). The CH$_4$ mole fraction of 943.36 ppb at the same depth (946.5 after gravitational correction) corresponds to an effective age of 96 years (MacFarling Meure et al., 2006) (Figs. 3a, 3b). The second method provides younger ages because the growth rate in the atmospheric mixing ratios of these gases has increased over time, biasing the effective ages towards younger values (Trudinger et al. 2002). Table 1 lists effective CO$_2$ ages in the deepest firn air sample for several sites; we here compare the effective CO$_2$ age between sites rather than the modeled mean age, as it is purely empirical and does not rely on model assumptions.

Only few firn air sites have effective CO$_2$ ages around 93 years or older: 91 years from the South Pole (Battle et al., 1996) and 129 years from Megadunes (Severinghaus et al., 2010; Table 1). These sites are located in
interior Antarctica and have low annual mean temperatures and low snow accumulation rates (Table 1). Firn densification takes a long time if snow accumulation and/or temperature are low, therefore firn air can be preserved for a long time without being trapped. In contrast, the Styx site is located near the coast and has relatively high snowfall, and therefore the age of 93 years is very unusual. Sites of comparable climate characteristics typically have an oldest firn air age of around 40 years. This indicates that there may be other factors that can permit preservation of the old firn air at Styx Glacier.

3.3 Density layering and its influence on bubble trapping

Firn density is the primary control on the bubble close-off process. Density layering leads to staggered bubble trapping, with high-density layers closing off before low-density ones (Stauffer et al., 1985; Etheridge et al. 1992, Mitchell et al. 2015, Rhodes et al. 2016).

Because the mole fractions of atmospheric greenhouse gases (CO$_2$, CH$_4$, N$_2$O) have increased during the last century, we may obtain information on the timing of the closure of the bubbles from the greenhouse gas mole fractions of the air trapped in closed bubbles. In this study, we used the CH$_4$ concentration in closed bubbles ($[\text{CH}_4]_c$) and the total air content of the firn ice as indicators of the close-off process. The density and $[\text{CH}_4]_c$ show an anti-correlation (Fig. 5). Our results confirm the CH$_4$ concentration-total air content relation observed in West Antarctic Ice Sheet (WAIS) Divide firn ice (Mitchell et al., 2015). High-density layers reach the lock-in and close-off densities at shallower depths than low-density layers. Thus, air bubbles are trapped at shallower depths in high-density layers. Early trapped bubbles preserve older air with lower greenhouse gas mole fractions. Higher air content is expected in the high-density layers, in which open porosity is small and closed porosity is large (Fig. 5). However, we cannot entirely exclude the possibility of some post-coring bubble close-off (Aydin et al., 2010). High open porosity in low-density layers may have more chances to trap modern ice storage air, which has higher mole fraction of CH$_4$ than atmospheric background levels.

Figure 5a shows $[\text{CH}_4]_c$ and total air content in the LIZ of the Styx firn. $[\text{CH}_4]_c$ generally decreases with depth and the cm-scale variability is reduced in the deep layers, while the total air content generally increases with depth. The $[\text{CH}_4]_c$ greater than CH$_4$ mole fraction in neighboring firn air (green line in Fig. 5a-d) indicates part of bubbles formed after coring and increased the $[\text{CH}_4]_c$, as previous studies also observed (Mitchell et al.,
Most of $[\text{CH}_4]$ data show large cm-scale variations (Fig. 5). The highs and lows of $[\text{CH}_4]$ repeat with cycles of 6 to 24 cm (Fig. 5e). Note that the layering observed in the snow pit likewise shows irregular intervals (Fig. 2b). From the layer spacing, we conclude that bubble trapping at Styx is not controlled by annual layers (Section 4), as was observed at Law Dome (Etheridge et al. 1992).

The evolution of $\text{CH}_4$ in the closed porosity may give information on how the snow layers can induce inhomogenous records and help constrain the gas age distribution in ice (Fourteau et al. 2017). However, the details are beyond the scope of this study and we will focus on the firn air age in the open porosity.

### 3.4 High-resolution firn density measurements

The X-ray measurements show highly variable density on cm scales. We converted the high-resolution density to total porosity using the following equation:

$$\Phi_{\text{total}} = 1 - \frac{\rho}{\rho_{\text{ice}}} \quad (2)$$

where $\rho = \text{density of porous ice}$; $\rho_{\text{ice}} = \text{density of bubble-free ice}$ (919 kg m$^{-3}$); and $\Phi = \text{porosity}$. We test the idea that the lock-in zone corresponds to the depth range bounded by the first closed layer (porosity below 0.1) on the shallow side, and the last open layer (porosity above 0.1) on the deep side.

At Styx Glacier, the shallowest depth where the running mean of total porosity with a 1 cm-thick window reaches below 0.1, is 48.1 m (Figs. 6a and 6b). It is approximately 4.3 m shallower than the LID of 52.4 m defined by the modelled firn air $\delta^{15}$N-N$_2$ profile. Meanwhile, the deepest point, where the running mean (with a 1 cm-thick window) exceeds 0.1, is at 63.7 m (Figs. 6a and 6c), which is shallower than the $z_{\text{COD}}$ of 64.8 m defined by the deepest successful firn pumping depth. Although the LID and $z_{\text{COD}}$ from the density data are different from those defined by firn air data, the thickness of LIZ from density data (between the two orange lines in Fig. 6a) is comparable to that from firn air analysis (between two blue lines in Fig. 6) (15.6 vs. 12.4 m).

The offsets of the LIZ about 1-4 m between those from total porosity and the firn air measurement may be due to for example small calibration offsets in the density data set, the fact that actual critical porosity may be variable and depend on the study site, or on horizontal snow density variations and the horizontal extent of diffusion-impeding layers. The similarity in the LIZ thicknesses from the two methods support the idea that the
large variations of density can increase the LIZ thickness by shallowing LID and/or deepening the zCOD. The thick LIZ eventually permits storing old firn air at Styx (Table 1). Usually, the LIZ thickness increases with a snow accumulation rate (Witrant et al., 2012), presumably because at high accumulation sites density variability in the lock-in zone tends to increase (Horhold et al. 2011). Refrozen melt layers may also act as high density, diffusion-impeding layers allowing for older firn air to be sampled as observed in Devon Island (Witrant et al., 2012). We demonstrate here that the snow density variability is an important factor in determining the firn air age. We suggest that sites with higher density variations at the LIZ have a high possibility of a thick LIZ and therefore old firn air, even in warm, relatively high-precipitation coastal climates.

4 Discussions

To quantitatively compare density variability of Styx snow with those at other glacier sites, we may use the standard deviation of densities (σρ) near the mean air-isolation density (Hörhold et al., 2011; Martine et al., 1992). The mean density at the mean air-isolation depth (ρcrit) can be related to mean annual temperature (T in Kelvin) using the following equation, which is empirically obtained from air content measurements (Martinerie et al., 1992):

\[
ρ_{\text{crit}} = \left(\frac{1}{ρ_{\text{ice}}} + 7.6 \times 10^{-4} \times T - 0.057\right)^{-1}
\]

where \(ρ_{\text{ice}}\) is the density of bubble-free pure ice.

Although this equation cannot provide exact \(ρ_{\text{crit}}\), we can take advantage in estimating the density at LIZ without gas chemistry data (Hörhold et al., 2011). We note that Martinerie et al. (1994) suggested slightly different coefficients for the equation based on a different set of data; however, the results do not significantly change our conclusions. We also note that Bréant et al. (2018) used an equation relating ice density at LID to snow accumulation rate; however, we prefer to use the relation of temperature-ice density at LIZ by Martinerie et al. (1992) because the latter is more relevant to the ice density at LIZ. Using the Styx high-resolution X-ray density data at depth interval of 43.13-66.97 m, we calculated the standard deviation of densities (σρ). For each \(σ_p\), we used 1000 density data points (Fig. 7) as Hörhold et al. (2011) did for the other sites listed in Table 2. At Styx,
\(\rho_{\text{crit}}\) is 821.68 kg m\(^{-3}\) according to equation (4), and the standard deviation of densities at \(\rho_{\text{crit}}\) \((\sigma_{\rho}, \rho_{\text{crit}})\) is 19.33 \(\pm 1.87\) kg m\(^{-3}\), which is greater than those in the other previously studied sites (Hörhold et al., 2011; Fig. 7, Table 2). The high \(\sigma_{\rho}, \rho_{\text{crit}}\) at Styx likely facilitates the thick LIZ and old firn air.

A high-density (low-density) layer at surface may become a low-density (high-density) layer (Freitag et al., 2004; Fujita et al., 2009) at density of 600-650 kg m\(^{-3}\), which occurs at shallower depths than LIZ (Hörhold et al., 2011). Thus, vertical snow layering at surface may not directly give information about density variability at LIZ (Hörhold et al., 2011). However, conditions for snow layering at the surface still may give us clues on the density variability at LIZ. The conditions may include redistribution of snow by wind and formation of wind and/or radiation crusts (Martinerie et al., 1992; Hörhold et al., 2011). To test the possibility of seasonal causes, we analyzed stable isotopes of surface snow (\(\delta^{18}\)O) because the surface \(\delta^{18}\)O generally follows seasonal variation (depleted in winter and enriched in summer). Figures 2e and 2f show the stable isotope profiles of snow (\(\delta^{18}\)O) at Styx Glacier, which are apart by \(~100\) m; one is from a snow pit made in 2014 and the other is from the main ice core drilled in 2014. The \(\delta^{18}\)O profiles commonly show cycles with intervals of \(~40\) cm per year, given that local maxima of \(\delta^{18}\)O indicate summer, and minima winter layers. Meanwhile, the repetition of the density layers has twenty cycles (high and low density layer pairs) in the top 180 cm depth at the snow pit (Fig. 2b). Using a snow accumulation rate of \(~40\) cm y\(^{-1}\) in recent years, the density layers have 4~5 cycles y\(^{-1}\), indicating that the formation of snow density layers is mainly controlled by non-seasonal factors.

A blizzard occurred during the ice coring campaign in December of 2014. We observed that the blizzard strongly reworked the surface snow. The Automatic Weather System (AWS) installed within 10 m from the borehole site show that blizzard events (wind speed > 15 m s\(^{-1}\)) took place on December 29 in 2015, May 23, June 26, August 17, and September 7 in 2016 (Fig. S1). The number of blizzard events in a year is similar to the mean density layer cycle of 4~5 y\(^{-1}\). Although Blizzards occur more frequently in winter, the frequency of 5 yr\(^{-1}\) is comparable to the number of the density layer cycles of 4~5 yr\(^{-1}\). During the blizzard events, westerly wind prevailed, and snow particles may have been redeposited with a sorted-size distribution (large grains in the bottom and small grains on the top) similar to winnowing seen in sedimentary records (Sepp Kipfstuhl, personal communication). Between the blizzards, the solar radiation and temperature gradient may have
facilitated the diagenesis of the snow layers (Alley, 1988; Fegyveresi et al., 2018). During the diagenesis processes, fine and coarse flake layers may form high-density and low-density layers, respectively. In summary, blizzard events may have played a major role in forming snow density layers.

5 Conclusions and implications

About 93-year-old firn air (effective CO$_2$ age) was found at Styx Glacier, East Antarctica, located near the Ross Sea coast. This is of great scientific interest because such old firn air is commonly only found in the inland sites such as the South Pole and Megadunes. The thickness of Styx LIZ is relatively greater than those in other sites where snow accumulation and temperature are similar. The thicker LIZ made the Styx firn layer preserve old firn air because the age of stagnant firn air rapidly increases with depth in the LIZ as air exchange with the atmosphere has stopped. We hypothesized that the high snow density variations in the LIZ of Styx Glacier made the thick LIZ and old firn air. To test the hypothesis, we conducted high-resolution X-ray density measurements. We argue that the thick LIZ is related to the high density variations at Styx Glacier. We also examined why high snow density variability developed at Styx site. The effect of strong wind (e.g., blizzards) may facilitate the density layer formation. It is likely that old firn air (>55 years) can be found in areas where climatological conditions are favorable for high snow density variations at LIZ even when the sites are located near the coast. We may take advantage in sampling and transportation from the coastal sites, because logistics is easier for those sites. Theoretically, the oldest firn air should be available in a site that has both strong layering and a low accumulation rates. Older firn air, perhaps as old as 150 years, may still be found under such suitable conditions on the Antarctic continent.

Acknowledgements. We thank Jeff Severinghaus and Ross Beaudette at Scripps Institution of Oceanography for accurate $\delta^{15}$N-N$_2$ analysis, and Jacob Schwander at University of Bern for a kind advice in constructing the SNU firn air sampling device. We also thank Mauro Rubino and an anonymous reviewer for their constructive comments. This study was supported by Korea Polar Research Institute (PE 18040) and National Research Foundation of Korea (NRF-2018R1A2B3003256).
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Environmental Sciences, Seoul National University, Seoul, 188 pp., 2019.
Figure 1. (a) Location map of study site, Styx Glacier, Antarctica and (b) a photo of surface snow density layers. The thickness of the snow density layers varies horizontally. The top boundaries of the high-density layers are sharp (horizontal red-dashed line). A hole on a high-density layer surface is indicated by a red-dashed circle. The length of the black sharp pencil in (b) is 14.3 cm.
Figure 2. Snow-pit photos at Styx Glacier. (a) The snow pit with dimensions of 280 × 65 × 220 cm (length × width × height). (b) The illustration of qualitatively defined hard (dark blue) and soft (pale blue) layers observed in the top 180-cm-depth interval. Progressive blue color changes indicate a gradual density decrease with depth. Red line is a 10-cm-resolution density profile. (c) Coarse grains observed in a soft layer. The grains were placed on a black glove. (d) Enlarged snow layers. Dashed red lines indicate top boundaries of fine-grained hard layers. (e) and (f) Stable isotope ratio ($\delta^{18}$O) of snow profiles at the main core and a snow pit 100 m away from the main ice core borehole, respectively.
Figure 3. CO$_2$, CH$_4$, SF$_6$ mole fractions and $\delta^{15}$N of N$_2$ measurements (circles), and model results (solid line) for the Styx firn air (air in open porosity). Black lines are modeled ages for the gas species.
Figure 4. Comparison of model-simulated CO$_2$ age distributions at Styx (this study), South Pole (Battle et al., 1996), and Megadunes (Severinghaus et al., 2010).
Figure 5. (a-d) CH$_4$ mole fraction in closed pores ([CH$_4$]$_{cl}$) (red line) and total air content (air volume per ice weight) (blue line) in the lock-in zone of Styx Glacier. (e) Comparison of density with [CH$_4$]$_{cl}$ and total air content near z$_{COD}$. A small dashed-box in (d) indicates the depth interval of Fig. 5e.
Figure 6. X-ray high-resolution density data obtained from the lock-in zone. (b) and (c) are enlarged portion of (a). Black lines show individual density data, while the red lines are 1-cm running means. Blue and orange lines represent the boundaries of the LIZ estimated from the gas compositions (between two vertical blue lines) and the critical porosity thresholds (between two orange vertical lines), respectively (see section 3.4).
Figure 7. Density variability calculated from 1000 depth points and their average density. The standard deviation at the critical density (821.68 kg m$^{-3}$) calculated from the approximate second order polynomial ($R = 0.84$) is $19.33 \pm 1.87$ kg m$^{-3}$. The blue and red areas are the density ranges near the LID (52.38 - 52.48 m) and the $z_{COD}$ (64.91 - 65.01 m), respectively.
Table 1. Glaciological characteristics of Styx Glacier and other firn air sampling sites.

<table>
<thead>
<tr>
<th>Site</th>
<th>T (°C)</th>
<th>A (cm ice yr⁻¹)</th>
<th>Effective CO₂ age (year)</th>
<th>LID (m)</th>
<th>COD (m)</th>
<th>LIZ thickness (m)</th>
<th>References</th>
</tr>
</thead>
<tbody>
<tr>
<td>Styx</td>
<td>-31.7</td>
<td>10</td>
<td>93</td>
<td>52.4</td>
<td>64.8</td>
<td>12.4</td>
<td>This study, Yang et al. (2018)</td>
</tr>
<tr>
<td>Megadunes</td>
<td>-49</td>
<td>~0</td>
<td>129</td>
<td>64.5</td>
<td>68.5</td>
<td>4</td>
<td>Severinghaus et al. (2010)</td>
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<tr>
<td>South Pole</td>
<td>-51.0</td>
<td>8</td>
<td>91</td>
<td>115</td>
<td>125</td>
<td>10</td>
<td>Severinghaus et al. (2001)</td>
</tr>
<tr>
<td>Siple Dome</td>
<td>-25.4</td>
<td>13</td>
<td>59</td>
<td>49</td>
<td>58</td>
<td>9</td>
<td>Severinghaus et al. (2001)</td>
</tr>
<tr>
<td>Dome C</td>
<td>-54.5</td>
<td>2.7</td>
<td>33</td>
<td>97</td>
<td>100</td>
<td>3</td>
<td>Landais et al. (2006)</td>
</tr>
<tr>
<td>WAIS Divide</td>
<td>-31</td>
<td>22</td>
<td>39</td>
<td>~67</td>
<td>76.5</td>
<td>9.5</td>
<td>Battle et al. (2011)</td>
</tr>
<tr>
<td>NEEM</td>
<td>-28.9</td>
<td>22</td>
<td>50</td>
<td>63</td>
<td>78</td>
<td>15</td>
<td>Buizert et al. (2012a)</td>
</tr>
<tr>
<td>NGRIP</td>
<td>-31.1</td>
<td>19</td>
<td>45</td>
<td>67.5</td>
<td>78</td>
<td>11.5</td>
<td>Kawamura et al. (2006)</td>
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<tr>
<td>Summit</td>
<td>-32</td>
<td>23</td>
<td>26</td>
<td>70</td>
<td>80.8</td>
<td>10.8</td>
<td>Witrant et al. (2012)</td>
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<tr>
<td>DE-08</td>
<td>-19</td>
<td>120</td>
<td>13</td>
<td>71.8</td>
<td>88.5</td>
<td>16.8</td>
<td>Etheridge et al. (1996)</td>
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</table>
Table 2. Comparison of standard deviation of density ($\sigma_\rho$) at critical density ($\rho_{\text{crit}}$). For data from all other sites, except the Styx, refer to Hörhold et al. (2011).

<table>
<thead>
<tr>
<th>Campaign/Region</th>
<th>Core name</th>
<th>$\rho_{\text{crit}}$ (kg m$^{-3}$)</th>
<th>$\sigma_\rho$, $\rho_{\text{crit}}$ (kg m$^{-3}$)</th>
<th>T (℃)</th>
<th>A (cm ice yr$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Styx</td>
<td>Styx</td>
<td>821.68</td>
<td>19.33±1.87</td>
<td>-31.7</td>
<td>10</td>
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<tr>
<td>NGT B16</td>
<td></td>
<td>819.27</td>
<td>12.26</td>
<td>-27</td>
<td>15.5</td>
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<tr>
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<td></td>
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<td>12.81</td>
<td>-30</td>
<td>11.3</td>
</tr>
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<td>820.81</td>
<td>12.91</td>
<td>-30</td>
<td>11.8</td>
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<tr>
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<td>13.23</td>
<td>-30.6</td>
<td>20</td>
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<tr>
<td>NGT B29</td>
<td></td>
<td>821.32</td>
<td>10.50</td>
<td>-31.6</td>
<td>16.7</td>
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<tr>
<td>Berkner Island</td>
<td>B25</td>
<td>819.16</td>
<td>14.57</td>
<td>-27</td>
<td>15</td>
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<tr>
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<td>10.27</td>
<td>-42</td>
<td>6.9</td>
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<tr>
<td>DML B32</td>
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<td>11.28</td>
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<td>8.12</td>
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<td>7.3</td>
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<tr>
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<td>815.00</td>
<td>16.59</td>
<td>-18.1</td>
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<tr>
<td>Pre-IPICS B39</td>
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<td>814.91</td>
<td>17.11</td>
<td>-17.9</td>
<td>84</td>
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<td>Pre-IPICS DML95</td>
<td></td>
<td>815.51</td>
<td>13.42</td>
<td>-19.2</td>
<td>60</td>
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<tr>
<td>Pre-IPICS DML97</td>
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<td>816.07</td>
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<td>-20.4</td>
<td>53</td>
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<td>Dome C</td>
<td>EDC2</td>
<td>832.02</td>
<td>4.59</td>
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<td>WDC06A</td>
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