



## Quantifying bioalbedo: A new physically-based model and critique of empirical methods for characterizing biological influence on ice and snow albedo

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### 15 **Abstract.**

The darkening effects of biological impurities on ice and snow have been recognized as a significant control on the surface energy balance of terrestrial snow, sea ice, glaciers and ice sheets. With a heightened interest in understanding the impacts of a changing climate on snow and ice processes, quantifying the impact of biological impurities on ice and snow albedo ('bioalbedo') and its evolution through time is a rapidly growing field of research. However, rigorous quantification of bioalbedo has remained elusive because of difficulties isolating the biological contribution to ice albedo from that of inorganic impurities and the variable optical properties of the ice itself. For this reason, isolation of the biological signature in reflectance data obtained from aerial/orbital platforms has not been achieved, even when ground-based biological measurements have been available. This paper provides the cell specific optical properties that are required to model the spectral signatures and broadband darkening of ice. Applying radiative transfer theory, these properties provide the physical basis needed to link biological and glaciological ground measurements with remotely sensed reflectance data. Using these new capabilities we confirm that biological impurities can influence ice albedo then identify ten challenges to the measurement of bioalbedo in the field with the aim of improving future experimental designs to better quantify bioalbedo feedbacks. These challenges are: 1) Ambiguity in terminology, 2) Characterizing snow or ice optical properties, 3) Characterizing solar irradiance, 4) Determining optical properties of cells, 5) Measuring biomass, 6) Characterizing vertical distribution of cells, 7) Characterizing abiotic impurities, 8) Surface anisotropy, 9) Measuring indirect albedo feedbacks, and 10) Measurement and instrument configurations.

### **1 Background:**

35 The presence of biological impurities in the cryosphere has been known for more than a century with scientific interest dating back to 1676 when Van Leeuwenhoek described microbial life in snow samples. In the late 19<sup>th</sup> Century, Nordenskiöld (1875) coined the term 'cryoconite' to describe the distinct biological aggregates that were visible on ice surface of the Greenland Ice Sheet and noted the potential for biological activity to amplify melt, stating that it was the 'greatest enemy to the mass of ice' and a potential accelerator of deglaciation. Since the late twentieth century, understanding snow and ice melt processes has gained new significance because loss of terrestrial ice is increasingly contributing to sea level rise and driving climate-amplifying positive feedbacks, many of which are linked to snow and ice reflectivity (Budyko 1969; Flanner et al. 2011). The



direct impact of biological material on snow and ice albedo has recently become known as ‘bioalbedo’, following an abbreviation of the term “biological albedo” presented by Kohshima et al., (1993).

There are two distinct areas of research into the biological darkening of the cryosphere: the first examines snow algal blooms (green and red patches) in seasonal snowpacks, and the other has focused on ice algal blooms on ice ablation zones. Ground reflectance and aerial remote sensing have been used to quantify algae in snow in the North American High Sierra (Painter et al., 2001). The same research group has also developed techniques for deriving grain size and specific surface area from spectral reflectance data that has proven useful for satellite observations of aging snow (Painter et al., 2007; 2013). Satellite retrieval of snow algal discoloration in the Harding Ice Field has also been achieved using similar techniques by Takeuchi et al. (2006). Both groups used different methods for quantifying microbial abundance in snow from remotely sensed reflectance data: Painter et al. (2001) used the integral of a chlorophyll-specific absorption feature that relies upon the availability of data with sufficiently high spectral resolution, whereas Takeuchi et al. (2006) used a broader carotenoid absorption feature as a diagnostic biomarker. Other groups have examined snow algae from a biological perspective. Remias et al. (2005; 2010) showed that snow algae produce pigments in response to environmental stresses that changes their colour, and thereby influences snow albedo. Hodson et al. (*in press*) also documented how both surface and sub-surface algal populations influenced the optical properties of maritime Antarctic snow.

Glacier albedo can also be modified by ice algae, which are taxonomically distinct from the algae that inhabit snow (Uetake et al., 2010; Yallop et al., 2012). Ice algae are dominated by several species of green algae and cyanobacteria. The relative abundance of different species of ice algae and cyanobacteria vary spatially: the green alga *Mesotaenium berggrenii* has been shown to dominate near the termini of Qaanaaq Glacier (north-west Greenland: Uetake et al. 2010) and Tyndall Glacier (Patagonia: Takeuchi and Kohshima, 2004) and is also present in western Greenland (Yallop et al. 2012). *Ancylonema nordenskiöldii* was found to dominate the upper ablation areas of Gulkana Glacier (Alaska: Takeuchi et al. 2009) and Qaanaaq Glacier (Uetake et al. 2010) and spread extensively over the western Greenland Ice Sheet along with *Cylindrocystis* spp. (Yallop et al. 2012). An abundance of both *Ancylonema* and *Cylindrocystis* on the Greenland Ice Sheet was reported in the nineteenth century (Berggren, 1871; Nordenskiöld, 1875) and mid twentieth century (Gerdel and Drouet, 1960).

The albedo reducing effect of ice algae arises from their dark pigmentation. This was discussed by Yallop et al. (2012) who linked algal pigmentation to ice albedo in south-west Greenland, but did not explicitly identify the pigments responsible. Ice algae produce a range of pigments to facilitate both light harvesting across the visible spectrum and energy dissipation under conditions of excessive irradiance (Remias et al., 2009), via ,for example, the interconversion of xanthophyll pigments to quench excess excitation energy as heat (Niyogi et al., 1997; Goss and Jakob, 2010). In addition, ice algae inhabiting surface habitats produce specialist UV-absorbing pigments (purpurogallins) that are distinct from the typical UV-absorbing pigments (mycosporine-like amino acids) identified in other algal lineages (Remias et al., 2012). Yallop et al., (2012) found these pigments to obscure the primary light harvesting and protective pigments and to lower the albedo of the ice sheet surface.

The wide spatial coverage of algae on the Greenland Ice Sheet combined with their strong absorption of visible light (which accounts for approximately half of the solar energy incident upon the ice surface: Liou, 2002) makes them important agents of albedo reduction and therefore melt accelerators. However, empirical studies that link ice algal biomass to albedo are scarce. To date, published measurements have been limited to broadband albedos measured using opposed pyranometers without quantification of the physical properties of the ice surface, nor its inorganic impurity content (Yallop et al., 2012; Lutz et al., 2014). Therefore, the biological contribution to ice albedo (‘bioalbedo’) and its spatiotemporal variability remains



unquantified. In this paper, we will show that it cannot be isolated from non-biological albedo reducing processes without refinements to current empirical techniques and without the use of a spectrally resolved radiative transfer model.

Weak, often semi-quantitative correlations have been reported between cell numbers and albedo (e.g. Lutz et al., 2016) which is unsurprising since the majority of the variation in albedo is due to physical characteristics of the snow or ice including internal factors such as ice crystal size and water content, as well as external factors such as solar angle, atmospheric effects, shading and multiple reflections between surface roughness elements (Gardner and Sharp, 2010). Takeuchi (2002), Takeuchi and Kohshima (2004) and Takeuchi et al., (2015) showed the albedo reduction resulting from biotic and abiotic impurities to vary between different glaciers, signifying that the mass and optical properties of abiotic impurities such as dusts are also crucial determinants of surface albedo. This may be particularly relevant on the western Greenland Ice Sheet, where high concentrations of dust may be outcropping from melting Holocene ice (Bøggild et al., 2010; Wientjes et al., 2010; 2011). Because ice particles have low absorption and high scattering in the visible spectrum (Warren, 1982), snow and ice albedo is not only sensitive to the concentration and optical properties of the biological impurities but also to their vertical distribution and the changing optical properties of the ice grains themselves. These complexities make the interpretation of previous results difficult. For these reasons, simple empirical models linking albedo to microbial abundance on glacier ice will have small coefficients of determination and leave a large portion of the variance in albedo unexplained. To understand bioalbedo, a physical modelling approach is required to link ground measurements to spectral data and to integrate effects of both biological and non-biological impurities being present in and on the ice. Two decades ago, an equivalent transition from a paradigm of empirical to physical modelling occurred in the field of remote sensing of terrestrial vegetation, stimulated by the difficulty in extracting sufficient information empirically (Verstraete et al., 1996). Physical modelling establishes a functional relationship between snow and ice properties and surface radiance, enabling both predictions of albedo change given changes in optical properties and surface parameter retrieval from spectral reflectance, and therefore allows us to quantify the bioalbedo of a complex mixture of ice, biotic and abiotic impurities.

Aoki et al. (2013) and Cook et al. (2017) have previously used radiative transfer models to suggest algal cells can directly influence the melt rate of a snowpack due to the bioalbedo effect. Here we improve upon their models by creating a library of optical properties for algal cells using Mie theory and coupling this to an adapted form of the SNICAR radiative transfer scheme (Flanner et al., 2007). This new physical bioalbedo model is then used to support a critique of empirical bioalbedo studies with the overall aim of guiding future experimental design. In the first section, a brief review of radiative transfer modelling of snow and ice will precede the presentation of the model. Then, the critique of empirical bioalbedo studies will proceed by detailed examination of ten distinct challenges whose effects are quantified using the model where possible.

## 2 Radiative Transfer modelling

Several radiative transfer models (RTMs) exist for snow, the most well-known being the two-stream radiative flux models SNICAR (Flanner et al., 2007) and TARTES (Libois et al., 2013), and the plane-parallel multi-stream DIScrete Ordinates Radiative Transfer scheme (DISORT, Stamnes et al., 1988) used by the MODTRAN (Berk et al., 2014) and SBDART (Ricchiuzzi et al., 1998) atmospheric transmission models. The two stream approach is computationally efficient and able to accurately predict spectral reflectance hemispherically; however, a multi-stream approach is required to predict directional reflectance. These models are well-validated for clean snow (Grenfell et al., 1994) and snow contaminated with black carbon and dust (Painter et al., 2007, Flanner et al., 2007; Gardner and Sharp, 2010, Brandt et al., 2011). Equivalent radiative transfer schemes for glacier ice are scarce, perhaps due to the diverse range of glacier surface types. Smooth ice is best represented as a bulk medium of ice with air bubbles and cracks of known size distribution and specular reflection from the upper surface



(Gardner and Sharp, 2010; Dadic et al., 2013). In contrast, weathered ice is best described as a collection of ice grains in a bulk medium of air (or water in the case of saturated ice) and is therefore more optically similar to very coarse snow. We suggest that, since algal blooms are a feature of weathered, ablating ice in summer, an adapted snow model is more appropriate for most bioalbedo studies. RTMs are often simplified by assuming ice grains (or air bubbles) to be spherical and homogeneous, such that Mie theory can be used to determine absorption and scattering in the snowpack. This requires a priori knowledge of spectral absorption and scattering coefficients of ice, which have been well known since the 1980's (Warren, 1982) and recently up dated (Warren and Brandt, 2008; Picard et al., 2013). The effects of impurities have been studied many times with particular emphasis on black carbon and certain dusts (see review by Gardner and Sharp, 2010). Aoki et al. (2013) and Cook et al. (2017) have attempted to incorporate biological impurities into radiative transfer schemes. Aoki et al.'s (2013) model was not fully predictive as it required field spectra to back-calculate absorption coefficients while Cook et al.'s (2017) model lacked field validation. Both models characterized biological impurities by updating the absorption term in an inorganic impurity model, which can be a poor approximation given the size of cells compared to the wavelength of visible light. To date no published datasets include both reflectance at high spectral resolution and sufficient biological and physical metadata to properly validate radiative transfer models for algal snow or ice. Here we address this knowledge gap by coupling a bio-optical model that populates a lookup library of biological impurities created using Mie theory with an adapted form of the two-stream radiative transfer model SNICAR (Flanner et al., 2007). We refer to these coupled models as 'Bio-SNICAR'.

### 3 Bio-SNICAR

BioSNICAR is composed of two discrete units: a) a model that calculates optical properties of algal cells; b) a two-stream radiative transfer scheme that accesses a library of biological impurities populated using a). The library comprises NetCDF files containing the following cell-specific information (dimensionless if no unit reported):

1. Cell radius (m)
2. Mass extinction cross section ( $\text{m}^2 \text{kg}^{-1}$ )
3. Single-scattering albedo
4. Scattering asymmetry parameter
5. Effective surface area-weighted cell radius (m)
6. Number-mean radius (m)
7. Geometric standard deviation of lognormal cell distribution
8. Particle density ( $\text{kg m}^{-3}$ )

There are two stages to obtaining the optical properties of cells. First, the complex refractive index of the cells must be known. The real part of the refractive index (representing refraction) is assumed to be constant for all cells at a value of 1.5 (Pottier et al., 2005) whereas the imaginary part of the refractive index ( $k$ , representing absorption) varies according to the pigmentation of the cells. This is modelled using a bio-optical scheme developed using Pottier et al.'s (2005) study of algal cells in bioreactors (also used by Cook et al., 2017) from user-defined values of pigment abundance (as % total cellular dry mass) according to:

$$k_{\lambda} = \frac{\lambda}{4\pi} \rho_{dm} \frac{1-x_w}{x_w} \sum_{i=1}^N E a_i(\lambda) w_i \quad (\text{Eq. 1})$$

where

$$x_w = 1 - \frac{C_x}{N_T V_{32} \rho_{dm}} \quad (\text{Eq. 2})$$

and  $k$  = imaginary part of refractive index,  $\lambda$  = wavelength,  $E a$  = in vivo mass absorption coefficient ( $\text{m}^2 \text{kg}^{-1}$ ) of pigment ( $i$ ),  $\rho_{dm}$  = density of cellular dry mass ( $\text{kg m}^{-3}$ ),  $x_w$  = water fraction in cell,  $C_x$  = dry mass concentration ( $\text{g L}^{-1}$  in culture suspension),  $N$  = number of pigments,  $N_T$  = number particle density ( $\text{m}^{-3}$ ),  $V_{32}$  = mean efficient volume for the particle ( $\text{m}^3$ ),  $w_i$  = mass fraction of pigment ( $i$ ) as percentage of total cellular dry mass. The water fraction ( $x_w$ ) and density of cellular dry mass ( $\rho_{dm}$ )



are assumed constant throughout the simulations presented in this paper ( $\rho_{dm}=1400 \text{ kg m}^{-3}$ ,  $x_w = 0.8$ ; Dauchet et al., 2015); however, we provide Equation 2 to enable derivation of  $x_w$  from empirical data (see Pottier et al., 2005; Kandilian et al., 2016). It should be noted that there are limits of applicability for Equation 1, since  $k_\lambda$  will become infinite when  $x_w = 0$ . Equation 1 is therefore applicable for  $0 < x_w < 1$ , where realistic values are likely to be close to 0.8 (Dauchet et al., 2015). The model does not currently include UV absorbing pigments.

Information about the size of the cells is obtained using a second model that creates particle size distributions from the mean cell radius and standard deviation. This information can be obtained from microscopy of field samples or varied for model experimentation (e.g. Takeuchi et al., 2006). From this information the model calculates all the size information required for the optical properties of the biological impurities to be obtained using Mie theory. Mie theory utilizes the size information and the complex refractive index to predict the spectral mass absorption cross section, single scattering albedo (defined as the ratio of scattering efficiency to total extinction efficiency) and scattering asymmetry parameter. We used this approach to generate a library of single cell biological impurity optical properties for a range of sizes and pigment compositions that can be updated as required.

The library of biological impurities (Appendix 1) is coupled to an adapted form of the two-stream radiative flux model SNICAR (Flanner et al., 2007). The mass concentration of cells ( $g^{algae}/g^{ice}$ ) is prescribed in each of  $n$  layers in a vertical ice column, which is assumed infinite and homogeneous in the horizontal dimension (plane-parallel). The model allows mixing of any number of biological impurities, dusts and black carbon species within each layer provided the relevant optical properties have been calculated and added to the impurity library. By default, SNICAR derives incoming spectral irradiance using an atmospheric RTM (Zender et al., 1997) for a mid-latitude site during clear sky conditions in winter (constant solar zenith angle of  $60^\circ$ ). For simulating albedo at a particular location, the incoming spectral irradiance can be provided to SNICAR as a dataset derived from field measurements or modelling using an atmospheric radiative transfer scheme such as SBDART (Ricchiuzzi et al., 1998) or COART (Jin et al., 2006). For this paper, we used the default incoming spectral irradiance as our simulations are not for a particular field site. We have also produced a lookup library of Mie optical properties for a range of snow/ice grain radii, from  $10 \mu\text{m}$  to  $5 \text{ cm}$ , and the model is able to simulate interstitial melt water using two methods: (a) spheres of liquid water interspersed in the ice matrix (we include this for completeness but note that this introduces additional scattering surfaces and therefore does a poor job of simulating interstitial water) or (b) as liquid water coatings around the ice grains (the preferred method).

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#### 4 Empirical studies

Modelling enables quantification of changes in albedo resulting from variations in individual components of the ice-atmosphere system. BioSNICAR facilitates this for ice and snow contaminated by biological and non-biological impurities. In this section, we use this capability to identify ten specific challenges to empirical studies of bioalbedo and, where applicable, quantify their impact upon ice albedo using Bio-SNICAR.

##### 4.1 Ambiguity in albedo terminology:

There is an abundance of terminology related to surface reflectance properties in the literature, some of which is ambiguous or used erroneously (Schaeppmann-Strubb et al., 2006). Here, we define the most important concepts in an effort to standardize the terminology used in bioalbedo studies. The albedo of a surface is the survival probability of a photon incident on the surface from above (Grenfell, 2011). For snow and ice where emission at long wavelengths ( $> 4 \mu\text{m}$ ) is negligible, albedo can usually



be well approximated as the ratio of upwelling irradiance from a surface to the downwelling irradiance incident upon that surface. This can be for a specific wavelength (“spectral albedo”) or integrated over visible and IR wavelengths, weighting by the incoming spectral irradiance (“broadband albedo”). Irradiance is the radiative flux incident upon, emitted from or passing through a surface measured in  $\text{W m}^{-2}$ , whereas radiance is a directional measure of the radiative flux per unit area over a particular solid angle, measured in  $\text{W m}^{-2} \text{sr}^{-1}$ . Spectral irradiance is irradiance measured per unit wavelength with units  $\text{W m}^{-2} \mu\text{m}^{-1}$ . The term ‘bioalbedo’ is frequently used to describe the impact of biological activity on the albedo of snow and ice.

For an isotropically scattering (Lambertian) surface, incoming solar radiation is scattered equally in all directions in the overlying hemisphere such that reflected radiance in any one direction is equal to that in all other directions. However, all natural surfaces exhibit some degree of anisotropy (Schaeppman-Strub et al., 2006). Snow and ice preferentially scatter light in the forward direction, although the angular distribution depends upon lighting geometry, ice optical properties, surface roughness and wavelength (Aoki et al., 2000; Dumont et al., 2010; Hudson et al., 2006). This means that albedo measurements must be made using instruments that sense over a full hemisphere, or alternatively directional measurements must be integrated over a hemisphere weighted by the cosine of the nadir angle (Milton et al., 2009), taking into account the surface bidirectional reflectance distribution function (BRDF). The BRDF describes the amount of radiation reflected in each direction in the overlying hemisphere. Technically, the BRDF is also a conceptual quantity because it is defined by infinitesimal illumination and viewing angles. However, the BRDF can be well-approximated using measurements or simulations at fine angular resolution (Schaeppman-Strub et al., 2006) although this is difficult under natural illumination because even in clear sky conditions a large fraction of the UV and blue light is diffuse. Snow and ice albedo is also sensitive to the direction of incoming solar irradiance. For example, at oblique angles a photon can be thought of as ‘skimming’ the shallow subsurface of a snowpack and is therefore more likely to scatter back out above the snow/ice, whereas a photon entering perpendicular to the surface will experience downward forward scattering deeper into the snowpack, reducing the likelihood of the photon scattering back across the ice-atmosphere interface (Gardner and Sharp, 2010). It is also important to consider whether incoming irradiance is direct or diffuse. Direct irradiance arrives at a surface having travelled in a straight line from its source, without scattering. Diffuse incoming light, having been scattered at least once from atmospheric particles and molecules, arrives at a surface from all directions within the overlying hemisphere. Truly diffuse light is therefore isotropic. These considerations lead to several possible interpretations of reflectance measurements that have been arranged into a taxonomy by Nicodemus et al. (1977: Fig. 1A). In brief, field measurements using pyranometers or spectral radiometers with cosine-collectors can usually be best described as bihemispheric reflectance (i.e. ‘broadband albedo’ if the waveband is sufficiently broad). Spectral measurements made under natural illumination using a collimated fibre optic or multispectral imaging system are usually best described as hemispherical conical reflectance factor (HCRF, which approximates to hemispheric directional reflectance factor if the viewing lens collimates to a small solid angle, or to bidirectional reflectance distribution function if both the incoming and reflected radiance collimate to a small solid angle). Under natural conditions HCRF is similar to BRDF at wavelengths  $> 0.8 \mu\text{m}$ , but they are increasingly different at shorter wavelengths due to the strong wavelength dependence of atmospheric Rayleigh scattering (Hudson et al., 2006). Many satellite reflectance measurements are hemispherical-conical because they are the integrated signal of radiance upwelling from a naturally illuminated surface over some finite solid viewing angle (Grenfell, 2011) and therefore significant post-processing is often required to provide satellite albedo products. For conversion to albedo, these directional or conical reflectance values must be integrated over the entire overlying hemisphere, taking surface anisotropy into account (see section 4.8).

Certain spectral conditions must also be met for a measured reflectance to be termed ‘albedo’. Since broadband albedo is a crucial determinant of ice and snow energy balance, it is essential that measurements take into account a large fraction of the incident solar irradiance. Solar irradiance extends well beyond the visible wavelengths, meaning opposed photosynthetically



active radiation (PAR) sensors that typically only measure between about 0.4 and 0.7  $\mu\text{m}$  are insufficient for measuring broadband albedo. Since UV and visible wavelengths account for only about half of the incoming solar energy (Painter et al. 2012; Liou, 2002, Table 2.3), studies that report values of bioalbedo using PAR sensors (e.g. Lutz et al., 2014, 2016) likely overestimate the albedo reducing effect of biological impurities because the influence of the impurities are restricted to the visible and near-UV wavelengths. Over 99% of incident solar energy lies in the wavelength range 0.3 - 4  $\mu\text{m}$ , which should be the standard range for characterizing the broadband albedo of a sunlit surface.

#### 4.2: Characterizing snow or ice optical characteristics

Snow or ice physical characteristics are the primary driver of albedo (Fig.1B). This is because scattering occurs at ice-air interfaces. The greater area of interfaces per unit volume of snow/ice, the more likely a photon is to scatter back out of the snowpack before being absorbed (Warren, 1982). A high specific surface area (i.e. small effective grain size) is therefore associated with higher albedo, where specific surface area (SSA) is defined as:

$$SSA = \frac{S}{M} = \frac{S}{\rho_{ice} * V} = \frac{3}{\rho_{ice} * r_{eff}} \quad (\text{Eq. 3})$$

where  $S$  = surface area,  $M$  = mass,  $V$  = volume,  $\rho_{ice}$  = density of ice ( $917 \text{ kg m}^{-3}$ ) and  $r_{eff}$  = effective radius (i.e. radius of ice spheres with the same SSA as the snow or ice).

For pure snow, grain size has negligible impact upon spectral reflectance at 0.3 – 0.45  $\mu\text{m}$  wavelengths where ice is virtually non-absorbing (Fig. 1B; Warren and Brandt, 2008), however it has a major impact at the NIR wavelengths. Infilling of interstitial spaces by meltwater replaces air-ice interfaces with water-ice interfaces and promotes forward scattering of light deeper into the snow or ice, lowering the albedo (Gardner and Sharp, 2010). This is a key albedo-melt feedback that explains why wet snow has a lower albedo than dry snow (with the albedo reduction occurring mostly at NIR wavelengths). Radiative transfer schemes tend to parameterize the geometry of the ice grains using the SSA (or effective grain size) under the assumption that the real snowpack can be well described by a model snowpack of homogeneous spheres with the same total surface area and total mass as the real snowpack (Warren, 1982; Grenfell et al., 1994). However, determining these values empirically for real ice and snow is challenging. For snow, the integral of an absorption feature centred around 1.03  $\mu\text{m}$  scales with grain size and can be measured using contact-spectroscopy (Fig. 1C; Nolin and Dozier, 2000; Painter et al., 2007). However, this may not hold for large grain radii or high impurity concentrations, and the absorption feature can be shifted to shorter wavelengths by liquid water (Green et al., 2002, 2006; Gallet et al. 2014). The same may be true for NIR (0.7 – 1.3  $\mu\text{m}$ ) photographic methods (Yamaguchi et al. 2014). Other options include microcomputed tomography (Micro-CT: Dadic et al., 2013) although this requires cores to remain completely frozen through extraction, transport and analysis. Grain sizes could be determined using a polarizing hand lens (Hubbard and Glasser, 2005; Negi et al., 2010) - although the error can be high (Painter et al., 2007) - or semi-automated analysis of photographs of thin/thick sections (Dadic et al., 2013) or the sample surface.

#### 4.3: Characterizing the incoming solar irradiance

Broadband albedo is an apparent property of a surface that depends upon the spectral and angular distribution of incoming irradiance as well as the inherent optical properties of the surface itself. Incoming solar irradiance is modified by the solar zenith angle, atmospheric composition, cloud cover and cloud optical thickness (Gardner and Sharp, 2010). Clouds increase the diffuse portion of the incoming irradiance and preferentially absorb near-IR wavelengths and scatter near-UV and visible



wavelengths upwelling from the surface, resulting in a spectral shift towards shorter wavelengths. Cloud cover therefore increases the broadband albedo of the surface relative to clear-sky conditions (Grenfell and Maykut, 1977; Grenfell and Perovich, 1984; Gardner and Sharp, 2010). Furthermore, scattering and absorption by atmospheric gases and aerosols modify the spectral quality and attenuate the top-of-atmosphere solar irradiance.

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Incoming solar irradiance is best measured using a spectral radiometer with a cosine-collector oriented to look upwards (Grenfell, 1981). The cosine collector enables spectral measurements to be made over the entire overlying hemisphere, and can then be rotated to look downwards to detect upwelling irradiance, thereby enabling a surface albedo measurement. Sometimes, instruments with limited field of view (FOV) might be used (discussed in section 4.10). In this case, the spectral distribution of incoming solar irradiance is best measured using a white-reference panel (Schaeppman-Strub et al., 2006; McCoy, 2005; Milton et al., 2009). Because the panel is as close as possible to an isotropic scattering surface, light reflected from it approximates the spectral characteristics of incoming solar irradiance. Common materials used for white reference panels are Spectralon™ and Barium Sulphate (Jaglarz et al., 2006). The white reference material used must approximate a Lambertian scattering surface as closely as possible so that anisotropic reflectance and spectral absorption effects are minimized in the measurement standard, although even Spectralon™ surfaces exhibit a degree of anisotropy (Bonnefoy, 2001; Voss and Chang, 2006). For occasions where field data are unavailable, incoming solar irradiance can be modelled using the latitude, longitude, time of day and site specific information or approximations regarding the atmospheric profile and an atmospheric radiative transfer model such as SBDART or COART. In our simulations, we use the default SNICAR incoming irradiance, but direct the reader to our README file that accompanies BioSNICAR where the protocol for updating the incoming irradiance for specific sites using atmospheric RTMs is described. In areas of high relief a surface model that includes multiple reflections and site-specific radiative processes may be required.

#### 4.4: Characterizing the optical properties of individual cells

25 The optical properties of individual cells are difficult to measure in the field for several reasons. Spectral measurements resolved at the scale of a single cell require equipment that cannot be readily utilized in the field. Yallop et al. (2012) successfully employed microspectrophotometry in the laboratory to identify some pigments in ice algae. However, the freezing, transportation, storage, thawing and filtering of the samples required for this analytical technique limit its utility to the more stable pigments. Furthermore, the concentrations of each pigment (specifically those associated with xanthophyll cycling) are sensitive to changes in the light field experienced during sampling (Demers et al., 1991; Grant and Louda, 2010). Changes in the light field can stimulate or relax non-photochemical quenching mechanisms and open or close photosynthetic reaction centres in the algal cells, leading to rapid changes to their pigment profiles (Grant and Louda, 2010). With knowledge of biomass concentration and distribution and the surrounding medium, an inverse modelling approach may allow pigment mass fractions to be estimated (e.g. Bigidare et al., 1990; Bricaud et al., 2004; Moisan et al., 2011) which could in turn enable inferences to be made about environmental stresses and stress-responses in the ice surface algal ecosystem.

The simulations shown in Figure 2 demonstrate that increased secondary carotenoid concentration results in lower broadband albedo but also changes the shape of the spectral reflectance curve. Carotenoids lower the albedo between 0.4 – 0.6  $\mu\text{m}$  and have the effect of masking the spectral signature of chlorophyll and extinguishing the characteristic ‘chlorophyll bump’ (e.g. Seager et al., 2005). This effect may hinder attempts to determine cell concentrations by remotely sensing chlorophyll absorption features, despite this being feasible for other algal blooms (e.g. lakes and oceans) and vegetation mapping. We simulated four pigment scenarios: high (1.5% chlorophyll a, 10% primary carotenoids, 10% secondary carotenoids), medium (1.5% chlorophyll a, 5% primary carotenoids, 5% secondary carotenoids), low (1.5% chlorophyll a, 1% primary carotenoids,



1% secondary carotenoids) and chlorophyll-only (1.5% chlorophyll a). We were guided by published values for algal pigment content, although they vary widely because pigmentation changes in response to environmental stresses and between species (e.g. Lamers et al., 2008). For example, Griffiths et al. (2011) found pigments to comprise 0.5 – 5% of cellular dry weight in the microalga *Chlorella vulgaris*, Christaki et al. (2015) suggested between 0.5 - 8% for phycobiliproteins in algal cells and up to 14% for carotenes specifically, while 1-2% is commonly suggested for chlorophyll a (Christaki et al., 2015; Kirk et al., 1975). Secondary carotenoid content between 7.7 and 10% for secondary carotenoids was reported for light-stressed microalgae by Lamers et al. (2008). Each pigment has a unique effect on the cell optical properties. For our simulations the pigments were mixed according to Equation 1. In each scenario there was no interstitial water, the ice grains had a constant radius of 1500  $\mu\text{m}$  and the biomass of the biological impurities in a 3 mm surface layer was varied (here, 0.01, 0.05, 0.1, 0.5, 1, 1.5 and 2  $\text{mg}^{\text{alg}}/\text{g}^{\text{snow}}$ ). Increasing the biomass from 0.01 to 2  $\text{mg}^{\text{alg}}/\text{g}^{\text{snow}}$  reduced the broadband albedo from 0.68 to 0.38 for the high carotenoid scenario, 0.39 for the medium carotenoid scenario, 0.41 for the low carotenoid scenario and 0.50 for the chlorophyll only scenario (Fig. 2). This highlights the importance of pigmentation for albedo reduction but also suggests that once a threshold concentration of carotenoid pigments is present in the cell further increases have diminishing impact on their albedo-reducing efficiency. The presence of carotenoids masks the spectral signature of the chlorophyll (Fig. 2). Previous empirical studies have reported biomass concentrations as high as  $5.2 \times 10^4$  cells  $\text{mL}^{-1}$  (Lutz et al., 2014) for melting snow in eastern Greenland and  $29.5 \times 10^4$  cells  $\text{mL}^{-1}$  for ice algae in western Greenland (Yallop et al., 2012) although it is not possible to convert this into  $\text{g}^{\text{algae}}/\text{g}^{\text{snow}}$  without knowing precisely the dimensions and physical characteristics of the algal cells and snow/ice sampled.

The optical properties of algal cells also depend upon cell size. Bio-SNICAR calculates the Mie optical properties for algal cells of different radii before appending them to the lookup library. In Figure 3 we show results for small (radius = 5  $\mu\text{m}$ ), medium (radius = 15  $\mu\text{m}$ ) and large cells (radius = 25  $\mu\text{m}$ ). The model assumes the cells to be spherical and homogeneous; however, ice algal cells are often irregularly shaped. We apply an effective spherical radius approach (again a collection of spherical cells with equivalent SSAs, as was done for nonspherical snow grains) justified by the random orientation of diverse cells. Decreasing cell size for the same total biomass is shown to enhance the albedo reducing effect (the broadband albedo for the highest simulated biomass is 0.22 for small cells, 0.39 for medium cells and 0.48 for large cells). This albedo reduction is concentrated at wavelengths < 1.1  $\mu\text{m}$ , whereas at wavelengths > 1.1  $\mu\text{m}$  the albedo can actually increase with additional biomass, especially for smaller cells. These observations are consistent with inorganic impurities including soot and mineral dusts (Figs 4 and 5 in Warren and Wiscombe, 1980; Flanner et al., 2007) and results from the total cross section of many smaller particles being greater than fewer larger particles for an equal impurity mass. There is uncertainty in the biological literature concerning the impact of cell size on the optical properties of algae arising from the spatial distribution of pigments within cells (Haardt and Maske, 1987), variable intracellular architecture (Haardt and Maske, 1987), packaging effects (Morel and Bricaud, 1981) and non-linear relationships between cell growth and pigment mass fraction (Alvarez et al., 2017). These simulations confirm that empirical determination of cell size is important for radiative transfer modelling of bioalbedo, and further emphasise the need for studies focussed on the links between photophysiology and optical properties of ice algae. We expect that the range of cell sizes modelled here represent a likely range for real ice algal cells, although ice-dwelling cyanobacteria can be filamentous and several tens of microns in length (Uetake et al., 2010; Yallop et al., 2012).

#### 4.5: Measuring biomass

The biomass concentration of cells ( $\text{g}^{\text{algae}}/\text{g}^{\text{ice}}$ ) in the snow or ice medium is a crucial determinant of surface albedo (Fig. 2, 3, 4). Impurity loading is required per unit volume or mass of ice or snow rather than per unit area of surface although the distribution of cells within the volume at the time of sampling is also crucial (see Section 4.6). Cells can be counted in field samples using traditional microscopy or flow cytometry, although the two methods may produce different results (Stibal et al.,



2015). The chlorophyll content of a sample has traditionally been used as a proxy for biomass; however, the amount of chlorophyll in an individual cell is species-specific and changes over time as a mechanism of photoacclimation (e.g. Felip and Catalan, 2000) making the use of chlorophyll-based biomass estimates questionable, especially in bioalbedo studies where both the pigmentation and abundance of cells are crucial variables. For snow, Painter et al. (2001) developed a method for  
5 retrieving snow algal biomass from spectral reflectance. This is not directly applicable to algal blooms on ice for several reasons: ice algae produce additional pigments that can obscure the diagnostic chlorophyll absorption feature used by Painter et al. (2001), and dust loading on ablating ice is often more prevalent than on seasonal snow. Heavy dust loading can also reduce the reflectance in the visible wavelengths in such a way as to be indistinguishable from the chlorophyll-based ‘red-edge’ biosignature (Fig. 4C) that is commonly used to remotely detect terrestrial vegetation and ocean plankton blooms (e.g.  
10 Roling et al., 2015; Joint and Groom, 2000) prohibiting the use of this biomarker for remotely sensing ice algae unless dust can be accurately accounted for (i.e. there is a high risk of biomarker false positive). In addition, melting glacier surfaces generally have large surface roughness features (e.g. Smeets and van den Broecke, 2008) and less efficient scattering, both of which increase the anisotropy of scattering relative to snow. Increased surface and shallow-subsurface heterogeneity, combined with greater inorganic impurity loading make biological signatures more difficult to distinguish from spectral noise. The  
15 chlorophyll absorption feature, already subtle due to the presence of additional pigmentation in ice algae compared to snow algae can be further attenuated by this noise.

#### **4.6: Characterizing the depth-distribution of cells**

20 The vertical distribution of cells within snow or ice is an important factor in determining their impact on albedo (Fig. 4A). This has an important influence upon biomass measurement. When field studies report number of cells per unit volume of melted snow or ice, the depth from which the volume of ice was removed must also be reported. For two samples of equal volume from ice with biological impurities concentrated into the upper surface, sampling very shallow ice over a wider area will result in a high cell concentration compared to a narrow, deep sample where more clean subsurface ice is included. Ice  
25 algae concentrate into a very thin surface layer of the order of 1 mm (Lutz et al., 2014; Yallop et al., 2012). Best practice would involve removing samples at well-defined depth intervals. There is a practical limit of approximately 2 cm for removing weathered ice, because it is often a rough surface. However, careful visual inspection or photographic records may provide further detail about vertical distribution at spatial resolution < 2cm to support modelling. As shown in Figure 4A, the albedo impact of distributing the same biomass concentration over 1 mm and 1 cm is dramatic.

30

#### **4.7: Characterizing abiotic impurities**

Inorganic light-absorbing impurities include mineral dusts and black carbon. These impurities have an albedo reducing effect (Warren and Wiscombe, 1980; Warren, 1984; Gardner and Sharp, 2010) and can be difficult to distinguish from biological  
35 albedo reduction, having similar spectral absorption signatures in the visible wavelengths (Fig. 4B, C). Isolating the inorganic impurity loading is therefore necessary for quantifying bioalbedo. By incorporating measured inorganic impurities into a radiative transfer scheme, and knowing the ice properties, the effects of biological impurities can be determined to first order by differencing the real spectrum from a modelled inorganic impurity-only spectra (although where impurities are mixed vertically more complex non-linear unmixing techniques are required). To quantify inorganic particulates, ice samples can be  
40 melted and filtered onto quartz fibre filters and total inorganic mass measured using thermogravimetry, which can be combined with X-ray diffraction analysis to characterize the composition of the minerals (e.g. Smith et al., 2016). Scanning electron microscopy with energy dispersive x-ray spectroscopy (SEM-EDX) can provide information on the geochemistry of samples and particle size distributions. There are also many optical microscopy techniques that can be used to determine sample



mineralogy (Winchell, 1951; Shelley, 1985). It is not unreasonable to assume that the mineralogy of cryoconite is representative of the surface ice, since cryoconite acts as a store for mineral grains from all sources in the supraglacial zone (Stibal et al. 2012; Cook et al. 2016). Therefore, the same analyses can be applied to cryoconite grains with the benefit of easy sample collection and no need to filter onto papers, although the inorganic components must be extracted from the organic matter chemically or by combustion. Dust can also be delivered to the surface by melt-out from underlying glacial ice. Dust loading on ice obscures the spectral signature of algal cells in the visible to NIR wavelengths (Fig. 4C). While the effect of dusts may be negligible in many cases, the possibility of very high concentrations in some areas (e.g. Wientjes et al., 2011) and the similarity of spectral signatures resulting from biotic and abiotic impurities make abiotic impurities a crucial component of bioalbedo investigations, especially those aiming to detect biological impurities remotely or quantifying biological albedo reduction empirically. The reflectance of some minerals is difficult to distinguish from that of microbes (Seager et al., 2005) and interactions of microbes with minerals can obscure or modify biological reflectance spectra (Röling et al., 2015). Therefore, if a chlorophyll or carotenoid based biomarker is designed to be universally applicable (e.g. detection of photosynthetic life across Earth's ice sheets and glaciers as well as extraterrestrial ice) then it must distinguish biotic from abiotic impurities. Our modelling shows that for equal mass concentrations, algal cells are generally more effective albedo reducers than mineral dusts (Figs 2 and 4) although this depends upon the specific optical properties of the impurities.

#### 4.8: Anisotropic scattering

Anisotropic scattering, characterized by the BRDF, causes reflectance measured over a particular solid angle to differ from the hemispheric albedo. This occurs because ice scatters light preferentially in the forward direction (Winther, 1993). The degree of anisotropy may exhibit significant spatiotemporal variability since enhanced forward scattering is generally associated with smoother surfaces and enhanced backscattering is generally associated with rougher surfaces (Nolin and Payne, 2007) and the scattering direction is dependent upon the directionality of incoming irradiance. The correction for this is known as the anisotropic reflectance factor (ARF) which is derived from empirical angular reflectance measurements approximating to the bidirectional reflectance distribution function (Hudson et al., 2006). The ARF varies according to the optical properties of the medium and impurity content (Warren, 2013). Several studies have estimated BRDFs and HCRFs of snow and sea ice surfaces (e.g. Hudson et al., (2006); Marks et al., (2015); Arnold et al., (2002), but equivalent measurements are scarce for glacier ice in various states of melt and impurity loading, although Naegeli et al., (2015) measured HCRF for various surface types on an alpine glacier and Gruell and de Ruyter de Wildt (1999) measured anisotropic reflection from melting glaciers in two Landsat TM bands. Naegeli et al., (2017) recently showed that ice surfaces with higher impurity loads scatter more anisotropically, possibly due to locally enhanced melt. ARF's for ice and snow containing biological impurities have not yet been reported and the ARF's may change depending upon cell size, shape and pigment content.

#### 4.9: Indirect biological albedo feedback

In addition to the direct effects of algae on lowering ice albedo there are also indirect effects such as increased meltwater generation and modified ice-grain evolution. A bioalbedo-induced increase in the net radiation balance at the surface can change the rate of ice grain evolution, preferentially melt smaller grains with larger SSA's, and produce meltwater that fills interstitial pore spaces or coats ice grains. All of these indirect effects tend to reduce the number of scatters near the surface, reducing the albedo across the visible spectrum (Fig. 1A, Fig4D). These same effects are true of abiotic impurities. Biological impurities on ice are likely to be algal or cyanobacterial cells which can be prone to attachment to mineral fragments and/or forming biofilms. This may capture mineral fragments and prevent their removal by meltwater, thereby increasing the total



abiotic impurity loading on an ice or snow surface. Because of the complex interplay between these processes direct and indirect quantification of the effect of biological impurities on ice is challenging to quantify.

#### 4.10: Measurement and instrument configurations

5

To measure the albedo of a surface, a cosine collector can be used (Grenfell, 1981). This should be fixed in position perpendicular to the surface and rotated through 180° to look upwards and downwards (Fig. 5). Working under overcast skies allows for the quantification of albedo under diffuse illumination. To avoid error resulting from variable cloud conditions, downwards irradiance measurements can be made before and after upwards irradiance measurements and interpolated to the  
10 precise measurement time to calculate the albedo (S. Warren, personal communication). Depending upon the frame used to support the sensor, a shadowing correction should be applied (Nicolaus et al., 2010; For reference, Brandt et al., 2011 used a 1.7% shadowing correction). The user should maintain sufficient distance from the sensor to avoid shading or reflection from clothing.

15 For characterizing small sample surfaces on heterogeneous ablating ice (e.g. measuring the albedo of an algal bloom) the cosine-collector method suffers from several limitations. Ablating ice is often highly heterogeneous at scales of centimetres to metres meaning that albedo measurements made using cosine collectors with wide viewing-angles will necessarily be the integrated signal from a variety of surfaces. This can be mitigated to some extent by reducing the height of the sensor, but to determine the reflectance properties of a surface component with dimensions of the order of decimetres (e.g. an algal bloom),  
20 a hemispheric-conical reflectance measurements (HCR: Fig. 1A) made using a sensor with a limited field of view are likely more appropriate. The spectral radiometer fibre optic is collimated using a lens with a field-of-view that restricts the ground viewing area. To achieve this, measurements are made from a constant viewing angle (usually nadir) at a fixed distance above a sample surface (Fig. 5). The sensor alternately measures light reflected and scattered from the sample surface and a white reference panel. With knowledge of ARFs for the sample surface and reflectance panel, these measurements can be used to  
25 approximate albedo (Grenfell, 2011). This measurement also requires correction for instrument shading (Nicolaus et al., 2010; Grenfell, 2011). Finally, for reflectance-factor measurements the white reference panel must be clean and flat (Schaeppman-Strub et al., 2006; McCoy et al., 2005; Milton et al., 2009) and a white reference measurement should be taken before and after each new measurement so that changes in solar angle, cloud cover or atmospheric effects can be accounted for.

30 For both albedo and HCR measurements, there are additional instrument and procedural considerations. Prior to use, the spectral radiometer should be allowed time to initiate, due to the different warm-up rates of three internal spectral radiometer arrays. Failure to do so can introduce ‘step’ artefacts into the measured spectra. For the commonly used ASD FieldSpec Pro (Analytical Spectral Devices), optimum warm up times of 90 minutes are recommended (NERC Field Spectroscopy Facility, 2007); however, this may not be feasible in the field when battery power in low temperatures is an issue. In these cases it is  
35 possible to warm up the spectral radiometer in advance using external power or a ‘sacrificial’ battery before switching to a new battery for deployment in the field. The instrument set-up also involves selecting the most appropriate fore-optic. For measurements integrating over a given area of sample surface the height the sensor should be held can be calculated using the instantaneous field of view (IFOV) of the fore-optic. Since there are uncertainties related to the true IFOV of a spectral radiometer, even when collimated with a fore-optic (MacArthur et al., 2012), the sample surface should be a) approximately  
40 homogeneous, and b) significantly larger than the area observed by the spectral radiometer. Some degree of spatial integration is desirable for many bioalbedo studies. To avoid saturating the sensors under changing illumination, the spectral radiometer must also be optimized before each measurement to allow the sensors to adjust to new lighting conditions.



Ideally, measurements should be taken within 2 hours of solar noon to minimize changes in illumination due to solar angle. The measurement itself must be made at a constant angle (typically nadir). This can be achieved using a bubble-level on a pistol-grip (e.g. ASD's pistol-grip) or by setting the sensor on a tripod. Sensor tilt can introduce significant measurement error. Errors due to surface slope are greatest under clear skies with large solar zenith angles and can be minimised by making albedo  
5 measurements close to solar noon and corrected for in post-processing (Grenfell et al., 1994). Several pseudo-replicate measurements should be averaged to minimize the instrument error for a particular measurement. Unless specific angular measurements are being made, the user and instrument should be oriented towards the incoming solar irradiance to prevent shading the sample surface. For more specific information regarding instrument and measurement configurations for field spectroscopy, we suggest the NERC FSF instrument guide (<http://fsf.nerc.ac.uk/resources/guides/>) the ASD User Guide  
10 (<http://support.asdi.com/Document/FileGet.aspx?f=600000.PDF>), a critique of field spectroscopy by MacArthur et al. (2015) and the overview provided by Grenfell (2011).

For some studies, a measure of broadband - rather than spectral - albedo is sufficient. In this case the measurement is more straight forwards. Opposed pyranometers have commonly been used to measure broadband albedo (e.g. Cutler and Munro,  
15 1992; Yallop et al., 2012; Lutz et al., 2014) with the ratio of irradiance measured by the downwards-looking and upwards-looking pyranometers providing a measure of surface albedo. This requires that each pyranometer has a viewing geometry of 180° zenith over 360° azimuth (i.e. hemispheric) and that the spectral range of the pyranometer is at least 300-2500 nm, representing 95% of the total incoming solar energy (Kopp et al., 2005) (see Section 4.1). Tilt error is very important, so some measure of tilt angle (e.g. bubble level) is crucial for the opposed pyranometer method. The pyranometers must also be held a  
20 sufficient distance from the user or fixed to a tripod with a lever-arm to avoid shading or reflection from the operator. A consistent height above the surface must also be maintained to ensure the same size of sampling area between measurements. The guideline that measurements be taken within 2 hours of solar noon also applies. Since the change in solar angle over time will vary seasonally and with latitude and the acceptable range of solar zenith angles will vary between experiments, the best approach is to identify an appropriate sampling window for each specific study. Opposed pyranometers allow albedo  
25 measurements to be determined directly but requires the two sensors to be cross-calibrated. Using a single pyranometer and rotating upwards and downwards negates the need for cross calibration.

For albedo measurements by any method, metadata collection is crucial. The minimum required metadata can be divided into three categories: a) instrument configuration; b) illumination conditions; c) surface conditions. This metadata is specific to  
30 each measurement and is required for transparency and appropriate interpretation of the data. We provide data booking sheets for bioalbedo studies in Appendix 2. We suggest that these booking sheets, or versions thereof, can be used to standardize bioalbedo measurements for the purpose of integrating spectral reflectance, glaciological and biological measurements, and facilitate effective communication between bioalbedo modellers and empiricists.

## 35 **5 Conclusion:**

Bioalbedo is a significant component of the energy balance of glaciers and ice sheets that is yet to be quantified. The lack of understanding of the physical mechanisms driving albedo reduction by biological impurities and difficulty in isolating the biological albedo signal from inorganic impurities and ice optical properties remain significant hurdles for bioalbedo research.  
40 In this paper, a new physical model was presented, providing a framework for studying biological albedo reduction from first principles. The model was used to quantify the effects of biological variables such as biomass, distribution in an ice column, pigmentation, and cell size on the surface albedo. Biological impurities were confirmed to be potentially significant components of the surface albedo, and the optical properties of the cells were found to be crucial determinants of the magnitude



of their albedo lowering effect. However, it was emphasised that empirical knowledge of bioalbedo is currently lacking. Ten specific challenges to integrating theoretical and empirical studies or linking ground measurements to remotely sensed spectral data were identified and quantified using the model, leading to suggestions for improved protocols for field studies.

## 5 6. Code Availability

The code and all dependencies are available in our repository (<https://bitbucket.org/jmcook/biosnicar>) along with an instruction manual and README.

## 10 8 Data availability

The data used in our modelling experiments are provided in our repository (<https://bitbucket.org/jmcook/biosnicar>).

## 9 Author contributions

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JC developed the code, ran the experiments, wrote the paper and produced the figures. AJH, AG, MF, AT, CW, RB, JN, TDI and MT all provided useful comments and guidance regarding the content of the paper. AG and MF provided specific guidance on the model development and implementation. All authors refined manuscript drafts.

## 20 10 Competing interests

The authors cite no conflict of interest

## 11 Disclaimer

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N/A

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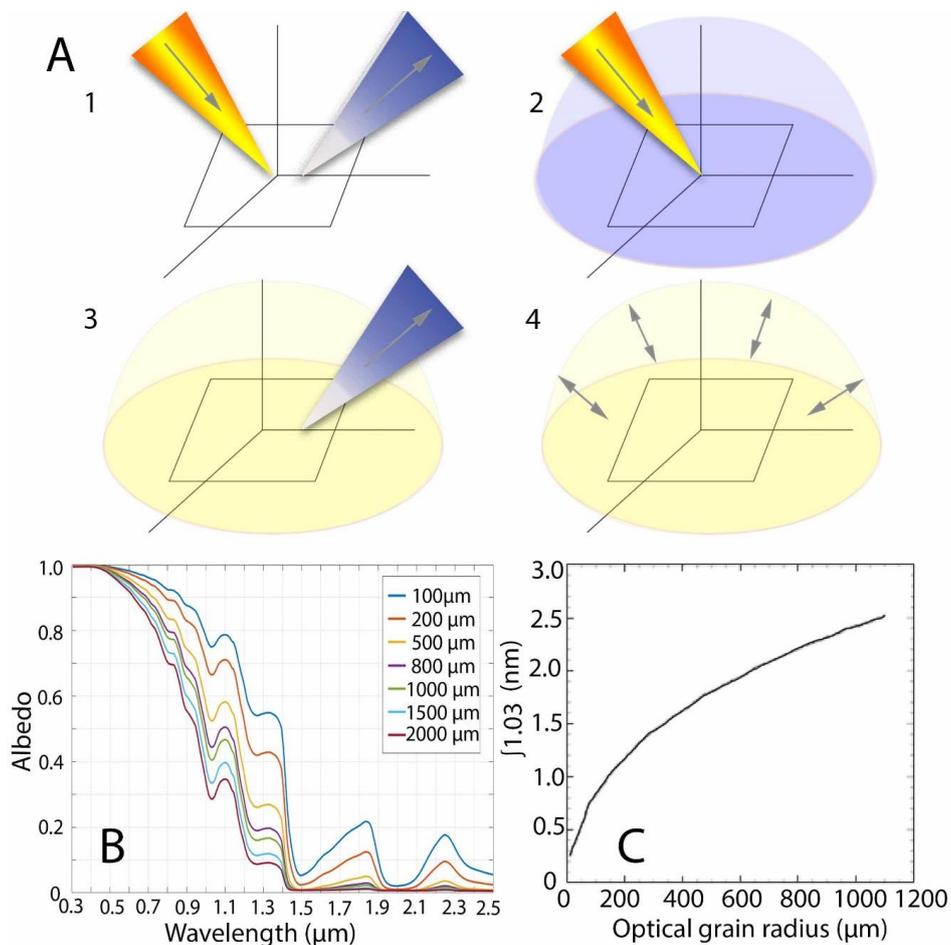
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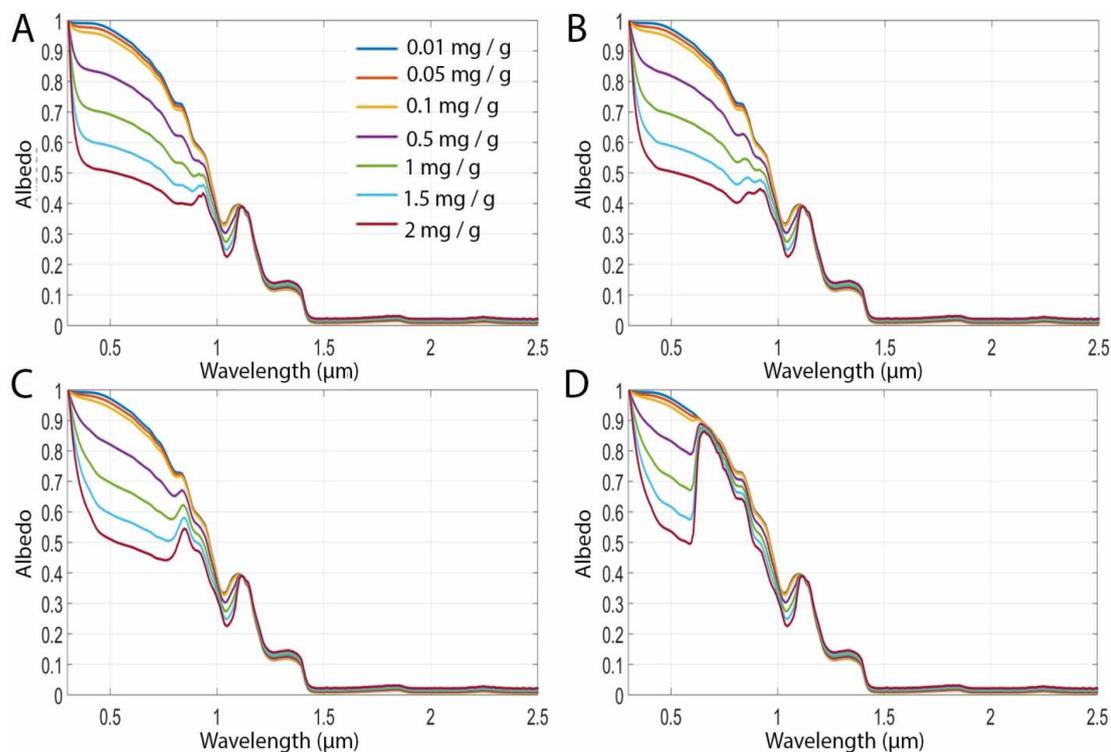
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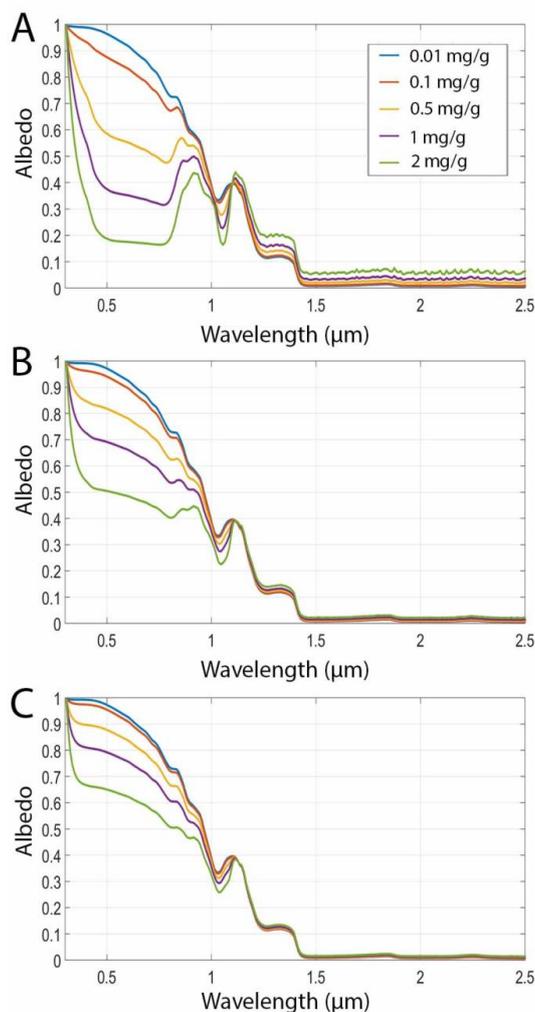
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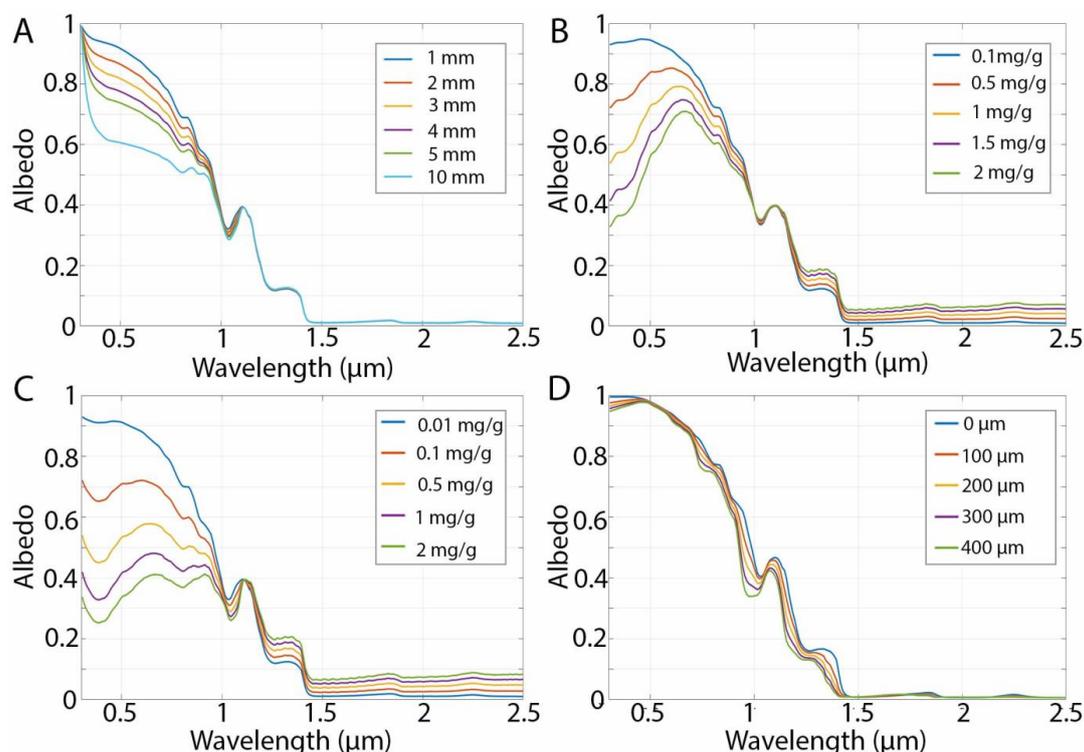
**Figure 1:** A) Diagram of the four measurable reflectance scenarios described by Nicodemus et al. (1977). 1 = biconical  
 5 reflectance, 2 = conical hemispheric reflectance (blue hemisphere represents outgoing radiance), 3 = hemispheric-conical  
 reflectance (yellow hemisphere represents incoming radiance), 4 = bihemispheric reflectance. B) Spectral albedo for clean  
 snow with grain radii 100 - 2000  $\mu\text{m}$ , solar zenith = 60°, no impurities; C) Relationship between optical grain radius and the  
 integral of the 1.03  $\mu\text{m}$  absorption feature, redrawn from Painter et al. (2007; Journal of Glaciology).



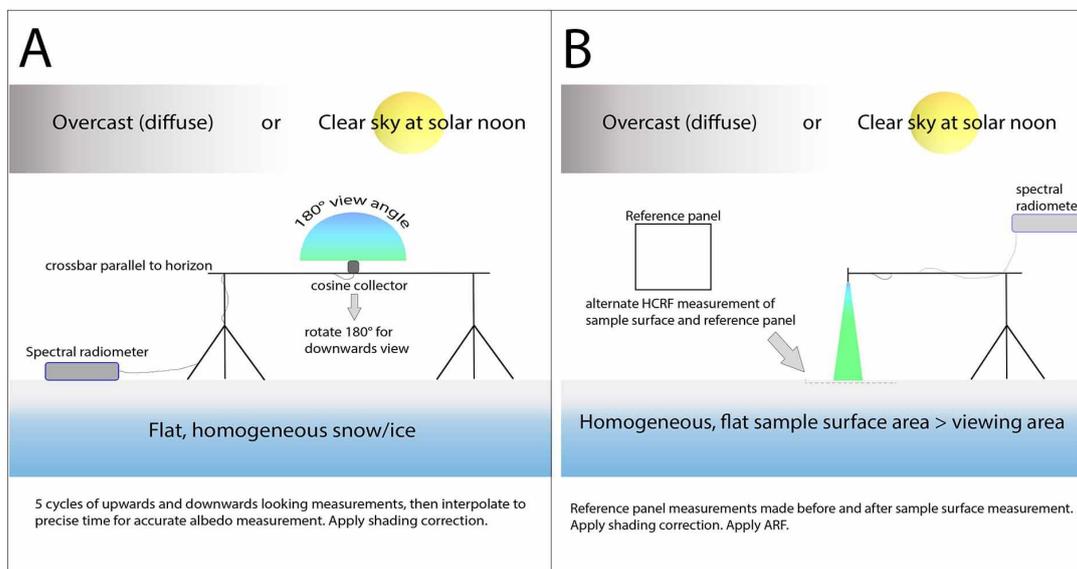
**Figure 2:** Spectral albedo of snow (grain radius 1500  $\mu\text{m}$ ) with equal biomass loading of algal cells with varying pigmentation. In all four simulations, chlorophyll a = 1.5% total cell dry weight. In A) Primary and Secondary carotenoids = 10% dry weight each. In B) primary and second carotenoids = 5% dry weight each. In C) primary and secondary carotenoids = 1% dry weight each. In D) no carotenoids are present, the cell contains chlorophyll only. In all simulations the solar zenith was 60°. The legend applies to all four subplots.



**Figure 3:** Simulations of 1500 μm radius ice grains with no interstitial water or inorganic impurities and biomass concentrations 0.01, 0.1, 0.5, 1 and 2 mg<sup>alg</sup>/g<sup>snow</sup> confined to a thin (3 mm) surface layer. The mass fraction (% dry weight) of pigments in the cells was 1.5% for chlorophyll a and 5% for each of primary and secondary carotenoids. In A) the cell radius was 5 μm, in B) the cell radius was 15 μm and in C) the cell radius was 25 μm. In all plots the solar zenith was 60°. Legend applies to all three subplots.



**Figure 4:** A constant biomass ( $0.5 \text{ mg}^{\text{algae}}/\text{g}^{\text{ice}}$ , pigment mass fractions (% total cell dry mass) = 1.5% chlorophyll a, 5% primary and secondary carotenoids,  $15 \mu\text{m}$  cell radius) distributed vertically in layers of ice ( $1500 \mu\text{m}$  grain radius) of varying thickness (1, 2, 3, 4, 5, 10 mm). B) Varying concentrations of mineral dust in a 3 mm surface layer ( $0.1, 0.5, 1, 1.5, 2 \text{ mg}^{\text{dust}}/\text{g}^{\text{ice}}$ ) on otherwise clean ice (grain radius  $1500 \mu\text{m}$ ). The dust used was SNICAR's 'dust 4' which has grain radii  $2.5 - 5 \mu\text{m}$ ; C) Equal mass concentrations ( $0.01, 0.1, 0.5, 1, 2 \text{ mg}^{\text{impurity}}/\text{g}^{\text{ice}}$ ) of algal cells (pigment mass fractions (% total cell dry mass) = 1.5% chlorophyll a, 5% primary and secondary carotenoids,  $15 \mu\text{m}$  cell radius) and mineral dust (SNICAR's 'dust 4' which has grain radii  $2.5 - 5 \mu\text{m}$ ) in a 3 mm surface layer in otherwise clean ice ( $1500 \mu\text{m}$  grain radius); D) Albedo of a dry snowpack (grain radii =  $1000 \mu\text{m}$ ) and snowpacks with liquid water as a coating around the ice grains. The legend indicates the thickness of water layer around a  $1000 \mu\text{m}$  ice grain.



**Figure 5:** Schematic diagram of albedo measurement configurations for A) cosine collector, and B) HCRF measurement modes.

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## Appendices

### A1: Lookup Library for Biological Cells

For all cells:

- 5 Real Part of Refractive Index = 1.5
- Water fraction = 0.8
- Density of dry cell = 1400 kg m<sup>-3</sup>

Name	Size (µm)	Chlorophyll a	Chlorophyll b	1 Carotenoids	2 Carotenoids	Phycocyanin	Phycoerythrin
Biological Impurity 1	30	1.5	0	10	10	0	0
Biological Impurity 2	30	1.5	0	5	5	0	0
Biological Impurity 3	30	1.5	0	1	1	0	0
Biological Impurity 4	30	1.5	0	0	0	0	0
Biological Impurity 5	10	1.5	0	5	5	0	0
Biological Impurity 6	50	1.5	0	5	5	0	0

10 Table 1: Overview of cell size and pigment mass fractions (% total cell dry weight) for each biological impurity

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