Description of changes made to ‘Attribution of sea ice model biases to specific model
errors enabled by new induced surface flux framework’

In this document, changes made to the manuscript arising from the reviews are described.
The original point-by-point response to the reviews, and a version of the manuscript with
changes tracked, are appended.

1. Fundamental changes to the analysis
In response to a suggestion by Reviewer 2, fields of surface flux estimated by the formulae
described in Section 2.3 of the original paper were compared to modelled actual surface flux
fields. While the estimated fields captured well the seasonal and spatial variation in surface
flux, large discrepancies in the absolute values were apparent (30-40 Wm^{-2} in some
months). As a result of these, the methods of estimating surface flux were refined in the
following ways:

- Surface flux contributions were estimated for each ice thickness category separately, and then summed (as opposed to using the mean ice thickness across all
categories);
- The reference temperature about which the Stefan-Boltzmann relation is linearised was allowed to vary in space and time (it was previously uniformly 0°C). For each
grid cell, this reference temperature was set to the monthly mean surface
temperature.
- A representation of the turbulent fluxes was added to the formulae.
The biases in estimated surface flux were greatly reduced by this method. The remaining biases, and their implication for the results, are discussed in the new Appendix A, which examines potential errors in the induced surface flux analysis.

Although the results of the analysis were not qualitatively changed by the new methods, the
contribution of the ice thickness bias during the winter to the surface flux bias was somewhat
increased, probably because the greater efficiency of ice production for thin ice is captured
more effectively by using the full thickness category information.

Because the new surface flux formulae use much more detailed information about the model
diagnostics, the approach of calculating surface flux dependency on each variable has
changed. Whereas in the previous version, the model-observation mean was used to
calculate partial derivates (i.e. surface flux dependency on variable), in the new version only
the model state is used to calculate this. Implications of this are also discussed in the new
Appendix A.

2. Major additions and corrections by section
The structure of the paper has been altered slightly. The description of the induced surface
flux (ISF) analysis has been moved later in the paper, to a new Section 4, after the model
sea ice and surface radiation evaluation in Section 3. The derivation of the formulae is now
described briefly within this section, rather than in an appendix. Hence the new structure is
In these sections, the presentation of the logic behind the ISF method has been altered, in part inspired by the changes described in Section 1 above. It is hoped that as a result the presentation is improved and clarified.

The argument is stated again here for ease of reading: at any model point in space and time, we approximate model surface flux as a function of climate variables that affect the surface flux on timescales shorter than that on which they affect each other. In this way, by taking partial derivatives, we characterise at each point in space and time the rate at which surface flux depends on any climate variable. Multiplication of the resulting field by an estimate of model bias in that variable therefore produces a field of estimated surface flux bias induced (on a near-instantaneous timescale) by the model bias in that variable. These fields can be averaged in space and time to give large-scale estimates of surface flux bias, bypassing nonlinearities present in the relationship between surface flux and climate variable. This allows the proximate causes of surface flux bias (and hence bias in sea ice growth and melt) to be directly quantified.

The argument above is presented fully in the new Section 4, but is summarised in the Introduction, and more briefly summarised in the abstract. It is stated more carefully that the ISF method can diagnose only proximate causes of surface flux bias, as pointed out by Reviewer 1.

In the model description, the sub-gridscale thickness distribution of HadGEM2-ES is now described more carefully, as it is now key to the analysis. The snow thermodynamics is also described in more detail, as requested by Reviewer 2. The motivation for the use of the period 1980-1999 is also described more fully, also as requested by Reviewer 2.

No major corrections are present in the observational data description section.
2.3 Sea ice state and surface radiation evaluation

For this section, and all subsequent sections, the Arctic Ocean region has been refined to exclude more of the Barents Sea, where processes are considered sufficiently different to the rest of the Arctic Ocean as to render the assumptions used in the ISF analysis questionable (for example there is high oceanic heat convergence here, and ice concentrations are low or zero year-round). Figure 1, which maps this region, has been changed accordingly.

In addition, unless stated otherwise, all model results are now given for the ensemble mean of the 4 historical simulations of HadGEM2-ES (rather than for the first historical simulation only, as had been the case for the previous version).

As a result of these changes, many of the sea ice thickness and surface radiation biases quoted in Section 3 are different. The previous Figures 3-5 (presenting the evaluations) have also been overhauled accordingly, and are now numbered Figures 2-4 (as discussed below, it was decided that the previous Figure 2 was superfluous and has been removed).

The qualitative conclusions of the evaluation are unchanged: sea ice thickness is too low in the annual mean and too amplified: net SW is too high and upwelling SW too low in the summer, net LW and downwelling LW are too low in the winter.

2.4 ISF method description

This section has been completely rewritten, reflecting the changes in the methodology described above. Firstly, the use of the word ‘bias’ is defined (as requested by Reviewer 2, all instances of ‘anomaly’ the paper have been replaced with ‘bias’). This is followed by a discussion of the relationship between sea ice mass balance, surface flux and oceanic heat convergence, clarified along the lines suggested by both reviewers, and with additional evidence cited that in both HadGEM2-ES and in the real world, oceanic heat convergence is a minor contributor to sea ice mass balance in the Arctic Ocean interior.

With the link between surface flux and sea ice mass balance established, the approach of the study (see 2.1 above) is set out in detail. The formulae by which surface flux values are estimated at each point in model space and time are described, and the assumptions behind their derivation are discussed (as requested by Reviewer 1, Thorndike (1992) is cited at this point, as it contains many of the key assumptions used).

Following the request by Reviewer 2 for a more detailed description of the method application, the calculation of induced surface fluxes is then described in turn for three climate variables (melt onset, downwelling LW and ice thickness), with the aid of a new figure (Figure 5), which shows for each variable the fields of variable bias, surface flux dependence, and induced surface flux bias. The calculation of the ice thickness component is described in particular detail, as it now involves the estimation of thickness biases by category which is nontrivial.

It is noted that in the updated analysis the treatment of the turbulent fluxes is changed; their contribution is no longer neglected. However, they are treated in such a way that their...
dependence on any climate variable examined cannot be evaluated, save for the ice area. As a result of this change in emphasis, their omission is no longer justified in this section (and the corresponding Figure 2 is removed); instead, the implications of their treatment is discussed in the new Section 6.

2.5 ISF results
The presentation of the results of the ISF analysis has been changed in several ways. Firstly, the ISF biases (with CERES as radiation reference dataset) are presented in the new Table 1, following the request by Reviewer 1 for more systematic quantification of the results. The table also shows total ISF bias, total net radiation bias relative to CERES, the residual between the two, and the CERES-ERAI net radiation bias, in preparation for a discussion of the relationship between observational uncertainty and ISF uncertainty, inspired by the request by Reviewer 2 for the ISF residuals to be discussed.

The results are also presented in Figure 6 (equivalent to Figure 6 in the previous version). This figure is unchanged in essence, but now shows the ensemble mean induced surface fluxes for the newly-refined Arctic Ocean region (see 2.3 above). Total ISF fluxes are also shown more clearly, using black bars.

Following the presentation of Table 1 and Figure 6, the significant ISF biases are described in turn (June surface melt onset, August ice area, early winter ice thickness, winter downwelling LW). In each case, the bias is now quantified, and an equivalent figure for bias in sea ice growth and melt is described, following the request by Reviewer 1 to make the surface flux-sea ice mass balance link more explicit. Internal variability in the ISF biases is then described, using all 80 ensemble years (Reviewer 2 request). Residuals between total ISF and total net radiation biases are quantified and compared to observational uncertainty in net radiation (Reviewer 2 request). The direct effect of observational uncertainty on ISF biases is described. Potential errors in the ISF biases arising from method assumptions, discussed in detail in Appendix A, are discussed (Reviewer 1&2 requests).

The discussion of ISF spatial patterns that follows has been left largely unchanged. A new paragraph has been appended to this, comparing the spatial pattern in total ISF bias to that in net radiation bias, with the aid of a new figure (Figure 7).

2.6 Discussion
Following the suggestion by Reviewer 1, the diagnosis of the thickness-growth and ice albedo feedbacks from the ISF analysis has been justified in a new paragraph. The resulting relative contributions (over the course of a year) of the feedbacks and forcings has been quantified (both in terms of surface flux and sea ice melt/growth bias), thereby justifying a statement in the abstract that Reviewer 1 had rightly questioned. The word ‘forcing’ has also been defined more carefully, as it is different from the most commonly used meaning of this word in a climate context.

Most of the ensuing discussion is unchanged, except for minor rewordings discussed in Section 3 below, and a mention of Holland et al (2010), requested by Reviewer 2, at the point at which it is concluded that the June surface albedo simulation problems are the
principal cause of the low annual mean thickness of HadGEM2-ES. However, two paragraphs have been appended, discussing the implications for the results of firstly the imperfect treatment of turbulent fluxes, and secondly the omission of oceanic heat convergence as an additional potential source of model bias in sea ice growth and melt. In this second paragraph, HadGEM2-ES Arctic Ocean heat convergence is compared to observational estimates, following requests by Reviewers 1&2.

2.7 Conclusions
This section has not undergone substantial changes as it is considered that the methodological changes have not resulted in any changes to the conclusions of the study. However, two paragraphs have been reworded for improved clarity.

2.8 Appendix A (analysis of ISF errors)
This section is equivalent to Appendix B in the original study, but is substantially larger and is based on several suggestions by Reviewers 1&2 for ways in which errors in the ISF analysis should be quantified. Its conclusions are quoted at the appropriate point in Section 5.
Firstly, the error in quantifying dependence of surface flux on climate variable is discussed. This discussion is motivated by comparison of estimated surface flux fields to actual, as suggested by Reviewer 2. As discussed above, although the new methodology has improved the correspondence some differences remain, which can be categorised into three types. In each case, a likely cause of the difference is stated, and its impact on the ISF biases estimated.
Secondly, the error in characterising induced surface flux bias as a product of variable bias with surface flux dependence is described. In this discussion, higher order derivatives of the surface flux are evaluated. The error inherent in using the model state to evaluate surface flux dependence is calculated using the mixed partial derivative terms (this follows from a point made by Reviewer 2).

3. List of minor corrections and rewordings
In the following list, the page and line number references refer to the non-tracked changes version of the manuscript. Each alteration is followed with an indication of whether it was requested by Reviewer 1, 2 or neither (in the last case with additional justification). Alterations which are directly related to the new structure, the new methods, or any of the enhancements listed above, are not listed.
1-16 (abstract) ‘Counteracting’ removed (R2)
1-22 (Intro) ‘-’ replaced with ‘to’ (R2)
1-27 Stammerjohn et al cited (R2)
1-28 ‘whoseloss’ replaced with ‘the loss of which’ (R2)
Appendix A: Original point-by-point reviewer response

A1. Reply to Reviewer Comment 1 (Anonymous)

We thank the reviewer for taking the time to read our manuscript, and for his/her useful suggestions for its improvement, which we address inline below. The reviewer’s comments are quoted in italics.

West et al propose a new analysis framework to understand model biases in Arctic sea ice which they apply to HadGEM2-ES, a model with known biases in sea ice characteristics.

The attribution of climate model errors in the sea ice zone is a very important open topic and the paper provides original and likely efficient means to evaluate such errors. The main problem I think is writing, which I found often imprecise, and renders a proper evaluation of the paper difficult.

In particular, the methods absolutely require clarification and should use better and simpler terminology. Because I did not fully get the methods, it was thereafter really complicated to follow, in particular the discussion and conclusions.
A second requirement to make this paper acceptable is to early on in the result section to explain that the induced surface flux method works - eg. to describe how well the different methods to compute surface flux biases converge. Now this is done here and there, and I have constantly been doubting of the quality of the methods, because of the absence of such evaluation.

The reviewer is right that the convergence of the different methods, demonstrated in Figure 6, deserves better discussion, and probably quantification, which we propose to carry out by more thorough analysis of the errors of the induced surface flux method in Appendix B, and discussion of these in Section 4, as described below.

Our view is that the spread amongst the different estimates is caused predominantly by observational uncertainty, with errors introduced by the induced surface flux method assumptions relatively small in magnitude by comparison. This is supported by the fact that the difference between the sum of the induced surface flux contributions and each surface flux anomaly is comparable in magnitude to the differences between surface flux anomalies wrt the different datasets (ERAI, ISCCP-FD, CERES).

We think that it would be difficult to show that the induced surface flux (ISF) method works purely by comparison with the direct surface radiation evaluation, because the difference between the different estimates are dominated by the observational uncertainty. We think a better way would be by a more thorough evaluation of the impact of assumptions made by the ISF method in Appendix B. For example:

- Evaluation of errors introduced by the simple model by comparing modelled fields of net radiation to those predicted by the formulae, as suggested by Reviewer 2
- Evaluation of errors caused by ignoring higher-order derivatives, by calculating these terms

The magnitude of these errors would then be compared to the observational uncertainty, as estimated by the difference between the direct radiation evaluations shown in Figure 6.

There is a more fundamental point: the ISF method is not just a way of calculating surface flux anomalies due to a particular process, but also of characterising them, because their definition is to some degree subjective. For example, suppose for the month of May a model shows mean downwelling SW of 300 Wm\(^{-2}\), and albedo 0.8, given net SW of 60Wm\(^{-2}\), but observational estimates shows mean downwelling SW of 250 Wm\(^{-2}\), and albedo of 0.7, giving net SW of 75 Wm\(^{-2}\), with a model anomaly of 15 Wm\(^{-2}\). Clearly the downwelling SW anomaly induces a positive surface flux anomaly, the albedo anomaly induces a negative one, and the total surface flux anomaly is -15 Wm\(^{-2}\). But the exact induced anomalies are subjective. The approach used in the paper is equivalent to multiplying the anomaly in one process by the mean in the other, giving contributions of +12.5 Wm\(^{-2}\) and -27.5 Wm\(^{-2}\) by the downwelling SW and albedo anomalies respectively.

Hence while it would in theory possible to say whether the sum of the ISF contributions was correct (if we knew the exact actual surface flux error), it would not be possible to say whether each individual contribution were correct. The main requirement is that each contribution is physically realistic, and provides useful information.
A third thing I would have enjoyed to see is a specific discussion of how the ice-albedo and growth-thickness feedbacks can be diagnosed from the method. It is claimed in the abstract that your method can separate these effects, and I am in trouble to see how that statement is presently supported in the text. I can guess feedbacks are acting from Fig. 6, but I think this topic deserves a bit more to support the claim made in the abstract.

This does require greater justification. The sentence in which the relevant anomalies are identified with the surface albedo feedback and thickness-growth feedback (page 10, line 31) will be expanded accordingly.

The thickness-growth feedback, for example, ostensibly acts by altering the energy balance at the base of the ice; as the ice thickens, the temperature gradient decreases, basal conduction decreases, and the energy balance becomes less strongly negative, so the ice thickens more slowly. But by energy conservation, this process must also have some manifestation in the external fluxes: as the ice is losing energy less quickly, some external energy flux must also have changed. Under the assumptions of the simple model used (similar to Thorndike 1992 as you note below), which ignores sensible heat storage, it is the upwelling LW term that changes: as the ice thickens, its top surface also cools to maintain flux continuity.

Hence any change to the ice energy balance resulting from the ice thickening, and therefore conducting less efficiently, can be diagnosed as the contribution of the ice thickness to the change in upwelling LW radiation.

I have also not understood why energetic errors of oceanic origin have been ignored from the discussion, especially in the North Atlantic sector of the Arctic - where there is a low bias.

Energy passing to the ice through the ocean-to-ice heat flux has two main sources: solar input to the ocean, and oceanic heat convergence. The first is implicitly taken account of through the analysis of the effect of ice fraction anomalies on net SW radiation. We make the case that Arctic-wide, the importance of the oceanic heat convergence in driving summer basal ice melt is small by comparison to direct solar input, and this will be more thoroughly justified in the revision, referencing model results by e.g. Steele et al 2010 in addition to the observational references originally included.

However, you are right that the contribution of the oceanic heat convergence should be properly quantified. In the revised version of the paper, we will include estimates of Arctic-wide ocean heat convergence (for model and observations), and set the results shown in Figure 6 in this context.

Finally, the authors claim in the conclusions that they can “quantify” the origin of errors, but apart from Fig. 6 (which I liked a lot), I did not really see a quantification of the errors. Is that quantification the main point - or is it the consistent comparison of the different sources of error?

We do quantify individual contributions to the surface flux biases as shown in Figure 6 for a few illustrative months, in section 4. However, we will examine whether there is scope for a
more systematic approach, for example quoting the annual average flux contribution for
each state variable in a table.

Also, it was difficult to ultimately figure out whether biases in external forcings or in the sea
ice model are the ultimate cause of the biases. Is your method capable to tell after all?

Briefly, the answer is no: the method cannot tell the ultimate cause of the surface flux biases.
It is designed to diagnose the proximate cause of the biases.

The induced surface flux (ISF) method, alone, can determine only the first-order cause of the
net surface flux bias. The state variables examined (downwelling SW, downwelling LW, melt
onset occurrence, ice fraction, ice thickness) affect the surface flux on very short timescales,
and are unambiguously properties of the atmosphere (radiation) and sea ice (melt onset,
fraction, thickness). Hence the ISF method allows the short-term causes of surface flux bias
to be decomposed into those arising from the atmosphere, and from the sea ice.

It’s recognised that the causes of biases in the state variables themselves may lie in different
systems. For example, ice thickness biases will have some ultimate cause in the
atmosphere – indeed, that is one finding of the paper. Conversely, while the case is made in
Section 5 that cloud errors are to blame for downwelling LW biases, it is likely that the sea
ice simulation will nevertheless have some influence on how this is manifested. However, all
these effects act on relatively long timescales, compared to the almost instantaneous
timescale on which the state variables affect the net surface flux.

A last general comment - the logics of the arguments should be better presented.

We apologise that the presentation of logic is unsatisfactory. This paper has been through
several rounds of restructuring as the analysis has developed, and the coherence of
argument has probably suffered due to this. We will try to significantly improve this in the
revision.

I am pretty confident that - if these presentation issues are seriously addressed by the team
of coauthors, this will make an excellent contribution to their favourite cryospheric journal.

A few specific comments.

* I have tried to understand what the generic approach is. Here is what I have understood.
The present presentation is too lengthy, misses the essential elements and overdiscusses
details. A synthetic view is missing. There are three means to evaluate errors in surface
energy budget (I have understood two of them)

1) The direct computation of surface flux bias, i.e. the difference between simulated and
observed surface flux (or one of its components)
This is correct although we would add the caveat that this is still only an estimate of the actual surface flux bias – the observational uncertainty is very large.

2) The induced surface flux bias, which is the contribution of bias in a specific variable to surface flux bias, namely calculated as \( \Delta F_x = \frac{dF_x}{dx} \Delta x(\text{mod-obs}) \).

To evaluate derivatives, the SEB is simplified using two different approximations during the cold and warm seasons, based on ideas from Thorndike et al 1992.

I don’t think there is a need to calculate those derivatives in the body of the paper.

We broadly agree but note the concern by Reviewer 2 that some of the calculations by which the induced surface fluxes are arrived at are incompletely explained. It may be necessary to show at least one derivative, for illustration, with the additional detail that is planned for the revision.

If the derivatives are well calculated and if the non-linearities are not too important, the sum of \( \Delta F_x \) should hopefully approach the surface flux bias.

3) The third diagnostic is “the sea ice latent heat flux uptake anomaly implied by the ice volume anomalies relative to PIOMAS”.

I have tried to figure out what the authors mean, but I did not really managed. The wording is not precise enough for the reader to what is meant by this and what is gained by comparing that to the surface flux biases. I guess “latent heat flux” is confusing in the context of the surface energy budget. But whether that thing is a heat storage anomaly divided by time or something else, I don’t know. Maybe an “ice thickness bias converted to Joules” or “an energetic equivalent ice thickness bias”?

This is correct. For each month, the modelled field of rate of change of ice thickness is calculated as half the difference between the following and the preceding month. A similar field is calculated for the reference dataset, PIOMAS. The reference field is subtracted from the modelled field to create a model anomaly of ice thickness change. This is then multiplied by ice density, specific latent heat of fusion, reversed in sign, and divided by the number of seconds in a month, to create an equivalent sea ice latent heat storage anomaly in Wm-2.

For the reasons discussed in Section 2.3, the surface heat flux is viewed as the main source of this latent heat storage. It’s noted however that it would be useful to provide an estimate of modelled and observed oceanic heat convergence which provides an additional input to the latent heat storage.

We will clarify this point in the revised manuscript.

Besides an explanation of what it means, we would need an explanation of what should be taken from that diagnostic.
It is important to clarify this point because a lot of the argumentation was based on that.

* The two methods to compute the surface flux derivatives is called "a model". I think it is a "computation method". It is actually inspired from Thorndike et al (1992) – which should be acknowledged - and maybe from earlier works in EBM. What you are doing is to derive the surface energy budget wrt anything.

The ‘model’ versus ‘computation method’ is an interesting distinction – our interpretation is that it hinges on whether the formulae are viewed as a way of calculating induced surface fluxes (model) or characterising them (computation method), as discussed above. As in our view both are valid interpretations, either phrase might be more appropriate depending on the circumstances.

Thorndike et al will be cited.

A2. Reply to reviewer comment 2 (Francois Massonnet)

We thank Francois Massonnet for his helpful and thorough review of our manuscript. Below, we address in turn his major and minor comments inline, which are quoted in italics.

Major comments

1) It is not always easy to follow the authors’ methodology. I see the general idea behind the approach: expressing the net flux to the ice as a function of state variables, then linearizing around a reference state to obtain the flux bias resulting from the bias in one of the model components. However, I could surely not reproduce the results myself, just based on the text. I appreciate the efforts to publish the code in Supplementary Material, but the text itself should have all elements. For example, I do not understand how the bias in $F_{sfc}$ attributed to error in melt onset is derived (i.e., from eq. A6 to 4). Furthermore, the attribution of flux error to melt onset error seems to not be a function of the melt onset itself, but rather a function of concentration difference. This is confusing: melt onset is defined by the time of the day where surface melting commences, and the right hand side of Eq. 4 does not display surface melting terms. At some point it came to my mind that the authors were perhaps using "melt onset" for "ice retreat", but I’m not sure. In all cases, this is confusing.

The definition of the state variable ‘melt onset occurrence’, and its relation to the net surface flux, is not very clearly explained in the paper, and certainly requires expansion. This will be altered in the revised version of the paper, and we will ensure more generally the replicability of the calculations of the induced surface fluxes (for example, noting the use of local ice and snow thicknesses as you suggest below). However, a brief explanation of this particular component, melt onset occurrence, is also provided here. Very simply, its purpose is to capture the effect of meltpond formation on the surface flux.
HadGEM2-ES parameterises the effect of meltponds by reducing surface albedo linearly from 0.8 to 0.65 as the surface temperature goes from -1°C to 0°C, after Curry et al. (2001). Because we have daily surface temperature fields, we can judge for each modelled year which day ‘melt onset’ — defined as the day surface temperature first goes above -0.5°C — occurs. Comparison of these dates to the observational SSMI estimates referenced in the paper shows that modelled melt onset occurs, on average, 20-25 days earlier across most of the Arctic Ocean in the model than in observations.

In the induced surface flux analysis, we examine the effect on modelled surface flux of the melt onset process occurring at the wrong time of year. The relevant state variable here is ‘melt onset occurrence’, which takes the value 0 or 1 depending on whether a grid cell on a particular day has yet exceeded -0.5°C (model definition) or whether a liquid water microwave signature has been detected (observational definition). In a similar way to the other state variables used in the paper, the observations are averaged over the period 1980-1999 to obtain a daily climatology of melt onset occurrence. For each modelled year, this climatology is subtracted from the modelled melt onset occurrence fields to obtain a modelled melt onset anomaly. This anomaly is then multiplied by the relevant partial derivative — in this case, (downwelling SW) * (cold snow albedo – melting snow albedo) to produce the induced anomaly in net SW.

2) Besides the need for clarity in the methodology, a key question is to what extent the assumption of linearity holds, in particular for what bias range Eqs 3 and 4 would be valid. Will the methods work for models with very large biases? Another point is that this linearisation involves the use of Eqs 1 and 2, that are themselves derived using linearity assumptions. I trust that the approach is valid, because the sum of individual contributions (Fig. 6) seems to match the flux errors from datasets and from volume estimates, but a quantification of this match should be done (perhaps by calculating residuals). Overall I find that the authors have not discussed the validity of this assumption, and this is critical given how non linearly the ice behaves.

You are right that this needs to be quantified. For some of the state variables the dependence of surface flux is linear, and the second partial derivatives go to zero (e.g. all state variables in the melting season, and downwelling LW in the freezing season). However the dependence of surface flux on ice and snow thickness is nonlinear and it would be useful to examine the circumstances in which higher derivatives are important.

As you mention, additional assumptions are made in deriving equation (1): linearity of upwelling longwave dependence on surface temperature, and uniform conduction of heat within the ice. We will try to quantify the impact of these also by comparing actual net surface flux fields to predicted fields, as you suggest below in point 5).

3) I would also like to see if the method is robust to internal variability. Could the authors take one or several of the four other members of the HadGEM2-ES model and run the same analysis? In other words, is the 1980-1999 period long enough to identify and attribute the biases?
Analysis of other ensemble members would be a valuable enhancement of the study. For the revised version of the paper, we plan to carry out the same analysis on the other three ensemble members, and to quantify the consistency with the results from the first member.

4) The authors have not cited an important study: Holland et al., 2008 (doi:10.1007/s00382-008-0493-4). In that study, the inter-model scatter in the sea ice mass budget (present-day conditions) is shown to be explained by the way models absorb shortwave radiation. This is directly relevant to the study here, and I think the authors should go through the Holland et al. study to position their results with respect to theirs. In particular the claim that turbulent fluxes are of relatively minor importance relative to radiative fluxes in setting the surface energy balance (Fig. 2, and p. 5 line 11-16) should be put in perspective with that study. As the Arctic sea ice mean state changes, turbulent fluxes appear to be of increasing importance.

Holland et al (2008) show annual sea ice melt rates to be strongly correlated with summer net SW across the CMIP3 ensemble. The causality here could go in either, or most likely both, directions. This appears to be consistent with the finding in the current study that the excessive sea ice melt in HadGEM2-ES is driven by surface albedo and net SW issues. This will be referenced.

The neglecting of turbulent fluxes is a shortcoming of our study. As you point out, while they are comparatively small in an absolute sense, they may nevertheless be important in driving future changes. In a similar way, model anomalies in turbulent fluxes may be of comparable size to those in radiative fluxes even if the model absolute values are much smaller. We will expand on this point in the Discussion.

5) The authors should prove, with a figure, that the model developed in the appendix is good enough to do the investigations. Could they plot, for one or several grid cells and one or several freezing seasons, the reconstructed flux $F_{sfc}$ (Eq. A4) and the actual flux from the model? A quantification of the correspondence would be a plus.

This would also be a valuable exercise. Actual and calculated modelled fields will be compared for a few sample years, and the correspondence described, quantified and illustrated. The most appropriate place for this would probably be the discussion of errors in Appendix B.

6) Nothing is said about the treatment of snow in the HadGEM2-ES model. How many snow layers are there, what is snow conductivity, etc.? There is only one snow layer, and the conductivity is 0.33 Wm-1K-1. Like the ice, the snow has no heat capacity. It should be made clear, however (here and in the revised version of the paper) that sensible heat storage is parameterised in the top 10cm of the snow-ice column during surface exchange calculations, to aid stability.
7) The authors repeatedly use the word "anomaly" to describe the difference between modeled and reference quantities, but I would avoid this word and use "bias" or "error" instead. To me, an "anomaly" is used to describe the deviation of a signal with respect to its own mean. We have some concern that use of the word 'bias' might suggest that HadGEM2-ES is being evaluated with respect to the 'truth', but most datasets used are only very rough approximations to this. We will change 'anomaly' to 'bias', but clearly define at the outset the meaning of the word 'bias' for the purposes of the paper – the difference of the model relative to a particular observational estimate.

8) I’m unclear about whether ocean surface temperature biases are accounted for in the analysis. From Fig. 6, it looks like they are not. On the other hand, p. 5 lines 6-10 seem to suggest that the Arctic Ocean is critical in setting the ice energetic balance (and this is also seen in Keen et al (https://link.springer.com/article/10.1007/s00382-013-1679-y, their Fig. 4). So, I’m puzzled: is the contribution of oceanic surface temperature bias taken into account or not in the analysis?

It is true that the ocean contributes a significant amount of heat to the ice in the summer. However, we make the case in our study that the major part of this heat comes from direct solar heating of the ocean, an effect which is taken account of through analysis of the effect of ice fraction on the net SW bias. This case will be strengthened in the revision by citing evidence from models (e.g. Steele et al, 2010) in addition to the evidence from observations already referenced.

This is also relevant to the minor comment at 5-6/10 below.

Minor comments
1-18 - "countered by a counteracting" is a bit odd. ‘Counteracting’ is superfluous and will be removed.


Change will be made as suggested.

1-28 - Along with an earlier melt onset date, you can mention that freeze up has been delayed: Stammerjohn et al., 2012, their Fig.2 (doi:10.1029/2012GL050874)

This will be done.

1-29 - "whoseloss" –> "the loss of which"
Change will be made as suggested.

2-1 - Evaluation against observations of volume is quite impossible (even extent observations are not direct observations), so I would use "observational or reanalysis reference datasets"

Change will be made as suggested.

2-18 Instead of "anomalies" I would use "biases"

See response to point 7) above.

3-24 The period 1980-1999 is used for evaluation, because it "predates the rapid sea ice loss". Why is it a problem to have a period with strong trend in the analysis? Is it expecting that the SEB would change too rapidly during a period with strong trends? Would the analysis be robust if the model output was evaluated on a distinct and later period (2000-2015 for instance, using historical + RCP8.5 runs). Please elaborate.

The trend itself is not problematic. Our motivation for using this period was to evaluate the model on a ‘reference’ time period at least partially independent of the time period that is usually used to evaluate sea ice trends. We will make this clear in the revised version.

3-35 Reanalysis data also suffer from biases because of errors in atmospheric forcing, this could be stated as well.

This will be stated.

5-11/16 Can the authors explain exactly what they mean by "Heat flux due to snowfall". The presence of snow affects heat conduction fluxes and acts to reduce bottom growth, is that what the authors are talking about?

The presence of snow matters only because it takes energy to melt it; snow falling on ice changes the enthalpy of the snow-ice system. A 3m column of bare fresh ice at 0°C, for example, will take \(-9.2 \times 10^8\) Jm\(^{-2}\) to melt it. If 50cm fresh snow falls on the ice, the combined snow-ice column will take \(-9.8 \times 10^8\) Jm\(^{-2}\) to melt it. Hence the falling of snow on ice represents a transfer of negative latent heat from the atmosphere system to the snow-ice system, which must be taken account of in calculating the total surface flux. We will try to explain this more fully in the revised version of the paper.
5-26 I assume \( h_I \) and \( h_s \) refer to in-situ / actual thickness (this is the one that matters for vertical thermodynamics). It would be good to mention that here, as there is usually a lot of confusion between that quantity and the grid cell average thickness.

Yes, this is the case and that will be clarified.

5-24 Eq. 1: Maybe I missed it, but what is the value for ice albedo? Does albedo depend on the ice state?

The ice albedo used is 0.61. This does not depend on ice thickness, but falls linearly to 0.535 as ice surface temperature rises from -1°C to 0°C. This information will be added to the text.

5-27 The subscript for snow thickness is "S" here while it is "s" elsewhere

Change will be made as suggested.

5-28 The symbol for albedo is \( \alpha_I \) while in the equation it’s \( \alpha_{ice} \)

Change will be made as suggested.

5-30 Eq. 2: the big "dot" is a bit disturbing, it makes me think at a scalar product. I would use a simple dot or no dot at all.

Change will be made as suggested.

5-6/10 The sentence "Because of this, although advection-derived ocean heating..." is unclear to me. First, can you demonstrate the oceanic heat convergence (that is not accounted for in your framework) is a small contributor to volume changes? Second, I do not follow the logical articulation with the next sentence "Hence the surface energy...". Please clarify. In the same line, reading the recent paper by Lei et al. (2018, doi:10.1002/2017JC013548) could be useful to add up to the discussion.

Regarding your first question, Reviewer 1 also suggested that it would be sensible for the oceanic heat convergence to be properly quantified, and this will be done in the revised version.

Regarding your second question, see first our response to your major comment 8) above. We agree our wording here is confusing. The case we are making is probably best illustrated by this schematic:
The point made is that the source energy for the ocean-to-ice heat flux derives, in the main, from the surface heat flux (specifically, solar heating in summer), and not from oceanic heat convergence, over most parts of the Arctic Ocean (clearly there are regions where this is not true, e.g. near the ice edge in the Atlantic sector). Therefore the surface flux analysis (specifically, the effect of ice fraction anomalies on net SW) implicitly accounts for a large part of the ocean-to-ice heat flux.

As we mention above, additional evidence will be cited for this in the revision, as well as rewording this sentence.

Thank you for drawing our attention to Lei et al (2018), which draws together a very wide range of observational data to investigate mechanisms of sea ice growth and melt. If we have understood this study correctly, it deduces a strong role for direct solar heating in driving summer sea ice basal melting by noting an association between areas of low summer sea ice concentration and high early autumn oceanic heat fluxes (as measured by ice mass balance buoys). Hence this would also be a valuable study to quote in this context.

6-1 Please give the albedo values used.
These will be provided in the revised version: 0.535 for melting ice, 0.61 for cold ice, 0.65 for melting snow, 0.8 for cold snow.

6-3 "summarises" --> "summarise"
Change will be made as suggested.

6-10 Eq.3: please describe the meaning of the terms of the equation. In particular, what is $h_{I_{eff}}$? It is necessary to have this information in the text somewhere.
These will be described fully.
6-23 The partial derivatives are to be evaluated at a reference state, and I understand here that a mid-point between observations and model is taken ("Where observational datasets were available, the reference quantities in the partial derivative fields were calculated as model-observation means"). The authors should explain why it was done this way. I assume that the reconstructed flux error would be mathematically closer to the actual error than if the reference was taken as either the model or the observed value.

I don’t think that’s the case. For any function $f$, evaluated at two values $x_1$ and $x_2$, if we try to approximate $(x_2) - f(x_1)$ using the first term of the Taylor series evaluated at some point $\lambda x_1 + (1-\lambda)x_2$, where $0 \leq \lambda \leq 1$, the coefficient of the second Taylor series term is minimised when $\lambda=1/2$, i.e. at the midpoint of $x_1$ and $x_2$. Hence evaluating the partial derivatives at the model-observation midpoint provides our best guess. We will briefly note our reasoning for using this in the revision.

6-30 The paragraph starts by saying that 4 ensemble members were run, but Fig. 3 only shows one. Can you clarify?

Only one ensemble member is used for the analysis – this will be clarified, although as indicated above reference will be made in the revised version to results for the other three members.

8-28 The word "save" should be removed, I think.

It would probably be best replaced by ‘except for’.

10-25/28 Can you go a bit more quantitative here? From Fig. 6, the residual of the analysis can be calculated as the sum of individual contributions (the stars) and the actual flux error.

But what is the ‘actual flux error’? The surface flux anomalies wrt ERAI, ISCCP-FD, CERES cannot be regarded as such because of the observational uncertainties – this is clear, because the difference between the sum of the contributions and each surface flux anomaly is comparable in magnitude to the differences between surface flux anomalies wrt the different datasets. Each is only an estimate of the actual flux error, which cannot be exactly known – it is equally possible that the sum of the contributions is a more accurate estimate than any.

We can provide the residual of the analysis for completion with respect to ERAI, ISCCP-FD or CERES – but the actual numbers will differ greatly depending on which dataset is used.

10-31/32 The surface albedo feedback is not just a sea ice concentration thing. The melting of snow, the thinning of the ice are also key players in the surface albedo reduction, even
though ice concentration remains unchanged. Wouldn't it make more sense to include these factors as well in the definition of surface albedo feedback?

It is true that the surface albedo feedback also includes these effects. However, it would not be possible to include the effect of snow melting because of the lack of reference dataset. The effect of ice thickness on albedo is not actually modelled by HadGEM2-ES – the albedo switches abruptly to the open ocean value when the sea ice thickness falls to zero. Hence there would be two separate effects to estimate here: the direct effect of ice thickness anomalies on albedo, and the effect of HadGEM2-ES not modelling this. Given this, as well as large uncertainties in observations of the link between albedo and ice thickness for thin ice, we think this effect is outside the scope of the study. However, it will be mentioned as a possible additional contributing factor to the summer net SW bias.

11-15 The sentence "Hence the melt onset anomaly, acting alone, would induce a seasonal cycle of sea ice thickness both lower, and more amplified, than that observed..." is unclear, especially regarding the "lower" part. Can you please rephrase?

How would the following be:

'Hence the melt onset anomaly, acting alone, would induce a seasonal cycle of sea ice thickness lower in the annual mean, but also more amplified, than that observed, because the surface albedo and thickness-growth feedbacks act to translate lower ice thicknesses into faster melt and growth. For similar reasons, the downwelling LW anomaly, acting alone, would induce a seasonal cycle of sea ice thickness higher in the annual mean, and also less amplified, than that observed.'

12-1 "concludethat" –> conclude that

Change will be made as suggested.

Fig. 3. I'm puzzled by panel (a). Sea ice extent seems small. Is that because the domain "Arctic Ocean" is restricted to the seas of Fig. 1? In other observational records, like NSIDC, winter sea ice extent is more in the 14-16 million km2 range.

Yes, the extent is calculated over only the Arctic Ocean domain, like other variables in this paper, and so the winter extent appears much lower than in the well-known NSIDC figures. We think that this is appropriate, because much of the winter variability in whole-Arctic sea ice extent is due to processes in the subpolar seas, which are not relevant to the Arctic Ocean process analysis in this study.

References
Attribution of sea ice model biases to specific model errors enabled by new induced surface flux framework

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Abstract. A new framework is presented for analysing the proximate causes of model sea ice biases, demonstrated with the CMIP5 model HadGEM2-ES. Arctic sea ice extent has decreased over recent decades and has reached historic minima in late summer in recent years. Climate models project an ice-free Arctic in late summer during the 21\(^{\text{st}}\) century, with wide-ranging implications for global climate and geopolitics. However, substantial spread remains in climate model projections of the rate of sea ice decline, with drivers poorly understood. In the framework described, in this framework the sea ice volume is treated as a consequence of the integrated surface energy balance. A system of simple models allows the local dependence of the surface flux, as a function of time and space, on specific model variables (ice area, ice thickness, surface melt onset and downwelling longwave and shortwave radiation) to be described. When these are combined with reference datasets of the variable in question, it is possible to estimate the surface flux bias induced by the model bias in that variable. Specific portions of the surface flux anomaly to be attributed to individual processes by calculating for each process an induced surface flux anomaly. The method allows detailed quantification of the role played by the surface albedo and ice thickness-growth feedbacks in causing anomalous sea ice melt and growth to the role played by other forcings which can be viewed as external to the sea ice state on short timescales. It shows biases in the HadGEM2-ES sea ice volume simulation to be due to a bias in spring surface melt onset date, partly countered by a counteracting bias in winter downwelling longwave radiation. The framework is
applicable in principle to any model and has the potential to greatly improve understanding of the reasons for ensemble spread in modelled sea ice state.

1. Introduction

The Arctic sea ice cover has witnessed rapid change during the past 30 years, most notably with a decline in September extent of 1.05 x 10^6 km^2 / decade from 1986 to 2015 (HadISST1.2, Rayner et al 2003). In association with the changes in extent, evidence of declining Arctic sea ice thickness has been observed from submarine and satellite data (Rothrock et al, 2008, Lindsay and Schweiger, 2015). Arctic sea ice is also thought to have become younger on average as reserves of older ice have been lost (Maslanik et al, 2011), and the beginning of summer melt has been observed to become earlier in the year (Markus et al, 2009) and the onset of winter freezing has been observed to become later (Stammerjohn et al, 2012).

The changes have focussed much interest on model projections of Arctic sea ice, whose loss influences the climate directly through increased absorption of shortwave (SW) radiation during summer and through greater release of heat from the ocean to the atmosphere during winter (Stroeve et al, 2012b). However, substantial spread remains in model simulations of present-day Arctic sea ice, and of the long-term rate of decline under climate change (Stroeve et al, 2012a). The causes of this spread are at present poorly understood, resulting in considerable uncertainty in future projections of Arctic sea ice.

Evaluating sea ice extent or volume with respect to observations-reference datasets shows that some models clearly reproduce present-day sea ice state more accurately than others (e.g. Wang and Overland, 2012; Massonet et al, 2012; Shu et al, 2015). However, an accurate simulation of sea ice extent and volume under the present-day climate does not necessarily imply an accurate future projection of sea ice change, as a correct simulation can be obtained by accident due to cancelling model errors, or internal variability. Sea ice extent in particular is known to be a very unsuitable metric for diagnosing model performance due to its very high internal variability (Notz, 2015; Swart et al, 2015). Hence there is a need to better understand the drivers which lead a model to simulate a given Arctic sea ice state.

Ice volume is, to first order, proportional to the heat required to melt the ice, and therefore acts to integrate the surface and basal energy balance. Basal melting in the interior ice pack has been shown to derive in the main from direct solar heating of the ocean (e.g Maykut and McPhee, 1995), while basal freezing derives principally from conduction of energy upward through the ice (Perovich and Elder, 2002); this implies that the surface energy balance (SEB) contains the principal sources and sinks of energy for the sea ice on an Arctic-wide scale. However, a complex two-way relationship exists between sea ice thickness and surface energy balance, via the surface temperature and surface albedo, giving rise to the thickness-growth feedback (Bitz and Roe, 2004) and the surface albedo feedback (Bitz, 2008), both of which exert first-order control on the sea ice state. Hence many components of the SEB cannot be viewed as independent of the sea ice state in any meaningful sense.

This study, which presents a new framework to investigate the causes of modelled sea ice bias anomalies, is motivated by a desire to separate, to first order, many external drivers of the SEB (and hence the sea ice state) from the thickness-growth and albedo feedbacks, and thereby better understand the processes that result in a
particular modelled sea ice state being simulated. The analysis uses as a case study the four members of the historical ensemble of the coupled model HadGEM2-ES, a member of the CMIP5 historical ensemble. CMIP5 model which simulates anomalously low annual minimum ice extent, and which simulates an ice volume which is both too low in the annual mean and too amplified in the seasonal cycle, a similar behaviour to that identified by Shu et al (2015) in the CMIP5 ensemble mean.

In the framework presented, the total surface flux principal components of the SEB, shortwave and longwave radiative fluxes, are expressed in terms of key Arctic climate variables as functions of space and time using two simple models, for the freezing and melting seasons respectively, which are shown to capture well the large-scale spatial and seasonal variation of the surface flux. With the use of the simple models, the local dependence of modelled surface flux on key variables can be described. Hence, using and reference datasets for the climate variables, the model biases in surface flux induced by each climate variable can be estimated. In this way, biases in ice growth and melt caused by the sea ice albedo feedback, the ice thickness-growth feedback, and of various external factors, to be separately quantified. In this way it can be seen how model biases in the external drivers are able to produce a particular sea ice state, offering a valuable tool for setting sea ice state biases in context, and for understanding spread in sea ice simulation within multimodel ensembles.

In Section 2, the HadGEM2-ES model and the observational datasets used are described in turn, and the induced surface flux method is introduced. In Section 3 the sea ice and surface radiation simulations of HadGEM2-ES are evaluated. In Section 4, the induced surface flux method is introduced, and in Section 5 the induced surface flux analysis is applied to HadGEM2-ES, allowing quantification of the role played by biases in specific Arctic climate variables in causing anomalous ice growth and melt. In Section 6 the implications of the results are discussed, in particular the mechanisms by which the identified external drivers identified cause determine the modelled sea ice state, and the likely drivers behind the corresponding model biases. Conclusions are presented in Section 7.

2. Model and observational data and methods

2.1 The HadGEM2-ES model

HadGEM2-ES is a coupled climate model employing additional components to simulate terrestrial and oceanic ecosystems, and tropospheric chemistry (Collins et al, 2011). It is part of the ‘HadGEM2’ family, a collection of models that use the HadGEM2-AO coupled atmosphere-ocean system. HadGEM2-AO is developed from HadGEM1 (Johns et al, 2006), a coupled atmosphere-ocean model whose sea ice extent simulation was recognised as being among the closest to observations out of the CMIP3 ensemble models (Wang and Overland, 2009). While the atmospheric and ocean components of HadGEM2-ES contain a large number of improvements relative to HadGEM1, many of these targeted at improving simulations of tropical weather, the sea ice component is very similar to that of HadGEM1 except for three minor modifications (Martin et al, 2011, table A4).
A fundamental feature of the sea ice component of HadGEM2-ES, which is important for the analysis described in Section 4 below, is that it includes a subgridscale sea ice thickness distribution (Thorndike et al., 1975). In this formulation, ice in each grid cell is separated into five thickness categories with boundaries at 0, 0.6m, 1.4m, 2.4m, 3.6m and 20m, each with its own area, thermodynamics and surface exchange calculations. It also includes elastic-viscous-plastic sea ice dynamics (Hunke and Dukowicz, 1997) and incremental remapping (Lipscomb and Hunke, 2004). The thermodynamic component is a zero-layer model, with no heat capacity, described in the appendix of Semtner (1976). The insulating effect of snow is modelled by means of a single layer with conductivity 0.33Wm⁻¹K⁻¹, also with no heat capacity (although sensible heat storage is parameterised in the top 10cm of the snow-ice column during surface exchange calculations, to aid stability). Most processes are calculated in the ocean model, but the surface energy balance (SEB) calculations are carried out in the atmosphere model, which passes top melting flux and conductive heat flux to the ocean model as forcing for the remaining components. A more complete description of the sea ice component can be found in McLaren et al (2006).

This study uses the first four ensemble members of the CMIP5 ‘historical’ experiment of HadGEM2-ES, forced with observed solar, volcanic and anthropogenic forcing from 1860 to 2005. The period 1980-1999 is chosen for the model evaluation, as a period which predates much of the recent rapid Arctic sea ice loss, and is hence at least partially independent of the period normally used to evaluate sea ice trends. It has the added advantage of being with its associated climatic changes, and which is recent enough to allow the use of a reasonable range of observational data. All analysis is carried out with data restricted to the Arctic Ocean region, shown in Figure 1.

### 2.2 Observational data

Uncertainty in observed variables tends to be higher in the Arctic than in many other parts of the world. There are severe practical difficulties with collecting in situ data on a large scale over regions of ice-covered ocean. While satellites have in many cases been able to produce Arctic-wide measurements of some characteristics, most notably sea ice concentration, the relative lack of in situ observations against which these can be calibrated means knowledge of the observational biases is limited. Reanalysis data over the Arctic is also more subject to the reanalysis model errors than in other regions, due to errors in atmospheric forcing, and the existence of there are fewer direct observations available for assimilation (Lindsay et al, 2014). The approach of this study is to use a wide range of observational data to evaluate modelled sea ice state and surface radiative fluxes, and to use as reference datasets for the induced surface flux framework, using the small number of in situ validation studies to set results in context as far as possible.

To evaluate modelled sea ice fraction, we use the HadISST1.2 dataset (Rayner et al, 2003), derived from passive microwave observations. To evaluate modelled sea ice thickness Arctic-wide, we use the ice-ocean model PIOMAS (Schweiger et al, 2011), which is forced with the NCEP reanalysis and assimilates ice concentration data. Laxon et al (2013) and Wang et al (2016) found PIOMAS to estimate anomalously low winter ice thicknesses compared to satellite observations in some years. In particular, Wang et al (2016) found PIOMAS to have a mean bias of -0.31m relative to observations from the ICESat (Ice, Cloud and land Elevation Satellite) satellite laser sensor. To set the PIOMAS comparison in context, we use two additional datasets to
evaluate the model over smaller regions; measurements from radar altimetry aboard the ERS satellites from 1993-2000 (Laxon et al, 2003), limited to latitudes below 82°N; and estimates compiled by Rothrock et al (2008), derived from a multiple regression of submarine transects over the Central Arctic Ocean from 1975-2000, constrained to be seasonally symmetric.

To evaluate modelled surface radiative fluxes across the whole Arctic Ocean, three datasets are used. Firstly, we use the CERES-EBAF (Clouds and Earth’s Radiant Energy Systems – Energy Balanced And Filled) Ed2.7 dataset (Loeb et al, 2009), based on direct measurements of top-of-atmosphere radiances from EOS sensors aboard NASA satellites, available from 2000 – present. Secondly, we use the ISCCP-FD (International Satellite Cloud Climatology Project FD-series) product derived from the ISCCP-D cloud product described below using a radiative transfer algorithm (Zhang et al, 2004). Lastly, we use the ERA-Interim (ERAI) atmospheric reanalysis dataset, which provides gridded surface flux data from 1979-present using a reanalysis system driven by the ECMWF (European Centre for Medium-range Weather Forecasts) IFS forecast model and the 4D-Var atmospheric data assimilation system (Dee et al, 2011).

To the authors’ knowledge, in-situ validation of these datasets in the Arctic has been quite limited, but Christensen et al (2016) found CERES to perform quite well relative to other products, albeit underestimating downwelling LW fluxes from November – February by 10-20 Wm⁻² relative to in situ observations at Barrow (Alaska). Liu et al (2005) found ISCCP-FD to simulate SW radiative fluxes fairly accurately relative to observations from SHEBA, but to underestimate downwelling SW fluxes in spring by over 30 Wm⁻², also overestimating downwelling LW fluxes in winter by around 40 Wm⁻². Finally, Lindsay et al (2014) identified ERAI as producing a relatively accurate simulation of surface fluxes compared to in situ observations at Barrow (Alaska) and Ny-Ålesund (Svalbard), although tending to underestimate downwelling SW fluxes in the spring by up to 20 Wm⁻² and overestimate downwelling LW fluxes in the winter by around 15 Wm⁻².

In addition to the datasets above, in section 4 we make use of satellite estimates of date of melt onset over sea ice (Anderson et al, 2012), also derived from passive microwave sensors; and in section 5, the CERES-SYN dataset (Rutan et al, 2015), similar to CERES-EBAF but available at higher temporal resolution, is used to examine modelled surface radiation evolution during May in more detail.

2.3 Calculating induced surface flux anomaly
Because the latent heat of sea ice is an order of magnitude greater than the sensible heat required to raise the ice to the melting temperature, ice volume is very nearly proportional to the heat required to melt the ice. Ice volume therefore acts to integrate the surface and basal energy balance, and is largely determined by the fluxes at these interfaces. Across much of the Arctic the sea ice is insulated from the main source of heat energy from beneath, the warm Atlantic water layer, by fresh water derived mainly from river runoff (e.g. Serrée et al, 2006; Stroeve et al, 2012b). Because of this, although advection-derived ocean heating can be important in setting the wintertime ice edge (Bitz et al, 2005), direct solar heating of the ocean is likely to be an order of magnitude higher in accounting for basal melting of the sea ice, as observed by Maykut and McPhee, 1995, McPhee et al,
The surface energy balance is composed of four radiative fluxes (downwelling and upwelling SW and LW), two turbulent fluxes (sensible and latent) and of an additional flux due to snowfall. The radiative fluxes dominate in terms of magnitude in both observations (e.g. Persson et al., 2002) and models (the relative magnitude of radiative and turbulent fluxes in HadGEM2-ES is demonstrated in Figure 2). For this reason, and because no large-scale observations of the turbulent fluxes or snowfall exist, in this study we analyse processes affecting only the radiative fluxes.

We express the net radiative flux \( F_{SW} + F_{LW} \), where \( F_{SW} \) and \( F_{LW} \) are fluxes of net surface SW and LW radiation respectively, as functions of key Arctic climate variables using two formulae derived in Appendix A, valid for the ice freezing and melting seasons respectively (as different processes are important depending on the season). The formulae are applied to monthly means of data in model grid cells measuring tens of km across, within which the relevant variables are not necessarily constant or uniform. Errors associated with this spatial and temporal extrapolation are briefly discussed in Appendix B.

During the ice freezing season

\[
F_{SW} + F_{LW} = \sigma_T F_{SW} \frac{S_{SW} + S_{LW} + C + BT_s}{1 - B \cdot R_{SW}} \tag{1}
\]

where \( F_{SW} \) and \( F_{LW} \) are fluxes of downwelling SW and LW radiation respectively, \( R_{SW} = \frac{h_i + h_s}{k_i + k_s} \) is the thermal insulation of the ice and snow column (\( k_i \) and \( k_s \) being ice and snow conductivity respectively, \( h_i \) and \( h_s \) ice and snow thickness), \( B \) is the linearised rate of upwelling longwave dependence on surface temperature at 0°C, \( T_s \) ice base temperature in Celsius, \( C \) upwelling longwave at 0°C and \( \sigma_T \) ice albedo.

During the ice melting season

\[
F_{SW} + F_{LW} = F_{SW} + F_{LW} + C + F_{SW} \frac{1 - \alpha_{sea} - \alpha_{snow} - \alpha_{sea} - \alpha_{sea} - \alpha_{sea} - \alpha_{sea} - \alpha_{sea} - \alpha_{sea}}{1 - B \cdot R_{SW}} \tag{2}
\]

where \( \alpha_{sea} \), \( \alpha_{snow} \) and \( \alpha_{cold} \) indicate albedo of open sea, melting snow on sea ice and cold snow on sea ice respectively, \( \alpha_{sea} \) and \( \alpha_{cold} \) indicate the corresponding area fractions.

(1) and (2) summaries the dependence of surface radiative flux, assumed to account for the major part of the surface flux, on ice thickness, snow thickness, downwelling longwave and shortwave fluxes, ice fraction, snow fraction and melting surface fraction, during the freezing and melting seasons respectively. Hence, at each point
in space and time, if a model bias in one of these quantities is known, an associated surface flux anomaly can be estimated to first order by multiplying the bias in the quantity by the partial derivative of (1) or (2) with respect to that quantity, calculated at some reference state. For example, given a model anomaly in ice thickness \( \left( h_{\text{MODEL}} - h_{\text{OBS}} \right) \) during the freezing season, a resulting induced surface flux anomaly can be calculated as

\[
\Delta F_{\text{sw}} = \left( h_{\text{MODEL}} - h_{\text{OBS}} \right) \cdot \frac{\partial F_{\text{sw}}}{\partial h_{\text{eff}}} = \left( h_{\text{OBS}} - h_{\text{MODEL}} \right) \cdot \frac{B}{k} \cdot \frac{L_{\text{W,REF}}}{\left(1 - R_{\text{eff}}^2\right)} + B T_{\text{b}}
\]

(Here the SW flux is neglected for clarity).

In a similar way, given a model anomaly of melting surface fraction for any point in space and time, an induced surface flux anomaly can be calculated in a similar way to above, by multiplying the model anomaly in melt onset occurrence by the partial derivative of (7) with respect to melt onset occurrence:

\[
\Delta F_{\text{sw, melt onset}} = \left( f_{\text{cold}} - f_{\text{OBS, cold}} \right) \cdot \text{SW,REF} \cdot \left( T_{\text{cold}} - T_{\text{cold, obs}} \right)
\]

Using equations (1) and (2), an induced surface flux anomaly due to downwelling SW, downwelling LW, ice thickness, ice fraction and melt onset occurrence was calculated for all model grid cells, for all months in the period 1980-1999. Firstly, a model anomaly in the relevant climate variable was calculated by bilinearly regressing the chosen observational dataset to the model grid, and then averaging this across all dataset years into a climatology. The model bias in the relevant quantity for each grid cell and model month was then calculated by subtracting this climatology. Secondly, to convert the modelled bias in the relevant quantity to an induced surface flux anomaly, the fields were multiplied by the partial derivative of equation (1) (where modelled \( T_{\text{ct}} \leq -1^\circ \text{C} \)) or equation (2) (where modelled \( T_{\text{ct}} \geq -1^\circ \text{C} \)). Where observational datasets were available, the reference quantities in the partial derivative fields were calculated as model-observation means. Where no such datasets were available (e.g. for snow thickness), the model field was used. No flux anomalies due to snow thickness or snow fraction were computed, due to a lack of observational datasets with which to calculate model anomaly in these quantities.

### 3. Evaluating sea ice and surface radiation in HadGEM2-ES

From 1980-1999, the four members of the HadGEM2-ES historically-forced ensemble simulate a mean September sea ice extent of 5.78 x 10^6 km^2, with ensemble standard deviation of 0.24 x 10^6 km^2. By comparison, the mean observed September sea ice extent over this period was 6.88 x 10^6 km^2 according to the HadISST1.2 dataset. Over the reference period, therefore, modelled September sea ice extent is systematically lower than that observed (Figure 2a).

Mean ice thickness is consistently lower than that estimated by PIOMAS for the Arctic Ocean region (Figure 2b), with the highest bias anomalies of -0.4m occurring in October, close to the minimum of the annual cycle, and a near-zero bias anomaly in May, close to the maximum. Modelled ice thickness is also biased low relative to Envisat the ERS satellite measurements (Figure 2c), with thickness bias anomalies ranging from
1.06 in November to -0.22 in April, and relative to the submarine data (Figure 23d), with thickness biases ranging from -1.5 m in August to -0.8 m in January and May. Hence it is very likely that ice thickness in HadGEM2-ES is biased low in the annual mean, with biases tending to be higher when ice thickness is lower. In other words, the ice thickness annual cycle of HadGEM2-ES is likely to be too amplified, with both anomalously high ice melt during the summer and ice growth during the winter.

Maps of the ice thickness bias in April and October (Figure 23b-d) show agreement that the low ice thickness bias is smaller on the Pacific side of the Arctic than on the Atlantic side of the Arctic, becoming very small or even positive in the Beaufort Sea. There is also striking agreement in the spatial pattern of the amplification bias of the seasonal cycle, as diagnosed by April-October ice thickness difference (Figure 34). All three ice thickness datasets show the HadGEM2-ES ice thickness seasonal cycle to be too amplified across much of the Arctic, by up to 1 m in the Siberian shelf seas; in addition, all show that in the Beaufort Sea, the amplification is nonexistent or even negative. There is clear association between areas where modelled annual mean ice thickness is biased low, and areas where the modelled seasonal cycle is overamplified, and vice versa.

In the following discussion of radiative fluxes, the convention is that positive numbers denote a downwards flux. Fluxes of downwelling SW radiation are higher in HadGEM2-ES than in all observational estimates during the spring (Figure 45a-c), with May biases of 226, 434 and 53 Wm$^{-2}$ relative to CERES, ERAI, and ISCCP-FD respectively. We note that as ERAI and ISCCP-FD have been found to underestimate downwelling SW during spring at specific locations, the true model bias is perhaps more likely to lie towards the lower end of these estimates. During the summer, upwelling SW radiation is consistently lower in magnitude than in HadGEM2-ES, with June biases of 164, 374 and 440 Wm$^{-2}$ with respect to ERAI, CERES and ISCCP-FD respectively (a positive bias in an upward flux demonstrates that the model is too low in magnitude). There is no consistent signal for a low bias in downwelling SW during the summer, suggesting a model surface albedo bias. The effect is that modelled net downward SW flux is too large with respect to all observational datasets in May and June, and with respect to some in July and August. Relative to CERES, the May downwelling SW bias displays no clear spatial differentiation over the Arctic Ocean (Figure 45a), but the June upwelling SW bias, and hence the net SW bias, tend to be somewhat higher in magnitude towards the central Arctic (Figure 45b-c).

Fluxes of longwave (LW) radiation are lower in magnitude in HadGEM2-ES throughout the winter than in all observational datasets (Figure 45d-f). For downwelling LW, the mean model biases from December-April are -165, -22 and -2240 Wm$^{-2}$ for ERAI, CERES and ISCCP-FD respectively; for upwelling LW, the biases are 11, 158 and 186 Wm$^{-2}$ for CERES, ERAI and ISCCP respectively. Because the downwelling LW biases vary more than the upwelling LW biases, there is uncertainty in inferring a model bias in net downwelling LW; ISCCP suggests a large model bias of -224 Wm$^{-2}$, CERES a smaller bias of -11 Wm$^{-2}$, while ERAI suggests no bias at all. As in situ studies have shown both underestimation (by CERES) and overestimation (by ERAI and ISCCP-FD) of downwelling LW in winter, there is no clear indication as to where the true model bias lies in this quantity may lie. Maps of the downwelling and net down LW bias relative to CERES in February (Figure 45d-f) show the bias tends to be somewhat higher towards the North American side of the Arctic, and lower on the Siberian side.
In summary, there is evidence of a low bias in net downward LW during the winter, and a high bias in net downward SW during the summer, each of order of magnitude ~10 Wm$^{-2}$. This is consistent with surface radiation fluxes being the likely first-order cause of the amplified sea ice thickness seasonal cycle. In the next section we describe the process by which surface radiation biases can be attributed to particular model processes by calculating induced surface flux biases. We attempt to attribute the surface radiation biases to particular processes using the methodology described in Section 2.

4. Calculating induced surface flux bias: Methods

In this section, and throughout the rest of the paper, a difference between a model simulation of a particular variable, and any reference dataset for that variable, is referred to as a ‘bias’. In a similar way, the difference in model surface flux judged to arise from the difference in a particular variable relative to a reference dataset is referred to as an ‘induced surface flux bias’. Attention is drawn to the fact that, due to observational inaccuracy, true model bias relative to the real world may be somewhat different from the biases described in this way.

Due to the latent heat of sea ice being an order of magnitude greater than the sensible heat required to raise the ice to the melting temperature, ice volume is very nearly proportional to the heat required to melt the ice. Ice volume therefore acts to integrate the surface and basal energy balance, and is largely determined by the fluxes at these interfaces. Across much of the Arctic the sea ice is insulated from the main source of heat energy from beneath, the warm Atlantic water layer, by fresh water derived mainly from river runoff (e.g. Serreze et al, 2006; Stroeve et al, 2012b). Because of this, in the Arctic Ocean interior direct solar heating of the ocean is likely to be an order of magnitude higher in accounting for basal melting of the sea ice, as observed by Maykut and McPhee, 1995, McPhee et al, 2003 and Perovich et al, 2008, and modelled by Steele et al (2010) and Bitz et al (2005). In particular, it has been found that in HadGEM2-ES oceanic heat convergence is of negligible importance to the sea ice heat budget (Keen et al, 2018). Hence the surface energy balance in the Arctic Ocean is of primary importance in controlling the evolution of sea ice volume.

We use a system of well-understood simple models, similar to those used in Thorndike (1992), to estimate, for each model grid cell, and month within the period, the rate at which the surface flux would be expected to change with a particular model variable. For each model grid cell and month, we construct a function $F_{i,j} = g_{i,j}(v_1, ..., v_i) F_{i,j}$ being surface flux, where the $v_i$ are climate variables that affect the surface flux on timescales shorter than that on which they affect each other, and can therefore be said to be independent for the purposes of this analysis. In this way, at each model grid cell and month the rate at which the surface flux depends on variable $v_i$ can be approximated by $\partial F_{i,j}/\partial v_i$. Given a reference dataset for variable $v_i$, it then becomes possible to estimate, for each point in time and space, the surface flux bias induced by the bias in $v_i$ as $\partial F_{i,j}/\partial v_i \left( v_{\text{MODEL},i,j} - v_{\text{REFERENCE},i,j} \right)$. The chief advantage of this method is that the resulting fields of induced surface flux bias can then be averaged in time or space to determine the large-scale effects of particular model biases, effectively bypassing nonlinearities in surface flux dependence.
The functions \( g_{x,i,j} \) are constructed as follows. Firstly, a model grid cell in a particular month is classified as freezing or melting depending upon whether the monthly mean surface temperature is greater or lower than -2°C. If the grid cell is classified as freezing, the surface flux is approximated as

\[
F_{\text{ice}} = g^{\text{u}}_{x,i,j} = \alpha_{\text{ice}} (F_{\text{atmos-ice}} + BT_{\text{ice}}) \sum_{\text{cat}} \gamma_{\text{cat}}^\text{ice} (1 - BR_{\text{cat}}^\text{ice})^{-1} + (1 - \alpha_{\text{ice}})F_{\text{atmos-ocean}} \tag{1}
\]

Here \( F_{\text{atmos-ice}} = F_{\text{LW}} - \varepsilon_{\text{ice}} \sigma T_{\text{ice}}^4 + F_{\text{SW}} + (1 - \alpha_{\text{ice}})F_{\text{ice}} \), where \( \alpha_{\text{ice}} \) is ice area, \( F_{\text{LW}} \) and \( F_{\text{SW}} \) downwelling longwave, shortwave and sensible heat flux respectively, the latter evaluated over only the ice-covered portion of the grid cell, \( \varepsilon_{\text{ice}} \) ice emissivity, \( \sigma = 5.67 \times 10^{-8} \text{Wm}^{-2}\text{K}^{-4} \) the Stefan-Boltzmann constant, \( T_{\text{ice}} \) is monthly mean surface temperature and \( \alpha_{\text{ice}} \) is mean surface albedo over ice.

Bolzmann constant, \( T_{\text{ice}} \), is monthly mean surface temperature and \( \alpha_{\text{ice}} \) is mean surface albedo over ice.

\[
F_{\text{atmos-ocean}} = F_{\text{LW}} - \varepsilon_{\text{ice}} \sigma T_{\text{ice}}^4 + F_{\text{SW}} + (1 - \alpha_{\text{ice}})F_{\text{ice}} \quad \text{where} \quad \varepsilon_{\text{ocean}} \quad \text{is ocean surface emissivity,} \quad T_{\text{ocean}} \quad \text{ocean surface temperature (assumed to be -1.8°C),} \quad \text{and} \quad F_{\text{SW}}, \quad F_{\text{LW}}, \quad F_{\text{heat}} \quad \text{sensible and latent heat flux respectively over the ice-free portions of the grid cell,} \quad B = 4\varepsilon_{\text{ice}} \sigma T_{\text{ice}}^4 \quad \text{approximates the local rate of dependence of surface flux on surface temperature,} \quad T_b \quad \text{is ice base temperature,} \quad \gamma_{\text{cat}}^\text{ice} \quad \text{the area of ice in ice thickness category cat as a fraction of total ice area, where cat ranges from 1 to 5, and}
\]

\[
R_{\text{cat}} = \frac{h_{\text{cat}} + h_{\text{snow}}}{k_{\text{cat}}} \quad \text{is the thermal insulation of the snow-ice column in category cat, where cat ranges from 1 to 5,} \quad k_{\text{ice}} \quad \text{and} \quad k_{\text{snow}} \quad \text{being ice and snow conductivity respectively,} \quad h_{\text{cat}} \quad \text{and} \quad h_{\text{snow}} \quad \text{local ice and snow thickness,}
\]

If the grid cell is classified as melting, the surface flux is approximated as

\[
F_{\text{mel}} = g^{\text{u}}_{x,i,j} = F_{\text{LW}} - \varepsilon_{\text{ice}} \sigma T_{\text{ice}}^4 + F_{\text{SW}} + (1 - \alpha_{\text{ice}})F_{\text{ice}} \quad \text{where} \quad T_b = 0^\circ \text{C} \tag{2}
\]

The ice surface albedo \( \alpha_{\text{ice}} \) is further expressed as

\[
\alpha_{\text{ice}} = (\alpha_{\text{melt-ice}} - \alpha_{\text{cat}}) + I_{\text{melt}} (\alpha_{\text{melt-snow}} - \alpha_{\text{melt-ice}}) + (1 - \gamma_{\text{melt}}) (1 - I_{\text{melt}}) (\alpha_{\text{cold-ice}} - \alpha_{\text{melt-ice}}) + (1 - \gamma_{\text{melt}}) I_{\text{melt}} (\alpha_{\text{cold-snow}} - \alpha_{\text{melt-snow}}) \tag{3}
\]

Here \( \alpha_{\text{cat}} = 0.06, \alpha_{\text{melt-ice}} = 0.535, \alpha_{\text{melt-snow}} = 0.65, \alpha_{\text{cold-ice}} = 0.61 \) and \( \alpha_{\text{cold-snow}} = 0.8 \) denote the parameterised albedos of open water, melting ice, melting snow, cold ice and cold snow respectively, and
\( \gamma_{\text{melt}} \) denotes melting surface fraction as a fraction of ice area, while \( I_{\text{snow}} \) is an indicator for the presence of snow that is set to 1 or 0 depending on whether monthly mean snow thickness exceeds 1 mm.

The derivation of the formulae is briefly described. The surface flux is composed of four radiative fluxes (downwelling and upwelling SW and LW), two turbulent fluxes (sensible and latent) and of an additional flux due to snowfall (which affects the surface flux as it represents a transfer of negative latent heat, since snow lying on ice changes the enthalpy of the snow-ice system). Hence

\[
F_{\text{sfc}} = \left(1 - \alpha_{\text{ice}} F_{\text{SW}} \right) + F_{\text{LW}} + F_{\text{snow}} + F_{\text{lat}} + F_{\text{snowfall}}
\]

is used as a starting point from which the derivation of (2) follows in the melting season, assuming a surface temperature of 0°C and neglecting the snowfall contribution. (3) is designed to mimic the calculation of ice albedo in HadGEM2-ES, which parameterises the effect of meltponds after Curry (2001), reducing albedo linearly as surface temperature rises from -1°C to 0°C. (1) is derived by considering separately the contributions to \( F_{\text{sfc}} \) from the area of the grid cell covered by each ice category (and by open water). For each ice category, the conductive flux through the ice is assumed to be uniform; the dependence of \( F_{\text{sfc}} \) on surface temperature is then linearised, using monthly mean surface temperature at each grid point, \( T_{\text{sfc-REF}} \), as a reference about which to take the linearisation. By setting the conductive flux equal to \( F_{\text{sfc}} \), the variable \( T_{\text{sfc}} \) is eliminated. Finally, the contributions to \( F_{\text{sfc}} \) are multiplied by category ice area and summed. In deriving (1), the contributions of the snowfall flux and of the latent heat flux over ice are neglected.

In this way, using equations (1)-(3), we construct the functions \( g_{s,\ell} \), which depend on downwelling LW, downwelling SW, sensible heat flux, category ice thickness, category ice area (freezing cells), total ice area (melting cells), snow thickness, snow area and surface melt onset, variables which have the required property of tending to affect the surface flux on timescales shorter than that on which they affect each other. Hence at each point in space and time the rate of dependence of surface flux on each variable can be approximated by

\[
\frac{\partial g_{s,\ell}}{\partial \nu_i}.
\]

We describe for the case of three variables how this process can be used to estimate the surface flux bias induced by biases in that variable, firstly for the variable of melting surface fraction (for simplicity, we describe only the process over grid cells judged to be melting). Model daily surface temperature fields are used to judge, for each month of the year, the average melting surface fraction in each grid cell. The satellite-derived observational estimates of surface melt onset described in Section 2.2 are used to produce a climatology of melting surface fraction for each month and grid cell, and this is subtracted to produce a model bias. This bias is then multiplied by the partial derivative of equation (2) with respect to melting surface fraction, evaluated with monthly mean fields of \( F_{\text{SW}} \) and \( I_{\text{snow}} \), to produce a monthly mean field of surface flux bias induced by the model bias in melting surface fraction.
By a similar method, the effect of downwelling LW radiation on surface flux can be estimated, illustrated here using CERES as a reference dataset (in section 4 below the analysis is performed using multiple datasets) to produce fields of model bias in downwelling LW radiation. For freezing grid cells, these are then multiplied by

$$\sum_{\text{cat}} a_{\text{cat}} \left(1 - BR_{\text{cat}}^\text{ice} \right)^{-1} \left| \sum_{\text{cat}} a_{\text{cat}} \right|$$

the partial derivative of (1) with respect to downwelling LW, to produce fields of surface flux bias induced by model bias in downwelling LW. For melting grid cells, the induced surface flux bias is equal to the downwelling LW bias, as the surface temperature does not change in response to the bias.

The most complex variable to analyse in this way is the ice thickness. Ice thickness strongly affects the surface flux in the freezing season; thicker ice is associated with less conduction, a colder surface temperature and a weaker negative surface flux, and hence reduced ice growth. However, it appears in equation (1) only implicitly, in the form of the individual category mean thicknesses $\bar{h}_{\text{cat}}$. To use this equation to estimate the effect of ice thickness biases on surface flux, a method of estimating the way biases are distributed amongst thickness categories is needed. Given an estimated model bias in mean thickness $\bar{h}_{\text{ice}}$, it can be argued that the least arbitrary approach is to estimate the model bias in each thickness category to be $\bar{h}_{\text{cat}}$ also (i.e. the thickness distribution is uniformly shifted to higher, or lower values). However, this leads to unphysical results at the low end of the distribution; in the case of a negative bias, it implicitly assumes the creation of sea ice of negative thickness; in the case of a positive bias, it assumes that no sea ice of thicknesses between 0m and $\bar{h}_{\text{ice}}$ m exists.

Hence we use a slightly modified approach. The model bias in the lowest thickness category is estimated to be $\bar{h}_{\text{ice}}/2$, equivalent to translating the top end of the category by $\bar{h}_{\text{ice}}$, but allowing the lower end to remain at 0.

The model biases in the other four categories are then estimated to be $\bar{h}_{\text{ice}}^\text{cat} - \bar{h}_{\text{ice}}/2$, i.e. the translation is increased to ensure that the mean ice thickness bias remains correct. Following this, we iterate through the categories, identifying grid cells where the bias is such that a negative category sea ice thickness in the reference dataset is implied; in these cells, the bias is reduced such that the reference thickness in that category becomes 0, and the bias in the remaining categories is increased proportionally to ensure the mean sea ice thickness bias remains correct.

Hence we create, for each category, fields of sea ice thickness bias. These are multiplied by the partial derivative of equation (1) with respect to category ice thickness, $(A + BT)^\text{ice} a_{\text{cat}} \left(1 - BR_{\text{cat}}^\text{ice} \right)^{-1} \left( \sum_{\text{cat}} a_{\text{cat}} \right)^{-1}$, to create fields of induced surface flux bias for each category. These are then summed to obtain the total induced surface flux bias due to ice thickness bias.

The process of calculating induced surface flux bias is illustrated in Figure 5 for example months for these three variables. Figure 5a-c illustrates the melt onset analysis. Figure 5a shows the HadGEM2-ES bias in melting
surface fraction for the month of June 1980, relative to the NSIDC climatology; the bias is generally positive, reflecting melt onset modelled earlier than observed during this month. Figure 5b shows the field of rate of change of surface flux with respect to melt onset occurrence (effectively downwelling SW multiplied by the difference in parameterised albedos); this tends to be higher in the Central Arctic, reflecting a greater tendency to clear skies here. Finally Figure 5c shows the product of these two fields, the modelled surface flux bias induced by the model bias in melt onset. This is also generally positive, by up to 25 Wm\(^{-2}\) in the central Arctic, reflecting the greater absorption of SW radiation induced by the early melt onset.

Figures 5d-f demonstrate the same process for the downwelling LW radiation in January 1980, using CERES as reference dataset. Modelled downwelling LW radiation is seen to be considerably lower in magnitude than that observed by CERES for the 2000-2013 period, by up to 30 Wm\(^{-2}\) in many parts of the Central Arctic (Figure 2d). The rate of change of surface flux with respect to downwelling LW is shown to be higher (closer to 1) in regions of thinner ice (Figure 2e). This has the result that the induced surface flux bias is greatly reduced relative to the downwelling LW bias in regions of thicker ice (Figure 2f), reflecting the lower efficiency of ice creation in regions of thicker ice; the bias is below 10 Wm\(^{-2}\) over much of the Arctic, only approaching 20 Wm\(^{-2}\) in the Barents and Kara seas.

Figures 5g-i demonstrate the process of calculating surface flux bias induced by the bias in ice thickness in model category 1 (0-0.6m) for the month of January 1980, using PIOMAS as reference dataset. Modelled ice thickness tends to be thinner than estimated by PIOMAS over much of the Arctic for this month, except for an area on the Pacific side of the Arctic; as described above, the bias in category 1 is assumed to be half the total bias. Figure 5h shows the rate of change of surface flux with respect to category 1 ice thickness, which tends to be high in regions where category 1 ice covers higher fractions of the grid cell, generally near the ice edge.

## 5. Calculating induced surface flux bias: Results

### 4. Induced surface flux anomaly

We calculate fields of surface flux anomaly for each month in the model period 1980-1999 induced by model anomalies in downwelling LW, downwelling SW, ice thickness, ice fraction and melt onset occurrence, using in turn CERES, ISCCP-FD and ERAI as reference datasets. The resulting fields are integrated over the Arctic Ocean and averaged over the model period 1980-1999 to produce, for each climate variable, a seasonal cycle of average induced surface flux anomaly. In Figure 6, the seasonal cycles of each component of the surface flux anomaly are shown as a set of stacked barplots, with the results using CERES, ISCCP-FD and ERAI as reference shown from left-right for each month. These are compared, firstly to the net radiative flux anomalies implied by the direct radiation evaluation above (using again CERES, ISCCP-FD and ERAI in turn), and secondly to the sea ice latent heat flux uptake anomaly implied by the ice volume anomalies relative to PIOMAS.

Consistent with the direct radiation comparison and the ice volume anomalies, the induced surface flux anomalies usually sum to positive values during the summer and negative values during the winter. This implies that model biases in the processes investigated tend to cause anomalous ice melt during the summer and anomalous ice growth during the winter. Despite the large discrepancies between the radiative flux estimates of
CERES, ISCCP-FD and ERAI, the choice of radiation dataset does not significantly affect any of the induced surface flux anomalies save those due to downwelling SW and LW radiation in the summer, where very large spread is apparent.

Induced surface flux anomalies are generally small in May. In June an anomaly of 8 Wm$^{-2}$ is induced by melt onset occurrence; the SSMR observations do not show melt onset to occur in the central Arctic until late June on average, while HadGEM2-ES allows the entire Arctic Ocean to reach the melting temperature by the end of May, inducing a large model anomaly in surface albedo. June also displays a small but growing surface flux anomaly due to ice fraction anomaly, consistent with the anomalously fast ice melt of HadGEM2-ES. By August, although there is no remaining melt onset occurrence anomaly, the surface flux anomaly due to ice fraction is 9-10 Wm$^{-2}$, allowing the positive surface flux anomaly to be maintained throughout the summer.

As noted above, the surface flux anomaly due to downwelling SW and LW radiation varies greatly with the use of observational dataset throughout the summer. Existing in situ validation studies of ERAI (Lindsay et al., 2014) and of CERES (Christensen et al., 2016) show both datasets to model downwelling LW accurately to within 15 Wm$^{-2}$, although CERES is shown to be biased high in June and low in August, while ERAI is biased high in July. Even if these results were representative of the broader Arctic Ocean, it would be hard to interpret the true net effect of the combined LW and SW anomalies, as there will tend to be opposite in sign and of similar orders of magnitude. It is concluded that it is not possible to determine the net effect of downwelling radiative anomalies on surface flux during the summer with current observational data.

As the surface of the Arctic Ocean begins to cool in early autumn, a growing negative anomaly due to the now large deficit in ice thickness begins to appear, reaching a maximum of -6 Wm$^{-2}$ in November. The anomaly reduces sharply through the winter as model anomalies in ice thickness with respect to PIOMAS become smaller. From November – April, the downwelling LW induces an additional negative surface flux anomaly, with a mean value of -4 Wm$^{-2}$ indicated from December – February. It is superficially surprising that the choice of radiation dataset does not greatly affect the total induced surface flux, given the large spread in net LW radiation during winter estimated by each dataset. This occurs because in the induced surface flux analysis, the upwelling LW flux anomaly is calculated from other variables, and is therefore strongly anti-correlated with the downwelling LW flux anomaly estimate via the surface temperature.

The sum of the induced surface flux anomalies is of a similar shape and order of magnitude to the sea ice latent heat anomalies implied by the ice thickness anomalies with respect to PIOMAS (Figure 6), with positive anomalies of order 10 Wm$^{-2}$ during the ice melt season, and negative anomalies of order 5 Wm$^{-2}$ during the ice freeze season. The most obvious discrepancy occurs in July, when the sum of the induced surface flux anomalies is small and of indeterminate sign, while a large positive anomaly is implied by the sea ice thickness simulation. This may be due to the ‘missing process’ of surface albedo anomaly due to the presence of snow on sea ice. Early surface melt onset, and sea ice fraction loss, as modelled by HadGEM2-ES, would be expected to be associated also with early loss of snow on sea ice, with an associated surface albedo anomaly, with this process reaching its maximum influence at a time between that of the surface melt onset (June) and that of the sea ice fraction loss (August).
In winter the sum of the induced anomalies is consistently lower in magnitude than the sea ice latent heat flux anomaly, indicating that the tendency of model anomalies ice thickness and downwelling LW to cause additional ice growth does not fully account for the anomalous growth seen relative to PIOMAS. This may be due to the use of ice thickness reference dataset; it was noted in Section 2.3 that PIOMAS may underestimate ice thickness in late winter, which would cause the sea ice latent heat flux anomaly, and the surface flux anomaly induced by ice thickness anomalies, to be overestimated and underestimated respectively. It may also indicate that model anomaly in snow thickness, not investigated here, plays a part in inducing additional surface flux anomaly. A third factor in play may be sub-grid scale variation in ice thickness, which as discussed in Appendix B causes an underestimate in the freezing season surface flux anomalies of the order of 10%, or ~1 W m⁻². It is noted that errors due to sub-monthly scale time covariance in the relevant climate variables, discussed further in Appendix B, are an order of magnitude lower than most absolute values discussed here, and in the context of the large observational uncertainties are unlikely to be important.

Using the methods described in Section 4 we calculate surface flux biases induced by model biases in downwelling SW, downwelling LW, ice area, local ice thickness and surface melt occurrence. The resulting fields are averaged over the model period and over the Arctic Ocean region, to produce for each variable a seasonal cycle of surface flux bias induced by the bias in that variable. The induced surface flux (ISF) biases are displayed in Figure 6, together with total ISF bias, radiative flux biases estimated by the direct radiation evaluation relative to ISCCP-FD, CERES and ERAI, and also sea ice latent heat uptake biases implied by the ice thickness biases relative to PIOMAS. The ISF biases are also shown in Table 1, using CERES as reference dataset for the radiative terms.

ISF biases tend to sum to negative values during the winter (indicating anomalous modelled energy loss and ice growth) and to positive values during the summer (indicating anomalous modelled energy gain and ice melt), consistent with the radiation and ice thickness evaluation. Major roles are identified for particular processes in certain months. Firstly, in June a bias in surface melt onset induces a surface flux bias of -13.6 W m⁻², equivalent roughly to an extra 11 cm of melt. This is associated with the meltpond parameterisation of HadGEM2-ES lowering the surface albedo at the end of May as the surface reaches the melting point, in contrast to SSMI observations which show surface melting to commence on average in mid to late June in the 1980-1999 period. Secondly, in August a bias in ice fraction induces a surface flux bias of 9.6 W m⁻², equivalent to an extra 8 cm of melt. This is associated with the overly fast retreat of sea ice in HadGEM2-ES, and the low extents in late summer, as noted in Section 3.

Thirdly, the large model biases in downwelling LW present throughout the freezing season induce substantial surface flux biases, ranging from -6.5 to -3.8 W m⁻² from October-March (the surface flux biases are considerably lower than the original downwelling LW biases because of the increasing inefficiency by which surface heat loss is converted to sea ice growth as ice thickens). Throughout this period, the total extra heat loss estimated by this process is roughly equivalent ice growth ranging from 20-33 cm. Fourthly, the negative biases in ice thickness present at the end of summer also induce substantial surface flux biases which tend to decrease throughout the freezing season as the thickness biases decrease, with an induced surface flux bias of -8.3 W m⁻² in November reducing to -2.0 W m⁻² in March. This effect is roughly equivalent to an extra 24 cm of ice growth.

It is noted that while large ISF biases due to downwelling SW and LW are evident during summer, there is very
large spread in these values between observational datasets, to the extent that the sign of the biases are
uncertain. It is concluded that it is not possible to determine the net effect of downwelling radiative biases on
surface flux during the summer with current observational data.

Internal variability in the ISF biases is measured by taking the standard deviation of the whole-Arctic ISF bias
for each process and month across all 20 years in the model period, and all four ensemble members used.
Variability is highest in the ice area term, reaching 4.0 Wm\(^{-2}\) in July. Variability reaches considerable size in
some other terms in some months, for example 1.1 Wm\(^{-2}\) for surface melt onset in June, 1.9 Wm\(^{-2}\) for ice
thickness in November, but is otherwise mainly under 1 Wm\(^{-2}\) in magnitude. In each case, therefore, the ISF
biases noted above are persistent features of the model.

Residuals between the total ISF bias and the directly evaluated radiative flux biases (demonstrated using CERES
as radiation reference dataset in Table 1) are comparable in magnitude to the differences between the three
different evaluations of the radiative flux biases, indicating that observational uncertainty is likely to dominate
uncertainty in the ISF biases themselves. For example, the residual between total ISF bias and net radiation bias
varies from -15.4 Wm\(^{-2}\) in June to 8.1 Wm\(^{-2}\) in November, while the difference between net radiation bias as
evaluated by CERES and ERAI respectively varies from -16.9 Wm\(^{-2}\) in July to -2.4 Wm\(^{-2}\) in September. As
discussed in Section 3 above, evidence from in situ validation studies is inconclusive as to the true size of the
modelled downwelling LW bias, and hence as to the magnitude of the surface flux bias induced by downwelling
LW. On the other hand, the evidence of PIOMAS underestimating winter sea ice thickness suggests that the
magnitude of this bias, and the associated ISF bias, may be underestimated. It is also noted that there is a high
uncertainty of the order ±10 Wm\(^{-2}\) in the ice area contribution during the winter. This is because, the rate of
dependence of surface flux on ice area is very high in freezing grid cells (generally 100-200 Wm\(^{-2}\)) due to the
large differences between turbulent fluxes over sea ice and open water.

In Appendix A potential errors in the ISF analysis are discussed and are found to be quite small in magnitude
relative to the difference between observational datasets. Firstly, due to sub-monthly variation in the component
variables, the winter downwelling LW component may be underestimated in magnitude by around 0.6 Wm\(^{-2}\) on
average, and the ice area component in August may be overestimated by around 1.6 Wm\(^{-2}\). Secondly, due to a
separate effect by which the ISF biases do not exactly sum to the total surface flux bias, the total bias in October
is likely to be overestimated in magnitude by 3.6 Wm\(^{-2}\). Thirdly, due to nonlinearities in the surface flux
dependence on ice thickness, the ice thickness component is overestimated in magnitude by 0.7 Wm\(^{-2}\) on
average from October-April, with a maximum overestimation in November of 1.9 Wm\(^{-2}\). We note that it is
possible that in some months the sum of the ISF biases may be a truer representation of the actual surface flux
bias than any of the individual evaluations, as the method combines observational estimates with physical
relationships between the various flux components. For example, in the satellite datasets observational errors in
the different components are not constrained to correlate in a physically realistic sense.

The most obvious discrepancy between the total ISF bias and the net radiation bias occurs in July, when the sum
of the induced surface flux biases is small and of indeterminate sign, while a large positive bias is implied by the
sea ice thickness and surface radiation simulations. This may be due to the 'missing process' of surface albedo
bias due to the presence of snow on sea ice. Early surface melt onset, and sea ice fraction loss, as modelled by
HadGEM2-ES, would be expected to be associated also with early loss of snow on sea ice, with an associated surface albedo bias, with this process reaching its maximum influence at a time between that of the surface melt onset (June) and that of the sea ice fraction loss (August). We note also that the direct effect of thinning ice on ice albedo could induce an additional flux bias relative to the real world, despite the fact that this effect is not modelled in HadGEM2-ES.

An annual mean total ISF bias of -3.6 (CERES) and -4.5 Wm\(^{-2}\) (ERAI) is present when the satellite datasets are used as reference (the annual mean total ISF bias for ERAI is -0.1 Wm\(^{-2}\)). It is noted that given a negligible contribution of oceanic heat convergence to the sea ice heat budget in HadGEM2-ES or in the real world, as is argued in Section 4, the annual mean surface flux bias would be expected to be substantially smaller than these figures, as a surface flux bias of -4.5 Wm\(^{-2}\) is equivalent to a relative thickening of the model sea ice cover by 9m over the 1980-1999 period. Analysis of potential sources of error in the ISF calculations in Appendix A does not produce evidence of a systematic bias that could explain these large annual mean negative biases, although the early-winter errors in the ice thickness component could explain a small portion (0.4 Wm\(^{-2}\)). Given the large discrepancy amongst observational datasets, therefore, it is likely that observational inaccuracy plays a significant part in introducing this annual mean bias.

Spatial patterns in the ISF biases are now discussed. Consistent with the pattern of net SW bias anomaly identified in section 3, the spatial pattern of surface flux bias induced by melt onset occurrence is characterised by a weak maximum in the central Arctic, with values falling away towards the coast. A more sharply-defined pattern is produced by the ice fraction bias anomaly in August, with high values across the shelf seas and the Atlantic side of the Arctic falling to low or negative values in the Beaufort Sea; the pattern displayed by the ice thickness-induced bias anomaly in November is almost a mirror image. Finally, the surface flux bias induced by downwelling LW in February displays slightly higher values on the Siberian side of the Arctic than the North American side, the reverse pattern to that displayed by the downwelling LW itself in Figure 5d. The contrast is due to the role the effective ice thickness scale factor plays in determining the induced surface flux bias; thicker ice, such as that which tends to be found on the American side of the Arctic in both model and observations, tends to greatly reduce the flux bias. This represents the thickness-growth feedback, the reality that thicker ice will grow less quickly than thin ice under the same atmospheric conditions.

The spatial patterns of total ISF bias shows many similarities to total net radiation bias evaluated by CERES in most months of the year (Figure 7), notably a tendency in July and August for positive surface flux biases to be concentrated on the Atlantic side of the Arctic, and a tendency throughout the freezing season for negative surface flux biases to be least pronounced in the Beaufort Sea, where the ice thickness biases are likely to be lowest. We note that the spatial pattern of amplification of the ice thickness seasonal cycle displayed in Figure 3 is very similar, with amplification most pronounced near the Atlantic Ocean ice edge, and least pronounced in the Beaufort Sea.

The surface flux biases produced by ice fraction bias anomaly in August, and ice thickness bias anomalies in November, provide reasons for the spatial variation in amplification of the ice thickness seasonal cycle seen in Figure 4, as well as the close resemblance of this pattern to the model bias anomalies in...
annual mean ice thickness. Ice which is thinner in the annual mean will tend to melt faster in summer, due to the net SW biases associated with greater creation of open water (the ice albedo feedback), and to freeze faster in winter, due to greater conduction of energy through the ice (the ice thickness-growth feedback).

5.6. Discussion

The calculation of the surface radiative flux biases induced by various key processes in the Arctic Ocean produces results qualitatively consistent with the surface radiation evaluation, and with the surface flux biases implied by the sea ice simulation. Melt onset occurrence and sea ice fraction biases tend to cause anomalous surface warming, and sea ice melt, during the summer, in the HadGEM2-ES historical simulation: downwelling LW and ice thickness biases tend to cause anomalous surface cooling, and hence sea ice growth, during the winter. It is recognised that it would not be expected that the induced surface fluxes would sum to values exactly consistent with either the radiation evaluation, or the sea ice volume evaluation, due to the main to observational inaccuracy, but also due to the approximations made when deriving the simple models of Section 2.3.

It is helpful to divide the processes examined into feedbacks (surface flux biases induced by biases in the sea ice state itself) and forcings (those induced by downwelling radiative fluxes and melt onset occurrence). In this sense, a ‘forcing’ refers to a variable which is independent of the sea ice volume on short timescales, rather than being used in the traditional sense of a radiative forcing.

The surface flux bias induced by biases in ice fraction during the melting season can be identified with the effect of the surface albedo feedback on the sea ice state. This is because during the melting season the ice area affects the estimated surface flux only through the surface albedo, and the surface flux biases induced in this way cause associated biases in ice melt. On the other hand, the surface flux bias induced by biases in ice thickness during the freezing season can be identified with the effect of the thickness-growth feedback on the sea ice state. This is perhaps less obvious, as the ice thickness affects the estimated surface flux via the surface temperature and upwelling LW radiation, while the thickness-growth feedback is usually understood to result from differences in conduction. However, the assumption of flux continuity at the surface in constructing the estimated surface flux means that the cooler surface temperatures, and shallower temperatures gradients occurring for thicker ice categories are manifestations of the same process. Slower ice growth at higher ice thicknesses has a manifestation in a smaller negative surface flux, and the surface temperature is the mechanism by which this is demonstrated. Hence the effect of the thickness-growth feedback is described by the ice thickness-induced component of the surface flux bias.

In this way, the ISF analysis allows the effect of the surface albedo and thickness-growth feedbacks on the sea ice state to be quantified, and compared to the effect of other drivers. Arctic-wide, the surface albedo feedback, diagnosed as the ice area-induced component of the surface flux bias, contributes an average of $5.2 \text{ Wm}^{-2}$ to the surface flux bias over the summer months, equivalent to an extra 13cm of ice melt; this is very similar to the effect of the surface melt onset-induced component, which contributes an average of $5.3 \text{ Wm}^{-2}$, equivalent also to an extra 13cm of ice melt. In the freezing season, meanwhile, the thickness-growth feedback, diagnosed as the ice thickness-induced component of the surface flux bias, contributes an average of $-4.4 \text{ Wm}^{-2}$ to the surface flux bias.
flux bias from October-April, equivalent to an extra 26cm of ice freezing, while the downwelling LW-induced component (using CERES as reference dataset) contributes an average of -4.9 Wm\(^{-2}\), equivalent to an extra 29cm of freezing over this period.

The biases of the HadGEM2-ES sea ice state can be understood by considering in turn the separate ISF components, their magnitudes, and the times of year when they are important. The surface flux anomaly induced by anomalies in ice fraction can be identified with the surface albedo feedback; that induced by anomalies in ice thickness can be identified with the thickness-growth feedback. The other processes examined—downwelling SW, LW, and melt onset occurrence—can be viewed as external forcings. Demonstrated in Figure 7, the anomalous summer sea ice melt is initiated by the early melt onset occurrence, and maintained by the surface albedo feedback, which acts preferentially in areas of thinner ice; the anomalous winter ice growth is maintained both by the thickness-growth feedback (occurring mainly in areas of thinner ice, of greater importance in early winter) and by the downwelling LW bias anomaly (more spatially uniform, in late winter). It is unclear that any significant role is played by the downwelling SW bias anomaly, as at the only time of year when the radiation datasets agree that this bias anomaly is of significant value (May), the induced surface flux bias anomaly is more than balanced by that induced by downwelling LW. However this may have a role in causing the later melt onset anomaly, as discussed below.

The means by which the external forcings—anomalous LW winter cooling, and early late spring melt onset—cause an amplified seasonal cycle in sea ice thickness are clear. It is also possible to understand how, in the absence of other forcings, these combine to create an annual mean sea ice thickness which is biased low, as seen in Section 3. The melt onset forcing, by inducing additional ice melting through its effect on the ice albedo, acts to greatly enhance subsequent sea ice melt through the surface albedo feedback. The downwelling LW, on the other hand, by inducing ice freezing, acts to attenuate subsequent sea ice freezing through the thickness-growth feedback. The effect is that surface flux bias anomalies induced by melt onset occurrence are enhanced, while those induced by downwelling LW are diminished.

Acting together, the ice thickness-growth feedback and surface albedo feedback create a strong association between lower ice thicknesses and amplified seasonal cycles, because ice which tends to be thinner will both grow faster during the winter, and melt faster during the summer. Hence the melt onset bias, acting alone, would induce a seasonal cycle of sea ice thickness lower in the annual mean, but also more amplified, than that observed, because the surface albedo and thickness-growth feedbacks act to translate lower ice thicknesses into faster melt and growth. For similar reasons, the downwelling LW bias, acting alone, would induce a seasonal cycle of sea ice thickness higher in the annual mean, and also less amplified, than that observed. Hence the melt onset anomaly, acting alone, would induce a seasonal cycle of ice thickness both lower, and more amplified, than that observed, while the downwelling LW anomaly acting alone would induce a seasonal cycle of ice thickness higher and less amplified. The bias seen in HadGEM2-ES is a result of the melt onset bias anomaly ‘winning out’ over the downwelling LW, due to its occurring at a time of year when the intrinsic sea ice feedbacks render the ice far more sensitive to surface radiation. The anomalously low ice cover in September arises as a consequence of the low annual mean ice thickness, and in particular of the anomalously severe summer ice melt. The finding that the low annual mean ice thickness is driven by surface albedo biases is
consistent with the finding by Holland et al (2010) that variance in mean sea ice volume in the CMIP3 ensemble was mostly explained by variation in summer absorbed SW radiation.

The feedbacks of the sea ice state explain the association between spatial patterns of annual mean ice thickness bias and ice thickness seasonal cycle amplification. However, the external forcings (melt onset and downwelling LW bias) cannot entirely explain the spatial patterns in the mean sea ice state biases, because on a regional scale effects of sea ice convergence, and hence dynamics, become more important. The annual mean ice thickness bias seen in HadGEM2-ES is associated with a thickness maximum on the Pacific side of the Arctic, at variance with observations which show a similar maximum on the Atlantic side. It was shown by Tsamados et al (2013) that such a bias could be reduced by introducing a more realistic sea ice rheology.

The study would be incomplete without a discussion of possible causes of the two external drivers identified by this analysis as causing sea ice model biases. Underestimation of wintertime downwelling LW fluxes in the Arctic is known to be a widespread model bias in the CMIP5 ensemble (e.g. Boeke and Taylor, 2016). Pithan et al (2014) showed that this bias was likely to be a result of insufficient liquid water content of clouds forming in subzero air masses, resulting in a failure to simulate a particular mode of Arctic winter climate over sea ice; the ‘mild mode’, characterised by mild surface temperatures and weak inversions, whose key diagnostic is observed to be a net LW flux of close to 0 W m\(^{-2}\) (Stramler et al, 2011; Raddatz et al, 2015 amongst others). HadGEM2-ES was not one of the models assessed by Pithan et al (2014), but its winter climate simulation displays many of the characteristic biases displayed by these, notably a tendency to model very low cloud liquid water fractions during winter compared to MODIS observations (Figure 8a) and a failure to simulate the milder mode of Arctic winter climate as demonstrated in SHEBA observations, diagnosed by 6-hourly fluxes of net LW (Figure 8b).

Here we conclude that a similar mechanism is likely to be at work in HadGEM2-ES, and that insufficient cloud liquid water is the principal driver of the anomalously low downwelling LW fluxes.

The causes of the early melt onset bias of HadGEM2-ES are harder to determine. For most of the spring, comparison of daily upwelling LW fields of HadGEM2-ES to CERES-SYN observations (not shown) shows the Arctic surface to be anomalously cold in the model, as during the winter. During May, however, upwelling LW values rise much more steeply in the model, and surface melt onset commences during mid-to-late May, far earlier than in the satellite observations. A possible cause of the overly rapid surface warming during May is the zero-layer thermodynamics approximation used by HadGEM2-ES, in which the ice heat capacity is ignored.

Comparing fields of surface temperature in HadGEM2-ES between the beginning and the end of May shows a ‘missing’ ice sensible heat uptake flux of 10-30 W m\(^{-2}\) over much of the central Arctic, which would in turn be associated with a reduction of flux into the upper ice surface of 5-15 W m\(^{-2}\). Examination of modelled and observed daily timeseries of downwelling LW and net SW fluxes in late May and early June suggests that a surface flux reduction of this magnitude could delay surface melt by up to 2 weeks, a substantial part of the modelled melt onset bias anomaly seen. Another cause of the rapid warming may be the increasing relative magnitude of the downwelling SW response to cloud bias anomalies as May progresses (compared to the downwelling LW response). Comparison of 5-day means of HadGEM2-ES radiative fluxes during May to those from the CERES-SYN product (not shown)
support this hypothesis; a modelled bias anomaly in downwelling SW grows quickly during early May, from ~0Wm$^{-2}$ to ~30Wm$^{-2}$, while the modelled bias anomaly in downwelling LW remains roughly constant.

The ISF analysis as presented does not comprise an exhaustive list of processes affecting Arctic Ocean surface fluxes. The missing processes of the effects of snow fraction and ice thickness bias on the surface albedo have already been noted; the effect of snow thickness bias on winter conduction and surface temperature is another process which cannot be included due to inadequate observations. Model biases in the turbulent fluxes may also be significant; while the process which is likely most important in determining these during the winter is captured (ice fraction in the freezing season), a more detailed treatment of turbulent fluxes would also examine the effect on these of the overlying atmospheric conditions. It is also noted that snowfall itself is a component of the surface flux which could in theory be evaluated directly given a sufficiently reliable observational reference.

Finally, it is noted that a complete treatment of model biases affecting the sea ice volume budget would also examine causes of bias in oceanic heat convergence. For the reasons discussed in Section 4 these are likely to be small in the Arctic Ocean interior in HadGEM2-ES and observations, but the model bias could nevertheless conceivably be of considerable size in the context of the surface flux biases shown in Figure 6. The total Arctic Ocean heat convergence modelled by HadGEM2-ES for the period 1980-1999 is 4.4 Wm$^{-2}$, although this figure shows high sensitivity to the location of the boundary in the Atlantic sector, suggesting that most of this heat is released close to the Atlantic ice edge. This figure is slightly higher than the 3 Wm$^{-2}$ found by Serreze et al (2007) in their analysis of the Arctic Ocean heat budget, but is broadly consistent with observational estimates of oceanic heat transport through the Fram Strait (likely to be the major contributor to Arctic Ocean heat convergence) from 1997 to 2000 by Schauer et al, 2004. This suggests that errors in oceanic heat convergence are unlikely to contribute significantly to sea ice volume biases in HadGEM2-ES. However, for a hypothetical model which simulated greater oceanic heat convergence in the Arctic Ocean interior, the surface flux analysis presented here would fail to adequately describe the model bias in the sea ice volume budget.

Conclusions

HadGEM2-ES simulates a sea ice cover which is not extensive enough at annual minimum. Comparison to various ice thickness datasets shows that it also has too low an annual mean ice thickness, and that its ice thickness seasonal cycle is likely to be overamplified. Evidence of a positive net SW bias during the ice melt season, and a negative net LW bias during the ice freezing season is apparent from evaluations using multiple radiation datasets.

An evaluation of processes influencing surface radiation, combined with simple models to estimate their effect, produces results consistent with the evaluation of the sea ice state and surface radiation; processes tend to cause anomalous ice melt during the melting season, and anomalous ice growth during the freezing season. Consequently model bias anomalies in sea ice growth and melt rate can be attributed in detail to different causes; in particular, the roles played by the sea ice albedo feedback, by the sea ice thickness-growth feedback, and by external forcings, can be quantified. The analysis reveals how the melt onset bias anomaly of HadGEM2-
ES tends to make model ice thickness both low in the annual mean, and too amplified in the seasonal cycle, with
the downwelling LW bias anomaly acting to mitigate both effects. The result is consistent with the prediction of
DeWeaver et al (2008) that sea ice state is more sensitive to surface forcing during the ice melt season than
during the ice freeze season. The analysis also suggests that through an indirect effect on surface albedo at a
time when sea ice is particularly sensitive to surface radiation, bias anomalies, the zero-layer approximation,
which was until recently commonplace in coupled models, may be of first-order importance in the sea ice state
bias of HadGEM2-ES.

The analysis also makes explicit the link between the spatial pattern of anomalies in annual mean ice thickness,
and anomalies in the April-October ice thickness difference. Regions where ice thickness tends to be biased
particularly low in the annual mean also display higher amplification in the seasonal cycle, due to the direct
action of the thickness-growth and ice albedo feedbacks, despite the initiating factors of melt onset occurrence
and downwelling LW anomaly being comparatively spatially uniform. However, the reasons for the underlying
spatial distribution of the annual mean ice thickness anomalies in HadGEM2-ES are likely to lie in ice dynamics
rather than thermodynamics.

A clear link has been demonstrated between the spatial pattern of biases in annual mean ice thickness, likely
driven by ice dynamics, and that of biases in the April-October thickness. Where ice thickness is biased low in the
annual mean, an enhanced seasonal cycle is apparent. This is due to the thickness-growth and ice albedo
feedbacks, initiated by melt early melt onset and downwelling LW bias, both of which are spatially uniform.

The method is limited by the current inability to evaluate the impact of anomalies in modelled snow cover, as
well as by the large observational uncertainties in summer surface radiation, underlying the importance of
reducing uncertainty in large-scale observations of Arctic climate variables. Adding in the ‘missing processes’
of freezing season snow thickness, and melt season snow fraction, would represent useful extensions to the
analysis presented. Other potential causes of SEB anomalies not investigated in this study include processes
casing anomalies in the turbulent fluxes, and in particular the effects of anomalies in sea ice fraction during the
freezing season.

Large observational uncertainties for snow cover and summer surface radiation limit the overall accuracy of the
methodology presented here. The addition of freezing season snow thickness, and melt season snow fraction,
would represent useful extensions to the analysis presented. An additional caveat regarding this analysis is that it
does not consider factors influencing turbulent fluxes (with the exception of the ice area, but this contribution is
subject to particularly high uncertainty). It also does not consider the influence of oceanic heat convergence on
sea ice state; in HadGEM2-ES the latter is small (~10%), but might be more significant in other models.

In the case study presented here, the analysis provides mechanisms behind a model bias in sea ice simulation.
However, the analysis could also be used to investigate a sea ice simulation that was ostensibly more consistent
with observations, to determine whether or not the correct simulation was the consequence of model biases that
cause opposite errors in the surface energy budget; a negative result would greatly increase confidence in the
future projections of such a model. The analysis could be also used to investigate a whole model ensemble, to
attribute spread in modelled sea ice state to spread in the underlying processes affecting the SEB, focussing
attention on ways in which spread in modelled sea ice could be reduced. It is noteworthy that Shu et al (2015)
found the CMIP5 ensemble mean Arctic sea ice volume to be biased low in the annual mean, and overamplified
in the seasonal cycle, relative to PIOMAS (albeit over the entire Northern Hemisphere), suggesting that the
behaviour exhibited by HadGEM2-ES may be quite common in this ensemble.

Finally, it is suggested that the ISF method, as well as being used to compare a model to observations, could
also be used to understand the reasons for the biases of one model with respect to another. Such a comparison
would avoid the issues of observational uncertainty discussed above, enabling the contributions of the different
model variables to the surface flux biases to be evaluated more accurately.

Appendix A: Derivation of formulae used for surface flux analysis

The surface energy balance equation on an area of sea ice and open water, in which fluxes arriving at the
surface from above are equated with fluxes of energy into the ice and ocean below, can be expressed as follows,
ignoring the contribution from snowfall occurring on both sides:

\[ F_{SW} + F_{LW} + F_{sens} + F_{lat} = F_{cond} + F_{melt} + F_{sublim} + F_{wat} \] (A1)

Here \( F_{SW} \), \( F_{LW} \), \( F_{sens} \), \( F_{lat} \), \( F_{cond} \), \( F_{melt} \), \( F_{sublim} \), and \( F_{wat} \) refer to fluxes of net SW, LW, sensible and
latent heat flux, conductive heat flux into sea ice, top melting flux into sea ice, sea ice net sublimation flux, and
flux of energy directly into seawater. Here and below, the convention used is that positive numbers denote
downward fluxes. For the reasons given above, we neglect the turbulent heat fluxes. By implication we also
neglect the sea ice net sublimation flux. Hence the surface flux is equated with \( F_{SW} + F_{LW} \).

During the ice freezing season, surface temperatures are below freezing, conduction into the ice is substantial,
and \( F_{melt} = 0 \). In addition, ice fraction tends to be close to 1 over much of the Arctic. We assume for the
purposes of this analysis that ice fraction during the freezing season is equal to 1 (equivalent to neglecting \( F_{wat} \))
while recognizing that due to turbulent fluxes being observed to be very large in regions of open water, \( F_{wat} \)
may in reality be of significant size. In the following analysis, as a result, the effects of model anomalies in ice
fraction during the freezing season on the surface flux are not estimated.

Linearising the dependence of \( F_{SW} + F_{LW} \) on the surface temperature \( T_{sfc} \) about the freezing point as
\( A + BT_{sfc} \), and assuming uniform conduction within the ice (thereby ignoring sensible heating of ice), we have

\[ A + BT_{sfc} = \frac{T_{sfc} - T_0}{R_{ice}} \] (A2)

where the ice thermal insulation \( R_{ice} = h_I/k_I + h_S/k_S \), \( h_I \), \( k_I \), \( h_S \), and \( k_S \) being ice and snow thickness, and
ice and snow conductivity, respectively.
Solving for $T_{SW}$ and re-substituting, we have

$$F_{SW} + F_{LW} = \frac{A + BT_{0}}{1 - B \cdot R_{ice}} \quad (A3)$$

$A$ represents the surface temperature-independent part of the surface flux, and can be identified with

$$F_{SW} + F_{LW} + C$$

where $F_{SW}$ is the downwelling LW flux and $C$ is the upwelling longwave flux associated with a surface temperature of $0^\circ C$. Hence

$$F_{SW} = \frac{\alpha F_{SW} \cdot F_{LW} + C + BT_{0}}{1 - B \cdot R_{ice}} \quad (A4)$$

Equation (A4) summarises the dependence of the net radiative flux on downwelling radiative fluxes and on ice and snow thickness, and is equivalent to equation (1), used to calculate induced surface flux anomaly during the freezing season.

During the ice melting season, conductive flux is near zero, surface temperature is close to $0^\circ C$ throughout, and ice fraction is no longer necessarily close to 1. Hence the surface flux is

$$F_{SW} = F_{LW} + C + \alpha F_{SW} \quad (A5)$$

where $\alpha$ is surface albedo and $C = e^{\sigma T_f^4}$ is upwelling longwave radiation at $T_f = 0^\circ C$. In HadGEM2-ES, surface albedo is parameterised after Curry et al (2001), which approximates the effect of sea ice meltponds by reducing albedo as the surface temperature approaches the melting point. Surface albedo is therefore affected by ice fraction, the presence or otherwise of snow, and whether or not melt onset has occurred. The effect can be summarised as

$$F_{SW} = F_{LW} + C + F_{SW} \cdot \left(1 - \alpha_{sea} - \alpha_{ice} - \alpha_{snow} - \alpha_{cold} \right) \quad (A6)$$

where $\alpha_{sea}$, $\alpha_{ice}$, $\alpha_{snow}$ and $\alpha_{cold}$ indicate the modelled albedos of open water, bare ice, melting snow and cold snow respectively, and $I_{sea}$, $I_{ice}$, $I_{snow}$ and $I_{cold}$ indicate the presence of sea ice, snow cover and cold snow (that is far from the melting point) respectively. Integrated over large areas, $I_{ice}$ becomes ice area fraction $a_{ice}$ (and similarly for $I_{snow}$ and $I_{cold}$), giving equation (2), used to calculate induced surface flux anomaly during the melting season.

Appendix B: Error associated with spatial and temporal extrapolation
The application of equations (A4) and (A6) to monthly means of data valid for grid cells tens of km across introduces two potential sources of error in estimating the surface flux anomaly induced. Firstly, subgrid-scale variation in ice thickness, which is not accounted for by the simple model, tends to increase the efficiency of heat loss and hence ice creation during the freezing season via the scale factor \( \frac{1}{1 - Bh_{\text{eff}}} \), rendering ice growth more sensitive to errors in the forcing variable. In an analysis of a field of category ice thickness produced by HadGEM2-ES for January 1994, this was found to cause an average underestimate of 8% in the scale factor value, and hence the induced surface flux anomaly.

Secondly, error will occur due to covariance in time between variables being multiplied. As both melt onset occurrence and ice fraction observations were available on daily timescales, the effect of this covariance was assessed by calculating induced surface flux anomaly due to melt onset occurrence and ice fraction on both daily and monthly timescales. The daily covariance was found to cause a maximum error of 0.1 Wm\(^{-2}\) in the melt onset-induced anomaly, in June, and of 0.5 Wm\(^{-2}\) in the ice fraction-induced anomaly, in August, with most other monthly anomalies being below 0.1 Wm\(^{-2}\).

Appendix A: Analysis of potential errors in ISF bias calculation

Due to observational uncertainty, it is difficult to directly evaluate the ISF bias calculations. Instead, we examine in turn the two principal sources of error in the method; firstly, error in correctly characterising the dependence of surface flux on a climate variable, and secondly, error in approximating the surface flux bias induced by this as the product of the surface flux dependence with the model bias in that variable.

To analyse the first source of error, we begin by comparing fields of the approximated surface flux \( g_{\text{s,t}} \) to those of the real modelled surface flux \( F_{\text{sfc}} \). The \( g_{\text{s,t}} \) are found to capture well the large-scale seasonal and spatial variation in surface flux, but are prone to systematic errors which vary seasonally, indicated in Figure A1; firstly, a tendency to underestimate modelled negative surface flux in magnitude from October-April by 13% on average; secondly, a tendency to overestimate modelled positive surface flux from June-August by up to 10 Wm\(^{-2}\); thirdly, during May, a underestimation varying from 5-20 Wm\(^{-2}\).

Examining first the winter underestimation (demonstrated in Figure A1 a-c), it is found that for each model month the relationship between estimated and actual surface flux is strongly linear, with underestimation factors ranging from 6 ± 1% in December to 17 ± 2% in April. This suggests that the cause lies in systematic underestimation of the scale factor \( \sum_{\text{cat}} \gamma_{\text{s,t}, \text{REF}} \left( 1 - B R_{\text{s,t}} \right)^{-1} \). A possible cause is covariance in time between \( \gamma_{\text{s,t}, \text{REF}} \) and \( R_{\text{s,t}} \) within each month, particularly in the first ice category; during the freezing season, occurrence of high fractions of ice in category 1, the thinnest category, would be expected to be associated with formation of new ice, and correspondingly lower mean thicknesses of ice in this category, lower values of \( R_{\text{s,t}} \).
and higher values of $(1 - BR_{cat})^{-1}$. A calculation using daily values of $\gamma_{cat-REF}$ ranging from 0.1 – 0.5, and daily values of $R_{cat}^{out}$ ranging from 0.2 – 0.5m, predicts that this effect would in this case lead to an underestimation of 9% in the magnitude of the surface flux, sufficient to explain all of the underestimation in October, December and January, and most in November, February and March. This effect would produce a corresponding underestimation of the rate of dependence of surface flux on downwelling LW radiation and ice thickness throughout the freezing season. It was estimated that the downwelling LW component of the ISF bias is underestimated by 0.6 Wm$^{-2}$ for the freezing season on average due to this effect.

Secondly, we examine the tendency to overestimate surface flux during the summer (Figure A1d-f), an effect that displays a spatially uniform bias rather than a spatially uniform ratio, ranging from 5-15 Wm$^{-2}$ in July and August; the bias is smaller, and in the central Arctic negative, during June. A possible contributing factor to this bias is within-month covariance between ice area and downwelling SW; during July and August, both downwelling SW and surface albedo fall sharply, an effect that would tend cause the monthly mean surface flux to be overestimated. To estimate this effect, monthly trends in these variables were estimated by computing half the difference between modelled fields for the following and previous month. For July, an overestimation in surface flux of magnitude 5-15 Wm$^{-2}$ was indeed predicted in the Siberian seas, as well as the southern Beaufort and Chukchi Seas; however, in the central Arctic no overestimation was predicted, due to near-zero trends in ice area in the summer months. It is possible that some covariance between ice area and downwelling SW is nevertheless present in these regions, due to enhanced evaporation and cloud cover in regions of reduced ice fraction.

However, this effect would have no direct impact on the ISF biases because these are computed from monthly means of the model bias in one variable by the model mean in the other; hence, it is covariance between bias and mean that would induce inaccuracy in this case. By similarly approximating the trend in monthly mean model bias as half the difference between model bias in the adjacent months, the error in downwelling SW and ice area contributions were evaluated. Error in the downwelling SW term was found to be significant early in the summer, with an error of -2.7 Wm$^{-2}$ in June; error in the ice area term was found to be significant later in the summer, with errors of -1.7 Wm$^{-2}$ and -1.6 Wm$^{-2}$ in July and August respectively. However, the August error is small relative to the total ISF bias identified.

Thirdly, we examine the reasons for the underestimation of surface flux in May (Figure A1g-i), a pattern unique to this month which is seen to be small in the central Arctic but to approach 20 Wm$^{-2}$ at the Arctic Ocean coasts. A likely cause of this inaccuracy is the classification of grid cells as ‘freezing’ or ‘melting’ for entire months. During May, as has been seen, most model grid cells in fact cross from one category to the other; however, virtually all Arctic Ocean grid cells are classified as freezing for the month as a whole. The difference field between estimated ‘freezing surface flux’ and ‘melting surface flux’ is similar in magnitude and in spatial pattern to the underestimation field, being near-zero in the central Arctic but rising to 25 Wm$^{-2}$ close to the Arctic Ocean coasts. It is concluded that the actual model mean surface flux is much higher than that estimated near the coast due to these grid cells experiencing melting conditions from relatively early in the month. Although this error is not directly relevant to the results of this paper, as no unequivocal ISF biases were
identified for May, it would have the potential to lead to overestimation of the dependence of surface flux on ice thickness, and underestimation of dependence on all other variables, as the upwelling LW flux is unable to counteract changes in surface forcing once the surface has hit the melting point.

Having examined potential causes of error in estimating dependence of surface flux on individual variables, the validity of estimating ISF biases as the product of these with model variable biases is now discussed. Even if the dependence of monthly mean surface flux on variable $v_i$ at a model grid cell is perfectly described by

$$\frac{\partial g_{s,f}}{\partial v_i}$$

that dependence changes as the realisation varies from the model state to the real-world state. As a simplified example, a component of the surface flux, net SW, is equal to $F_{\text{SW}}(1 - \alpha_{\text{de}})$, and induced surface flux biases due to model biases in $\alpha_{\text{de}}$ would be calculated as $F_{\text{SW}}(1 - \alpha_{\text{de}}^{\text{mod}})$ and $F_{\text{SW}}^{\text{mod}}(1 - \alpha_{\text{de}}^{\text{mod}})$ respectively. However, the sum of the two induced surface flux biases will not be exactly equal to the true surface flux bias $F_{\text{SW}}(1 - \alpha_{\text{de}})$, but will differ from it by $F_{\text{SW}}^{\text{mod}}(1 - \alpha_{\text{de}}^{\text{mod}})$. This is due to the dependencies being evaluated on model states which are themselves biased.

This apparent problem can be resolved only by viewing the ISF method as a way not simply of estimating model biases due to a particular variable, but of characterising them, i.e. by accepting that the quantity that we are trying to estimate is itself somewhat subjective. Instead of requiring the ISF method to be correct, it is required that it gives useful, physically realistic results. In the case given above, a sufficient condition is that $F_{\text{SW}}(1 - \alpha_{\text{de}}^{\text{mod}})$ is small relative to $F_{\text{SW}}(1 - \alpha_{\text{de}})$ and $F_{\text{SW}}^{\text{mod}}(1 - \alpha_{\text{de}}^{\text{mod}})$, i.e. that the model bias in both downwelling SW and in surface albedo is small relative to the absolute magnitudes of these variables.

More generally, the difference between the surface flux bias $F_{\text{de}}$ and the sum of the induced surface flux biases can be approximated by

$$\sum v_i \frac{\partial g_{s,f}}{\partial v_i}$$

a term that can be calculated relatively easily as many of the derivatives go to zero. Averaged over the Arctic Ocean this term was small in most months of the year, but of significant size in October (3.6 Wm$^{-2}$), due to co-location of substantial negative biases in downwelling LW and category 1 ice thickness in this month, indicating that the true surface flux bias in this month may be substantially smaller (in absolute terms) than the -11.5 Wm$^{-2}$ obtained from summing the ISF biases.

Finally, the induced surface flux calculation implicitly assumes a linear dependence of surface flux on each climate variable. However, this is not the case for the ice thickness, where higher-order derivatives do not go to zero, and in some regions of thinner ice actually diverge. It is possible to quantify the error introduced by the assumption of linearity by comparing the partial derivative $\frac{g_{s,f}(1 - B R_{\text{cat}})}{(1 - B R_{\text{cat}})^2 \sum q_{\text{cat}}}$ to the
quantity \( (A + BT_s)h_{cat}^{-1} (1 - BR_{ice}^{-1})^{-1} (1 - BR_{ice-REF}^{-1})^{-1} \left( \sum_{cat} a_{cat} \right)^{-1} \), where \( R_{ice-REF}^{cat} = h_{cat}^{OBS} / k_j + h_{cat}^{OBS} / k_S \), being climatological ice thickness in the reference dataset, in this case PIOMAS, and all other terms defined as in Section 4. It can be shown that multiplying this quantity by the model bias produces the exact bias in estimated surface flux that is being approximated by \( \partial g_{s,f} / \partial h_{j} (h_{j}^{MODEL} - h_{j}^{OBS}) \). Hence the bias in the ice thickness component induced by the nonlinearity can be calculated directly. It is found that the nonlinearity causes the ice thickness component to be overestimated in magnitude by 0.7 Wm\(^{-2}\) on average from October-April, with a maximum overestimation of 1.9 Wm\(^{-2}\) in November.

Code availability

The code used to create the fields of induced surface flux bias is written in Python and is provided as a supplement (directory 'ISF'). The code used to create Figures 1-8, as well as Figure A1, is also provided written in Python and is provided as a supplement (directory 'Figures'). In addition, the routines used to estimate errors in the ISF analysis are provided create the monthly induced surface flux anomaly fields used in Figures 6 & 7 from the basic model and observation fields is provided (directory 'Analysis'). Finally, the code used to create Table 1 is provided (directory 'Tables'). A set of auxiliary routines used by most of the above are also provided (directory 'Library'). Most routines make use of the open source Iris library, and several make use of the open source Cartopy library.

Data availability

Monthly mean ice thickness, ice fraction, snow thickness and surface radiation, as well as daily surface temperature and surface radiation, for the first historical member of HadGEM2-ES, is available from the CMIP5 archive at https://cmip.llnl.gov/cmip5/data_portal.html.

NSIDC ice concentration and melt onset data can be downloaded at http://nsidc.org/data/NSIDC-0051 and http://nsidc.org/data/NSIDC-0105 respectively.

PIOMAS ice thickness data can be downloaded at http://psc.apl.uw.edu/research/projects/arctic-sea-ice-volume-anomaly/data/.

ERAI surface radiation data can be downloaded at http://apps.ecmwf.int/datasets/data/interim-full-daily/levtype=sfc/.

ISCCP-FD surface radiation data is available at https://isccp.giss.nasa.gov/projects/browse_fc.html.


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References


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Table 1. Surface flux biases induced by model bias in 5 different variables in HadGEM2-ES (Wm⁻²), with CERES used as reference dataset for the radiative components. Total ISF bias and total net radiative flux bias relative to CERES are shown for comparison, as well as their residual; the difference between net radiative flux bias as evaluated by CERES and ERAI is also shown. A positive number denotes a downwards flux, and vice versa.
Figure 1. The Arctic Ocean region used in the analysis.
Figure 2. 1980-1999 mean surface fluxes over the Arctic Ocean region for the first historical run of HadGEM2-ES.
Figure 23. (a) HadGEM2-ES 1980-1999 mean Arctic Ocean ice extent, compared to HadISST1.2 1980-1999, with September ice fraction bias anomaly map; (b-d) HadGEM2-ES 1980-1999 ice thickness compared to (b) PIOMAS, (c) Envisat and (d) submarine datasets over respective regions of coverage, with April and October ice thickness bias anomaly maps. For each seasonal cycle plot, the model is in black and observations in red. In (c), data is not plotted from May-September due to the region of coverage being very small.
Figure 34. HadGEM2-ES 1980-1999 model bias anomaly in ice thickness change from October-April compared to (a) PIOMAS 1980-1999; (b) Envisat 1993-2000; (c) submarine regression analysis 1980-1999. Differences are taken as model-observation so that areas of green (purple) correspond to areas where the HadGEM2-ES model simulates too much (not enough) sea ice growth through the winter.
Figure 5. (a) Downwelling SW, (b) upwelling SW, (c) net down SW, (d) downwelling LW, (e) upwelling LW, (f) net down LW, for HadGEM2-ES 1980-1999 over the Arctic Ocean region, compared to CERES 2000-2013, ISCCP-D 1983-1999 and ERAI 1980-1999. For all fluxes, a positive number denotes a downward flux and vice versa. Maps of flux bias anomalies relative to CERES are shown for downwelling SW in May, upwelling and net down SW in June, and downwelling and net down LW in February.
Figure 5. Demonstrating the calculation of fields of surface flux bias due to model bias in melting surface fraction (a-c), downwelling LW (d-f), category 1 ice thickness (g-i) and category 5 ice thickness (j-l). The left-hand column shows model bias in each variable; the middle column the local rate of dependence of surface flux on each variable as calculated above; the right column the induced surface flux bias, calculated as the product of these two fields.
Figure 6. Surface flux anomaly induced by model anomalies in ice fraction, melt onset occurrence, ice thickness, downwelling SW and downwelling LW respectively, for the Arctic Ocean region in HadGEM2-ES, 1980-1999, as estimated by the simple models described in Section 2.3. For each month, induced anomalies are estimated using in turn CERES, ISCCP-FD and ERAI as radiation reference datasets, from left-right. Sea ice latent heat flux uptake anomaly relative to PIOMAS is indicated in black. Net radiative flux anomalies relative to CERES, ISCCP-FD and ERAI are indicated in brown. Spatial patterns of induced surface flux anomaly for four processes in key months, with CERES as reference dataset, are displayed beneath.
Figure 6. Surface flux bias induced by model biases in ice fraction, melt onset occurrence, ice thickness, downwelling SW and downwelling LW respectively, for the Arctic Ocean region in HadGEM2-ES, 1980-1999, as estimated by the simple models described in Section 2.3. For each month, induced surface flux biases are estimated using in turn CERES, ISCCP-FD and ERAI as radiation reference datasets, from left-right. Sea ice latent heat flux uptake bias relative to PIOMAS is indicated in black. Net radiative flux biases relative to CERES, ISCCP-FD and ERAI are indicated in brown. Spatial patterns of induced surface flux bias for four processes in key months, with CERES as reference dataset, are displayed beneath.
Figure 7. Surface flux anomalies caused by anomalies in external forcings, and to feedbacks due to anomalies in the sea ice state, represented as stacked filled regions. All values shown are means across radiation datasets shown in Figure 4; summer radiative flux anomalies are not plotted due to very large spread among datasets.
Figure 7. Comparing fields of total ISF bias to net radiation bias relative to CERES for each month of the year, for the four historical members of HadGEM2-ES, 1980-1999.
Figure 8. Frequency distributions of (a) October-April cloud liquid water percentage in HadGEM2-ES compared to MODIS observations, for the Arctic Ocean region; (b) December-February surface net downwelling LW in HadGEM2-ES in the SHEBA region, compared to the values observed at SHEBA.
Figure A. Illustrating approximated (left) and actual (centre) model net surface flux, as well as the approximation error (right), in (a–c) February; (d–f) May; (g–i) July, for the period 1980-1999 in the first historical run of HadGEM2-ES.