#### Response to the Interactive comment on

"Basal drag of Fleming Glacier, Antarctica, Part A: sensitivity of inversion to temperature and bedrock uncertainty"

by Chen Zhao et al.

Anonymous Referee #1 Received and published: 14 Feb 2018

We are grateful to Reviewer 1 for the positive and constructive suggestions to improve our paper. We have addressed all comments below. The line numbers in the responses are based on the revised manuscript without track changes.

Please note that Mathieu Morlighem created the ice thickness data for the Fleming Glacier system using mass conservation method, which is very important for most experiments done in this study. We do value his contribution to this paper, so we add him as the co-author in the revised text.

In response to the reviewer 2's question about our choice of enhancement factor, we implemented a new sensitivity test. This was more thorough than our original test, and with a more up-to-date setup. And in fact, it reveals that our original choice was not optimal. So we added the sensitivity tests for various E values (0.5, 1.0, 2.0, 4.0) in Sect. 3.6 and Sect. 4.2, and the optimal value of 1.0 was chosen as the E in the CONTROL experiment. We redid all the simulations and modified the text and figures accordingly. Our conclusions have not changed.

## General comments

This paper presents results from a series of Elmer/Ice simulations of the Wordie Ice Shelf-Fleming Glacier system in West Antarctica. It aims to demonstrate the sensitivity of model inversion to englacial temperature, bedrock topography and ice front boundary, as well as provide a realistic basal shear stress field. It uses a similar multi-step inversion process to Gladstone et al. (2014), where surface relaxation is followed by an inversion for basal friction coefficient (C); then a steady-state temperature simulation using this C and velocity; and another inversion using the steady-state temperature. This process is applied iteratively in three cycles, which they show helps remove the dependence on the initial temperature distribution. They argue this is particularly important to Fleming Glacier given the sensitivity of the system to englacial temperatures, due to the dominance of internal deformation over basal sliding. Using one of the initial temperature distributions, they run the inversion process several more times, testing the sensitivity of the inverted basal traction coefficient to bed geometry (e.g. bedmap2 versus mass-conserved), and the ice front boundary condition.

Overall this manuscript is well structured and clearly written, although some of the description of figures and discussion of results are fairly laborious and may benefit from being reduced in length. The conclusions are clearly supported by the results presented. I recommend this manuscript is published in The Cryosphere, provided the authors address the following comments.

#### Specific comments

Line 47: "especially for small-scale glaciers." Not sure this is relevant, or are there

papers that show greater sensitivity of small- over large-scales systems?

To our knowledge, no study has shown that. We removed "especially for small-scale glaciers".

Line 45 - 50: These two sentences appear to be contradicting each other – firstly you say that these uncertain quantities pose a challenge for modelling basal shear stress, and then you say they are not important (to that particular ice cap). I wonder if it's worth holding off on discussing the results of the Vestfonna studies until the discussion.

The first sentence is a general statement for most glaciers, which we quote Vaughan and Arthern (2007). But, the Vestfonna Ice Cap is mentioned as a case showing the less sensitivity to the basal topography, which is contrasting to what we find for the Fleming Glacier in this study. The Fleming Glacier is the main focus of this paper, but we think it is good to mention the Vestfonna here.

Line 132: Why do you make this assumption? I know it is discussed further on that the ice shelf is effectively only 1.5 km long by 2008, but before knowing this, this statement seems strange, especially given that an ice shelf is mentioned previously.

We did not have a clear way to provide the ice thickness for a short fringing ice shelf left, which is detected from the DEM data in Jan 2008 (we clarify this in Sect. 2.1 and Fig. S1). This small ice shelf disappeared in Apr 2008, as shown in Fig. 1c. To discuss the sensitivity to different ice front position, we expanded remarks in Sect. 3.6 and Sect. 4.3, and relevant results and discussions have been added to the text.

#### Line 163: What is your justification for using a linear sliding law?

Different sliding laws in inverse modeling will not change the inversed basal shear stress distribution, and it will just lead to different basal friction coefficients based on different sliding law. In diagnostic studies that invert to find the basal shear stress which gives the best agreement with observed surface velocities, the choice of sliding "law" is not relevant provided that the required stress can be generated by adjustments of the parameters in the sliding law – in this case the coefficient C. The inversion procedure modifies C to modify stress – adjusting the momentum balance. That solution of the Stokes equation provides an updated estimate of basal velocity – which enters the next cycle of the inversion search. The question does remain whether this is physically suitable relationship to apply when the system is evolving, but this is not relevant here. So we adopted the simplest sliding law here following Gagliardini et al. (2013); Gillet-Chaulet et al. (2012). We clarified this in the text (Line 190).

## Line 263: What do you mean by "imposed by a neigboring glacier"?

We made a hypothesis that the ice front of the Fleming Glacier had a continuation of the advancing glacier by exerting a normal stress on the ice front. Here we modified into "imposed by a hypothetical undeforming continuation of the advancing glacier".

## Line 274 - 277: This seems out of place here, and the related discussion in Section 4.3 is not obvious.

Now that we have adopted the E of 1.0 as the CONTROL setup, we find that the surface lowering near the ice front during the surface relaxation was <25 m in each cycle. But we still need to know whether the small changes in surface elevation at the ice front will affect the basal friction deduced from inversion, which is discussed in Sect. 4.4. So we modified this sentence to a separate paragraph and modified the

relevant discussions in Sect. 4.4 (Line 333-337).

Line 334: add ", than CONTROL" to end of sentence? The similarity between BEDZC and BEDMC compared to CONTROL seems unsurprisingly, i.e. the two surfaces are more similar than the two thicknesses.

We added ", compared with CONTROL".

Line 352: Possibly worth mentioning Sun et al. (2014) here as another study that demonstrates the sensitivity of grounding line dynamics to bedrock topography.

#### Added.

Line 356 – 357: This seems unsurprising seeing as BEDZC makes use of surface (and mass conserved thickness?) from 2008, the same year as the velocity observations.

Yes, we agree. To clarify this, we clarified this in Line 438-440) "Both BEDMC and BEDZC use the 2008 surface DEM and this improvement over the Bedmap2 surface DEM in CONTROL appears significant, even before turning to the matter of ice thickness."

## Line 387 – 88: Why doesn't altering the sea level affect the grounding line position?

We ran all the experiments with the grounding line fixed. The sea level adjustments are meant as a convenient tool for altering the force applied at the ice front, including the influence of uncertainties in ice thickness (and hence bed depth) at the ice front/grounding line. We have clarified this in Line 302-304.

Note that the height above buoyancy calculations for 2008 in the companion paper (Zhao et al., companion paper) indicate that the glacier – as described by our datasets – would have remained grounded at the ice front for all but the largest sea level forcing.

#### Line 420 – 422: Not sure "spreading" is the right word: spreading in which direction?

By "spreading" we meant longitudinal extensional flow. We modified this sentence into "The lowered surface at the ice front in experiments IFBC1 and CONTROL is apparently the consequence of rapid deformation due to its own weight (longitudinal extension with locally high vertical shear) of an ice cliff, which is over 100 m higher than the control sea level." (Line 526-529).

## Technical corrections

Line 21 - 23: unnecessary repetition of "temperature-dependent" deformation, combine to one sentence

## Modified.

Line 67: add comma at end of line

## Added.

Line 108: Here FG is used for Fleming Glacier, whereas previously FGL is used. I suggest you use FG consistently (to me GL is grounding line).

## Modified "FGL" into "FG" in whole text.

Line 184: Inconsistent use of basal sliding/drag/friction coefficient, as well as inconsistent use of boldface C. Discuss results in the present tense

Modified all these terms "basal sliding/drag/friction coefficient" into "basal friction

coefficient". The font in equations is unchangeable so we could just make sure all the C in the main text shares the same font. We have adjusted the tenses used in the paper for consistency.

Line 295: remove quotations from CONTROL"

Deleted

Line 353: "most accurate", rather than "best"?

Modified into "more accurate"

Line 403 – 410: remove quotations from simulation names, e.g. "IFBC3".

Deleted

Figure 8: Could the 1500 m/yr contour be included in the other plots too, to help with comparisons?

Added

## References

- Gagliardini, O., Zwinger, T., Gillet-Chaulet, F., Durand, G., Favier, L., de Fleurian, B., Greve, R., Malinen, M., Martín, C., Råback, P., Ruokolainen, J., Sacchettini, M., Schäfer, M., Seddik, H., and Thies, J.: Capabilities and performance of Elmer/Ice, a new-generation ice sheet model, Geosci. Model Dev., 6, 1299-1318, 2013.
- Gillet-Chaulet, F., Gagliardini, O., Seddik, H., Nodet, M., Durand, G., Ritz, C., Zwinger, T., Greve, R., and Vaughan, D. G.: Greenland ice sheet contribution to sea-level rise from a new-generation ice-sheet model, The Cryosphere, 6, 1561-1576, 2012.
- Vaughan, D. G. and Arthern, R.: Why Is It Hard to Predict the Future of Ice Sheets?, Science, 315, 1503-1504, 2007.
- Zhao, C., Gladstone, R., Zwinger, T., Warner, R., and King, M. A.: Basal friction of Fleming Glacier, Antarctica, Part B: implications of evolution from 2008 to 2015, The Cryosphere, companion paper. companion paper.

Response to the Interactive comment on

"Basal drag of Fleming Glacier, Antarctica, Part A: sensitivity of inversion to temperature and bedrock uncertainty"

by Chen Zhao et al.

Anonymous Referee #2 Received and published: 07 Mar 2018

We are grateful to reviewer 2 for the positive and constructive suggestions to improve our paper. In particular, we now explore the effect of moving the location of the ice front, as suggested. We have addressed the comments below. The line numbers in the responses are based on the revised manuscript without track changes.

Please note that Mathieu Morlighem created the ice thickness data for the Fleming Glacier system using mass conservation method, which is very important for most experiments done in this study. We do value his contribution to this paper, so we add him as the co-author in the revised text.

In response to the question about our choice of enhancement factor, we implemented a new sensitivity test to enhancement factor (E). This was more thorough than our original test, and with a more up-to-date setup. In fact it reveals that our original choice was not optimal. We added the sensitivity tests to various E values (0.5, 1.0, 2.0, 4.0) as described in Sect. 3.6 and discussed in Sect. 4.2. The optimal value E =1.0 was chosen as the enhancement factor in all the other experiments. We redid all the simulations and modified the text and figures. We retain the sensitivity tests for the multi-cycle inversion scheme as the first results presented, since in all other cases only the third cycle results are discussed. Our conclusions have not changed.

#### General comments

This paper from Zhao and colleagues evaluates the sensitivity of the inversion of the basal friction coefficient of Fleming glacier, Antarctica, to (i) initial (i.e., before the inversion) temperature, (ii) different bed topographies and (iii) ice front boundary conditions. The simulations are performed with a control inverse method (MacAyeal, 1993) implemented in the Elmer/Ice ice sheet model and uses the full Stokes version of the Elmer/Ice model.

The novelty here is the use of a three-cycle spin-up (initially proposed in Gladstone et al, 2014, but for one cycle) scheme to avoid the influence of initial temperature field on the final inversion results.

The paper is quite long compared to what it could be. There is a substantial number of repetitions, which should be avoided when possible. The figures are not very clear, some differences pointed out by the authors between experiments being barely visible, thus I was not always sure by how much the three cycle methods improved the inversions results. In many places in the text I was often doubtful about the assertions. Moreover, I am not an English native speaker, but I am sure that the English could be improved. Related comments are written down below.

I have a concern with the Bedmap2 data. Since this is not written in the paper, I would

like to be sure that the authors removed the difference between the Geoid and Ellipsoid height, as they did for the other DEM used, which led to have 15m of sea level. If no mistake was made with the Bedmap2 data, could you please adapt your figures to a sea level at 0, which is more common?

All data used in the study are self-consistent which is the key concern here. In this study we adopted an ellipsoidal height references for all datasets (surface and bedrock elevation data) (WGS84 ellipsoid). To clarify this, we added a few words in Sect. 2.2 (Line 96-98) "The first is from the Bedmap2 dataset (Fretwell et al., 2013), with a resolution of 1 km (hereafter bed\_bm; Fig. 2b), which is converted from the EIGEN-GL04C geoid to WGS84 ellipsoid heights."

We don't think the issue of the reference value of sea level should cause confusion and we are sure all the elevation data is under the same height reference system. To be quite clear – the 15 m sea level elevation is determined from examining the 2008DEM used in the paper for the difference between elevations over the ocean and the glacier. But we agree to adapt my figures (Fig. 2b and Figs. 7g-i) to the meters above sea level with a sea level at 0 m.

I question the last experiment that consists in applying different sea level at the ice front in order to deal with the uncertainties linked to the potential presence of ice mélange, the proximity of icebergs that could push back the ice stream... First, this case need to be documented with literature, or, it needs to be strongly argued. Neither the former nor the latter is done here.

The mélange issue is not the main or only reason for exploring different force balances at the ice front – as stated in Sect. 3.6 (Line 301-305). Uncertainties in ice thickness/bed elevation are also a major consideration. The emergence of a curious sticky spot with high basal friction adjacent to the ice front further encouraged these sensitivity tests. Many previous studies have also argued that the ice mélange could suppress calving by exerting a buttressing force directly on the glacier terminus (Amundson et al., 2010; Krug et al., 2015; Robel, 2017; Todd and Christoffersen, 2014; Walter et al., 2017). We have added this in the main text (Sect. 3.6, Line 300).

Another thing is that the authors have an uncertainty on the position of the ice front, I think a better experiment would be to assess the sensitivity of the results to the position of the ice front, even though I don't think that changing it by 1.5 km (the uncertainty) would significantly change the results.

Thanks for this good point. We additionally conducted sensitivity tests to three different ice front positions in Sect. 3.6. It did not make a significant change, as expected by the reviewer, but different ice front positions affected the basal friction near the ice front. Relevant results and discussions have been added in Sect. 4.4.

I had issues understanding how you chose your experiments. For example, why choosing -20 C as an initial temperature pre-inversion? Is this number related to anything real, such as a yearly average temperature? In the paper from Schaffer 2012 that you cite, their cold and warm scenario were linked to observations, which is what you should do here, or at least explain how you chose those temperatures.

We don't have any observations for the temperature field except for the surface temperature from RACMO model, which ranges from -26 C to -7 C. The choice of -20 C or -5 C as an initial englacial temperature is not based on observations. In the Glen Flow law, the ice temperature is a function of pressure melting point via the Arrhenius law (Gillet-Chaulet et al., 2012):

 $A = A_0 e^{(-Q/[R(273.15+T)])}$ 

Here,  $A_0$  is the pre-exponential factor and Q is activation energy.  $A_0$  and Q have different values while the temperature T is lower or higher than -10 C. To test the sensitivity of inverse methods to the initial englacial temperature, we assumed two constant values, one is lower than -10 C and the other one higher.

The authors need to be consistent with the terms basal drag, basal friction coefficient, basal sliding coefficient, basal shear stress. They keep mixing up those terms all over the text to mostly talking about the basal friction coefficient.

Thanks for pointing this out. We have changed all the terms to use "basal friction coefficient".

Finally, I would recommend this paper to be merged with its companion paper, also in The Cryosphere Discussion, which deals with simulating the evolution of Fleming Glacier from 2008 to 2015. All those sensitivity analysis (the first two for me) that were done in this inversion are to me verifications that you start with a sufficiently good initial state. This is not my choice of course but the one of the editor.

This paper proposed the multi-cycle spin-up scheme to remove the effect of the plausible initial temperature assumption for the glaciers like the Fleming Glacier, which have strong, temperature-dependent, deformational flow in the fast-flowing regions. Sensitivity tests to various bedrock datasets and ice front boundary conditions for the Fleming system provided a good initial state and setting up for further simulations on this system. If we combine this paper with its companion paper, most of the above points would have to be put into the supplementary sections, which is not good for benefiting more researchers interested in the technical spin-up aspect. So we prefer keeping the two papers separate. In particular, with the addition of the ice front position sensitivity tests suggested by the reviewer this paper contains quite sufficient material to stand alone.

In all cases, this paper needs substantial rewriting before publication.

Specific comments

120: I don't think you have done a sufficient number of experiment to say so, at least to say it this way. Would you explore other glaciers with the same conclusion, this assertion would be more justified.

We gave this conclusion for glaciers like the Fleming system. To clarify this, we combined this sentence and next sentence into "This is particularly important for glaciers like the Fleming Glacier, which have areas of strongly temperature-dependent, deformational flow in the fast-flowing regions". We also modified "three cycle" into "multi-cycle" (Line 23-25).

122: Is it true ? Looking at your fig7 I see Vb/Vs=1 over a substantial area in the ice stream part ? Means that vertical deformation here is not significant...

Looking at Fig. S5b, there is a steep region between the 1000 m yr<sup>-1</sup> and 1500 m yr<sup>-1</sup>, where Vb is much smaller than the Vs. It means that the vertical deformation in the some parts of the fast flowing regions is significant.

124: You have done some sensitivity test, but I am not sure that those tests specifically show the importance of what you say. I go back into this below.

We respond to this later at the relevant point.

128: Here you put the glaciers of the AP and the WA ice sheet in the same category. The way those two parts of Antarctica are losing mass is fundamentally different and you should mention those differences.

We are aware that the ice shelf collapse in the AP is likely significantly driven by surface melting, and the ice shelves in the AP are more vulnerable to atmospheric warming. However, the Fleming Glacier in this study had nearly lost its ice shelf (the Wordie Ice Shelf) by 2008.

In recent studies on the Fleming Glacier (Friedl et al., 2018; Walker and Gardner, 2017), it is proposed that the glacier acceleration and thinning is likely to be triggered by the incursion of warm ocean water, associated with grounding line retreat, which has shown the possibility that some glaciers of the AP may lose mass in the same way with those in the WA.

We agree with the reviewer that it is important to consider both the similarity and difference between these regions, and we do extensively discuss this in our companion paper (Zhao et al., companion paper).

131: this sentence (mostly the same as in 114) is the kind you would find in an abstract but neither in the introduction nor in the main text.

We think this comment is a personal preference rather than a scientific critical argument. If the reviewer wishes to give a reason why it is not appropriate to put this sentence in the Introduction, we would consider removing it. Regarding the apparent duplication, our view is that an Abstract is a summary, not a substitute for aspects of the Introduction

133: Is this always the case ? Fast flowing outlet glaciers can have a small slope and be driven by basal sliding mostly, such as for the Siple coast glaciers... Could you rephrase.

We modified this sentence into "The high velocities of fast-flowing outlet glaciers arise from internal ice deformation or ice sliding at the bed or both." (Line 35-36).

135: This way, all those processes appear to have equal impacts onto the dynamics whatever the situation...Could you rephrase. And remove strongly.

We simply listed all the relevant factors regarding deformation here and we are not emphasizing the importance of each impact. We are happy to remove "strongly" since we do not discuss relative importance.

137: Same remark as above. What is disturbing is that you seem to put all those things in the same order in influence whatever the situation.

Same response as to 135.

140: Again, this kind of sentences should be in the abstract not here, at least to me.

Same response as to 131.

142: What you infer primarily with inverse methods is basal friction (or sliding) coefficient (sometimes ice rheology). Could you rephrase.

Modified "basal shear stress" to "basal friction coefficients", added "ice rheology". An inversion could produce basal velocities but it deduces basal shear stress by adjusting the basal friction coefficient in the description of basal shear stress inside a sliding law as a boundary condition to solving the momentum balance equations. So we don't agree that the basal shear stress is not the target of the inverse approach here.

144: In topography, do you put basal and surface topography ? I don't think so. Maybe use the term geometry or thickness and surface topography, because we need the thickness and one of the two surfaces... Please rephrase.

Modified "glacier topography" into "glacier geometry"

147: Why especially for small scale glacier ? We have major challenges for modeling temperature in the bigger glaciers as well. I understand you want to guide the reader to you specific case, but this comment is misleading.

Thanks for the reviewer's suggestion. We deleted "especially for small-scale glaciers".

148: I feel like your analysis mostly relies on those two publications dealing with the same glacier. from that you generalise things that should not be.

It is not our intention to generalize the Vestfonna case here. The Fleming case turns out to be a contrasting one. We are happy to address it further if there are specific concerns about it.

149: What type of inverse methods, did they use many ? Rephrase please.

We unintentionally suggested they used a range of techniques – they used the "Robin inverse method". We corrected this in the text (Line 51).

149: A lot of things here are not correct or need to be rephrased. 1) the results of Schafer2012 have a dependence to mesh resolution (you should read section 4.3). 2) this is not as simple as that for bed topography and velocity uncertainties. You should be less approximative in your assertions.

1) Thanks for pointing out this. Yes, the results of Schafer 2012 emphasized the importance of a finer mesh. So we delete "mesh resolution or".

2) The Sect. 4.4 of Schafer 2012 did show that the inverse method is not sensitive to the modification of the surface and bed elevation datasets.

151: This sentence is not clear, rephrase please.

Modified into "In their case, sliding dominated the flow regime, and the impact of internal deformation on ice velocity was relatively small compared to the important role of friction heating at the bed on the basal sliding" (Line 52-54).

152: And I don't think you are doing this generalisation in your paper. This is clearly overstating to me.

We just state that "No generalization on these findings to Antarctic outlet glaciers has been investigated", but we did not mean to do this generalization in this paper. To make it clearer, we changed this sentence into "It is unclear whether this property is specific to Vestfonna situation or if it also applies to other fast flowing glaciers." (Line 55).

154: Do you test this to all the inversion methods. please rephrase.

Modified into "to test the sensitivity of a variational inverse method (MacAyeal, 1993; Morlighem et al., 2010) for basal friction to basal geometry and to an assumed initial englacial temperature distribution for a different outlet glacier system" (Line 56-59).

156: What robust means here ? You will have tested on one single friction law, and almost the simplest one. You should rephrase.

"Robust" here means the robustness of simulated basal friction coefficient distribution to experiment design and the mismatch between the simulated and observed surface velocities. We don't want our simulated results to be dependent on our initial temperature assumptions. As discussed in the response to Reviewer 1, in diagnostic studies of the type we present here, the claimed physical character of the basal friction law is of little importance (assuming that it can produce the required range of basal shear stresses) so reliance on a single friction law is not a limitation. So we think it is appropriate to use "robust" here.

## 160: Maybe here you could add some figures, what are the velocities, the size, some more details about the glacier...

We added a sentence (Line 63-68) "The Fleming Glacier (FG) (Fig. 1b), as the main tributary glacier, has a current length of ~80 km and is ~10 km wide near the ice front (Friedl et al., 2018). This glacier has recently shown a rapid increase in surface-lowering rates (doubling near the ice front after 2008) (Zhao et al., 2017), and the largest velocity changes (> 500 m yr<sup>-1</sup> near the ice front) across the whole Antarctic over 2008-2015 (Walker and Gardner, 2017). "

165: You invert the basal friction (or sliding) coefficient. You need to be consistent over the text.

## Modified for whole text.

166: What you invert is the basal friction coefficient. Rephrase please.

## Modified.

180: Just a question here to be sure because you don't mention it after. Did you make sure you accounted for the Geoid-Ellipsoide difference for Bedmap2, which reference is the Geoid?

Yes, we adopted the bedmap2 data based on the WGS84 ellipsoid and we clarified this in Sect. 2.2 (Line 96-98). "The first is from the Bedmap2 dataset (Fretwell et al., 2013) with a resolution of 1 km (hereafter bed\_bm; Fig. 2b), which is converted from the EIGEN-GL04C geoid to WGS84 ellipsoid heights. " See also the discussion above under response to General Comments.

182: This is rather strange and unusual to use sea level of 15m. It would be much clearer to take the geoid as the reference.

As we stated above, the value of sea level will not make a difference in our experiments as long as we are sure all the elevation data is under the same height reference system. To be quite clear – the 15 m sea level elevation is determined from examining the 2008DEM used in the paper for the difference between elevations over the ocean and the glacier. But we agree to adapt my figures (Fig. 2b and Figs. 7g-i) to the meters above sea level with a sea level at 0 m.

186: Since you mentioned the Wordie ice shelf in the previous section, you should replace "This"

"this region" -> "the WIS-FG system"

187: shear stress - > friction coefficient

## Modified.

195: Could you break down this sentence in two parts, otherwise this is hard to read.

## Modified.

# 1100: To calculate the Hmc, did you use ElmerIce ? I think it needs to be mentioned since this would not be an official feature in Elmer.

No, we calculated Hmc using ISSM's mass conservation algorithm (Morlighem et al. 2011). We clarified the manuscript accordingly (Line 105-111) "H<sub>mc</sub> (where "mc" refers to "mass conservation") is the ice thickness data with a resolution of 450 m covering three regions shown in Fig. 2e. H<sub>mc</sub> for the yellow area is computed using the Ice Sheet System Model's mass conservation method (Morlighem et al., 2011; Morlighem et al., 2013), based on ice thickness measurements from the Center for Remote Sensing of Ice Sheets (CReSIS), using ice surface velocities in 2008 from Rignot et al. (2011b), surface accumulation from RACMO 2.3 (van Wessem et al., 2016) and 2002-2008 ice thinning rates from Zhao et al. (2017). The thickness data for the grey area is interpolated from Bedmap2 (Fretwell et al., 2013), while the data in the red area ensures a smooth transition between the two regions. The yellow area indicates the Fleming Glacier system with ice velocity >100 m yr<sup>-1</sup>."

## 1103: This is not really true for bedzc since Sbm has a resolution of 1000m. How did you interpolate Sbm from 1000m to 500m ?

We presume you meant to talk about bed\_mc here. We used a bilinear interpolation to downscale Sbm to 500 m. We have clarified this in the manuscript (Line 103).

## 1107: could you mention the fact that they are both part of the same basin.

Whether or not they are in the same "basin" depends on one's precise definition of a basin. What we mean is that each of these features has its own local minimum in bedrock elevation and a significant region of reverse bed slope. We have modified the text to make it clearer to the reader that both features are under the Fleming main trunk (Line 121-123).

## 1112: shear stress - > friction coefficient

Modified.

1124: basal drag - > basal friction coefficient

Modified.

1134: Here you need to mention the difference that you have between your reconstructed ice front and the grounding line of Rignot2011a.

Here we mentioned that the ice front position in 2008 was assumed to be same with the 1996 grounding line of Rignot et al. (2011a). So there is no difference here.

## 1144: My personal viewpoint is that the mesh resolution influence should always be checked beforehand... This is not such a strenuous task to do this.

Another experiment has been done with 20 vertical layers. The simulated C shows nearly the same distribution as the CONTROL experiment. So we modified this sentence into "In the current study an experiment with 20 extruded layers (not shown) gives very similar results as with 10 layers, confirming those findings also apply to the WIS-FG system." (Line 164-165).

## 1149: The temperature is fixed to what dataset ?

The surface temperature is fixed to the yearly average surface temperature over 1979-2014 computed from RACMO2.3/ANT27. We have moved the relevant paragraph

after this sentence (Line 173-179).

1152: You describe the BC and then you switch into something different, which should be more in the discussion section, not here. This way of writing just affect the reading in a bad way. Please consider not doing this in the text.

Thanks for your suggestion. We deleted this sentence. The uncertainties of ice thickness and bedrock topography, the low accuracy of ice front and grounding line locations, and the possible buttressing on the ice front by partly detached icebergs and ice mélange are now discussed in Sect. 3.6 and Sect. 4.4.

## 1159: Temporarily : what does it mean ?

Thank you for the query. We meant "temporally fixed" and have corrected accordingly.

1169: Ah here you talk about temperature data. It should be written in the same place as above.

This whole paragraph has been moved to the Line 173.

## 1178: Ok, Why 0.2 ? Did you check other values ?

Yes, we checked longer time and shorter times. Shorter time was not enough for Elmer/Ice to remove the non-physical spikes, which would lower the efficiency of following inverse running. If we relaxed the free surface for longer than 0.2 yr, the relaxed surface was much lower or higher than the observed one, since the simulated velocity close to the front was very high.

## 1186: drag - > sliding

## "drag" -> "friction"

1187: As there are many types of cost functions in the literature, you should define yours.

Added.

1193: Here I think you should cite Gillet2012 as it seems that you do exactly the same thing for the cost function

Added "(following for example Gillet-Chaulet et al. (2012))"

1200: You should add a figure showing the improvement made with E=2.5. I would also be very pleased to see the L-curve, for instance in a supplementary.

Thanks for your suggestions. The L-curve analysis figure has been added as Fig. S2 in the supplementary material.

As we mentioned above, we implemented a new sensitivity test to the enhancement factor E. This was more thorough than our original informal test, and with a more up-to-date setup. And in fact it reveals that our original choice was not optimal. So we added the sensitivity tests to various E values (0.5, 1.0, 2.0, 4.0) in Sect. 3.6 and Sect. 4.2, and the optimal value of 1.0 was chosen as the E value in the CONTROL experiment. We redid all the simulations and modified relative text and figures as required.

## 1207: Actuality: I am not sure we can use this word here, change please

"Actuality" -> "Reality"

1210: If you say so, you need to show that Greenland glaciers and the domain of your study can be similar to each other. Or you need to rephrase your sentence...

We guess you refer to 1209 in the original text? We delete "However" for a subtle shift of emphasis. The current temperature distribution in the Fleming Glacier cannot be accurately calculated or estimated in any way. Steady state is as good a guess as anything else.

1215: you mention Gong2016 (this is 2017 actually) for the spin up scheme or for Elmerice. For the latter, better to cite Gagliardini2013

Thank you for pointing this out. We modified this sentence into "Gong et al (2017) adopted the four-step spin-up scheme (Gladstone et al., 2014) in inverse modelling using Elmer/Ice (Gagliardini et al., 2013), without testing the effect of initial temperature assumption on the inversion results."

1219: There is a step here that is not common, surface relaxation with C at its initial chosen value. What is done usually is the inversion, then the relaxation over about 15 years. I wonder the effect of the surface relaxation using a C that is far from reality...

"For cycle 1, the surface relaxation and first inversion are implemented with an initial temperature assumption (described below) and uniform basal friction coefficient of  $10^{-4}$  MPa m<sup>-1</sup> a (following Gillet-Chaulet et al. (2012))." We clarified this in the text (Line 247-249).

Then we added another two cycles starting with surface relaxation from the initial geometry and simulated C from the previous cycle. Besides, the surface relaxation in each cycle was run for 0.2 yr, which is mentioned in Sect. 3.3. We also added a sentence in Sect. 3.3 (Line 200-202) "This is long enough to remove the non-physical spikes, but too short to significantly modify the geometry of the fast flowing regions of the Fleming Glacier"

## 1220: Basal sliding

As said above, we now use the consistent term "basal friction coefficient" in the whole text.

## 1225: Means you don't account for the modification of surface with relaxation at the beginning of the last two cycles ?

This seems to be a misunderstanding. Relaxation is carried out for each cycle, as stated. We point out that the relaxation of each cycle starts from the initial geometry. For each cycle, the modification of surface after relaxation (<25 m) is smaller than the uncertainty of the ice thickness based on the RMSE of difference between relaxed and observe surface elevations (see Table S1 in the supplementary material), which has been clarified in the Sect. 4.1 (Line 333-337). We feel this is quite clearly set out as it stands. This appears the sensible procedure to minimize the influence of any initial guess for *C* in the first cycle on the relaxation, as raised by the reviewer above.

## 1228: Basal friction

## Modified

1229: To your inverse method, not all of them

"inverse methods" -> "our inverse method"

1243: Don't say linear but rather Control

"linear" here is used to describe the way to generate the initial temperature field. The CONTROL experiment also contains a specific bedrock geometry (bed\_bm). For clarity, we have rewritten as (Line 277): "the linear initial temperature distribution described above."

1246 to 1265: I don't really understand the relevancy of this scenario. To me you should rather study the influence of the position of your ice front, since this is what you are not sure about with your hypothesis assuming ice front = grounding line.

The question is not as simple as ice front position, because division between intact ice shelf and iceberg/sea ice mélange is not clearly defined. Both ice front position and ice front pressure condition are relevant. The scenario here to adjust the external forcing on the calving front considers the uncertainties of ice thickness, bedrock depth, and backstress due to the ice mélange. But following the reviewer's suggestion, we have now added another sensitivity test to different ice front positions in Sect. 3.6 and Sect. 4.4. Note that we do not attempt to define a floating portion of the glacier.

## 1267: Results and discussion

#### Modified

#### 1270: what do you call robustness here ? Replace drag by sliding. Rephrase please.

As we responded above to the comments regarding 156, "robustness" here means selfconsistency. We think it is OK to use "robustness" here. To clarify it, we changed the sentence into "The evaluation criteria are the robustness of simulated basal friction coefficient distribution to experiment design and the mismatch between the simulated and observed surface velocities."

## "drag"->"friction"

#### 1273: There are only 3 TEMP experiments, be more clear

Modified. "the four TEMP experiments " -> "the CONTROL experiment and three TEMP experiments "

1275: Here what we need to have is a metric like the RMS, otherwise this is only a maximum value that is not representative of the rest of the data.

Thanks for the reviewer's suggestion. We have added this in Table S1 in the supplementary material. We calculated the root mean square difference (RMSD) of the difference between the relaxed and observed free surface for the fast flowing regions (>1500 m/yr). The RMSDs in elevation of all the experiments are all < 25 m.

#### 1277: I don't understand what you say here ?

As mentioned in Sect. 3.3, the surface relaxation was used to remove the non-physical spikes in the initial observed surface DEM, caused for example by observational uncertainties of the surface or bedrock data and/or by the resolution discrepancy between mesh and geometry data. However, the surface relaxation cannot avoid systematic coherent changes in the surface near the ice front. To discuss the sensitivity of inverse modeling to this systematic change, we adopted different ice front boundary conditions in Sect. 4.4, which led to different changes in glacier surface during the surface relaxation. We modified this sentence (Line 336-337) "However, the systematic changes generated at the ice front during the surface relaxation may have effect on the inverse modeling, and this is further discussed in Sect. 4.4."

1279: This is quite difficult to evaluate the differences between the different experiments in your maps. I would recommend to the relative differences with a reference experiment.

Thanks for your suggestion. We plotted the relative differences between TEMP1-3 and CONTROL in Fig. S4, but the differences were mainly dominated by the slow-flowing areas. So we computed the RMSDs for *C* (Table S2) and magnitudes of simulated basal velocity (Table S3) between TEMP1-3 and CONTROL for the fast flowing regions (> 1500 m yr<sup>-1</sup>) in each cycle to evaluate the consistency of these experiments. The RMSD of magnitude of observed and simulated surface velocity for each experiment is also computed (Table S5).

1281: Figures should be ordered differently, such vertically Control, Temp1, Temp2, Temp3, this is otherwise very difficult to follow.

We think it is alright to put the figures horizontally as long as we keep the consistency of all figures in this text. Forcing more than three columns into the plots will make them smaller and harder to distinguish features properly. We changed the vertical ordering of different experiments and put CONTROL at the first row for each figure as requested even though it can make the trends in sensitivity more difficult to discern.

1283: Looking at Schaffer2012, it does not seem to me that the dependence of their model to temperature scenario is smaller than yours... You do need to quantify your differences, because this is really not clear.

See comments to 1279.

1286: They showed a non influence onto the modelled surface velocity, not the friction coefficient, or I misread their paper... Their Fig8 shows the differences in terms of basal friction coefficient, but this slightly affect surface velocity as the inverse model tends to minimize the differences.

In Sect. 4.6 of Schäfer et al. (2012), they showed that the temperature scenario did not affect both surface and basal simulated velocities. So they made the conclusion that the obtained basal drag coefficients in their case did not depend strongly on the temperature.

1289 to 1291: Already said, please avoid repetitions. You are in the result and discussion, thus adding other unnecessary stuff is only distracting the reader.

The reviewer seems to have lost track of which parts of the figure we are discussing. Having discussed the differences between results after a first cycle, we are moving to discuss the extent to which an additional cycle (and in due course a third cycle) reduces the dependence on the assumed initial englacial temperature distribution. We think it is necessary and appropriate to explain here why we implement the further cycle. We could understand that the remarks about Vestfonna modeling seem being a little repetitive and we shortened them.

1283: I think this is normal to have different results if you choose a sort of outlier in your initial state, like -20 degrees everywhere for the initial state. I don't think you discussed this as a comparison with the final result? Is -20 in the range of this final result?

Thanks for the suggestion. We agree that in a single cycle it is normal to expect "outliers", that is to say a lack of robustness between the results. So we computed the RMSD of the difference between the simulated temperature and the initial

temperature assumption for each cycle (Table S4). It shows that the experiment TEMP1 (beginning with -20 C) still shows notable differences to other simulations, even after three cycles. "Given this choice of preferred temperature initialization (CONTROL), and the significant difference between this and the cold initialization (TEMP1), we argue that TEMP1 likely deviates furthest from an ideal temperature initialization, and that such a large initial deviation would require more than three cycles to converge on a basal friction coefficient distribution." This sentence has been added in the main text in Line 366-370.

## 1291: Drag - > friction coefficient

Modified.

## 1295: Could you quantify your sticky spots ?

Yes, we have modified this sentence into "However, for experiments CONTROL and TEMP2, the isolated sticky points ~3-5 km upstream of the ice front (with horizontal scale around ~1 km and peak basal friction coefficient of around  $6 \times 10^{-5}$  MPa m<sup>-1</sup> yr) mostly decrease or disappear from the first cycle (Figs. 5a, 5g) to the second cycle (Figs. 5b, 5h)" (Line 350-353).

## 1296: "therefore..." remove this as this was already written

This is actually a new point. Here we try to explain the motivation of running the third cycle. To clarify this, we modify this sentence: "Therefore, a third cycle was implemented to test whether a two-cycle spin-up scheme was enough to reduce the dependence on the initial temperature assumptions." (353-355).

## 1300: You should say Control instead of linear scenario

Here we are not talking about the CONTROL simulations rather the scenario with linear initial temperature.

1306 to 1308: Third time I see this in the paper, remove repetitions please.

We have deleted the earlier occurrence of a similar sentence in response to comment 1289-1291. But this is the appropriate place for the Vestfonna discussion.

## 1306: The low impact is on modelled surface velocity. There is an impact on basal friction coefficient (or basal drag as they say)

They said the low impact on both the modeled surface and basal velocity, and the basal drag coefficients does not strongly depend on the temperature (Sect. 4.6, Sect. 5, and Fig. 13 in Schäfer et al. (2012)). So we are not wrong here.

## 1313: No need to say "inside the yellow contour" in the text

We think it is helpful to guide the reader to a specific aspect of the figure without referring to the figure caption. If it is strongly against the Cryosphere's style we could remove the remark.

## 1318: "shows that internal deformation": you should vertical deformation here.

Modified to "vertical shear deformation" to avoid confusion with strain thinning.

1319: I don't agree with this assertion. Vb=Vs in the fastest flowing areas. In between those you have an area with Vb much lower than Vs, but this matches the places where driving stress is much higher. So this is the driven stress that may drive the vertical deformation, not only the ice internal temperature... You need to rephrase.

This comment does not contradict our statement. The reviewer points out that the high vertical shear rate in our domain is a product of both high driving stress, and deformable (i.e. warm) ice. This is clearly true. We state that the basal state is sensitive to ice temperature – we have made no statement about the relevance or not of driving stress. To make it clearer, we modified the text (Line 390) to emphasize that we are referring to the region of high slope between the upstream and downstream basins, where the driving stress is high. Actually, these are regions of local higher basal shear stress than the surrounding regions, which is more directly relevant to shear deformation near the bed.

## 1330 to 1332: not necessary because already mentioned

Moved to Sect. 4.1.

1334: Remove mentions to colors and rather explain with the physical parameters

See our response to comment 1313. We modified this sentence into "in the fast-flowing region (>1500 m yr<sup>-1</sup>, cyan contour in Fig. 7). The pattern in the region between the 1000 and 1500 m yr<sup>-1</sup> contours" (Line 422-423).

## 1340: I don't understand

We modified this sentence into "However, all three cases feature a low basal friction coefficient in the fast flow region (>1500 m yr<sup>-1</sup> in Fig. 7), which is approximately coincident with the FG downstream basin." (Line 425)

1345 to 347: Why mentioning the MISI in a paper that only deals with inversions, there is no point to me.

We agree that MISI cannot directly explain over- or under- estimation of velocities in an inversion. We deleted this sentence.

1347: basal friction

Modified.

1350: What is behind "it" ? The link with previous sentences is not quite clear. Rephrase please.

We deleted the sentence starting with "it means" and added one sentence before it. "One possible cause of the different basal friction coefficient distributions in these inversions might be the changed surface topography during the surface relaxation, especially near the ice front (Figs. S6)." (Line 431-433)

1355: Ok great, you calculated RMSEs. However, 1) you should have done it before (see previous comment above) and 2) please give us numbers.

Thanks for your suggestion. We have added the RMSD of each experiment in the text and Table S5 in the supplementary material.

1357: I guess this is justified by your RMSE. I think you should discuss more this result, as it suggest that using data taken over a short time range improve the results compared to Bedmap2, which is taken over larger time scales than a year... If I am not wrong.

Thanks for your point. We added few sentences here to clarify the reason why we chose BEDZC (Line 438-453)

1368 to 1371: remove this as already written in the methods section

In the Sec. 3.6, we only discussed the reason for setting different ice front boundary conditions. Here we are talking about the possible reasons for the high friction spots near the front. So we don't think this comment should be removed. This is not an exact copy of the earlier section, and it gives context for the current discussion, given the emergence of the high friction spots in the simulations of the previous sections.

1371: You did not really have investigated the sensitivity to uncertainty to me. You only have tested two datasets, one being more accurate than the other by the way. The Mass conservation based inversion for bedrock is quite an efficient method to infer the bedrock (see Morlighem2014 NG)

Modified "bedrock uncertainty" into "bedrock datasets".

Here we are presenting a sensitivity study - we are not aiming to explore the full range of uncertainty. We have chosen different bedrock datasets that can be justified, and we carry out a sensitivity experiment using these datasets. It is true that this does not quantify the full range of possible outcomes as a response to bedrock uncertainty, but we are not claiming to do that.

It is also not true to say that in general the mass conservation method is "more accurate" than interpolation of direct observations. It may often be preferable, but there are many factors.

1382: This is the kind of things you need to check really. You may have the answer in the paper by Mouginot 2012 in the Journal Remote Sensing. It seems to be a combination between 2007 to 2009 data.

The epoch we quote in the paper (Line 131) was taken from the published information about the various contributions to MEaSUREs velocity datasets we used. The velocity data for the Fleming system is derived from the PALSAR (see the supplementary information in Rignot et al. (2011b)). The PALSAR measurements used in that paper covers coastal sectors north of 77.5° S in "Fall 2007 and 2008". We did check the paper you mentioned, but it did not give us extra information.

We have modified this sentence into "Regarding velocities, Friedl et al. (2018) presented evidence that an acceleration phase occurred between Jan-Apr 2008, but the surface velocity data used in this study was extracted from measurements in Fall 2007 and 2008 (Rignot et al., 2011b)." (Line 468-470).

1387: I really question the relevancy of this experiment. Why doing so as it seems to me that more relevant experiment would be to adjust the ice front position, where you have your uncertainty, and check the sensitivity of inversion results. This latter experiment would not change much the results to me, because over 1.5 km of ice shelf, you don't have much buttressing, but it would be more relevant than what you propose to me.

We have added the experiments of adjusting the ice front position in the Sect. 4.4 to partly address this comment.

## 1421: I don't understand, in what context ?

Here we mean "The lowered surface at the ice front in experiments IFBC1 and CONTROL is apparently the consequence of rapid deformation due to its own weight (longitudinal extension with locally high vertical shear) of an ice cliff, which is over 100 m higher than the control sea level" (Line 526-529).

1429: This is still about this experiment. To test such an amplitude in the influence of

sea levels in inversion results, you need to cite literature about what buttressing could be added from ice mélange (see Krug2014 by the way)...

We guess you mean Krug et al. (2015). This sentence is about presences/absence of the ice front high basal friction being connected with the ice front boundary conditions and not about the local driving stress modification in the relaxation step. We are trying here to address the effect of uncertainty in bed elevation rather than buttressing and mélange. We emphasized that the experiments with different sea levels represent some small uncertainty in the actual sea level, but is also a proxy for pressure variations due to thickness and bed uncertainty and mélange back stress (Line 302-304).

We calculated that ice mélange back force ( $\sim 1.1e7 \text{ N m}^{-1}$ ) used to prevent the rotation of iceberg at the calving front (Krug et al., 2015) could account for the equivalent of up to  $\sim 2.3$  m sea level in terms of ice front boundary condition. We added this sentence in Line 513-514.

Figure 4: add relaxation time here

Added.

Figure 5 caption: Temp4 doesn't exist

Modified.

Figures in general: All the differences that you comment are not always visible. These are to me really tight differences so if you want to argue on this to underline the improvement that are brought by your 4 cycle spin up scheme, you should care more about the figures. Use relative differences between the Control and the other experiments.

Thanks for your suggestions. We hope it is understood that our study concerns the iteration of the original four step spin-up scheme of Gladstone et al (2014). We plotted the relative differences between TEMP1-3 and CONTROL in Fig. S4. We also computed the RMSDs of *C* (Table S2) and of the magnitudes of simulated basal velocity (Table S3) between TEMP1-3 and CONTROL for the fast flowing regions (> 1500 m yr<sup>-1</sup>) in each cycle to evaluate the consistency of these experiments. The RMSDs of magnitudes of observed and simulated surface velocity for each experiment is also computed (Table S5). We modified our analysis about the temperature simulations in Sect. 4.1 (Line 356-374).

Figures in general: please, for the readability order vertically your subplots like: Control, temp1, temp2, temp3

As we comment on l281, we changed the order of different experiments and put CONTROL at the first row for each figure.

Figure 7: Here is certainly a way to remove those zigzags discontinuity, I know Paraview is not user friendly for some stuff, but I don't think this is acceptable for a peer reviewed paper.

This figure has been moved into Fig. S4 in the supplementary material. We do not think the zigzag artefacts interfere with the interpretation of the figure, but can try to improve it if the editor regards it as important

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## **Basal** <u>drag friction</u> of Fleming Glacier, Antarctica, Part A: sensitivity of inversion to temperature and bedrock uncertainty

Chen Zhao<sup>1,3</sup>, Rupert M. Gladstone<sup>2</sup>, Roland C. Warner<sup>3</sup>, Matt A. King<sup>1</sup>, Thomas Zwinger<sup>4</sup>, <u>Mathieu Morlighem<sup>5</sup></u>

<sup>1</sup> <u>School of Technology, Environments and Design</u>School of Land and Food, University of Tasmania, Hobart, Australia

<sup>2</sup> Arctic Centre, University of Lapland, Rovaniemi, Finland

- <sup>3</sup> Antarctic Climate and Ecosystems Cooperative Research Centre, University of Tasmania, Hobart, Australia
- <sup>4</sup> CSC-IT Center for Science Ltd., Espoo, Finland

<sup>5</sup> Department of Earth System Science, University of California Irvine, CA 92697-3100, USA

## Abstract

5

10

- Many glaciers in West Antarctica and the Antarctic Peninsula are now rapidly losing ice mass. Understanding of the dynamics of these fast-flowing glaciers, and their potential future behavior, can be improved through ice sheet modeling studies. Inverse methods are commonly used in ice sheet models to infer the spatial distribution of a basal shear stressfriction coefficient, which has a large effect on the basal velocity and internal ice deformation. Here we use the full-Stokes Elmer/Ice model to simulate the Wordie Ice Shelf-
- 20 deformation. Here we use the full-Stokes Elmer/Ice model to simulate the Wordie Ice Shelf-Fleming Glacier system in the southern Antarctic Peninsula. With an control-inverse method, we model-infer the pattern of the basal dragfriction coefficient from the surface velocities observed in 2008. We propose a threemulti-cycle spin-up scheme to remove-reduce the influence of the assumed initial englacial temperature field on the final inversion. This is
- 25 particularly important for glaciers <u>like</u> with significant temperature-dependent internal deformation. We find that the Fleming Glacier, <u>which</u> <u>has have areas of strong</u>, <u>ly</u> temperature-dependent, deformational flow in the fast-flowing regions. Sensitivity tests using various bed elevation datasets, <u>and</u> ice front <u>positions and</u> boundary conditions demonstrate the importance of high-accuracy ice thickness/bed geometry data and precise location of the
- 30 ice front boundary.

35

## **1** Introduction

In response to rapid changes in both atmosphere and ocean, glaciers in West Antarctica (WA) and the Antarctic Peninsula (AP) have undergone rapid dynamic thinning and <u>increased</u> ice discharge over recent decades, which has led to a significant contribution to global sea level rise (Cook et al., 2016; Gardner et al., 2018; Wouters et al., 2015). Understanding the underlying processes, <u>especially for fast-flowing outlet glaciers</u>, is crucial to improve modeling of ice dynamics and enable reliable predictions of contributions to sea level change<sub>a</sub> especially for fast-flowing outlet glaciers.

- 40 The high velocities of the fast-flowing outlet glaciers are determined byarise from both internal ice deformation and or ice sliding at the bed or both. Internal dDeformation is strongly dependent on gravitational driving stress, englacial temperature, the development of anisotropic structure at the grain scale in polycrystalline ice (e.g. Gagliardini et al. (2009)) and larger scale weakening from fractures (Borstad et al., 2013). Basal sliding is strongly dependent on the gravitational driving stress, bedrock topography and the basal slipperiness,
- 45 which in turn is affected by the roughness of the bed, the presence of deformable till, or the

basal <u>subglacial</u> hydrology. Therefore, one of the keys to modeling fast-flowing glaciers is accurate knowledge of the basal conditions: the bedrock topography and the basal slipperiness (Gillet-Chaulet et al., 2016; Schäfer et al., 2012). Inverse methods are commonly used in ice sheet models to infer the basal <u>shear stressfriction coefficient</u>, and basal velocities, and ice rheology from the glacier topography geometry and observed surface velocities (Gillet-

50 <u>rheology</u> from the glacter topography geometry and observed surface velocities (Gillet-Chaulet et al., 2016; Gladstone et al., 2014; Morlighem et al., 2010).

However, pPoorly constrained quantities, like the basal topography, and the distribution of internal temperature, have provided major challenges for modeling the basal shear stress (Vaughan and Arthern, 2007), especially for small-scale glaciers. However, iIn a study

- 55 <u>studies</u> carried out on a fast-flowing outlet glacier draining from the Vestfonna ice cap in the Arctic (Schäfer et al., 2014; Schäfer et al., 2012), it was found that the <u>Robin</u> inverse methods did not depend strongly on the <u>mesh resolution or</u> uncertainties in the topographic and velocity data. In their case, sliding dominated the flow regime, and tThe impact of ice temperature internal deformation on ice internal deformation ice velocity was relatively small
- 60 compared to the important role of friction heating at the bed on the basal sliding (Schäfer et al., 2014; Schäfer et al., 2012). However, It is unclear whether this property is specific to the Vestforna situation or if it also applies to other fast flowing glaciersno generalization on these findings to Antarctie outlet glaciers has been investigated. The motivation of this paper is twofold: to test the sensitivity of <u>a variational inversion inverse</u> method\_(MacAyeal, 1993;
- 65 Morlighem et al., 2010)<u>s for basal friction</u> to basal geometry and <u>to an assumed initial</u> englacial temperature distribution for a different outlet glacier system, and to determine a robust basal <u>shear stressfriction coefficient</u> pattern for the Fleming Glacier, located in the southern AP, in 2008.

The Wordie Ice Shelf (WIS) (Fig. 1b) in the southern AP has experienced ongoing retreat and collapse since 1966, with its almost-complete disappearance by 2008 (Cook and Vaughan, 2010; Zhao et al., 2017). The Fleming Glacier (FGLFG) (Fig. 1b), as the main tributary glacier that fed the WIS, has a current length of ~80 km and is ~10 km wide near the ice front (Friedl et al., 2018). This glacier \_-has recently shown a rapid increase in surface-lowering rates (doubling near the ice front after 2008) (Zhao et al., 2017), and the largest velocity changes (> 500 m yr<sup>-1</sup> near the ice front) across the whole Antarctic ice sheet over 2008-2015 (Walker and Gardner, 2017).

In this study, we employed the Elmer/Ice code (Gagliardini et al., 2013), a three-dimensional (3D), finite element, full-Stokes ice sheet model, to invert<u>for</u> the basal <u>dragfriction</u> coefficient<u>distributions</u> over the whole WIS-<u>FGLFG</u> system <u>in-using</u> a parallel computing environment.

Here, wWe deduce the distribution of basal shear stress using a control inverse method to assess its sensitivity to bedrock topographies, assumptions about the initial temperature distribution, bedrock topographies, ocean boundary conditions and other constraint parameters in the model. We introduce the data in Sect. 2, present the ice sheet model, spin-up scheme and experiment design in Sect. 3, and discuss the results in Sect. 4 before we give

the conclusions in Sect. 5.

#### 2 Data

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#### 2.1 Surface elevation data in 2008

90 The surface topography in 2008 (Fig. 2a) is combined from two SPOT DEM products acquired on 21<sup>st</sup> Feb, 2007 (resolution: 240 m) and 10<sup>th</sup> Jan, 2008 (resolution: 40 m) (Korona et al., 2009) and an ASTER DEM product ranging from 2000 to 2009 (resolution: 100m) (Cook et al., 2012). The surface elevation data for the Fleming Glacier is mainly from the SPOT DEM product acquired on 10<sup>th</sup> Jan, 2008 (see masks of different DEM products in Fig. S1 in the supplementary material). Here we apply the SPOT DEM precision quality masks on

95 the raw data to extract the DEM data with correlation scores from 20% to 100%. Areas with low correlation scores were filled with the ASTER DEM data. To remove noise from the DEM data, the combined DEM (resolution: 40 m) is resampled to 400 m with a median filter and a window size of 10×10 pixels. The EGM96 geoid (Lemoine et al., 1998) is used to convert from the EGM96 Geoid values to WGS84 ellipsoidal heights. We extract a median 100 value of 15 m for the DEM data over Marguerite Bay (Fig. 1a) as the mean local sea level in the ellipsoid frame.

#### 2.2 Bed elevation data

The bed topography plays an very-important role in the basal sliding and distribution of fastflowing ice (De Rydt et al., 2013). However, high-resolution observations of bedrock 105 elevation for this region the WIS-FG system are still not available. To explore the sensitivity of the basal friction coefficientshear stress distribution to the uncertainty in the bedrock topography, we adopt three basal topographies. The first is from the Bedmap2 dataset (Fretwell et al., 2013) with a resolution of 1 km (hereafter bed bm; Fig. 2b), which is converted from the EIGEN-GL04C geoid to WGS84 ellipsoid heights. The other two are 110 derived using the equations below:

$$bed_zc = S_{2008} - H_{mc}$$

bed  $mc = S_{bm} - H_{mc}$ 

(1)(2)

where S<sub>2008</sub> is the 2008 surface DEM in 2008 mentioned described in Sec. 2.1, and S<sub>bm</sub> is the surface elevation data from Bedmap2 (Fretwell et al., 2013), again relative to the WGS84 115 ellipsoid. S<sub>bm</sub> is downscaled to 500 m with a bilinear interpolation method., and H<sub>me</sub> (where "mc" refers to "mass conservation") is the ice thickness data with a resolution of 450 m combined from covering three sources regions shown in Figs. 2e:-. Hmcdata for the yellow area is computed from the Center for Remote Sensing of Ice Sheets (CReSIS) ice thickness measurements using the Ice Sheet System Model's a mass conservation method (Morlighem 120 et al., 2011; Morlighem et al., 2013), based on ice thickness measurements from ;- the Center for Remote Sensing of Ice Sheets (CReSIS), using ice surface velocities in 2008 from Rignot et al. (2011b), surface accumulation from RACMO 2.3 (van Wessem et al., 2016) and 2002-2008 ice thinning rates from Zhao et al. (2017). The thickness data for the grey area is

- interpolated from Bedmap2 (Fretwell et al., 2013), while the data in the red area thickness 125 ensures a smooth transition between the two regionsis interpolated from CReSIS and Bedmap2 data. The yellow area is-indicates the Fleming Glacier system with ice velocity >100 m yr<sup>-1</sup>. The uncertainty of H<sub>mc</sub> (Fig. 2f) ranges from 10 m to 108 m. For the calculation of H<sub>mc</sub>, we assume the that the ice elevation changes over 2002 to 2008 (Zhao et al., 2017) were small compared to the uncertainty uncertainties in ice thickness (Fig. 2f) and could be 130 ignored in the ice thickness measurements which span a wider time frame. Both bed mc (Fig.
- 2c) and bed zc (Fig. 2d) have a higher resolution of 450 m while bed bm (Fig. 2b) has a resolution of 1 km. The uncertainty of bed bm for the fast-flowing regions of the Fleming Glacier (yellow and red area in Fig. 2e) ranges from 151 m to 322 m\_(Fretwell et al., 2013), while the uncertainty of bed mc and bed zc ranges from 10 m to 108 m (from uncertainties 135

in H<sub>mc</sub>).

The bed topography data (Fig. 2b) indicates the essentially marine character of the Fleming Glacier, showing two basins featuring retrograde slopes, both located underneath the main trunk of the Fleming Glacier's fast flow region. The region basin further upstream (hereafter "FG upstream basin") has a steeper retrograde slope than the one closer to the grounding line

140 of those basins (hereafter "FG downstream basin"). For the FG downstream basin, elevation differences between bed bm and the other two datasets (Figs 2c, 2d) show that bed bm has a generally steeper retrograde slope. The sensitivity of basal shear stressfriction coefficient distributions to the three bed datasets is discussed in Sect. 4.2.

#### 2.3 Surface velocity data in 2008

- The surface velocity data used for 2008 (Fig. 1b) were obtained from MEaSUREs InSAR-based Antarctic ice velocity (from the fall 2007 and/or 2008) produced by Rignot et al. (2011b) (version 1.0) with a resolution of 900 m and with uncertainties ranging from 4 m yr<sup>-1</sup> to 8 m yr<sup>-1</sup> over the study area. For the regions without data (grey area in Fig. 1b), we prescribe the surface speed to be 0. We do not use the finer (450 m) resolution MEaSUREs velocity here since the coarser (900 m) resolution data have been subjected to some post-
- processing, including smoothing and error corrections.

#### 3 Method

All the simulations are carried out using the Elmer/Ice model (Gagliardini et al., 2013). These simulations are used to solve the ice momentum balance equations with an control-inverse method to determine the basal dragfriction coefficients, and the steady state heat equation for to model the ice temperature distribution. The ice rheology is given by Glen's flow relation (Glen, 1955):

 $\tau = 2\eta \dot{\epsilon}$ 

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with where  $\tau$  is the deviatoric stress and  $\dot{\varepsilon}$  is the strain rate tensor. The viscosity  $\eta$  is computed as:

(3)

(4)

$$\eta = \frac{1}{2} (EA)^{-1/n} \dot{\varepsilon_e}^{(1-n)/n}$$

using where *E* is an overall flow enhancement factor, <u>A is a temperature-dependent rate factor</u> calculated using an *E*, and a function of the ice temperature relative to the pressure melting point according to the Arrhenius equation Law (Gagliardini et al., 2013),  $-\dot{\varepsilon}_e = \sqrt{tr(\dot{\varepsilon}^2)/2}$  is the effective strain rate, and *n* is the exponent in Glen's flow law. Table 1 lists the parameters used in this study.

#### 3.1 Mesh generation and refinement

We used GMSH (Geuzaine and Remacle, 2009) to generate <u>an initialthe</u> 2-D horizontal footprint mesh with the boundary defined from the grounding line data in 1996 (Rignot et al., 2011a) and the catchment boundary of the feeding glacier system (Cook et al., 2014), with the assumption that the ice front position in 2008 <u>was</u>-coincide<u>dnt</u> with the grounding line position in 1996 (Rignot et al., 2011a). <u>This assumption is tested as part of the sensitivity tests to various ice front positions.</u>

To reduce the computational cost without reducing the accuracy, we refined the mesh withusing the anisotropic mesh adaptation software YAMS (Frey and Alauzet, 2005) using the local Hessian matrix (second order derivatives) of the surface velocity data in 2008 from Rignot et al. (2011c) as a metric for the mesh density. The resulting mesh is shown in Fig. 3 with and has the minimum and maximum element sizes of approximately 250 m and 4 km, respectively. The 2-D mesh is then vertically extruded using 10 equally spaced, terrain following layers. Sensitivity tests have been done on the Vestfonna ice cap (Schäfer et al., 2014; Schäfer et al., 2012) to prove demonstrate the robustness of inverse simulations to the vertical mesh resolution. In the current study an experiment with 20 extruded layers (not shown) gives very similar results as with 10 layers, confirming those findings also apply to It would be useful to know whether the WIS-FGLFG system shows same robustness to the vertical resolution, but this is beyond the scope of current study. Experiments with various

.85 vertical resolution, but this is beyond the scope of current study. Experiments with various horizontal resolutions (1 km, 500 m, 250 m, and 125 m) show that 250 m are sufficient for the simulations on the WIS-FG system.

#### **3.2 Boundary Conditions**

For transient simulations (surface relaxation, section 3.3), the stress-free upper surface is allowed to evolve freely, with a minimum imposed ice thickness of 10 m over otherwise icefree terrain. For inverse and temperature simulations, the upper surface height and temperature are fixed.

The surface temperature is defined by the yearly averaged surface temperature over 1979-2014 computed from the regional atmospheric climate model RACMO2.3/ANT27 (van Wessem et al., 2014). The geothermal heat flux (GHF) at the bed is obtained from Fox Maule et al. (2005) using input data from the SeaRISE project, and the GHF is interpolated with bilinear interpolation method from the standard 5 km grid onto the anisotropic-mesh. A basal heat flux boundary condition combining GHF and basal friction heating is imposed for temperature simulations.

At the ice front, the normal component of the stress where the ice is below sea level is equal to the hydrostatic water pressure exerted by the ocean. The uncertainties of ice thickness and bedrock topography, the low accuracy of ice front and grounding line locations, and the possible buttressing on the ice front by partly detached icebergs and ice mélange (see Fig. 1c) would affect the calculation of ocean forcing there. Accordingly, wWe will discuss the sensitivity to the ice front boundary condition in Sect. 4.34. On the lateral boundary, which

- falls within glaciated regions, the normal component of the stress vector is set equal to the ice pressure exerted by the neighboring glacier ice while the tangential velocity is assumed to be zero.
- The bedrock is regarded as rigid, impenetrable, and temporarily temporally fixed in all simulations. The present-day solid Earth deformation rate in the Fleming glacier region (Zhao et al., 2017) is negligible compared to the uncertainty of the bedrock data. Assuming that basal melt is negligible under grounded iceSo, the normal basal velocity is assumed setto be zero at the ice/bed interfacehere. The sliding relation relates the basal sliding velocity  $u_b$  to basal shear stress  $\tau_b$ . Here, a simple linear sliding law following Gagliardini et al. (2013); Gillet-Chaulet et al. (2012) is applied on the bottom surface:

$$\tau_b = C u_b$$

(35)

where C is a basal sliding friction coefficient. During the initial surface relaxation, and at the start of the inversion, C is initialized to a constant value of 10<sup>-4</sup> MPa m<sup>-1</sup> yr (following Gillet-Chaulet et al. (2012)), which is replaced with the inverted C in subsequent following steps.
The surface temperature is defined by the yearly averaged surface temperature over 1979-2014 computed from the regional atmospheric climate model RACMO2.3/ANT27 (van Wessem et al., 2014). The geothermal heat flux (GHF) at the bed is obtained from Fox Maule et al. (2005) using input data from SeaRISE project, and the GHF is interpolated with bilinear interpolation method from the standard 5 km grid onto the anisotropic mesh. A basal heat flux boundary condition combining GHF and basal friction heating is imposed for temperature simulations.

#### **3.3 Surface relaxation**

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There may be non-physical spikes in the initial surface geometry, caused for example by observational uncertainties of the surface or bedrock data and/or by the resolution discrepancy between mesh and geometry data. To reduce these features, we relaxed the free surface of this domain during a short transient simulation of 0.2 yr length with a timestep of 0.01 yr. This is long enough to remove the non-physical spikes, but too short to significantly modify the geometry of the fast flowing regions of the Fleming Glacier.

#### 3.4 Inversion for basal shear stress friction coefficient

- 235 After the surface relaxation, we used <u>a variational the control</u>-inverse method (MacAyeal, 1993; Morlighem et al., 2010) implemented in Elmer/Ice (Gagliardini et al., 2013; Gillet-Chaulet et al., 2012) to constrain the basal <u>sliding friction</u> coefficient <u>C</u> in Eq. (35). To avoid non-physical negative values, we used a logarithmic representation of the basal dragfriction coefficient,  $\underline{C} \in = 10^{\beta}$ , where  $\beta$  can take any real value.
- 240 The inverse method is based on adjusting the spatial distribution of the basal dragfriction coefficient to minimize a cost function that represents the mismatch between the magnitudes of the simulated and observed surface velocities.

$$J_0 = \int_{\Gamma_s} \frac{1}{2} (|\boldsymbol{u}| - |\boldsymbol{u}^{obs}|)^2 d\Gamma_{-}$$

(6)

 $\frac{\text{where }\Gamma_s \text{ is the upper surface of the domain, } u \text{ and } u^{obs} \text{ are the simulated and observed}}{\text{surface velocities, respectively. We do not try to fit velocity directions.}}$ 

To avoid over-fitting of the inversion solution to non-physical noise in the observations, a regularization term is added to the cost function as:

$$J_{tot} = J_0 + \lambda J_{reg}$$

(4<u>7</u>)

where  $J_{\theta}$  represents the square of the magnitude of the mismatch between the simulated and observed surface velocities,  $J_{reg}$  is the <u>a</u> regularization term imposing a cost on spatial variations in the control parameter  $\beta$ ,  $\lambda$  is a positive regularization weighting parameter, and  $J_{tot}$  is the total cost (following for example Gillet-Chaulet et al. (2012)). Thus, the minimum of this cost function is no longer the best fit to observation but a compromise between fit to observations and smoothness in  $\beta$ . An L-curve analysis (Hansen, 2001) has been carried out for inversions in the current study to find the optimal  $\lambda$  by plotting the term  $J_{reg}$  as the function of  $J_{0}$  (Fig. S2 in the supplementary material). The optimal value of 10<sup>8</sup> is chosen for  $\lambda$  to minimize  $J_{0}$ .

With  $\lambda = 10^8$ , we compute the total cost  $J_{tot}$  with different values of flow enhancement factor E(0.7, 1, 2.5, 5, 10). It is found that inversions with smaller E gave a better simulated surface velocity for slow ice flow regions while greater E gave a better velocity for fast ice-flow regions. The optimal value of E = 2.5 is chosen for the current study.

#### 3.5 Steady-state temperature simulations

In the absence of a known englacial temperature distribution for the Fleming Glacier system, the steady state ice temperature is solved heat transfer equation is solved using an iterative method as described in Gagliardini et al. (2013) to provide temperatures for use in the inversion process. The ice velocity and geometry are held constant for this part of the simulation. Steady-state temperature simulations for a non-steady-state glacier system will result in the estimations of the temperatures that deviate from actualityreality. However, sSimilar experiments on the Greenland Ice Sheet indicated that the simulated steady-state temperature field could present-provide a reasonable thermal regime for calculation of basal conditions (Seroussi et al., 2013). Heat sources and internal energy transfer determine the temperature distribution within the ice. The heat transfer equation is solved using an iterative method as described in Gagliardini et al. (2013).

#### **3.6 Experiment design**

275 <u>Gong et al (2017) adopted a four-step spin-up scheme (Gladstone et al., 2014) in inverse</u> <u>modelling using Elmer/Ice (Gagliardini et al., 2013), without testing the effect of assumptions</u> <u>about the initial englacial temperature distribution on the inversion resultsThe four-step spinup scheme (Gladstone et al., 2014) has been adopted in inverse modeling using Elmer/Ice (Gong et al., 2017), without testing the effect of initial temperature assumption on the</u>

- inversion results. To explore the sensitivity of inverse modeling to initial temperature assumptions, we proposed a spin-up scheme with more cycles (three cycles in this study as presented in Fig. 4). For each cycle, we followed the spin-up scheme of from Gladstone et al. (2014):
  - 1. surface relaxation;

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- 2. inversion foof ther basal friction coefficient using the relaxed surface geometry;
- 3. a steady state temperature simulation using the simulated <u>velocity velocities</u> from that inversion;
  - 4. another inversion with <u>the previously obtained</u> steady-state temperature.
- The surface relaxation for each cycle starts from the same initial geometry described in Sect. 3.3. For cycle 1, the surface relaxation and first inversion are implemented with an initial temperature assumption (described below) and uniform basal dragfriction coefficient of  $10^{-4}$ <u>MPa m<sup>-1</sup> a (following Gillet-Chaulet et al. (2012))</u>. For cycles 2 and 3, the surface relaxation and inversion are implemented initiated with the simulated steady-state temperature and an initial distribution of basal dragfriction coefficient *C* from the final state of the previous cycle.
- To explore the sensitivity of <u>our</u> inverse methods to <u>assumed initial englacial</u> temperature distribution, <u>enhancement factor (*E*)</u>, basal topography, <u>ice front positions</u>, and the ice front boundary conditions, we <u>carried carry</u> out the experiments summarized in Table 2.

An assumed iInitial englacial temperature distribution is used in the first cycle of the scheme above and would affect the surface relaxation, would affect the modelled ice deformation and the ice velocity field, especially for fast-flowing regions, and consequently lead to a difference in the relaxed upper surfaceaffect the steady-state temperature calculation, which might affect the subsequent inversion process. To explore the impact of initial temperatures on inversion results with the three-cycle spin-up scheme, we proposed experiments with different initial temperature assumptions for the surface relaxation and initial inversion in

- 305 Cycle 1. TEMP1: a uniform temperature of -20 °C; TEMP2: a uniform temperature of -5 °C; CONTROL: a linearly increasing temperature from the upper surface values (see also Sect. 3.2) to the pressure melting temperature at the bed. To test the sensitivity of basal dragfriction to the relaxed geometry, we also added another experiment experiment "TEMP3": surface relaxation in the first cycle using the linear temperature, followed by inversion with a uniform
- 810 temperature of -20 °C. <u>Experiments TEMP1, TEMP2 and TEMP3 differ from CONTROL</u> only in the temperature fields imposed before the first temperature simulation.

Ma et al. (2010) tested the influence of ice anisotropy on the ice flow through various enhancement factors, and found that <u>ideal\_appropriate</u> *E*-values for the grounded ice are usually >1.0. To find out the most appropriate value of  $E_{\underline{(in Eq. 4)}}$  in this study, we evaluate inversion carried out with different values of *E* (EF1: *E* = 0.5, CONTROL: *E* = 1.0, EF2: *E* = 2.0, EF3: *E* = <u>4.0; Table 2</u>). Experiments EF1, EF2 and EF3 differ from CONTROL only in terms of the value used for *E*.

As described in Sec. 2.2, we generated three different bed topography datasets to explore the sensitivity of the inverse modelling. The three-cycle spin-up scheme is carried out for each bed dataset using the linear (described above) initial temperature distribution described above. These experiments are referred to as CONTROL, BEDZC, and BEDMC (Table 2). Experiments BEDZC and BEDMC differ from CONTROL only in terms of the bedrock data set used.

In our standard model domain we assume the 2008 ice front is coincident with the 1996 grounding line, which has an error of several km on fast-moving ice (Rignot et al., 2011a) and might have changed since 1996. The frontal surface elevation is from the SPOT DEM data in Jan 2008, which shows the ice front position is ~1.5 km downstream of the 1996 grounding line position. Since such a narrow residual ice shelf is considered unlikely to have a major influence, we construct the model geometry to have the ice front coincide with the 1996 grounding line for simplicity, i.e. all ice is considered grounded. To test the sensitivity of

inverse modelling to the ice front positions, we implement two further scenarios with different ice front positions: downstream (experiment IFP1) and upstream (experiment IFP2) of the 1996 grounding line position (CONTROL). In IFP1, we assume the ice front position is coincident with the frontal boundary of SPOT DEM data (~1.5 km downstream). In IFP2, we artificially put the ice front position ~1.5 km upstream of the 1996 grounding line position for ~1.5 km. IFP1 and IFP2 differ from CONTROL only in their ice front position.

In addition to the ice front position, there are other sources of uncertainty in the vicinity of the ice front: ice thickness, bedrock depth, and backstress due to the presence of ice mélange. These uncertainties from the ice thickness and bedrock datasets may have an

- 340 significant effect on the pressure boundary condition applied to the ice front, which conventionally balances the normal stress in the ice against the ocean water pressure. In view of the ice thickness uncertainty (ranging from 10 m to 100 m) and hence bedrock depth around the grounding line, and given the possibility of increased additional pressure buttressing force due to floating icebergs and ice mélange as indicated in many previous to the ice the stress of the ice to the stress of the ice to floating icebergs and ice mélange as indicated in many previous.
- 345 <u>studies (e.g.</u> Amundson et al. (2010); Krug et al. (2015); Robel (2017); Todd and Christoffersen (2014); Walter et al. (2017)<u>) and clearly seen in Fig. 1c</u>, we vary <u>the ocean pressure this</u> boundary condition by varying the sea level used to calculate <u>ocean water</u> pressure. This approach directly represents some small uncertainty in the <u>actual exact</u>-sea level-<u>itself</u>, but is also a proxy for pressure <u>errors-variations</u> due to <u>bedrock elevation/ice</u>
- thickness uncertainty and mélange back\_stress. <u>Firstly, in the CONTROL experiment, we assume an ocean pressure at the ice front computed using the observed sea level of 15 m, as mentioned in Sec. 2.1. We adjust the sea level by 10 m from hydraulic equilibrium to test the sensitivity of the inverse modeling to the ice front boundary condition. Firstly, we assume an ocean pressure at the ice front computed using the sea level mentioned in Sec. 2.1. We further simulate two alternative scenarios for the sea level used in the simulations to calculate ocean</u>
- simulate two alternative scenarios for the sea level used in the simulations to calculate ocean pressure: IFBC1 with <u>a</u> sea level of 5 m and IFBC2 with <u>a</u> sea level of 25 m. Another extreme scenario (IFBC3, Table 2) is adopted here by setting the ice front pressure to the ice overburden:

 $P_i(z) = \rho_i g(z_s - z)$   $860 \quad (57)$ 

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where  $P_i(z)$  is the pressure at the ice front as a function of height z,  $\rho_i$  is ice density (Table 1), g is the gravitational constant (Table 1), and  $z_s$  is the height of ice upper surface at the ice front. This is the pressure that would be imposed by a <u>hypothetical undeforming continuation</u> of the advancingneighboring glacier, and imposes zero normal strain rate at the ice front. The ice surface elevation  $z_s$  at the front is ~115 m, approximately 100 m above actual sea level. The total vertically integrated pressure imposed by this condition is equivalent to a sea level of ~60 m, although the vertical distribution of pressure is different todiffers from an ocean pressure condition. Experiments IFBC1, IFBC2 and IFBC3 differ from CONTROL only in their ice front boundary condition.

In our model domain we assume the 2008 grounding line is consistent with the 1996 grounding line, which has an error of several km on fast-moving ice (Rignot et al., 2011a) and might have changed since 1996. The frontal surface elevation is from the SPOT DEM data in Jan 2008, which shows the ice front position is ~1.5 km downstream of the 1996 grounding line position. Since such a narrow residual ice shelf was considered unlikely to have a major influence we constructed the model geometry to have the ice front coincide with the 1996 grounding line for simplicity, i.e. all ice is considered grounded.

#### 4 Results and discussions

The main focus of the current study is the sensitivity of the inversion to the <u>variations of three</u> <u>five</u> factors: temperature initialization, <u>enhancement factor</u>, bed topography, <u>ice front</u> <u>positions</u>, and ice front <u>stress balanceoceanic pressure boundary condition</u>. The evaluation criteria are the robustness of simulated basal <u>dragfriction</u> coefficient distribution <u>to</u> <u>experiment design</u> and the mismatch between the simulated and observed surface velocities.

#### 4.1 Sensitivity to initial temperature

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- We present the results for the inferred basal friction coefficients of from the CONTROL and four three TEMP experiments (Sect. 3.76, Table 2) for the WIS-FGLFG system in Fig. 5. The 2008 ice velocity contours are added as visual references for comparing the basal dragfriction coefficient patterns in the regions of fast flow, since the largest observed ice velocity changes occurred in fast flowing outlet flow-regions (Mouginot et al., 2014; Walker and Gardner, 2017).
- All those experiments showed that In each cycle, the absolute differenceroot-mean-square deviation (RMSD, sometimes also called root-mean-square error) between the relaxed and the observed surface was < 30-25 m (see Table S1 in the supplementary material), smaller than the ice thickness uncertainty (> 50 m) used in this study. However, the systematic changes generated at the ice front during the surface relaxation may have an effect on the inversion, and this is further discussed in Sect. 4.4.

However, we think some of the systematic changes generated by surface relaxation are not correcting real errors in the surface topography data, as discussed later in Sect. 4.3. After the first cycle (left column, Fig. 5), results showed different patterns of basal dragfriction coefficient for each experiment, especially in the fast-flowing regions with surface velocity higher thanexceeding 1000 m\_/yr<sup>-1</sup> (yellow contour in Fig. 5). The basal dragfriction is in the surface topography data is a surface velocity based of the surface topography data is a surface velocity higher thanexceeding 1000 m\_/yr<sup>-1</sup> (yellow contour in Fig. 5). The basal dragfriction is in the surface topography data is a surface velocity based of the surface velocity based of

- <u>coefficients</u> from TEMP2 (Fig. 5d5g) and CONTROL (Fig. 5g5a) share a very similar rib-like patternsticky spots around the ice front, and some isolated sticky spots ~3-5 km upstream of the ice front and another rib close to the yellow contour in the fast flow regions (> 1000 m yr<sup>+</sup>), but the\_TEMP1 (Fig. 5a5d) and TEMP3 (Fig. 5j) display different patterns, indicating
- 405 dependence on the initial temperature assumption. <u>The RMSDs of key properties are computed to evaluate the consistency of these experiments (Table S2-S5). This is in contrast to a similar inverse study on the Vestfonna ice cap (Schäfer et al., 2012), which showed little impact of temperature distribution on the basal sliding coefficient. That was due to a low contribution of ice deformation to ice motion compared to the basal sliding (Schäfer et al., 2012).</u>
- 410 2012). We return to this contrast after considering the effect of the second and third cycles of our spin-up.

To remove reduce the dependence on initial temperature and achieve a consistent equilibrium thermal regime with respect to the given slip-friction coefficient distribution-for surface relaxation, we carried out the second cycle shown in Fig. 4. The basal dragfriction 415 coefficients from the final step of Cycle 2 (the middle column in Fig. 5) shows greater similarity across all the temperature experiments. However, for experiments CONTROL and TEMP2, the isolated sticky points ~3-5 km upstream of the grounding lineice front (with horizontal scale around ~1 km and peak basal friction coefficient of around 6×10<sup>-5</sup> MPa m<sup>-1</sup> 420 yr) in the downstream basin show a trend of mostly decreasing decrease and or disappearing from the first cycle (left column of Figs. 5a, 5g) to the second cycle (middle column of Figs. 5b, 5h) for experiments CONTROL" and TEMP2. Therefore, a third cycle was implemented for all temperature assumptions to test whether a two-cycle spin-up scheme was enough to reduce the dependence on the initial temperature assumptions. After the third cycle, all the 425 scenarios depicted a similar basal dragfriction coefficient pattern (right column in Fig. 5). These differences in basal friction coefficients between the TEMP simulations can also be analyzed through Table S2 and Fig. S4. While these statistics and visualizations confirm the similarity between CONTROL, TEMP2 and TEMP3, it is evident that TEMP1 still shows notable differences to these simulations, even after three cycles (see also Table S3 for basal

430 <u>velocity RMSD</u>). The CONTROL simulation, starting with a linear interpolation of temperature from upper to lower surfaces, seems to be the best option for several reasons: the

choice of temperature value for upper and lower surfaces is physically motivated, which is not true for the other assumptions; it shows the lowest RMSD between simulated and observed upper surface velocity of the temperature sensitivity simulations (Table S5); and it shows the

- **435** least change in the temperature distribution over the three cycles (Table S4). Given this choice of preferred temperature initialization (CONTROL), and the significant difference between this and the cold initialization (TEMP1), we argue that TEMP1 likely deviates furthest from an ideal temperature initialization, and that such a large initial deviation would require more than three cycles to converge on a basal friction coefficient distribution. In other
- 440 words, we postulate that the three cycles are likely sufficient to provide a robust inversion only for initial temperatures moderately close to reality, with the linear interpolation in the vertical providing the optimal initial guess amongst our tests. HenceThe differences between the simulated and observed surface speed for the above experiments (Fig. 6) also prove that the three cycle scheme could provide relatively robust inversion results with little sensitivity
- to the initial temperature. Considering the linear temperature is likely closer to a realistic temperature distribution, we adopted the scenario with initial linear temperature for the experiments described hereafter.

The present study is focused on exploring the effects of uncertainties and their control, and while the dynamics of the FGLFG system will be discussed in more detail in a companion paper (Zhao et al., companion paper). However, a few comments are in order regarding the contrast with the previousan earlier study on the Vestfonna ice cap. The low impact of temperature distribution-profile on the basal sliding-friction coefficient distribution in that study was due to a lower contribution of ice deformational to ice-motion compared to the basal sliding (Schäfer et al., 2012). Internal ice deformation, and hence temperature, may be

- especially important for the WIS-FGLFG system due to steep surface slopes and corresponding high driving stresses in the region between the downstream and upstream basins (<u>~8-12 km upstream of the ice front in Fig. 7aS5a</u>). The patterns of basal dragfriction coefficient (right column of Fig. 5) all indicate substantial differences spatial variation in basal dragfriction over the fast flowing part of the FGLFG. For example, in the region flowing at overfaster than 1000 m yr<sup>-1</sup> (inside the yellow contour), we see very low
- dragfriction over the downstream basin, but higher dragfriction coefficients over the upstream bedrock high, and in a narrow band along the ice front. The nature of the basal shear stress is further complicated by substantial variations in the contribution of basal sliding velocity and of vertical shear deformation to the flow. A comparison between the simulated basal and
- surface velocities (Fig. 7b<u>S5b</u>) shows that <u>internal\_vertical\_deformation\_shear\_dominates</u> the ice dynamics in <u>the some of the fast-flowing regionsregion of high slope between the</u> <u>downstream and upstream basins, where the driving stress is relatively high</u>. This alone would suggest a high sensitivity of modelled sliding velocity and basal <u>dragfriction</u> to the englacial temperature.
- The three<u>multi</u>-cycle iterative spin-up scheme is suggested as an effective set-up for inverse modelling of fast-flowing glaciers that have high surface slopes and vertical shear strain rates and therefore are sensitive to the internal vertical ice temperature fielddistribution. In the present application to the Fleming system, three cycles were sufficient, except in the case of an unphysically cold initialization. In other cases, the inversion process is not so heavily dependent on the temperature field, for example for reproducing the shear margins of the
  - outlet glacier of Basin 3 on Austfonna ice cap, Svalbard (Gladstone et al., 2014).

## **4.2 Sensitivity to enhancement factor**

480 Sensitivity of inverse modelling to the flow enhancement factor has been explored by experiments EF1-3 and the results (after three-cycle procedure) are shown in Fig. 6. The simulated basal friction coefficients (left column in Fig. 6) show different patterns with different *E* values. Recall that from Eq. (4), smaller *E* means higher ice viscosity. The local high friction coefficient sticky spots near the ice front expanded both upstream and along the ice front with increased *E* values, forming a band across the ice front for E = 4.0 (EF3). 485 Conversely, inversions with smaller *E* give a better-simulated surface velocity at the ice front (middle column in Fig. 6), and also lead to smaller differences between the observed and relaxed surface elevation after the surface relaxation (right column in Fig. 6), whereas for EF3 the surface relaxation generates a considerable steepening of the surface slope towards the ice front (Fig. 6l). However, the computed RMSD of the surface velocity mismatch for the fast flowing regions (> 1500 m yr<sup>-1</sup>, middle column in Fig. 6 and Table S5) indicates that the experiment EF1 (*E* = 0.5) (Fig. 6e) shows greater underestimation of surface velocity than CONTROL (Fig. 6f). Therefore, the optimal value of *E* = 1.0 is chosen as the most suitable enhancement factor for the Fleming system.

#### 4.2-3 Sensitivity to bedrock topography

- Figure 8-7\_summarizes results from the three experiments using different bed topographies (Sect. 3.7, Table 2). The 2008 ice velocity contours are added as visual references for comparing the basal drag coefficient patterns in the regions of fast flow, since the largest observed ice velocity changes occurred in fast outlet flow regions (Mouginot et al., 2014; Walker and Gardner, 2017). As shown in Fig. 78, the simulated basal friction friction coefficient C varies <u>slightly</u> with bedrock geometry and its distribution shows greater similarity between BEDZC and BEDMC, compared with CONTROL. CONTROL (using
- similarly between BEDZE and BEDWe, compared with corvertor. Convertor (using <u>"bed\_bm"Bedmap2 bedrock data;</u> Fig. <u>8a7a</u>) shows slightly smaller basal <u>dragfriction</u> coefficients than BEDMC (Fig. <u>8b7b</u>) and BEDZC (Fig. <u>8e7c</u>) in the fast-flowing region (>1500 m yr<sup>-1</sup>, cyan contour in Fig. <u>87</u>). and tThe pattern in the region between the <u>yellow</u> 1000 and cyan-1500 m yr<sup>-1</sup> contours also differs <u>compared to in</u> the CONTROL case, which might be caused by the deeper bedrock of <u>bed\_bmBedmap2</u> in the FG downstream basinthis
   region (Fig. <u>8c7c</u>) compared to the other two detects (Fig. <u>8b7b</u>, <u>8i7i</u>). However, all three
- region (Fig. 8g7g), compared to the other two datasets (Figs. 8h7h, 8i7i). However, all three cases feature aindicate similar regions with low basal dragfriction coefficient in fast flow regions (>1500 m yr<sup>-1</sup>, cyan contour in Fig. 87), which is approximately consistent coincident with the boundary of the FG downstream basin.
- 510 The simulated surface velocities from BEDZC (Fig. <u>8e7e</u>) and BEDMC (Fig. <u>8f7f</u>) match the observed surface velocities better than those from CONTROL (Fig. <u>8d7d</u>) in the regions around the ice front/grounding line and more broadly for velocity exceeding 1000 m yr<sup>-1</sup>. <u>This point is supported by the computed RMSD of surface velocity mismatches (Table S5). One possible cause of the different basal friction coefficient distributions in these inversions might be the changed surface topography during the surface relaxation, especially near the ice front</u>
- (Figs. S6).

The deeper retrograde bed in the CONTROL simulation may indicate increased vulnerability to marine ice sheet instability, and more overestimation of surface velocity is found around the grounding line (Fig. 8d). One possible cause of the different basal shear stress in these

- 520 inversions might be the increased slope caused by the surface relaxation. However, we find the inversion process is not sensitive to the surface relaxation, and this is further discussed in Sect. 4.3. It means a high accuracy bedrock topography data is very important for inverse modeling owing to the fact that the bedrock resolution around the grounding line determines the ice dynamics. Comparisons of the distributions of velocity mismatch and of *C* between
- 525 BEDZC and BEDMC does not provide a <u>clear\_direct\_insight</u> into which is the <u>best\_more</u> <u>accurate\_basal</u> geometry for modelling <u>this-the Fleming</u> system. We <u>compute tThe\_computed</u> <u>root\_mean\_square\_errors\_(RMSED)</u> of the velocity mismatch for the regions with velocity >1500 m yr<sup>-1</sup> (<u>Table S5</u>) is only slightly higher for , and find the RMSE\_of\_BEDMC (62.60 m yr<sup>-1</sup>) is-than for marginally larger than the RMSE of\_BEDZC (61.78 m yr<sup>-1</sup>), and
- both are much lower than CONTROL. Both BEDMC and BEDZC use the 2008 surface DEM and this improvement over the Bedmap2 surface DEM in CONTROL appears significant, even before turning to the matter of ice thickness. While bBoth cases use the ice thickness extracted using the mass conservation mechanismapproach (which is independent of surface geometry) and the bed geometries are accordingly more similar to each other than they are to CONTROL (see Fig. 2 b-d). However, BEDZC maintains better internal consistency with the

<u>2008</u> surface elevation, since it results in the mass conserving ice thickness  $H_{mc}$  being employed, whereas, by the construction of bed\_mc (Eq. 2), the ice thickness in BEDMC is not entirely consistent with mass conservation, although still a more physically motivated interpolation than bed\_bm in CONTROL. The BEDMC and BEDZC ice thicknesses clearly differ by the difference between the Bedmap2 and 2008 DEMs, which should be greatest in areas of greatest lowering, and as we see BEDMC provides a useful sensitivity test case.

- Since bed\_zc is extracted from the accurate and contemporary DEM2008, it should also incorporate into the bed geometry (via H<sub>mc</sub>) more detail from the then current surface, compared to bed\_mc, extracted from Bedmap2's surface DEM, which was generated over a longer time range-data used in current study than BEDMC. Therefore, bed zc is suggested as
  - the best current bedrock elevation data for further ice sheet model<u>l</u>ing of the WIS-<u>FGLFG</u> system.

## 4.3-4 Sensitivity to ice front position and boundary condition

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- All the inversions presented so far feature both a bandsome sticky spots of with high basal 550 dragfriction coefficient near the ice front of the Fleming Glacier (right column of Fig. 5 and left column of Fig. 87) and a similar localized overestimate of upper surface velocities at the ice front (right column of Fig. 6 and middle column of Fig. 8). We now consider causes for possible uncertainties about in the force applied to the ice front, and whether the high basal friction near the ice front is likely to be a feature of the real system or emerges from the 555 inversion process as a compensating response to incorrect boundary forcing by the inversion process. These possible causes include uncertainty in local bedrock elevation (or equivalently ice thickness), uncertainty in observed sea level, uncertainty in exact ice front position and grounding line position, uncertainty in surface velocity, and uncertainty in potential backstress due to ice mélange and/-or grounded icebergs in contact with the ice front. The 560 sensitivity to various bedrock uncertainty datasets has been discussed in Sec. 4.23. In-our model domain we assume the 2008 grounding line is consistent with the 1996 grounding line, which has an error of several km on fast-moving ice (Rignot et al., 2011a) and might have changed since 1996. The frontal surface elevation is from the SPOT DEM data in Jan 2008, which shows the ice front position is ~1.5 km downstream of the 1996 grounding line
- 565 position. Since such a narrow residual ice shelf was considered unlikely to have a major influence we constructed the model geometry to have the ice front coincide with the 1996 grounding line for simplicity, i.e. all ice is considered grounded. In this framework By assuming the ice front position to coincide with the 1996 grounding line, uncertainty about the bedrock depth at the ice front feeds in-to significant uncertainty in the total restraining
- force from ocean pressure. <u>Regarding velocities</u>, Friedl et al. (2018) presented evidence that an acceleration phase occurred on the Fleming Glacier around between MarchJan-April 2008, but we are not sure the specific month of the surface velocity data used in this the current study was extracted from measurements in Fall 2007 and 2008 (Rignot et al., 2011b). <u>ThisH</u> means the surface velocity data, which is provide the target to be matched by the control inversione process, might not be consistent with the DEM data used here (acquired in Jan 2008).

To explore the influence of these different sources of uncertainty, we adopt different <u>ice front</u> <u>positions and effective</u> sea level heights within our vertical reference frame to apply a range of ocean pressures to the ice front as described in Sect. 3.6 (IFBC1-3 and IFP1-2, Table 2).

580 Experiments with different ice front positions (IFP1-2 in Table 2) directly affect the ice thickness and bed elevation at the ice front, which affects the ice front pressure condition. The simulated basal friction coefficients (left column in Fig. 8) show that the high sticky spots near the ice front migrate with the ice front position but with different patterns. The experiment IFP1 with a seaward shifted ice front position shows a decrease in magnitude of the high friction spots (Fig. 8b) and a better match with the observed velocity (Fig. 8e), while the IFP2 with a retreated ice front shows an increased C (Fig. 8c) and worse surface velocity match (Fig. 8f) compared with CONTROL experiment (Figs. 8a, 8d). In experiment IFP1,

thinner ice at the ice front leads to a relatively smaller ice velocity compared with CONTROL, so the model does not need to increase *C* to match the observed surface velocity. This does

590 not mean that ice front position in IFP1 is more accurate than CONTROL, since the time inconsistency of surface DEM data, ice front and grounding line position, and surface velocity data is the obstacle to obtaining a reliable basal friction pattern. Therefore, we speculate that some of the high basal friction spots near the ice front are artefacts. However, we do not exclude the possibility of high basal friction spots caused by the pinning points located at the 1996 grounding line, which is also proposed by Friedl et al. (2018). An accurate location of the ice front and grounding line is clearly important for inverse modelling of fast flowing

glaciers like the Fleming Glacier.

(~2.8×10<sup>11</sup> N).

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A higher sea level in the ice front boundary condition imposes a higher pressure at the ice front, i.e. a higher total retarding force, and we impose these different boundary conditions as a proxy for the sources of uncertainty discussed above.

Basal <u>dragfriction</u> coefficients *C* simulated from the IFBC<u>1-2 and CONTROL</u> experiments (Figs. 9a-c) present <u>different similar</u> patterns <u>but differ systematically</u> around the ice front regions of the FGL (within ~1 km of the grounding line). Experiments with higher sea levels display smaller *C* there (Fig. 9, left column) and provide a better match between modeled and

- 505 simulated surface velocities (Fig. 9, middle column), which is consistent with the computed RMSD of the surface velocity mismatch (Table S5). If the applied ice front boundary condition underestimates the real world forcing, the inversion process will compensate by increasing the basal dragfriction in this region. However, the large vertical shear strain rate imposes a limit to how much increasing basal drag can reduce the surface velocity, which could explain why the mismatch between the modeled and observed velocity is still large in the narrow band near the ice front (Fig. 9, middle column). For the fast flowing region (velocity > 1500 m a<sup>-1</sup>), the decreased basal shear force from IFBC2 to IFBC3 (~1.1×10<sup>+1</sup>-N) roughly matches the increased the ice front pressure over a 6 km length of ice front
- Experiment "IFBC3", with an extreme assumption of applying ice pressure corresponding to a neighbouring column of ice matching the ice front, shows very small basal dragfriction for the ice front area around the grounding line, and also resulted in lower drag over the downstream basin (Fig. 9d). However, "IFBC3" introduces a much greater mismatch to the observed surface velocities (Fig. 9h), with lower simulated underestimated velocities over a substantial region extending upstream from the ice front and greater overestimate of velocities further upstream. This is only a sensitivity test but implies a potentially suitable ice front
- pressure may lie between "IFBC2" and "IFBC3". This set of experiments also suggests that moderate changes influence only a limited area. It is hard to decide the best ice front boundary condition here owing to the lack of precise bedrock data (as seen above) and difficulty of estimating the additional pressure from the partly detaching icebergs and ice
- 625 difficulty of estimating the additional pressure from the partly detaching icebergs and ice mélange. As an indicator, the simulated ice mélange depth-integrated back stress (~1.1×10<sup>7</sup> N m<sup>-1</sup>) required to prevent the iceberg rotation at a calving front (Krug et al., 2015) would be comparable to an additional ~2.3 m in sea level in terms of ice front boundary condition for the Fleming Glacier. The thickness and density of mélange may affect this estimation. But it is cartainly clear that the ice front boundary conditions can have a cignificant effect on the second secon
- 630 is certain<u>ly clear</u> that the ice front boundary conditions can have a significant effect on the inversion results near the grounding line.

The different ice front boundary conditions also lead to <u>minor significant</u> differences in the surface relaxation at the ice front, with lower sea levels leading to <u>slightly</u> greater lowering and corresponding steepening of the surface adjacent to the ice front (Fig. 9, right column in

Fig. 9; for example, ~8 m lowering from IFBC2 to CONTROL and from CONTROL to IFBC1 at the ice front). The differences in surface elevation are localized to the ice front zone, with the relaxation over the rest of the domain essentially unaffected, even except for the most extreme forcing. This The lowered surface at the ice front in experiments IFBC1 and CONTROL is apparently the consequence of rapid deformation due to its own weight spreading(longitudinal extension with locally high vertical shear) of an ice cliff, which is over

100 m higher than the control sea level-at the ice front due to its own weight. Thus, the inversions are potentially influenced both by the ice front condition directly in the overall momentum balance and also by the increased local driving stress due to this artificially steeper surface slope. The band of higher basal drag near the ice front may be partly a

- 645 response to these issues. However, an additional simulation (not shown), in which a high sea level was used for the surface relaxation step and a low sea level for the inversion, gave a relaxed surface very close to observations (no steepening) and still<u>the sticky spot located ~1</u> <u>km upstream the ice front is a persistent feature except for the experiment IFBC3-showed the</u> high basal friction band near the ice front. This implies that the high friction near the ice front
- 650 is directly sensitive to the boundary condition at the ice front but not to associated artifacts in the surface relaxation.

#### Friedl et al. (2018)

Based on the experiments IFP 1-2 and IFBC1-3, we cannot be sure whethersuspect the high friction near the ice front is <u>likely a real feature</u>, an arteifact due to errors in the ice front boundary condition but we cannot rule out the possibility that this may be a real feature, or a combination of both. However, the impact diminishes rapidly with distance inland for moderate sea level shifts, which do not affect the general pattern of basal friction coefficients or the quality of the velocity matching more than ~2 km upstream of the grounding line.

#### **5** Conclusions

- 660 We have obtained a basal dragfriction coefficient distribution for the Wordie Ice Shelf-Fleming Glacier system in 2008, using an iterative spin-up scheme of simulations, observed surface velocities and a detailed surface DEM. We explored the sensitivity of the inversion for basal dragfriction to three-four inputs to the modelling process. Within the approximation of using simulated steady-state ice temperatures, we showed that three-multiple temperature-
- <u>inversion</u> cycles are necessary of iterationto removed the influence of initial englacial temperature assumptions, at least for plausible initial temperature assumptions, and that a poor initial assumption will lead to a requirement for a greater number of cycles. In contrast to the observed low sensitivity to the englacial temperature of outlet glaciers from the Vestfonna Ice Cap (Schäfer et al., 2014; Schäfer et al., 2012), the first cycle of our iterative process showed that the inferred basal stress of the Fleming Glacier system is highly sensitive
- to the englacial temperature distribution. This conclusion is expected to also apply to other fast-flowing glacier systems with a significant dependence on that feature high rates of the internal deformation. For such glacier systems, a multiple-cycle spin-up scheme is likely to be necessary.
- For oOur inversionmodel of the Wordie Ice Shelf-Fleming Glacier system, our is highly sensitiveity tests to different the choice of ice flow enhancement factors and basal elevation datasets indicate a high dependence of basal inversion on the accuracy of bed topography. The "bed\_zc" bed topography, which used ice thickness determined using the mass conservation method for the fast-flowing regions, using contemporary velocities and ice thinning rates, and applied to the then current DEM, is suggested as the best current bed topography for further simulations in this region.

For the Wordie Ice Shelf-Fleming Glacier system, which we treated as grounded adjacent to the ice front, the inferred basal <u>dragfriction</u> coefficient near that <u>grounding lineice front</u> is sensitive to the ice front <u>position and ocean pressure</u> boundary condition, emphasizing the

685 importance of the normal force on the ice front <u>and the accuracy of ice front positions</u>. <u>These factors have a very low impact on basal friction coefficients more than a few kilometers upstream of the grounding line, but may still be important when using inversion to initialize transient simulations, due to the high sensitivity of transient ice dynamic behavior to grounding line dynamicsThis finding, combined with the sensitivity of surface relaxation to ice front boundary condition, implies that an accurate representation of the ice front boundary
</u>

will be important for inverse modeling and transient simulations of the Wordie Ice Shelf-Fleming Glacier system.

#### **Author Contributions**

695 Chen Zhao and Rupert Gladstone designed the experiments together. Chen Zhao collected the datasets, ran the simulations, and drafted the paper. <u>Mathieu Morlighem generated the mass-conservation constrained ice thickness data.</u> All authors contributed to the refinement of the experiments, the interpretation of the results and the final manuscript.

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Figure 1. (a) The location of the Wordie Ice Shelf-Fleming Glacier system in the Antarctica Peninsula (pink polygon). (b) Surface speed in 2008 with a spatial resolution of 900 m 855 obtained from InSAR data (Rignot et al., 2011c) for the study regions. Colored lines represent the ice front position in 1947 (red), 1966 (brown), 1989 (green), Apr 2008 (blue), and Jan 2016 (magenta) obtained from Cook and Vaughan (2010), Wendt et al. (2010), and Zhao et al. (2017). The grey area inside the catchment shows the region without velocity data. (c) Ice front images acquired from ASTER L1T data on Feb 2<sup>nd</sup>, 2009. The dashed line in (b) and (c) is the 1996 grounding line position (Rignot et al., 2011a).

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Figure 2. (a) Surface elevation data in 2008 with black contours (interval: 200 m) representing the surface elevation. (b) bed elevation data from bed\_bm\_(meters above sea level, masl), (c) elevation difference between bed\_mc and bed\_bm (d) elevation difference between bed\_zc and bed\_bm. The black contours in (b-d) show the bed elevation with an interval of 200 m.

870 (e) The ice thickness data sources and (f) the uncertainty of the ice thickness data  $H_{mc}$  with black solid lines representing the observed ice surface velocity of 100 m yr<sup>-1</sup>.



Figure 3. (a) Mesh structure of the domain in the current study with surface velocity in 2008 (Rignot et al., 2011c) and the zoomed-in map for (b) the Fleming Glacier and (c) the Prospect Glacier.



Figure 4. Flow chart of simulation spin-up with three cycles.





Figure 5. Basal dragfriction coefficient C (MPa m<sup>-1</sup> yr) inferred from experiments: (a-c) <u>CONTROL</u>, (ad-f-e) TEMP1, (dg-fi) TEMP2, (g-i) CONTROL, and (j-l) <u>TEMP4TEMP3</u>.
The left (a, d, g, j), middle (b, e, h, k) and right columns (c, f, i, l) are the inferred basal dragfriction coefficients from Cycle 1, Cycle 2 and Cycle 3, respectively. The black, and yellow, and cyan solid lines represent observed surface speed contours of 100 m yr<sup>-1</sup>-and 1000 m yr<sup>-1</sup> and 1500 m yr<sup>-1</sup>, respectively.





Figure 6. Distribution of basal friction coefficient *C* (MPa m<sup>-1</sup> yr) (left column), mismatch between the observed and modeled surface velocity (observed minus simulated; middle column), and the difference between the observed initial surface and relaxed surface elevation (observed minus relaxed; right column) from experiments: (a, e, i) CONTROL, (b, f, j) EF1, (c, g, k) EF2, and (d, h, l) EF3Mismatch between the observed and simulated surface speed in 2008 (observed minus simulated) from experiments: (a-c) TEMP1, (d-f) TEMP2, (g i) CONTROL, and (j-l) TEMP3. The left (a, d, g, j), middle (b, e, h, k) and right columns (c, f, i, l) are the inferred basal drag coefficients from Cycle 1, Cycle 2 and Cycle 3, respectively. The black, and yellow and cyan solid lines represent observed surface speed contours of 100 m yr<sup>-1</sup>, and 1000 m yr<sup>-1</sup>, and 1500 m yr<sup>-1</sup>, respectively.



Figure 7. (a) The slope (degree) of the relaxed surface and (b) the ratio of magnitude of the modeled basal and surface velocity (basal over surface) after three cycle spin-up scheme from experiment: CONTROL. The maximum difference around the ice front is ~2600 m yr<sup>-1</sup>. The zigzag discontinuities in (a) are artefacts of the post-processing at partition boundaries only, and do not affect the simulations. The black and yellow solid lines represent surface speed contours of 100 m yr<sup>-1</sup> and 1000 m yr<sup>-1</sup>, respectively.





Figure 87. Distribution of basal friction coefficient C (MPa m<sup>-1</sup> yr) (left column) and mismatch between the observed and modeled surface velocity (observed minus simulated; middle column) from experiments: (a, d) CONTROL, (b, e) BEDMC, and (c, f) BEDZC with 930 bedrock data (meters above sea level, masl) from (g) bed\_bm; (h) bed\_mc; (i) bed\_zc, respectively. The black, yellow, and cyan solid lines represent observed surface speed contours of 100 m yr<sup>-1</sup>, 1000 m yr<sup>-1</sup> and 1500 m yr<sup>-1</sup>, respectively.



middle column), and the difference between the observed initial surface and relaxed surface elevation (observed minus relaxed; right column) from experiments: (a, d, g) CONTROL, (b, e, h) IFP1, and (c, f, i) IFP2. The black, yellow, and cyan solid lines represent surface speed contours of 100 m yr<sup>-1</sup>, 1000 m yr<sup>-1</sup>, and 1500 m yr<sup>-1</sup>, respectively. Black dotted line is the 1996 grounding line position (Rignot et al., 2011a).

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Figure 9. Left column: Distribution of basal friction coefficient C (MPa m<sup>-1</sup> yr) (left column), the mismatch between the observed and modeled surface velocity (observed minus simulated; middle column), and the difference between the observed initial surface and relaxed surface elevation (observed minus relaxed; right column) inferred from experiments: (a, e, i) CONTROLIFBC1, (b, f, j) IFBC1CONTROL, (c, g, k) IFBC2, and (d, h, l) IFBC3. Middle column: the mismatch between the observed and modeled surface velocity (observed minus simulated) from experiments: (e) IFBC1, (f) CONTROL, (g) IFBC2, and (h) IFBC3. The right column: the difference between the observed initial surface and relaxed surface elevation (observed minus relaxed) from experiments: (i) IFBC1, (j) CONTROL, (k) IFBC2, and (l) IFBC3. The black, yellow, and cyan solid lines represent surface speed contours of 100 m yr<sup>-1</sup>, 1000 m yr<sup>-1</sup>, respectively.

Parameters	Symbol	Values	Units	
Enhancement Factor	E	<del>2.5</del>		
Rheological parameter in	$A_0(T < -10 \ ^{\circ}C)$	3.985×10 <sup>-13</sup>	Pa <sup>-3</sup> s <sup>-1</sup>	
the Arrhenius law	$A_0 (T > -10 \ ^{\circ}C)$	1.916×10 <sup>3</sup>	Pa <sup>-3</sup> s <sup>-1</sup>	
Activation energy in the	$Q_0(T < -10 \text{ °C})$	-60	kJ mol <sup>-1</sup>	
Arrhenius law	$Q_0 (T > -10 \ ^{\circ}C)$	-139	kJ mol <sup>-1</sup>	
Gravitational constant	g	9.8	m s <sup>-2</sup>	
Exponent of Glen flow law	n	3		
Density of ocean water	$ ho_w$	1025	kg m <sup>-3</sup>	
Density of ice	$ ho_i$	900	kg m <sup>-3</sup>	

Table 1. List of parameter values used in this study.

Table 2 Experiment lists. n/a is short for "not applicable". <u>EF and SL are short for</u> <u>"enhancement factor" and "sea level", respectively. IF1 and IF2 represent the ice front</u> <u>positions located downstream and upstream of the 1996 grounding line position (Rignot et al.,</u> 2011a), respectively.

Experiment	<del>Des</del> <del>cript</del> ion <u>E</u> <u>F</u>	Bed topography used	Initial temperature in surface relaxation of Cycle 1	Initial temperature in first inversion of Cycle 1	Sea levelS L-used	Ice front position
CONTROL	Spin -up for thre e cycl es with initi al linea f tem pera ture 1.0	bed_bm	Linear temperature	Linear temperature	15 m	<u>GL1996</u>
TEMP1	Spin -up for thre e cycl es with initi al cons tant tem pera ture of_ 20 °C1.	bed_bm	-20 °C	-20 °C	15 m	<u>GL1996</u>

TEMP2	<pre> <u>U</u> <u>1.0S</u> pin- up for thre e eycel es with initi al cons tant tem pera ture of _5 °C 1.00 </pre>	bed_bm	-5 °C	-5 ℃	15 m	<u>GL1996</u>
TEMP3	1.08 pin- up for thre e cycl es with initi al linea f tem pera ture for surf ace rela xati on but a cons tam pera ture for the for thre e s with initi al linea f tem pera ture for thre e s with initi al linea f tem pera ture for thre e s with initi al linea f tem pera ture for ture f for ture f f f for ture f f f f f f f f f f f f f f f f f f f	bed_bm	-20 °C	Linear temperature	15 m	<u>GL1996</u>
<u>EF1</u>	0.5	bed bm	Linear temperature	Linear temperature	<u>15 m</u>	<u>GL1996</u>
EF2	2.0	bed bm	Linear temperature	Linear temperature	<u>15 m</u>	GL1996
EF3	4.0	bed_bm	Linear temperature	Linear temperature	<u>15 m</u>	GL1996
BEDZC	<u>1.0</u> \$- <del>pin-</del> <del>up</del> <del>for</del>	bed_zc	Linear temperature	Linear temperature	15 m	<u>GL1996</u>

	thre e cycl es with bed_ ze					
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BEDMC	es with bed_ me as the bed topo grap hy	bed_mc	Linear temperature	Linear temperature	15 m	<u>GL1996</u>
<u>IFP1</u>	<u>1.0</u>	bed_bm	Linear temperature	Linear temperature	<u>15 m</u>	<u>IF1</u>
<u>IFP2</u>	<u>1.0</u> <u>pin-</u> <del>up</del> <del>for</del> thre e	<u>bed_bm</u>	<u>Linear temperature</u>	Linear temperature	<u>15 m</u>	<u>IF2</u>
IFBC1	<del>cycl</del> es with low ice front pres sure	bed_bm	Linear temperature	Linear temperature	5 m	<u>GL1996</u>
IFBC2	1.08 pin- up for thre e eyel es with high ice front pres sure	bed_bm	Linear temperature	Linear temperature	25 m	<u>GL1996</u>
IFBC3	<u>1.0</u> <del>S</del> <del>pin-</del> <del>up</del>	bed_bm	Linear temperature	Linear temperature	n/a	<u>GL1996</u>

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