Dear Editor,

We are very grateful to you for accepting the successive deadline extensions. Thank to this, our revised manuscript has been largely modified to address all the reviewers’ concerns. We summarise below the main changes made to the manuscript.

First, Reviewer 2 asked us to focus primarily on Greenland rather than on the differences between the coupling techniques. We therefore started our Results section with a presentation of the results obtained with the two-way coupled experiment (i.e. 2W) which is more physically-based (compared to the NF and PF experiments, i.e. one-way-coupling and parameterised SMB-elevation feedback experiments). The results of the NF and PF simulations are then presented relatively to those of the 2W experiment.

Another important modification concerns the description of the spin-up method. We made our best to clarify the method, but we simply present the basic principles of the procedure. Indeed, since the first submission of the present paper, the description of the spin-up procedure has been published in details in a GMDD paper.

Section 3 is now devoted to the presentation of the three coupling techniques.

The role of ice-sheet dynamics has also been more thoroughly discussed in the Results section and the Discussion has been considerably extended.

Finally, we made our best to clarify all the points raised by the reviewers and to improve the English language.

Best regards,

Sébastien Le clec'h (on behalf of all co-authors)
We would like to thank the reviewer J. Fyke for the evaluation of our study and the constructive comments that helped us to improve the manuscript. Please find below the reviewer’s comments in black font and the author’s response in blue font.

Responses to J. Fyke (Reviewer 1)

Le clec’h et al present a study that assesses the strength of Greenland ice-sheet atmosphere feedbacks over the 21st century using a regional model that is coupled to an ice sheet model. I think this is a novel experiment and valuable study and has the potential to be cited extensively as ice sheets are increasingly incorporated into various climate model architectures. My suggestions for improvement, listed in ‘order of appearance’, are below. My primary general concerns, which I hope the authors can address adequately, involve some apparent inconsistencies in the coupling/spinup (e.g. use of topography anomalies, and uncertainty on how land surface types change in response to ice retreat, and what happens if the ice sheet wants to expand beyond present-day margins). Finally, please feel free to counter my suggestions if you think I’m in error.

Thank you for your constructive comments. We hope that we addressed your concerns in the following.

P1L1: “the projected Greenland sea level rise contribution is mainly controlled by the interactions between the Greenland ice sheet (GrIS) and the atmosphere”: while I tend to agree, relevant models can’t yet fully assess the ocean contribution, so I think this statement is overconfident. Please moderate.

We have considerably modified the text in the abstract and this statement has disappeared in the revised version:

“In the context of global warming, a growing attention is paid to the evolution of the Greenland ice sheet (GrIS) and its contribution to sea-level rise. Atmosphere-GrIS interactions, such as the temperature-elevation and the albedo feedbacks have the potential to modify the surface energy balance and thus to impact the GrIS surface mass balance (SMB). In turn, changes in the geometrical features of the ice sheet may alter both the climate and the ice dynamics governing the ice sheet evolution”.

P1L2: “in particular through the temperature and surface mass balance – elevation feedback”: no, the atmospherically-driven GrIS SLR contribution is controlled by radiative excess/warming. Feedbacks reinforce this effect but is do not control it.

Again, these lines have been modified (please see our previous comment).

P1L2: “fine scale processes”->“fine scale dynamical processes” ?
OK modified.

P1L15: “Furthermore, in 2150, using a fix ice sheet mask, as in the no coupling method, overestimates by 24 % the SLR contribution from SMB compared to the use of the ice sheet mask as simulated in the two-way method” this seems counter to the previous statement that SLR from two-way coupling is 9.3% larger than the uncoupled case. Is the difference due to dynamic
There is no contradiction but we acknowledge that the way this sentence was written was confusing. Actually, this sentence aims at quantifying the overestimation of the SLR projection inferred from changes in SMB only (and not from changes in simulated ice volume) when using a fixed ice sheet component. Therefore they ignore the albedo changes and the SMB-elevation feedbacks. By using such methods, we show that the use of a fixed ice sheet mask leads to an overestimation of the GrIS contribution to SLR of $\sim 6\%$ in 2150, and to an overestimation of $\sim 23\%$ of the SMB (with respect to the use of a time variable ice-sheet mask). These estimations are referred to as $\text{SMB}_{\text{MSK-NF}}$ (fixed ice-sheet mask) and $\text{SMB}_{\text{MSK-2W}}$ (time variable ice-sheet mask) and are both based on the SMB-integrated method, traditionally used in RCM-based studies that have no interactive ice-sheet component. Conversely, when considering the two-way and the one-way coupling experiments, we find that the GrIS contribution to sea-level rise (computed from ice volume changes simulated by GRISLI) is $9.3\%$ higher when GrIS-atmosphere feedbacks are accounted for (i.e. in the two-way coupled method). In the revised version, this has been better presented (see section 4.4) and reformulated in the abstract:

“As a result, the experiment with parameterised SMB-elevation feedback provides a sea-level contribution from GrIS in 2150 only 2.5% lower than the two-way coupled experiment, while the experiment with no feedback is 9.3 % lower. [...] In addition, we quantify that computing the GrIS contribution to sea level rise from SMB changes only over a fixed ice-sheet mask leads to an overestimation of ice loss of at least 6 % compared to the use of a time variable ice-sheet mask”.

P2L 4: “The atmospheric conditions control the variability” -> “Atmospheric conditions control variability and change”

This section has been completely re-written to provide clarifications on surface melting and snowfall drivers before dealing with atmosphere-GrIS feedbacks:

“The evolution of the Greenland ice sheet (GrIS) is governed by variations of ice dynamics and surface mass balance (SMB), the latter being defined as the difference between snow accumulation, further transformed into ice, and ablation processes (i.e. surface melting and sublimation). While surface melting strongly depends on the surface energy balance, snowfall is primarily controlled by atmospheric conditions (wind, humidity content, cloudiness...). However, various feedbacks between the atmosphere and the GrIS may lead to SMB variations that can therefore directly affect the GrIS total mass by impacting its surface characteristics, such as ice extent and thickness, with potential consequences on ice dynamics (e.g., due to change in surface slopes).”

P2L7: “SMB directly affect the GrIS total ice mass by impacting its characteristics such as thickness, ice volume and ice extent” - this can occur both directly and via impacts on ice dynamics. Explicitly state the latter (dynamics) for clarity.

Again this part has been drastically reformulated with clarity in mind (see previous comment).

P2L9: there are more foundational references regarding the dynamical GrIS impact on atmospheric flow. Suggest to use these in addition/instead. As just one arbitrary example: [link](http://onlinelibrary.wiley.com/doi/10.1034/j.1600-0870.1996.00014.x/abstract)

Thank you for the reference. We have added the following:

“These changes may in turn alter both local and global climate. As an example, changes in near-
surface temperature and surface energy balance may occur in response to changes in orography (temperature-elevation feedback) or in ice-covered area (albedo feedback; see Vizcaino et al., 2008, 2015; Lunt et al. 2004). On the other hand, topography changes may alter the atmospheric circulation patterns (Doyle and Shapiro, 1999, Petersen et al. 2003, Moore and Renfrew, 2005) causing changes in heat and humidity transports.”

P2L11: “different processes and feedbacks”→“different processes and feedbacks that regulate transient ice sheet change”
Thanks for the suggestion, we have modified the text accordingly.

P2L16: “The climate models usually represent” -> “For example, CMIP5 climate models unanimously represented”
We modified as: “For example, the CMIP5 climate models unanimously represent the ice sheet component with a fixed and constant topography, even under a warm transient climate forcing”.

P2L24: Suggest citing recent Lofverstrom et al. discussion study on resolution dependence of ice sheet conditions in GCMs: https://www.the-cryosphere-discuss.net/tc-2017-235/
The suggested reference has been added as well as the following paragraph:
“Using the AGCM NCAR-CAM3 run at different spatial resolutions (T21 to T85) and coupled to the SICOPOLIS ice-sheet model, Löfverström and Liakka (2017) investigated how the atmospheric model resolution influences the simulated ice sheets at the Last Glacial Maximum. They found that the North American and the Eurasian ice sheets were properly reproduced with the only T85 run. According to the authors, this is likely due to the inability of the atmospheric model to properly capture the temperature and precipitation fields (used to compute the SMB) at lower horizontal resolutions, as a consequence of the poorly resolved planetary waves and smooth topography”.

P2L35: “the authors only consider a strict linear relationship between topography and SMB changes” - please note more clearly either here or in next paragraph why this a handicap to these methods, leading to why your approach is better
We have added the following:
“However, in both parameterisations by Franco et al. (2012) and Edwards et al. (2014b), the authors only consider a strict linear relationship between topography and SMB changes. Although changes in temperature can be derived from a linear vertical lapse rate, other processes governing the SMB such as those related to energy balance, precipitation or atmospheric circulation do not follow a linear relationship with the altitude. While this approach may be valid at the local scale for small elevation changes, it may lead to a misrepresentation of the SMB-elevation feedbacks for substantial changes in altitude, especially at the ice-sheet margins.”

P2L9: “The second fundamental requirement is to represent the ice sheet topography changes in the atmospheric model by using an ISM instead of the fixed geometry usually used” This sentence is tautological since by definition a fixed geometry will not capture topography changes. Reword sentence.
The sentence has been reworded as: “The second fundamental requirement to describe the interactions between atmosphere and GrIS is to represent the ice sheet topography changes in the
atmospheric model by using an ISM (instead of the fixed geometry typically used) to take into account the effects of ice dynamics on the ice sheet topography changes”.

Throughout text: “developed” -> “developed”
OK, modified.

P4L7: 16 km high, from surface? Sea level?
This part of the text has been changed in “The MAR horizontal resolution is 25 km x 25 km covering the Greenland region (6600 grid points), from 60 °W to 20 °W and from 58 °N to 81 °N, and 24 vertical levels to describe the atmospheric column in sigma-pressure coordinates (Gallée and Schayes, 1994)”.

P4L12: “hydrological cycle” -> “atmospheric hydrological cycle”?
Yes, modified.

How does Crocus differ/integrate with SISVAT? Please clarify. In the case where the ice sheet expands or contracts, how is under-snow (or snow free) ice sheet surface exchanged for bare land surface (or vice versa)?
Crocus is a 1D snow model, while SISVAT is the surface model embedded in MAR. In SISVAT, each grid cell is assumed to be covered by at least 0.001% of two major surface types, namely tundra and snow (including ice sheet). Tundra is considered by SISVAT as a vegetation zone with an albedo ranging from 0.1 to 0.2 as a function of surface water and plant type. On the contrary, the Crocus snow model is used to compute the albedo of ice covered areas. In the 2W method, the percentage of tundra/snow evolves following the ice-sheet model advance and retreat. We now provide more information about the MAR model in Sec. 2.1 and we hope that the interplay between Crocus and SISVAT appears now clearer:

“MAR is a regional atmospheric model fully coupled with the land surface model SISVAT (Soil Ice Snow Vegetation Atmosphere Transfer model, see Gallée and Duynkerke, 1997) which includes the detailed one-dimensional snow model Crocus (Brun et al., 1992) which simulates fluxes of mass and energy between snow layers and reproduces snow grain properties and their effect on surface albedo [...]. Each grid cell is assumed to be covered by at least 0.001 % of tundra and snow. At each time step SISVAT computes the albedo of each surface type and the characteristics of the snowpack which are weighted and averaged as a function of the snow and vegetation coverage in each grid point, and then exchanged with MAR.”
In addition, we have included more details on the 2W coupling methodology (in Sec. 3.3):

“At the end of a MAR model year, MAR is paused and GRISLI is forced by the downscaled SMB and ST fields with the method of Franco et al. (2012) as in PF (Eq. 7). Then, GRISLI computes a new GrIS topography and a new ice extent at 5 km which are aggregated at the yearly time scale onto the 25 km MAR grid. The aggregated ice extent is used to update the fraction of tundra relative to ice/snow covered surface type for the subsequent MAR run. To account for the differences between MAR and GRISLI topographies, the surface elevation which is aggregated onto MAR is computed from GRISLI surface elevation anomalies added to the present-day observed topography (Eq. 7). It is then used as the updated surface elevation in MAR. As previously mentioned, topography changes are negligible before 2020. Hence, changes in ice-sheet geometry are fed to MAR only
after this date. Compared to the NF and PF approaches, this two-way coupled method is the most accurate to represent the GrIS-atmosphere feedbacks”.

P4L20: “The topography of the GrIS as well as the surface types (ocean, tundra and permanent ice) are provided by Bamber et al. (2013)” -> clarify this is for the NC is experiment (presumably)

We made this clarification in the text:
“Except for the experiment presented later in this study in which MAR is coupled to an ice-sheet model, the topography of the GrIS as well as the surface types (ocean, tundra and permanent ice) are taken from the Bamber et al. (2013) dataset aggregated on the 25 km grid.”

P5L10: “we have repeated the MIROC5 year 2095 (representative of the years 2090s) for 50 additional years” - this repetition is certainly not representative of this time period due to lack of continued change, and also lack of internal variability. While I don’t think this is a fatal flaw of the study, the authors should clearly note this caveat here and later during discussion of results, so readers clearly realize the effects of this artificial ‘extension’ (probably, fairly strongly reduced overall change, making the results presented here conservative).

We acknowledge the fact that, in our approach, we discard the role of interannual variability within the GCM after the year 2100. This could indeed result in conservative estimates due to non-linearities of SMB (in particular ablation). However, the GCM imprint of the year 2095 may also increase regional changes in term of GrIS response. We present these limitations in the revised version of the manuscript in the discussion section:

In section 3, we mentioned that the use of a constant forcing from 2100 to 2150 “implies that both climate changes and large-scale inter-annual variability are neglected beyond 2100”.

In the Discussion section: “A second question concerns the impact of a constant MIROC5 climate used to force MAR beyond 2100. As outlined in section 3, this results in discarding the continued change that the climate will likely undergo beyond 2100 suggesting that our SLR projections are underestimated. The second consequence is that inter-annual variability is neglected after 2100. This can lead to conservative estimates of Greenland melting contribution to sea level rise in the future due to non-linearities of the SMB. On the other hand, the imprint of the 2095 MIROC5 climate may amplify regional changes of the GrIS response”

P6L11: Why is the annual mean bottom snowpack temperature not used as the boundary condition for the ISM instead?

It would be indeed possible, but probably would have very low impact on the ice temperature profile simulated by GRISLI. Indeed, the annual temperature at the bottom of the snowpack is very similar to the annual mean surface temperature. Because GRISLI has a yearly time step, it can not see annual temperature variability simulated by Crocus. Therefore, we have used the annual mean temperature as a boundary condition for the ISM.

P6L19: also just due to the long timescale of ice sheet responses?

Yes, we agree with this comment. It seems more appropriate to only deal with the long time-scale response of the ice sheet. Our motivation has been reformulated as: “Due to the long time scale response of the ice sheet to a given climate forcing, a proper initialisation of the model is required before performing forward experiments”.

Moreover, our spin-up procedure also includes the calibration of unknown parameters (basal drag coefficient) and our inversion procedure can be seen as a way to correct model deficiencies. That being said, we have clarified and simplified the presentation of the spin-up procedure (see Section 2.2.2) as we now directly refer to the paper published in the discussion forum of the Geoscientific Model Development journal. In the revised manuscript, this paper refers to as: Le clec’h et al. (2018).

P7L1: what is meant by ‘vertical fields’? Please clarify.
We meant temperature and ice velocity profiles from Gillet-Chaulet et al., 2012. It is now specified in the revised manuscript.

Spin-up procedure: How does this procedure deal with ice growth outside the observed ice sheet extent? Figure 2 suggests this ice is simply removed? If so, how does this effective strong artificial sink of ice impact all subsequent sensitivity experiments? Please explain the impacts of this clearly in the text, if this is the case.
You are right, in our framework we apply an artificial strong negative SMB outside the observed present-day ice sheet mask. We do not think that it is a major flaw in our methodology as our spin-up procedure aims at reducing the mismatch between observed and simulated ice thickness. Assuming that MAR produces a realistic SMB on the ice sheet and because the simulated ice thickness is close to observations, we can hypothesise that our simulated ice flow is realistic. As such, in theory, the ice sheet should not grow outside the observed present-day ice sheet imprint. The artificial strong negative SMB outside the present-day ice sheet mask can be seen as a way to correct both the atmospheric model bias (e.g. positive / not enough negative SMB over the tundra) and the spin-up procedure bias (too strong ice export towards the margin). This has been explained at the end of the spin-up description (section 2.2.2).

P8L21: why not simply start the coupling at 2005 (i.e. the end point of the 1976-2005 initialization/spin-up period)?
Sure it would have been possible. However, as stated in the manuscript, the results would have been similar as the SMB changes through 2005-2020 does not produce any significant topography changes in GRISLI.

P9L13: The use of topography anomalies is concerning since it implies the SMB/ST field received by GRISLI is inconsistent with GRISLI’s height (for example, the ELA on the GRISLI grid would exist at a different elevation than if the GRISLI elevation was directly used). Can the authors comment on why this approach does not introduce problems with their experimental design? As it stands, this is not justified adequately. An alternate approach that would have avoided this problem would have been to use the spun-up GRISLI topography as the ‘fixed’ topography instead of the Bamber topography.
The issue here is that GRISLI tends to produce steeper slopes than what is observed. This has important consequences for the climate simulated by MAR due to, in particular to katabatic winds. This is why we made the choice to maintain the realism of the simulated present-day climate (computed on the Bamber et al. (2013) topography) and the consistency between the climate simulated by MAR and the climate used to force GRISLI, downscaled at the 5 km resolution using the method developed by Franco et al. (2012).
In Section 3.2, we mentioned that “Due to the topography differences between MAR and GRISLI, this approach has been chosen to avoid large inconsistencies between the SMB and ST fields computed by MAR and the ones corrected to account for the GRISLI topography”.

We also discussed the impact of the anomaly method in Section 5:

“A second limitation is related to the 2000-yr relaxation GRISLI experiment, run at the end of the spin-up procedure to reduce the model drift in terms of ice volume, that produces residual differences with the observed topography (Bamber et al. 2013) used in the MAR simulations. This has important consequences on the MAR simulated climate. In particular, the steeper slopes existing in the GRISLI topography (i.e. $S_{ctrl}$) tend to produce unrealistic katabatic winds. Therefore, we choose to use an anomaly method of the surface elevation onto which the SMB and ST fields are downscaled at the 5 km resolution grid (Eq. 7). The objective of this approach was first to maintain the realism of the simulated present-day climate computed on the observed topography (Bamber et al. 2013) and, secondly, to avoid inconsistencies between the climate simulated by MAR and that used to force GRISLI. However, this implies that the forcing climate is not fully consistent with the GRISLI topography. This should be taken into consideration in a future work to improve the quality of our results”.

Figure 2 and other figures: 5 years is likely not long enough to generate robust climatologies. Suggest using at least 10 years instead.

We have followed your suggestions and used 10 years to compute climatologies.

P11L10: the finding of very strong marginal cooling due to increased katabatics is very interesting and pertinent, and deserves a further explaining. It would be very useful the authors plotted overlaid near-surface wind anomaly vectors plus ST changes in ‘zoomed-in’ plot of a good illustrative portion of the margin.

We provided further explanations to justify the role of katabatic winds in the marginal cooling (see section 4.2.1):

“Over the ice sheet, the steeper surface slopes simulated in 2W in 2150 (discussed in Sec. 4.1.2) lead to a slight increase in katabatic winds (Fig. 9). However, at the ice sheet margin, i.e. where the ice mask in MAR is below 100%, there is a substantial decrease in surface winds. This is because the change in surface elevation as seen by the atmospheric model is computed from the aggregated changes in GRISLI at 5 km. As such, a non-zero fraction of tundra, which presents no change in surface elevation, results in smaller elevation changes compared to grid cell in the same region with permanent ice cover only. This induces artificially lower surface slopes at the margin with respect to the interior and a decrease in surface winds in these regions. Altogether, the slight increase in katabatic winds over the ice sheet and their reduction at the margin lead to a cold air convergence towards the ice sheet edge (Figs. 8b and 9 and Fig. S8-S9)”.

To support these explanations, we added a new figure (Figure 9) displaying the 2W near-surface wind vectors at the end of the 2W experiment as well as the wind strength anomaly between 2W for NF. A zoom-in plot showing near-surface wind anomaly vectors overlaid to ST changes is provided in the Supplementary Materials as Figure S7.

Similar to above point: it would be excellent to see a quiver plot of wind anomalies over the entire
ice sheet, given their importance. Also would it be possible to visualize the increased mixing in the boundary layer, leading to warming in the 2-W coupled case?

A similar plot as Figure 9 in the main text is also given in the Supplementary Materials (see Fig. S9).

**P11L23: do authors mean “Following the increase of the ST”?**

Yes, this is what we meant. However, due to modifications in the structure of the revised manuscript, this part of the text has been removed.

**P11L25: “, there is a decrease of 112 Gt yr⁻¹ 25 of ice ” -> “112 Gt/yr extra ice ablates”**

The sentence has been changed in: “This process is faster in 2W than in NF and PF. In 2150, the ablation zone is 14 % (resp. 11.7 %) larger in 2W than in NF (resp. PF) causing 112 Gt yr⁻¹ of extra ice ablation in 2W (w.r.t NF)”.

**P11L30: “14 % larger in 2-W” - can an estimate be made of the uncertainty in this value (and others) due to interannual variability? Put another way, can the authors confirm that the changes they see are significant in the face of background noise in ablation area (for example)?**

As specified in sections 2 and 5, the use of a constant climate forcing for MAR after 2100 (here the MIROC5 climate simulated for year 2095) implies that the inter-annual variability is neglected beyond 2100. As such, the relative changes in ablation areas after 2100 mentioned in the text are necessarily statistically significant, at least within the framework of our experimental setup. However, we acknowledge that a better approach would be to perform similar simulations with a prolonged RCP8.5 scenario (not available at the time of this study).

**P12L5: “lower surface temperature over these regions” - suggest reinforcing to readers once more here that this is *relative* to the NC experiment.**

Thanks for this remark. We paid attention to clarify the text when dealing with relative changes.

**P12L8/9: what does the +/- indicate here?**

This is the mean value +/- the root mean square error over the region. In the revised manuscript, the mean values do not longer appear with the +/- root mean square error. We chose to present the mean value results with the 5th and 95th percentiles when necessary.

**P12L13: “become ice or snow-free or snow free, exhibiting bare ice ” this is confusing. What happens if the entire GRISLI ice column disappears? Does tundra emerge?**

This part of the text has been modified and changes in ice-sheet extent are now only discussed in Section 4.4.

The point which which is addressed here has been clarified (see section 2.1):

“[In MAR], each grid cell is assumed to be covered by at least 0.001% of tundra and snow. At each time step SISVAT computes the albedo of each surface type and the characteristics of the snowpack which are weighted and averaged as a function of the snow and vegetation coverage in each grid point, and then exchanged with MAR”.


In both the NF and the PF experiments, the ice-sheet mask, as seen by GRISLI is not updated in the atmospheric model and MAR sees the present-day observed ice-sheet mask throughout the simulation. In the 2W experiment, the ice extent computed by GRISLI is then aggregated to MAR to update the fraction of tundra relative to ice/snow covered surface type for the subsequent MAR run. As a result, if the entire ice column disappears, MAR sees in each grid cell a fraction of tundra of 99.999% and modifies the albedo accordingly.

P12L25: Previous studies have highlighted a strong decrease in ice discharge across outlet glacier grounding lines as a consequence of increased surface melting. E.g. Gillet-Chaulet 2012, Goelzer 2013 and others. Is this same effect seen here?

At the end of the 2W experiment (2140-2150), there is a decrease of surface velocities compared to the 2000-2010 mean period (Figs. 6a, 7c), suggesting that ice discharge across outlet glaciers is reduced. Moreover, the negative anomaly of ice flux divergence (Fig. 5b) shows an upstream ice accumulation (i.e. ice accumulates faster than it discharges through outlet glaciers). These results strongly suggest a decrease of ice discharge across outlet glaciers, similarly to what was found by Gillet-Chaulet et al. (2012) and Goelzer et al. (2013).

P12L25: Is it completely correct to say the entire SLR contribution is caused by the ‘melting contribution’?

In our model, there is only a very few number of grid points in contact with the ocean. Therefore, calving is negligible and melting remains the dominant contribution to sea-level rise. However, to avoid confusions, we removed all expressions such as the Greenland melting contribution to SLR in the revised manuscript and simply use the Greenland ice sheet contribution instead.

P12L25: Can the authors quantify the reduction in marine margin extent in 2-W?

As explained in our previous response, the number of grid points in contact with ocean is negligible in our model. This is likely due to the too coarse GRISLI resolution (5 km) that prevents from properly resolving the complex topographic features of marine terminating glaciers. As a result, it is not possible to quantify accurately the marine margin extent. To illustrate the limitations induced by the coarse ice-sheet model resolution, we added the following paragraph in the Discussion section:

“Regarding the ice-sheet model, a 5 km horizontal resolution does not permit to capture the complex ice flow patterns of smallest outlet glaciers, whose characteristic length scale can be less than 1 km (Aschwanden et al., 2016) and to quantify accurately the ice discharge at the marine front. This may have large implications in the sea-level rise estimates. Using a 3D ice-sheet model with prescribed outlet glacier retreat, Goelzer et al. (2013) found an additional SLR contribution from outlet glaciers of 0.8 to 1.8 cm in 2100 and 1.3 to 3.8 cm in 2200, with the influence of their dynamics on SLR projections decreasing with time and with the increasing importance of the atmospheric forcing. This is in line with the fact that ice dynamics act to counteract ice loss from surface melting (see Section 4.2), as previously outlined by several authors (Edwards et al., 2014b, Goelzer et al., 2013, Huybrechts and de Wolde, 1999). However, despite the possible decreasing influence of marine terminating glaciers, at the centennial time scale, it seems to be preferable to evaluate more accurately the impact of ice dynamics and to better capture the complex geometry of fjords surrounding the marine-terminating glaciers”.
“This higher integrated SMB, obtained when using no updated ice sheet mask” - do the authors mean “lower”.? This sentence seems to directly contradict the previous sentence. If I’m mistaken here, a clearer description of the processes here is needed.

Yes, you’re right. In the revised manuscript (Section 4.4), we tried to better explain the issues related to the integrated-SMB method. We hope that the text has been clarified enough:

“A widely used method to estimate the projected GrIS to global sea-level rise is to compute the GrIS mass loss as the time-integral of the SMB computed by an atmospheric model over a fixed ice-sheet mask (Fettweis et al., 2013, Meyssignac et al., 2017, Church et al., 2013). In the present study, we go a step further since the ice mass variations related to SMB changes are computed over a changing ice-sheet mask as simulated by GRISLI. However, in both the NF and the PF experiments, the atmospheric model does not account for the variations in the ice-sheet extent simulated in GRISLI and the ice-sheet mask, taken from the observations (Bamber et al., 2013) is kept constant throughout the simulation. Taking the changes in ice-sheet mask into account may have strong impacts on the computed GrIS contribution to sea-level rise. To illustrate the influence of the ice sheet mask, we used the SMB outputs from the NF experiment at the MAR resolution and applied the integrated SMB method over the fixed observed ice-sheet mask (SMB_{MSK-NF}) and over the updated 2W mask (SMB_{MSK-2W}). Results reported in Table 2 indicate differences in SMB values exceeding 23 % in 2150. In the same way, compared to a time variable ice-sheet mask, the use of a fixed ice-sheet mask overestimates the sea-level rise by “6 % in 2150. Though a bit lower, this number is far from being negligible compared to the errors made when the SMB-elevation feedbacks are not taken into account (i.e. 7.6 %) and when all the feedbacks are ignored (i.e. 9.3 %). This strongly suggests that realistic SLR projections cannot neglect the evolution of the ice-sheet extent, only accounted for through the use of an ice-sheet model”.

General: The authors should consider quantifying actual feedback factors associated with the inclusion of elevation feedbacks (see Roe 2009, Reviews of Geophysics). This would be a good benchmark number to produce, for other works to compare to.

We agree that a formalised way to quantify the elevation feedback would be very interesting, in particular for inter-comparison exercises. However, the definition of such a metric has yet to be done. For now, we only compare our SLR projections with and without the elevation feedback to other papers available in the literature a similar approach has been followed (e.g. Vizcaino et al., 2015; Calov et al., 2018).

“As for the ISM, increasing the grid resolution of MAR” - do you mean “as for the regional climate model”.?

No, we think that an increase in both ISM and RCM resolutions could better constrain the SLR contribution from Greenland ice sheet. These aspects have been detailed in the Discussion section (in the revised manuscript).

“underestimated by simulating.” Unclear.

“By simulating” should be removed. This error was probably due to an improper “copy-paste”

“surface albedo and strength of katabatic winds.” -> “surface albedo and strength of katabatic winds, with a demonstrably strong return influence on SMB”

The Discussion section has been entirely re-written and this sentence has been removed from the
P15L27: “optimal resolution of the ice sheet and the atmospheric model, for ISM-RCM coupling.” While an interesting-sounding statement, I find it also a bit vague: by optimal, do the authors mean something like “of high enough respective resolutions to resolve both important atmospheric and important ice sheet dynamical processes”?

The point here is to find “high enough ISM and RCM resolutions to resolve both important atmospheric and important ice sheet dynamical processes”, while keeping a reasonable computational time. In the revised manuscript, the sentence has been modified by: “However, a compromise must be reached between the additional computing resources and the required degree of accuracy of sea-level projections”.

P15L30: “The next step of this study. . .” as described, this is extremely ambitious, with many challenges that outstrip the effort to implement atmospheric coupling. If it is truly a planned next step; great! But if not, I’d suggest not claiming to plan to do this.

In the revised manuscript (Section 5), we rather gave a few examples to illustrate the importance of having a description of the ocean-atmosphere-GriS coupled system describing the coupled ocean (see paragraph below), but we followed your recommendation and avoided expressions such as “the next step of this study”:

“There is a growing number of evidence for attributing the acceleration of outlet glaciers to the intrusion of warm waters from adjacent oceans in the fjord systems or in the cavity of floating ice tongues (e.g. Straneo et al., 2012; Johnson et al., 2011, Rignot et al., 2015) that can destabilise the glacier front and/or favour the ice-shelf breakup, decreasing thereby the buttressing effect and increasing the ice calving. In turn, the released freshwater flux in ocean may impact sea-surface temperatures, oceanic circulation and sea-ice cover. Moreover, atmosphere-ocean feedbacks also have an impact on the GrIS. As an example, Fettweis et al. (2013) showed that the disappearance of Arctic sea ice in summer induced by ocean warming enhances surface melting in northern Greenland through a decrease of surface albedo and the subsequent atmospheric warming. Thus, the absence of the oceanic component in our modelling setup appears as a limiting factor, although, the direct impact of ocean via sub-shelf melt at the ice sheet margin will likely be limited in the future as a result of inland retreat of GrIS”.

Note however, that MAR has already been coupled to a regional configuration of the oceanic model NEMO (e.g. Jourdain et al., 2011), but applied to the Ross Sea sector in Antarctica. We can therefore reasonably envisage that in the coming years, we will be able to develop a coupled atmosphere-ocean-ice-sheet model.


General: while the writing is 100% understandable and clear, a final proof-read by a native English speaker would be useful as a final stage, if possible, to clear up remaining small grammar issues.

We are aware of the fact that many English mistakes and syntax errors appeared in the submitted manuscript. We made a huge effort to improve English writing.
We would like to thank the reviewer for the evaluation of our study and the constructive comments that helped us to improve the manuscript. Please find below the reviewer’s comments in black font and the author’s response in blue font.

Responses to Reviewer 2

The paper claims to be focused on assessment of the future of the GrIS through 2150. But in fact, it seems more focused on assessment of a new technique for RCM-Ice Model coupling. Throughout the paper, focus shifts back and forth between the two. The experiment is to run a future simulation of the GrIS using MAR coupled with GRISLI in three different ways, and then compare/analyze the results from each other.

The coupling method is interesting, but the GrIS is more interesting. I believe the paper would be better if it would keep its focus firmly on the GrIS, while keeping the methods separate. I ultimately want to know, what do we learn about Greenland? Unfortunately, the figures do not really support that. Figures 1-4 do in a way; but the rest of the figures only really tell us about technical differences between coupling technique. In the revised version, we have largely restructured the manuscript. We now start the result section with a thorough analysis of the GrIS evolution simulated with the most comprehensive method, i.e. the two-way (2W) coupling. The primary focus being now the GrIS evolution, we hope that we have addressed your concerns on this point.

The experiments in the paper show that the different coupling techniques provide different answers. Unfortunately, it is hard to know which answers are closer to the truth, because there are no controls. I came into this believing that the most sophisticated coupler would produce the most melt and also be more accurate; but I had no proof on the accuracy part. This paper has reinforced my prior assumptions, without providing any additional evidence on accuracy. I am therefore hard pressed to say what it has added to my understanding of coupling technique.

All methods, and more generally all models, have their flaws. As stated in the manuscript, both the NF (No Feedback) and PF (parameterized altitude feedbacks) methods (corresponding respectively to NC and 1W in the first version of the manuscript) do not account for the change in atmospheric circulation induced by the change in ice-sheet orography and albedo. The PF method intends to represent the non-linearities of SMB changes with linear corrections based on vertical SMB gradients. Finally, it is fair to say that, compared to the NF and PF methods, the 2W method is the most physically based approach. The two approaches (NF and PF), are inaccurate by construction but have been widely used in the community because of the complexity of including a dynamical ice sheet model in RCMs. Related to your concern on accuracy, we acknowledge the fact that there is only minor constraints to test the validity of our projections: the satellite era covers a relatively short period for which the change in ice sheet topography is small. Thus, although we can state firmly that the 2W method has a stronger physical realism we cannot however guarantee the accuracy of the projections.
I did learn some things about the future Greenland itself, in spite of the figures not really helping with this. I learned:

1. Expect a steeper slope and stronger katabatic winds, in addition to the expected smaller ice sheet. This will result in colder (not warmer) temperatures near the coast.
2. In parts of Greenland, the ELA could be as high as 3000m by the year 2150. I find that idea astounding, at 77 degrees North latitude. Some discussion of this result would be really interesting.
3. Expected sea level rise contribution of Greenland in 150 years is 20cm; and the rate of melting will be continuing to rise at that point.
4. Ice loss and SMB are highly correlated over the next 150 years; so much so that plots of the two look highly similar. Unfortunately, the paper does not try to quantify the correlation.

Thanks for mentioning that. In the revised version of the manuscript, these points appear more clearly along with the description of the 2W results in Sec. 4.1. Katabatic winds are discussed in Section 4.2.1.

Concerning the ELA, we did not mention in the original text that it could be situated as high as 3000 m. In the revised manuscript, we added more information about the shift of the ELA towards higher altitudes (see Fig. 3 and section 4.1.1):

“The equilibrium line altitude (ELA, i.e. altitude for which SMB = 0) increases significantly between the beginning and the end of the 2W experiment, as a consequence of increased runoff for areas below 2000 m. As an example, at around 73.5 °N, on the eastern side of the ice sheet, the ELA moves from ~1000 m to ~2500 m (Fig. 3). In other regions, at the end of the 2W experiment, the ELA is generally situated between 1500 and 2000 m high, except in the northern part where it is between 1000 and 1500 m. This shift of ELA towards higher altitudes represents an increase of 24 % of the ablation area between the beginning and the end of the experiment”.

Concerning the correlation of SMB with total mass loss, we added more discussion on the role of ice dynamics (see Section 4.1.3). As now shown in this section, ice dynamics act to counteract ice loss from surface melting (see Figs 4 and 5). This was also noticed in previous studies (e.g., Goelzer et al., 2013, Edwards et al., 2014a). In turn, ice dynamics is impacted by changes in ice-sheet geometry (see Fig. 6a).

For the record, here’s what I learned about coupling techniques:

1. Integrating SMB over a fixed ice mask over time is a poor way to calculate total SLR contribution, due to the changing ice mask.
2. The 2w case melts more than the 1w or NC case in the RCP8.5 scenario.
3. Full Stokes solvers might yield better results.

Overall... I think this paper has done some interesting modelling runs, but so far has mostly failed to draw interesting conclusions from those runs, and to focus the reader’s attention on those conclusions. I would suggest the authors think through the question “What have we learned about Greenland;“ and then re-do the figures and commentary to support that learning, and focus the reader’s attention on it. The paper will also need significant discussion of these Greenland results, in comparison with other papers that have looked at the future of Greenland; for example, Vizcaino et al 2015. Especially interesting would be places where this
paper predicts something DIFFERENT from those other papers, and why? In this way, the reader needs to be drawn to focus on the most interesting things — the surprises! — first, without having to dig for them.

In the revised version of the paper we did our best to organise the ideas following your suggestion, emphasizing on the fate of the GrIS in our projections with the 2W method. We have also added a thorough discussion with existing literature (see in particular Section 4.4 and Section 5). These sections, as well as the Conclusions (Section 6) have been entirely re-written.

Once the paper has focused primarily on Greenland, I would then think about how to add discussion of a new coupling technique, without taking away from the main scientific focus of the paper. But in the absence of any solid provable way to prove that one coupling technique is better than another, I would avoid making too many claims about the 2w coupling; just that you think it is better, and it certainly melts more ice. In the parts (bulk) of the paper focused on Greenland, I would use whatever coupling technique you think is most realistic.

Here, we do not agree. The 2W method is definitively more physically based than the two other methods and explicitly represents feedbacks that are lacking in NF and PF. For example, the change in albedo in response to ice sheet retreat exerts a major control on local SMB changes that is completely discarded with the two simple coupling techniques. A simpler approach can provide similar estimates for GrIS melt but not always for physical reasons.

A secondary issue: the paper reports many numbers, and only a few of them have error bars. Where did those error bars come from, and why are error bars not reported for other numbers? Would it be possible to get error bars for other numbers?

Because we only have one scenario for each coupling technique we cannot assess statistically the uncertainty in our projections. The +/- signs that you saw for some numbers in the original manuscript stand for the spatial average of the standard deviation for a given variable. For example, in Sec. XX for deltaH=XX+/-YY, the YY is simply the standard deviation in ice thickness change (i.e. the XX value) from the initial condition for a given temporal snapshot. However, in the revised manuscript we now provide the 5th and the 95th percentile values to indicate the range of a given variable.

p.21.24: Studies by Vizcaíno et al (and also at GISS; see Fischer Nowicki 2014) use elevation classes to develop an SMB. Elevation classes are mathematically equivalent to custom-designed gridcells that follow elevation contours. They are therefore able to offer high resolution in the direction of the slope gradient, while continuing with low resolution perpendicular to the gradient.

We agree. Their technique is a way to downscale the SMB from their coarse GCM grid. The technique of elevation classes to downscale SMB is explained in Vizcaíno et al. (2013), not in Vizcaíno et al. (2015) that was cited in the first version of the manuscript. Following your suggestion, we also mentioned the study by Vizcaíno et al. (2013) in the revised manuscript in the Introduction section:
“To circumvent the low resolution, some authors have used the method of elevation classes and are therefore able to offer high resolution in the direction of the slope gradient (e.g. Vizcaino et al., 2013)”.

p.8 l.25: I have traditionally used different labels for the different coupling strategies described. Your "NC", I have traditionally called "1-way coupling." Your "2w", I would call "serial 1w coupling". Your "1w coupling," I would call "corrected 1w coupling." Given the differences in terminology, it’s probably best to describe what each of your schemes is (which you do), but don’t assume that others would use the same names. BTW, none of the coupling schemes here conserve energy, in the sense that two-way couplers (say) between the ocean and atmosphere typically do conserve energy. Therefore, I would be reluctant to call any of them true "two-way coupling."

Following your advice in agreement with the two other reviewers, we renamed the coupling experiments. The experiment with no feedback representation is now called NF (for no feedbacks). The experiment which represents the elevation feedbacks by correcting the MAR outputs is called PF (for parameterized feedbacks). The two-way experiment name remains identical (2W). Since the 2W method does not account for the ocean and since it is based on topography anomalies, we removed all the occurrences of “full two-way coupling” and “fully coupled” replaced them by “two-way coupling”.

p.9 l.7: Why is the 2w scheme more expensive? I see that you have to run the GCM and ice model together, rather than separately. But is any more expense actually involved?

In case you do not have an existing MAR simulation, it is true that the 2W is not drastically more expensive than the two other methods since the only difference is the additional time needed by the ice sheet model (negligible compared to the atmospheric model). However, the major advantage of the NF and PF methods is that we can use existing MAR simulations. In this case, we can run multiple sensitivity experiments since only the ice sheet model is run. We have clarified this point in the text:

 “[The PF] method offers the possibility to account artificially for the elevation feedbacks when using existing RCM simulations in which the topography is kept constant. As such, it is also transferable to any ice sheet model”.

p.9 l.21: Fig. 1 does not support the text. Now I see Fig. 1 is reporting anomalies; but I think it would be more interesting (and no less informative) if it would report actualy Temperature.

Following your suggestions to focus on what happens to GrIS, we start the Result section with a description of the results obtained with the 2W experiment. Therefore, Figure 1 has been changes. It now displays (in absolute values, not anomalies) the evolution of SMB and its components integrated over the whole ice sheet. The spatial distribution of the surface temperature anomaly (2140-2150 vs. 2000-2010) is now given in Fig. 2d.

p. 10 l.2: Cause-and-effect is backwards. Actually, the lower SMB is the CAUSE of the ELA shift.

We totally agree with the reviewer. More precisely, the ELA shift is mainly due to increased runoff (see Fig. 2c). This has been clarified:
“The equilibrium line altitude (ELA, i.e. altitude for which SMB = 0) increases significantly
between the beginning and the end of the 2W experiment, as a consequence of increased
runoff for areas below 2000 m”.

p. 10 section 4.1.2: This is the one section of the paper with error bars. How were those error
bars computed, it didn’t say? Unfortunately, some of the values reported are not statistically
significant; and many others are barely. A more clear way to report the reports in this section
would be something like "we saw no statistically significant change in the GRISLI ice sheet in
the years 2000-2050." This conclusion is already pretty apparent in the figure: the “interesting
stuff” happens further out in time, especially with the more advanced coupling.

As stated earlier in response to one of your comment, the values given in the section 4.1.2 of
the original manuscript were the spatial averages of the standard deviation. The idea behind
these numbers was to have an idea on how geographically different is the variable of interest.
However, we agree that these numbers, averaged over the entire ice sheet do not illustrate
statistically significant changes. In the revised manuscript, the results are most often discussed
as a function of the altitudinal locations. Therefore it does no longer make sense to provide
quantitative results averaged over the entire ice sheet. Instead, we often used the 5th and the
95th percentile values, as previously mentioned.

p.10 l.12-24: This looks like an explanation for the increased slope; but I’m not following it.
The increased slopes are simply due to larger and negative SMB changes at the margin relative
to the interior. Changes in surface slopes have consequences on ice dynamics with increased
slopes leading to increased velocities. We have made substantial text modifications in this
paragraph that now reads:

“The changes in local ice dynamics between the first and the last 10 years of the 2W experiment
are also related to changes in surface slope and ice thickness, particularly at the margins. To
investigate the ice dynamics changes at the local scale, we used the examples of the
Jakobshavn (western coast) and the Kangerlussuaq (eastern coast) glaciers for which the fine
scale structures of the ice velocity, obtained after the GRISLI initialisation procedure, are
relatively well reproduced compared to the observations (Figs 7ab).

For the Jakobshavn glacier, and for altitudes above 1500 m, the vertically-averaged ice
velocities increase by more than 15 m yr\(^{-1}\) (i.e. +10 %) as a result of increasing surface slopes,
and slow down by more than 200 m yr\(^{-1}\) (i.e. +29 %) for altitudes below 1000 m due to the
decreasing ice thickness (Fig. 7c). For altitudes above 500 m, the vertically-averaged velocity
is mainly driven by the SIA velocity (Figs. 7c-e). On the contrary, below 500 m, basal sliding
velocities are large due to low basal drag coefficient (see Fig. 3 in Le clec'h et al., 2018) and
the SSA velocity component dominates the ice flow (Figs 7c and 7g). However, while basal drag
is lower in locations below 500 m, the ice flow is limited by the strongly reduced ice thickness
(Fig. 4).

The Kangerlussuaq glacier is located in regions where the bedrock is characterised by a
succession of valleys surrounded by mountains merging in a canyon where the deepest part is
located 100 km away from the coast (Morlighem et al., 2017). The ice flow of the
Kangerlussuaq is therefore divided in different branches with increasing ice velocities towards
the ice sheet margin and becoming even larger when merging in the canyon (Fig. 7b). As for the Jakobshavn glacier, the ice flow accelerates at the end of the 2W experiment as a consequence of the increase in surface slope for high altitudes (~2000-2500 m, see Fig. 4). Conversely, a strong decrease of the ice flow is found in most of margin regions (Fig. 7d) directly related to the ice thinning (Fig. 4). Contrary to the case of the Jakobshavn glacier that presents large basal sliding velocities only below 500 m, the Kangerlussuaq shows low basal drag coefficients in the entire glacier (see Fig. 3 in Le clec’h et al. 2018) and thus the ice flow is mainly governed by the SSA component (Fig. 7h).
The high correlation between ice thickness anomaly and SMB anomaly shows that climate change due to the imposed RCP forcing is the major control on the Greenland ice sheet geometry change. However, we find important to keep two sections presenting on one hand the changes in SMB and, on the other hand, the changes in ice thickness because it allows to better constrain the role of ice dynamics. Indeed, in our revised version, we show that ice dynamics counteracts the SMB signal (see Section 4.1.3, Fig. 5 and the following paragraph):

“To quantify the role of ice dynamics on the GrIS geometry (Fig. 4), we plotted the ice flux divergence integrated over 150 years (2000-2150, see Fig. 5b). In particular, over the central plateau, the cumulated SMB (Fig. 5a) reaches about +50 m, 40 m of which are transported away by the ice dynamics (Fig. 5b). As a result, the ice thickness anomaly is reduced to only ~10 m in this region (Fig. 4). An opposite behaviour is found near the western coast, where the ice melting is partly compensated by ice convergence, resulting in a less negative ice thickness anomaly than that related to the SMB forcing. This shows that ice dynamics act to counteract ice loss from surface melting, as previously noticed by several authors (Huybrechts and de Wolde, 1999, Goelzer et al., 2013, Edwards et al., 2014b). As a consequence, it appears to be essential to account for ice dynamics to estimate accurately the mass balance of the whole ice sheet”.

p.12 l.30: I appreciate that doing wrong calculations will give the wrong answer. I’m glad that you are not doing that. But is this worth half a section to explain? It seems you are going out of your way because someone else did something fishy.

Our study is the first one to provide the GrIS melting projection that makes use of a RCM coupled to an ice sheet model. This means that all the previous studies based on RCMs, did
not consider the change in ice sheet mask. We therefore think that this section is particularly relevant for the ice-sheet surface mass balance community. To emphasize the importance of the results, this section has been re-written:

“A widely used method to estimate the projected GrIS to global sea-level rise is to compute the GrIS mass loss as the time-integral of the SMB computed by an atmospheric model over a fixed ice-sheet mask (Fettweis et al., 2013, Meyssignac et al., 2017, Church et al., 2013). In the present study, we go a step further since the ice mass variations related to SMB changes are computed over a changing ice-sheet mask as simulated by GRISLI. However, in both the NF and the PF experiments, the atmospheric model does not account for the variations in the ice-sheet extent simulated in GRISLI and the ice-sheet mask, taken from the observations (Bamber et al., 2013) is kept constant throughout the simulation. Taking the changes in ice-sheet mask into account may have strong impacts on the computed GrIS contribution to sea-level rise. To illustrate the influence of the ice sheet mask, we used the SMB outputs from the NF experiment at the MAR resolution and applied the integrated SMB method over the fixed observed ice-sheet mask (SMB\textsubscript{MSK-NF}) and over the updated 2W mask (SMB\textsubscript{MSK-2W}). Results reported in Table 2 indicate differences in SMB values exceeding 23 % in 2150. In the same way, compared to a time variable ice-sheet mask, the use of a fixed ice-sheet mask overestimates the sea-level rise by \~6 % in 2150. Though a bit lower, this number is far from being negligible compared to the errors made when the SMB-elevation feedbacks are not taken into account (i.e. 7.6 %) and when all the feedbacks are ignored (i.e. 9.3 %). This strongly suggests that realistic SLR projections cannot neglect the evolution of the ice-sheet extent, only accounted for through the use of an ice-sheet model”.

p. 13 l.2: the last sentence of this paragraph is the most important. Don’t "bury the leded"... put it up at the front.

We agree with you. In the revised manuscript, an entire paragraph is devoted to the role of the katabatic wind feedback as simulated in our model. We also added the new Figure 9 to support our findings (see Section 4.2.1):

“Over the ice sheet, the steeper surface slopes simulated in 2W in 2150 (discussed in Sec. 4.1.2) lead to a slight increase in katabatic winds (Fig. 9). However, at the ice sheet margin, i.e. where the ice mask in MAR is below 100%, there is a substantial decrease in surface winds. This is because the change in surface elevation as seen by the atmospheric model is computed from the aggregated changes in GRISLI at 5 km. As such, a non-zero fraction of tundra, which presents no change in surface elevation, results in smaller elevation changes compared to grid cell in the same region with permanent ice cover only. This induces artificially lower surface slopes at the margin with respect to the interior and a decrease in surface winds in these regions. Altogether, the slight increase in katabatic winds over the ice sheet and their reduction at the margin lead to a cold air convergence towards the ice sheet edge (Figs. 8b and 9 and Figs. S8-S9). Another consequence of the katabatic winds increase due to increased surface slopes in the GrIS interior, is to enhance the atmospheric exchanges along the slope of the ice sheet. The area with lower atmospheric pressure generated by the stronger katabatic winds is filled in by the warmer air coming from higher atmospheric levels in the boundary layer. The warming of the upper part of the boundary layer combined with the lower surface elevation, explains the ST increases in the interior of the GrIS”.
I don’t believe this argument on ice-ocean feedback. We know that tidewater glaciers retreat VERY quickly once they become imbalanced. How many tidewater glaciers will be left for us to simulate in the year 2050, 2100 or 2150? And what about going beyond that — when the REALLY interesting things start to happen? I just don’t believe that ocean coupling is very important for GrIS.

We agree with you concerning the high probability of having a decreasing influence of outlet glaciers in the future as a result of increased melting in margin areas. We have outlined this in the discussion section (see section 5). However, it remains difficult to accurately evaluate the time scale at which the influence of outlet glaciers on the whole Greenland ice sheet will be negligible. At the centennial time scale, it is therefore highly desirable to have a good representation of tide water glaciers because they have important consequences on inland ice dynamics. A strong change in ice dynamics could in turn strongly modify the SMB signal and the projected sea-level rise contribution. This process cannot represented in our model because of the too coarse GRISLI resolution. As an example, Goelzer et al. (2013) found an additional SLR contribution from outlet glaciers of 0.8 to 1.8 cm in 2100 and 1.3 to 3.8 cm in 2200, as mentioned in Section 5. In addition, ocean may exert a strong influence on ice dynamics through the intrusion of warm waters in the fjord system that can accelerate the destabilization of marine terminating glaciers and the subsequent ice discharge. This leads to a release of freshwater flux in the ocean, modifying oceanic circulation, sea-surface temperatures and sea-ice cover and the exchanges at the atmosphere-ocean interface, resulting in fine in SMB changes (due to changes in external forcings). These ideas have been developed in Section 5:

“There is a growing number of evidence for attributing the acceleration of outlet glaciers to the intrusion of warm waters from adjacent oceans in the fjord systems or in the cavity of floating ice tongues (e.g. Straneo et al., 2012; Johnson et al., 2011, Rignot et al., 2015) that can destabilise the glacier front and/or favour the ice-shelf breakup, decreasing thereby the buttressing effect and increasing the ice calving. In turn, the released freshwater flux in ocean may impact sea-surface temperatures, oceanic circulation and sea-ice cover. Moreover, atmosphere-ocean feedbacks also have an impact on the GrIS. As an example, Fettweis et al. (2013) showed that the disappearance of Arctic sea ice in summer induced by ocean warming enhances surface melting in northern Greenland through a decrease of surface albedo and the subsequent atmospheric warming. Thus, the absence of the oceanic component in our modelling setup appears as a limiting factor, although, the direct impact of ocean via sub-shelf melt at the ice sheet margin will likely be limited in the future as a result of inland retreat of GrIS”.

Any idea what happens beyond the year 2150? I know it’s outside the scope of this paper. But this paper opens up more tantalizing questions by simulating a non-steady-state process just a little bit of the way — to a point where the changes are continuing to accelerate. What does this simulation look like in 500 years? 1000 years? 5000 years? How important are the feedbacks on that timescale?
Unfortunately, running the MAR model over such long time scales is out of reach for the time being because of the considerable computational resources it would require. However, from the ice-sheet perspective, we can reasonably expect:

a/ an amplification of the SMB-elevation feedbacks, as suggested by the results presented in this paper (see Section 4.4);

b/ a smaller ice-sheet extent (possibly combined with a larger ablation area) with therefore a growing influence of the albedo effect amplifying warming and surface melting (see Section 4.4);

c/ increased surface slopes favouring thereby (see Section 4.1.3):
   i/ the convergence of cold air in margin areas through the effect of katabatic winds, acting therefore against warming;
   ii/ the increase of surface ice velocities in the interior regions.

d/ decreased ice thickness leading to a reduction of ice velocities (see Section 4.1.3)

e/ inland retreat of outlet glaciers resulting in their limited influence on ice dynamics (see Section 5);

f/ Multiplication of melt ponds at the surface of the ice sheet, possibly even in high altitude areas leading to:
   i/ surface albedo reduction;
   ii/ increased lubrication and basal sliding.

Of course all the processes listed above should be investigated with a coupled climate-ice-sheet model to investigate their relative influence at different time scales. In addition, atmosphere-ocean-ice sheet feedbacks should also be considered (see our response to your previous comment).

Fig 6A: Why is there a vertical-stripe pattern in western Greenland? That makes me suspicious of the model. Please explain...

This pattern is due to the interpolation method between the coarse MAR grid and the finer GRISLI grid.

** Figures: Please make sure of the following in figures:

a) Avoid the rainbow color scale in most cases (Fig 4). There are better choices.
   
   We kept the rainbow scales, but we paid attention to the choice of the colour to better illustrate our purpose.

b) If you do use the rainbow, avoid splitting green at zero (Fig 4A). One figure has green for both positive and negative numbers; not cool.
   
   We agree. This has been changed.

c) Avoid a color scale that’s read on one end and violet on the other; because then the smallest and largest values look almost the same.
   
   Once again, we agree with you. Colour scales have been changes accordingly.
d) When using color scales with red on one end and blue on the other, make sure that red always corresponds to places that are melting / getting warmer / losing mass; and blue corresponds to the opposite. Reverse the color scale if needed, in order to keep this consistent. All the colour sales are now consistent: Blue colours correspond to a decrease of the displayed variable, and red colours represent the opposite.

e) The figures in this paper all use different color scales and conventions, for no apparent reason. It looks like they don’t belong together. Please make them more uniform, unless there’s a good reason for the difference.
We followed your recommendation.

f) Please put a title on top of every plot, in font large enough to read. Make sure that every plot has units on every axis (either the color scale, or the x-y axis. Most fonts on most figures need to be larger.
We have put a title on most figures and the fonts are now larger.
We would like to thank the reviewer for the evaluation of our study and the constructive comments that helped us to improve the manuscript. Please find below the reviewer’s comments in black font and the author’s response in blue font.

Responses to Reviewer 3

The two-way coupling between a regional climate model and an ice sheet model is an important development that marks a clear step forward to improve the projections of the future contribution of the Greenland ice sheet to sea-level change. The manuscript compares results of the two-way coupling to former methods of representing the interactions between ice sheet and atmosphere and comes to important conclusions concerning the errors implicit to those simpler approaches. The manuscript is of clear interest to the readers of The Cryosphere but still needs to be substantially improved before being acceptable for publication. I recommend major revisions along the comments outlined below.

General comments

The language of the manuscript needs substantial improvement, because many formulations give rise to misinterpretations of the scientific content. While many mistakes could clearly be avoided with a better command of the English language, a large number of typographical errors and mistakes in the referencing suggest that the authors could have made a better effort to deliver a readable manuscript for the review process.

We apologize for the language of the submitted manuscript and the large number of typographical errors. In the revised version, we did our best to avoid English mistakes and typographical errors. We hope that the new version has become more readable.

The text itself reveals that the models are not actually fully coupled (use of anomaly method) and also gives indications why a full coupling is so much more difficult to achieve. I suggest to adjust the title and modify other occurrences of "fully coupled" in the text to "two-way" coupled to take this into account. A discussion item on this point and next steps that need to follow to work towards truly full coupling between RCMs and ISMs should be included.

Following your recommendation, we changed the title of the manuscript by replacing “fully coupled” with “coupled”. Now, the title is “Assessment of the Greenland ice sheet atmosphere feedbacks for the next century with a regional atmospheric model coupled to an ice sheet model”. For consistency reasons, we also removed the other occurrences of “fully coupled” in the text. The “fully coupled” experiment (in the first version of the manuscript) is now simply referred to as the “two-way” experiment (2W). In the discussion section (i.e. section 5), we also discussed the limits of this experiment with respect to a real fully coupling method between RCM and ISM.
The ice sheet initialisation procedure is somewhat non-standard and therefore requires a much better explanation. As it is heavily based on former work that is in part not well documented, an additional effort is required to describe the method in a way reproducible for other modellers.

Since the submission of the present paper, the ice sheet initialisation procedure used in this study has been published in *Geoscientific Model Development Discussions* (see https://doi.org/10.5194/gmd-2017-322). The GMDD paper describes in detail the different steps of the method, its sensitivity to various parameters as well as its limitations. We have therefore simplified the description of the ISM initialisation method in the present paper (Section 2.2.2) by only providing the basic principles of the procedure, and we propose to the readers willing to focus on the spin-up procedure to focus on the GMDD paper. However, in our response to reviewers, we tried to reply as clearly as possible to all questions related to the spin-up procedure.

Finally, the evaluation of the method appears to be based on an experiment that is not closely related to the model state actually used for the projections, which may be possible to resolve with an additional control experiment.

We apologize for this misunderstanding. The evaluation of the initialisation procedure is actually based on a 2000-year forward control experiment with constant forcing. In the revised paper, we have clarified the points related to this control experiment. Below, we also provide a detailed answer to the comments associated with this experiment.

The thermodynamic aspect of the model is not well represented, arguably because it plays a minor role for the present work. Nevertheless, substantial computing time is spent during initialisation to equilibrate the temperature and the role of bottom and surface boundary conditions is mentioned. Therefore, the model description in 2.2 requires at least a short description of this model component.

Following your recommendation, a short description of the thermodynamic aspect of the ice-sheet model has been added in section 2.2 with the following paragraph:

“Basal melting occurs when the basal temperature is at the pressure melting point. The ice temperature plays a crucial role in the dynamics of the ice sheet because it also affects the viscosity, and thus the ice flow in the entire ice column (Ritz et al., 1997, 2001). In turn, heat released by internal ice deformation and basal dragging over the bedrock modifies the temperature. The temperature field is computed by solving a time-dependent heat equation both in the ice and in the bedrock accounting for advection and vertical diffusion processes. At the surface, the boundary condition is provided by the prescribed surface temperature. At the base of the ice sheet, the boundary condition is given either by the geothermal heat flux or by the temperature melting point at the ice-bed interface.”

The experiment names are not specific enough and should be improved. For my understanding, what is presented as the method "no coupling" is in fact a one-way coupling, where the ice sheet is responding to changes computed by the RCM (with "no feedback"). 2-W is correctly described, but 1-W is somewhere between one-way and two-way coupling because it parameterises the feedback. Maybe you could use "no feedback", "parameterised feedback" and "two-way" instead.
We followed your recommendation and changed the name of the experiments referred to as NC and 1W in the first version of the manuscript. These experiments are now referred to as NF (for “No Feedback”) and PF (for “Parameterised Feedbacks”). NF corresponds to the experiment in which GRISLI is forced by the MAR climate, and PF corresponds to the experiment in which both SMB and ST fields simulated by MAR are corrected to account for topography changes simulated by GRISLI. In our revised manuscript the name of the two-way coupling method (2W) has remained unchanged.

The most important question in the comparison of results after 100 years and after 150 years is left open: why does the behaviour of 2-W suddenly change around 2010. For this it may be instructive to also look more detailed at around 2060, where a similar shift is possibly visible.

In the first version of the manuscript we insisted on the behaviour of the 2W experiment after year 2100. However, a closer examination of the results clearly shows that the evolution of the Greenland ice sheet to sea-level rise diverge from one experiment to the other as soon as 2025-2030, namely only a few years after GrIS-atmosphere feedbacks are accounted for in the 2W experiment and in the parameterized feedback experiments (SMB-elevation feedbacks in this case). As a result, while the effect of the feedbacks on sea-level rise becomes significant by the end of the 21st century, it starts to operate much earlier and is amplified over time. To illustrate this point, we added a –zoom-panel in the new Figure 12 displaying the evolution of the anomalies (2W-NF, PF-NF and 2W-PF) of the GrIS contribution to SLR.

Otherwise, I find the comparison redundant because the bottom line in most cases is ‘like after 100 years, only stronger’.

We fully agree with this remark. In the revised version of the manuscript, we mostly present the results at the end of the simulations (i.e. 2150) and we only discuss the temporal evolution, including the results by 2100, only in Sect. 4.4.

The integration of SMB anomalies already discussed in the manuscript could be added as an additional experiment, possible even two, if masking would be additionally taken into account. This would facilitate the comparison and place the discussion of the effect of masking on firmer ground.

Estimates of sea-level rise from the time integral of SMB anomalies were already discussed in the submitted manuscript. In the revised version they are referred to as SMB_{MSK-NF} and SMB_{MSK-2W}. Both are based on the SMB outputs from the NF experiment (at the MAR resolution), but the time integral of SMB anomalies is made either on the fixed present-day ice-sheet mask (SMB_{MSK-NF}) or on the time variable ice-sheet mask simulated in the 2W experiment. The results are discussed at the end of section 4.4. However, since these estimates are inferred from diagnostics of already performed experiments (i.e. NF and 2W), we think it is not appropriate to present them as additional experiments.

Title I would argue that the models are not "fully", but rather "two-way" coupled because an intermediate down-scaling step is necessary and, more importantly, an anomaly method is used. To avoid any confusion, we changed the title in:
“Assessment of the Greenland ice sheet - atmosphere feedbacks for the next century with a regional atmospheric model coupled to an ice sheet model”.

P1.L1 Better “the projected sea-level contribution from the *Greenland ice sheet*”. Also mention a typical time scale here to make clear this is about the centennial time-scale. The first sentence now reads as:

“In the context of global warming, a growing attention is paid to the evolution of the Greenland ice sheet (GrIS) and its contribution to sea-level rise at the centennial time scale”.

In the main text we use “the GrIS contribution to sea-level rise” or, following your suggestion, “the projected sea-level rise contribution from the Greenland ice sheet”.

P1.L2 Be more precise about the mechanisms and feedback(s). The next sentence ("these feedback*s*") suggests that "temperature and surface mass balance – elevation feedback" refers to at least two feedbacks. What are these precisely? "surface mass balance – elevation feedback" is clear, but what is the role of temperature? Note also that melting is clearly related to temperature increase, but the SMB is ultimately controlled by the energy balance.

The abstract has been extensively modified. The sentence you mentioned has been changes in: “Atmosphere-GrIS interactions, such as the temperature-elevation and the albedo feedbacks have the potential to modify the surface energy balance and thus to impact the GrIS surface mass balance (SMB). In turn, changes in the geometrical features of the ice sheet may alter both the climate and the ice dynamics governing the ice sheet evolution”.

P1.L5 A bit confusing to mention start date as 2020. It is understood later that before 2020 elevation changes are considered too small to make a difference. But at this place it may be better to give the period of the entire simulation (2006 - 2150). Note also that the RCP is not defined beyond 2100, so it is better to mention "prolonged RCP 8.5 scenario". Recommendation followed.

P1.L5 It seems confusing to call this simple method "no coupling", since it represents a one-way coupling. See also general comment on naming the experiments. As advised we have change the name of the experiment: No coupling (NC) becomes No Feedbacks (NF) experiment and one-way coupling (1-W) becomes Parameterised Feedbacks (PF) experiment.

P1.L6 Could mention that this one-way coupling methods attempts to incorporate or parametrise two-way interaction. It represents an intermediate method between one-way and two-way coupling. See also general comment on naming the experiments. In the abstract, we first present the two-way coupled approach, and then the one-way coupling experiment (i.e. NF). The parameterised feedback experiment is then defined as an “alternative one-way coupling approach in which the elevation changes feedbacks are parameterised in the ice-sheet model”.

P1.L7 I suggest to omit "offline". The correction may be offline to MAR, but it is online to the ice sheet model, as the correction is updated every time step and dependent on the current ice sheet
elevation. Could add what is happening with the extent, since it has been explicitly mentioned for the former method.
We agree, it’s offline to MAR but not for GRISLI. The ice sheet extent is not updated in the one way parameterised coupling method (PF). Only ice sheet topography changes computed by GRISLI relative to the observations are used to correct the SMB fields.
In the revised version of the manuscript we do not mention explicitly what happens with the ice sheet extent in the PF experiment, but we explain that only the surface mass balance - elevation feedbacks are parameterised, which implies that the ice-sheet extent in the atmospheric model is kept constant as in the NF experiment.

P1.L9 Clearer to replace "ice sheet elevation feedback" by "surface mass balance – elevation feedback".
"Ice-sheet elevation feedback” has been removed from the entire text in favour of SMB-elevation feedback, melt-elevation feedback or temperature-elevation feedback, depending on the context.

P1.L9 Maybe ", the one-way and two-way coupling methods ..." since the amplification occurs in both cases.
We agree. In the revised manuscript, this part of the abstract has been completely reorganised but we paid attention to make clear that SMB-elevation feedbacks are amplified over time both in the PF and the 2W experiments.

P1.L11 Some ice sheet margins are not in the coastal region. Replace by "ice sheet margins" or similar. This should be followed throughout the document for other occurrences.
Recommendation followed.

P1.L15 "52 400 km^2 smaller"
In the revised version of the manuscript, we discuss the relative changes (in %) in ablation area (rather than changes in ice-sheet extent) between NF vs 2W and PF vs. 2W.

P1.L16 "fixed ice sheet mask"
OK, modified

P1.L20 "always" is only true for the end of the simulation. In the first decades or so the volume loss difference cannot be significant. Maybe give an estimate for a time scale where this is true similar to the comparison one-way vs. two-way.
In the revised version of the abstract, we just mention that the effect of feedbacks is amplified over time. However, in Section 4.4, we specify that the feedbacks make the three simulations diverging from each other only a few years after taking into account the feedbacks (i.e. after 2020). This is illustrated in Figure 12. However, we also mention that the effect of the feedbacks becomes significant only after the end of the 21st century.

P2.L6 "the ablation" (singular) <-> "are processes" (plural). Reformulate
OK, this has been reformulated.
Some risk for confusion here. It is a bit simplistic for a paper discussing an RCM as an important component to reduce the interaction to changes in SMB and temperature. It is understood that these are the two variables used to force the ice sheet model, but a bit more detail is required. How does the change of ice extent change the albedo and therefore the energy balance?

In the revised version, our arguments have been a bit more developed. The first paragraph of the Introduction has been modified as follows:

“The evolution of the Greenland ice sheet (GrIS) is governed by variations of ice dynamics and surface mass balance (SMB), the latter being defined as the difference between snow accumulation, further transformed into ice, and ablation processes (i.e. surface melting and sublimation). While surface melting strongly depends on the surface energy balance, snowfall is primarily controlled by atmospheric conditions (wind, humidity content, cloudiness...). However, various feedbacks between the atmosphere and the GrIS may lead to SMB variations that can therefore directly affect the GrIS total mass by impacting its surface characteristics, such as ice extent and thickness, with potential consequences on ice dynamics. These changes may in turn alter both local and global climate. As an example, changes in near-surface temperature and surface energy balance may occur in response to changes in orography (temperature-elevation feedback) or in ice-covered area (albedo feedback; see Vizcaino et al., 2008, 2015; Lunt et al. 2004). On the other hand, topography changes may alter the atmospheric circulation patterns (Doyle and Shapiro, 1999, Petersen et al., 2003, Moore and Refrew, 2005) causing changes in heat and humidity transports”.

Does temperature enter the correction method and how?
Yes, both ST and SMB, used as forcings of the ice-sheet model are corrected in the PF experiment, using the Franco’s et al. (2012) method. In the 2W experiment, they are explicitly computed as a function of the evolving topography (computed by GRISLI), following the same procedure as in the PF experiment.

OK, accumulation and ablation are sensitive to ST, but why and how?
Processes have been clarified in the modified paragraph reported above

Also, what is the role of ST other than its influence on SMB, as boundary condition to ice thermodynamics? Does it have an impact on the simulations at all (I don’t expect it, but would be good to say something about why not and being able to exclude it).
The surface temperature (ST) applied as a boundary condition of the ISM allows to compute the vertical temperature profile (i.e. the surface conditions diffuse from the surface to the base of the ice sheet and modify the ice flow by changing the viscosity of the ice). Using ST as a climate forcing is therefore a pre-requisite to run the ice-sheet model. However, at the century timescale, ST has not time enough to diffuse farther than the surface layer. Thus, in the present study, changes in ST during the 150-yr experiment have only a very limited impact on ice dynamics.
Maybe already intended, but make really clear that the changes in ST have no direct effect on thickness volume and extent. Reformulate.

We acknowledge that this sentence was not clear. The overall paragraph has been reformulated (see above).

Replace "disrupt" by "modify"

OK, modified

More detail needed. Amplification of mass loss by what process under what forcing and compared to what other (control) experiment?

We clarified these points in the revised version:

“Compared to a control experiment in which the ISM is forced off-line by the atmospheric model run with the fixed present-day GrIS topography, they found an amplification of ice mass loss of 8–11 % and 24–31 % in 2100 and AD 2300 respectively, when the elevation feedbacks are taken into account (i.e. in the coupled experiment). This results from the combination of the positive elevation-SMB feedback in low lying areas, the negative feedback related to the elevation-desertification effect in accumulation areas, and the changes of surface slopes resulting from high mass loss in ablation areas and slight snowfall increase in the accumulation zone, enhancing the ice transport from the central regions to the ice margins”.

The beginning of this sentence suggests (and I agree) that increased resolution would help to improve the modelling compared to observations, while "more detailed physics" is at least for the ice sheet model typically associated with ‘less approximation’, i.e. higher order physics. Could you add some detail to distinguish these.

This sentence has been completely reformulated. In particular, we specified that the ISM (SICOPOLIS) used in Vizcaino et al (2015) is based on the shallow ice approximation and is therefore not able to properly capture fast flowing of outlet glaciers. As suggested by Reviewer 1, we also mentioned the study of Löfverström and Liakka (2017) who confirmed the importance of the spatial resolution in coupled climate – ice sheet experiments in a paleo-climatic context. They explain that ISM results are limited by the capacity of the climate model to simulate atmospheric temperature and precipitation at low spatial resolution as a consequence of the poorly resolved planetary waves and smooth topography.

Should introduce RCMs and add references to MAR, RACMO, HIRHAM ... already here, as that is the obvious choice to increased resolution. Introducing the Franco and Edwards methods is already a step further as it is based on RCM output.

As recommended we have firstly introduced RCMs with references for MAR (Fettweis et al. 2017), RACMO2 (Noël et al., 2015), Polar MM5 (Box et al. 2013) and HIRHAM5 (Langen et al. 2015). We then mentioned the altitude corrective methods of Franco et al. (2012) and Edwards et al. (2014).
We added the same references as those mentioned in our response to the comment P2.L26 (see just above).

P3.L9 Specify again for what it is a requirement.
The sentence has been modified as follows:
“The second fundamental requirement to describe the interactions between atmosphere and GrIS is to represent the ice sheet topography changes in the atmospheric model by using an ISM (instead of the fixed geometry typically used) to take into account the effects of ice dynamics on the ice sheet topography changes”.

P3.L10 Reformulate "usually used" to "typically used" or similar.
This has been reformulated (see the sentence reported just above in our previous answer).

P3.L11 Add reference to Goelzer et al. 2017 here, since it is specifically on GrIS models.
Sorry for this omission. The reference has been added.

P3.L18 Remove "high resolution" or specify explicitly at what resolution GRISLI is run.
We removed “high resolution” from the sentence and specified at what resolution MAR and GRISLI are run in section 2.

P3.L21 "two-way"
OK modified

P3.L25 I would consider the three methods part of the experimental setup and therefore name initialisation and experimental setup first.
The paragraph describing the organisation of the paper has been reformulated according to the new structure of our revised version. Section 2 (entitled Models) describes the atmospheric and the ice-sheet models together with their respective spin-up procedures and boundary conditions. Section 3 (entitled Coupling methods) is now focussed on the description of the three coupling methods.

P4.L4 "developed". Correct also throughout the manuscript.
OK corrected everywhere

P4.L4 "SISVAT" requires a reference and description of the acronym.
OK specified.

P4.L16 "ice albedo that has been improved by parametrising the impact of melt ponds on the albedo."
OK corrected.

P4.L19 Replace "provided by" by "taken from"
OK corrected
"forced with 6-hourly atmospheric fields". See also P8.L6
OK corrected.

Remove "forcing"
OK removed

P4.L27 Suggest reformulation to "... because it has been shown by Fettweis et al. (2013), to be the best choice from the CMIP5 data-base to reproduce the present-day climate compared to results of MAR forced by reanalyses."
OK, modified.

P5.L1 Heading "Climate model initialisation and experiment"
Section 2 (and its related subsections) has been reorganised following the recommendations of the three reviewers and heading is now “Models” This section is still divided in two subsections 2.1 and 2.2 devoted to the description of the MAR and GRISLI models respectively. Section 3 is devoted on the description of the three “coupling” experiments.

What is the difference between "spurious drifts" and "unwanted trends" or are they one and the same? Reformulate.
We apologise for this misunderstanding. We used two different expressions to deal with “unwanted trends”. The sentence has been reformulated as: “Before starting our experiments, MAR needs to be properly initialised to limit unwanted trends in the results”.

Replace "SISVAT requires more than 6 years", by "SISVAT requires less than 7 years" to make clear that the chosen 7-year period is long enough. Or otherwise explain why 7 years is considered OK.
OK modified.

Replace "provided by" by "taken from". Add explanation how the data was interpolated to the coarse MAR grid.
We replaced “provided” by “taken from” as suggested. In the revised version, we specified that the GrIS topography from the Bamber et al. (2013) dataset is aggregated on the MAR grid.

Be consistent in if SISVAT is written in italic or not.
OK corrected

Replace "following year 1976" by "from 1977 onward".
Sorry for this misunderstanding. In the revised version we clarified that MAR is initialized with MIROC5 climatic fields from 1970 to 1975 included. The MAR simulations start in 1976, but the results presented in this paper are for the period 2000-2150. This has been clarified in the revised manuscript. Therefore we changed the sentence in:
“Here, MAR is initialised with the atmospheric forcing fields from MIROC5 from 1970 until 1975 and the MAR simulations start in 1976. However, in this paper, the MAR results will be analysed for the period spanning from years 2000 to 2150”.

P5.L10 Need to explain in more detail why 2095 can be considered representative for the 2090s. Is it e.g. the year that is closest to the decadal mean? Are trends so linear that the middle of the decade are representative for the average? Typically one would use the decadal mean to represent the long-term average and not one individual year, unless it doesn’t matter for some reason.

We chose to force MAR with the 2095 climate from 2101 to 2150 because, averaged over the entire GrIS, the 2095 climate is one the closest to the decadal 2090-2100 mean climate. We acknowledge that, in the absence of a MIROC5 simulation run under a prolonged RCP8.5 scenario it would have been more appropriate to repeat the ten years (2090 -2100) until 2150, but it would have been more complex to set up.

P5.L12 Better to omit "coupled" here, since it is not clear what is coupled to what and it is further detailed later.
OK, “coupled” has been removed.

P5.L15 "the northern hemisphere ice sheet*s* (NH references) and the Greenland ice sheet (GrIS references)”. or "the northern hemisphere ice sheet*s* and the Greenland ice sheet (all references)"
OK modified according to the suggestion.

P5.L16 "... covering Greenland with ...", since the coverage extends outside of the ice sheet mask. Add information about the vertical.
We have also specified that GRISLI has 21 vertical evenly spaced levels.

P5.L17 Need to specify what "hybrid" means.
The word "hydrid” has been removed from this part of the text and introduced after having explained the basic principles of both the shallow-ice and the shallow-shelf approximations: “Using a hybrid model (i.e. based on both SIA and SSA approximations) allows to better represent the different deformation regimes found in an ice sheet”.

P5.L19 Need to add explanation on the thermodynamic aspect of the model.
We added new information on the thermodynamic aspects:
“Basal melting occurs when the basal temperature is at the pressure melting point. The ice temperature plays a crucial role in the dynamics of the ice sheet because it also affects the viscosity, and thus the ice flow in the entire ice column (Ritz et al., 1997, 2001). In turn, heat released by internal ice deformation and basal dragging over the bedrock modifies the temperature. The temperature field is computed by solving a time-dependent heat equation both in the ice and in the bedrock accounting for advection and vertical diffusion processes. At the surface, the boundary condition is provided by the prescribed surface temperature. At the
base of the ice sheet, the boundary condition is given either by the geothermal heat flux or by the temperature melting point at the ice bed interface”.

See also P5.L27 SIA velocity is even stronger controlled by ice thickness. Both the ice surface slopes and the ice thickness occur in the computation of the SIA and the SSA velocity with the same exponent. We therefore modified the sentence as: “The ice thickness and the ice-sheet surface slopes control the SIA and the SSA velocity components, but the SSA is also governed by basal dragging”

P5.L28 "SSA component is mainly controlled by the ice flux" is confusing because ice flux is velocity x ice thickness. Clarify! This has been clarified and corrected in the revised version (see our previous response).

P5.L29 "rheologies" is the wrong term here. Maybe "deformation regimes". Yes, you are right. We replaced by “deformation regimes”

P5.L30 Replace "ice melting point" by "pressure melting point" OK, corrected

P6.L3 Replace "floating criterion" by "floatation criterion" Ok, corrected

We acknowledge that this expression was too vague. It referred to the nature of the bedrock (i.e. water-saturated sediment or not). However, this part of the manuscript has been re-written and the sentence has been deleted.

P6.L5 Does that mean the enhancement factor differs for different regions? Explain
Alike most ice-sheet models, GRISLI considers the ice as a non-Newtonian viscous fluid that follows the Glen’s flow law (with the coefficient n generally fixed to 3). However, a particularity a GRISLI is also to account for a Newtonian contribution (i.e n =1) for low deformation rates leading to a polynomial Glen’s flow law in which we apply an enhancement factor in SIA areas to favour longitudinal deformations. In addition, a fixed ratio between the SIA and the SSA enhancement factor is used. The polynomial Glen’s flow law is expressed as:

\[
\frac{1}{\eta} = (E_1 B_1 (T) + E_3 B_3 (T) \cdot \tau^2 )\tau_{ij}
\]

where \(\eta\) is the ice viscosity, \(\tau\) is the shear stress tensor and \(\tau_{ij}\) is the deviatoric stress tensor, \(B1(T)\) and \(B3(T)\) are temperature-dependent coefficients following and Arrhenius equation for coefficients \(n= 1\) and \(n= 3\) respectively and \(E1\) and \(E3\) are the corresponding enhancement factors. As a result there are theoretically four enhancement factors (2 for the SIA component of the velocity with \(n= 1\) and \(n=3\) and 2 for the SSA component of the velocity). In practice, for the simulations presented in this paper, we used \(E1_{SIA} = E3_{SIA} = 1\) and \(E1_{SSA}/E1_{SSA} = 0.125\).
After a careful examination of the paper and reviewers comments, we do believe that any mention to the enhancement factor does not provide any added value to the manuscript. We therefore removed the corresponding sentence.

**P6.L6 "ice loading changes"
OK, corrected.**

**P6.L7 Add a reference for the used isostatic model.
OK, Le Meur et al. (1996) added for the ELRA model.**

**P6.L8 Add a reference describing the thermodynamic model.
As previously mentioned, we added a new paragraph to describe the thermodynamic aspects of the GRISLI model and added the references Ritz et al (1997, 2001).**

**P6.L10 This whole paragraph needs to be reworked. Be more specific. What is considered a boundary condition, what is input data and what is considered a forcing? What variables are concerned for ice flow, ice thermodynamics and isostasy?**
We acknowledge that this paragraph was very confusing. In the revised version, subsections 2.2.1 is now devoted to the description of the GRISLI ice-sheet model and section 2.2.2 is focused on the spin-up procedure. Following your recommendation, the paragraph concerning climate forcing, initial conditions and input data has been entirely re-written. Now it reads as: “The climatic forcing is given by the mean annual SMB and the mean annual ST. Because seasonal variations of surface temperature are rapidly dampened, ST is considered as a good approximation of the bottom snowpack temperature. The initial GrIS surface and bedrock topographies come from Bamber et al. (2013) and the geothermal heat flux is taken from Fox Maule et al. (2009)”.

**P6.L11 Is there a difference between "The annual mean near surface air temperature" and ST? If yes, explain, if not, use TS instead.
No, there is no difference: ST represents the mean annual near surface air temperature. This has been clarified in the new version of the manuscript.**

**P6.L13 What data are these ‘boundary conditions’ and which variables are taken from which data set? Surface elevation, bedrock elevation and ice thickness are not boundary conditions to the equations that GRISLI solves in the proper sense. You could call this "input data" instead.
As mentioned above, we clarified the text.**

**P6.L14 "The climatic forcings". Say what they are! TS and SMB?
The climatic fields used as GRISLI forcings are the SMB and the ST. This has been clearly specified in the new version.**

**P6.L15 If basal drag were a boundary condition, it could hardly be computed. Reformulate to make this clearer.**
The basal drag coefficient is only adjusted during the initialization. In forward experiments, it remains constant through time and its spatial distribution is fixed to that obtained at the end of the initialization. It can be thus considered as an input data, at least for transient experiments.

P6.18 Heading "Ice sheet model initialisation and experiments"
As previously mentioned this sub-section has been canceled and the text has been moved to the main section 2.2

P6.19 The motivation is not quite correct. I would argue that to equilibrate the model to a steady state is not a necessity given the approximations, but rather a choice. One could envision a transient spinup as initialisation with the exact same model.
Yes, we agree with this comment. It seems more appropriate to only deal with the long time-scale response of the ice sheet. Our motivation has been reformulated as: “Due to the long time scale response of the ice sheet to a given climate forcing, a proper initialisation of the model is required before performing forward experiments”

P6.20 Again, more precision needed. What the ice sheet model equilibrates to is rather the climate forcing held constant for this particular initial steady state experiment.
The text has been modified as follows:
“the aim of the initialisation is to start the simulations from a present-day ice sheet geometry as close as possible to the observed one while ensuring consistency between internal properties of the ice-sheet (e.g. basal sliding velocities and vertical profile of temperature) with the climate forcing”.

P6.20 Replace "sensitivity" by "forced" or "forward".
We replaced “sensitivity” by “forward”.

P6.20 Reference Le clec’h et al. (in prep) is not in the reference list. If you are referring to the present manuscript, say that instead of using an external reference.
No, we did not refer to the present manuscript. We simply omitted to add the reference Le clec’h et al. (in prep) in the reference list. This paper describes in details the initialization procedure. Since the submission of the present manuscript, Le clec’h et al. (in prep) has been published in GMDD. In the following (as well as in the revised manuscript) it is referred to as Le clec’h et al. (2018). As a result, in the revised version of the present manuscript, this reference appears as Le clec’h et al. (2018). Moreover, we made the choice to only present the basic principles of the initialisation procedure to avoid redundancy with the GMDD paper.

P6.22 Replace "avoid" by "reduce", since the method is not perfect. Also I’d suggest the formulation "reduce an initial adjustment of the model during the first years of the simulation due to factors not related to the climate forcing alone." or similar.
In the revised version, we no longer speak about “an initial adjustment of the model”. Instead, we explain that the aim of the initialisation procedure is to “reduce the difference between the observed and the simulated ice thickness”.
P6.L25 Reformulate "just over the bedrock". Maybe "basal conditions". We agree with this suggestion. However, as a result of the simplified description of the initialisation procedure, the corresponding sentence has been removed in the new version of our manuscript.

P6.L26 If basal conditions are "likely to change in time" your method to define spatially variable but *temporally fixed* basal drag coefficient could never be successful. Should add here that your method assumes them to be constant over the 150 years of your experiment. This is actually a limitation of the method and this is why it cannot be applied for long-term transient experiments. However, over the 150 years of the experiments, we assume that basal conditions do not change so much and that the best guess for the basal drag coefficient obtained at the end of the spin-up procedure is a good approximation of the basal dragging at the century time scale. Because of the simplified description of the spin-up procedure, this part of the text has been removed. However, we specified in the revised manuscript that "$\beta$ is a time constant but spatially variable basal drag coefficient".

P6.L27 Suggest to remove sentence "As a result any error in the basal velocity computation can spread vertically in the ice and generate slowdown or acceleration of ice sheet motion." In its present form this sentence is generally true in any case and doesn’t support your chose of assimilation method. We followed this suggestion and removed the sentence.

P6.L29 It is not clear to me at what point in the procedure observed velocities are actually used. Which observational data set is used? Reference needed. The observed velocities (Joughin et al., 2010) are only used as input data for the first iteration. The actual target is to reduce as best as possible the mismatch between the observed and the simulated ice thickness.

P6.L30 "three main steps:" Make a numbered list (possible with lists of sub-steps) to facilitate navigation of the different steps.
We acknowledge that is part of the text was not well written and contained several misleading formulations. As explained above, the presentation of the spin-up method has been reduced to its basic principles (see section 2.2.2) in this revised paper because the full description of the method can be found in Le chlec’h et al. (2018).

P6.L31 Replace "not necessary consistent between them" by "not necessarily mutually consistent". Recommendation followed.

P6.L32 It looks to me like the first guess of basal drag mentioned here is a very good first guess and further adjustment of the basal drag coefficient is very much based on it. At any rate, a full description of the procedure used to arrive at that stage should be included, otherwise the method is not reproducible with another model (and not even with GRISLI itself). See also general comment on initialisation.
The first guess of the basal drag coefficient comes from a preliminary version of the spin-up procedure summarized in the present paper and fully detailed in Le clec’h et al. (2018). This former procedure was set up for Ice2Sea simulations carried out with exactly the same GRISLI model version and the same initial Greenland ice-sheet topography (Bamber et al. 2013) as those used in the present study, but with a different climate forcing, implying the need for adjusting the basal drag coefficient. Furthermore Le clec’h et al. (2018) have shown that the final value of the basal drag coefficient (i.e. used for forward experiments) obtained after the spin-up procedure is very poorly dependent on the initial guess (see Fig. 3 in the GMDD paper).

P6.L32 Edwards et al is a multi-model intercomparison and does not give specific details on the assimilation technique for GRISLI. The model reference there is given as Quiquet et al., 2012), which does not provide information on spatially variables tuning. Again, the method to produce the first guess basal drag needs to be made transparent for other modellers to be able to reproduce the results.

The reviewer is right concerning the reference Edwards et al (2014a). This was cited to inform the reader that the initial guess of the basal drag coefficient was coming from the Ice2Sea project. We acknowledge this was not appropriate since this paper does not contain any detail about the assimilation technique. In the revised manuscript, we provide a piece of information about the method used to obtain this first guess (see our previous response).

P6.L32 "surface and bottom". I think you mean surface elevation and bedrock topography. Be more specific!

This is right. This part of the text has been reformulated (see our response related to the main steps of the initialization procedure).

P7.1.1 "vertical fields" Be more specific!

We dealt with the vertical and temperature profiles. Once again, this part has been reformulated.

P7.L2 If I understand correctly, you calculate something here in the first step to be used in the second step. Maybe you should say that. Confusing to mention already here "to have an ice flux as close as possible to observation" when diagnostically calculating something here will not have any influence on the match of the ice flux with observations in this step. This could be mentioned in the second step or as a general motivation for your method before.

Completely reformulated to make clearer the description. Indeed, the basal drag coefficient computed during the 1st step (see new description) is used in the 2nd step. This has been specified in the new version.

P7.L3 Not clear to me how to derive a factor (a/b) from a difference (H1-H0). Please provide an equation or better explanation what the underlying idea is, what is done here, and how it is calculated?

The revised manuscript includes equations supporting the spin-up description. We hope this will help to avoid any ambiguities.
You are mixing topography differences and ice thickness differences. Possibly similar or identical in absence of bedrock adjustment, but is it necessary to distinguish them?

Since bedrock adjustment is negligible in the present study (owing to the addressed time scales), surface elevation differences and ice thickness differences are supposed to be very similar. However, to avoid any confusion, we only use the term “ice thickness” throughout the revised manuscript.

Again "the factor allows to decrease (resp. increase) the surface ice velocity" is confusing, because this is not happening in this first step. Also "If *locally* the topography difference *is* positive ..."

The previous description was misleading. In the revised manuscript, the first step consists in computing a new value of the basal drag coefficient from the value obtained at the previous iteration and from the ratio of the sliding velocity over the corrected sliding velocity. This ratio represents the corrective factor to reduce the mismatch between observed and simulated ice thicknesses. The way the corrected sliding velocity is computed is now fully described in the revised manuscript. For the very first iteration (i.e. just after the 5-year relaxation), the first step is skipped because there is no difference between observed and simulated ice thicknesses and the procedure starts at the second step. This has been also specified in the new version of the paper.

How does deltaH translate into deltaV?
The relationship between the ratio of $H^G/H^{obs}$ and the ratio of the vertically averaged velocity is given by Equation 4.

How does the new velocity compare to observed surface velocities?
There was a confusion in the revised manuscript when dealing with surface ice velocity. Actually, surface ice velocity have to be replaced by “vertically-averaged velocity”. As a result, we do not compare the new surface ice velocities to the observed ones in the present paper. However, this comparison can be found in the GMDD paper (see Fig. 8 herein): we show that the overall patterns of the simulated ice surface velocities are generally in good agreement with observations (particularly in regions of fast ice flows), despite slight differences in the central plateau where the ice velocities are low.

How is the new coefficient calculated? Explain in detail.
The description of the method has been clarified in the revised manuscript and the way the new basal drag coefficient is calculated has been explained in detail (see the new section 2.2.2 in the revised paper).

This is reminiscent of the method of Pollard and DeConto 2012, could you describe the similarities and differences to their approach?
The reviewer is right. In the revised manuscript, we specified that our spin-up method is based on the same basic principles as that of Pollard and Deconto (2012) in that their basal sliding coefficient is adjusted so as to reduce the difference between simulated and observed ice-sheet
topography. We also mentioned the main differences between their method and ours. The new paragraph is:

“Based on the same basic principles as that of Pollard and DeConto (2012), our method consists in the adjustment of the spatially-varying basal drag coefficient (and thus of the basal sliding velocities, see equation 3) so as to reduce the difference between the observed and the simulated ice thickness. However, while the study by Pollard and DeConto (2012) requires long (multi-millennial) integrations for the method to converge, we suggest instead an iterative method of short (decadal to centennial) integrations starting from the observed ice thickness”.

Surprisingly your adjustment goes very fast (in total less than 2000 years). This makes me believe that the original basal drag was already a good guess and you only need minor adjustments. Is that correct? How different is the final basal drag field from the initial one? Can we see a figure for this comparison?

As previously mentioned, we have shown in the GMDD paper that the convergence of our spin-up method is only poorly dependent on the choice of the original basal drag coefficient. Sensitivity tests performed with a uniform β coefficient (β=1) and with the same spin-up parameters (i.e. 20 years for the duration of each iteration, 200 years for the free-evolving simulations and Nbcycle = 8) results in negligible differences in the final basal dragging compared to that inferred from our “standard method” (i.e. a first guess for β coming from Ice2Sea simulations), and in an ice thickness root mean square error (+ 62 m) fully comparable to that obtained in the present study after 8 cycles (+ 63 m). These results are illustrated in Figure 3 in the GMDD paper. Hence, they are not reported in the present study.

P7.L15 Replace "minimum gap" by "error".

This part has been reformulated. Throughout the manuscript we use “mismatch” or “differences” between simulated and observed ice thicknesses.

P7.L15 You additionally need to convince the reader here that this method is optimal in the parameter choices (adjustment time 20 y, relaxation time 200 years) and to make clear in how far the results are (not) dependent on these choices.

The results (in terms of time of convergence and ice thickness root mean square error) are obviously dependent on the choice of the adjustment and relaxation time and of the number of cycles. As explained below (see response concerning the stopping criterion) they have been chosen to minimise the ice thickness RMSE. This has been mentioned in the revised paper. Numerous sensitivity studies with different sets of parameters have also been carried out and presented in Le clec’h et al. (2018). In the revised paper, we specify that: “The overall process is stopped when the ice thickness root mean square error is not significantly improved. This ensures a good compromise between the reduction of the mismatch between observed and simulated ice thickness and the rapidity of the convergence of the spin-up method. In the present paper, the number of cycles that provides the best fit with observations (RMSE = + 63 m) is Nb_cycle = 8”.

P7.L15 You additionally need to convince the reader here that this method is optimal in the parameter choices (adjustment time 20 y, relaxation time 200 years) and to make clear in how far the results are (not) dependent on these choices.
After each step you have "a new set of initial conditions" for the next step. Maybe better to only name the final result of your initialisation your initial state as input for the forward experiments.

Completely reformulated

After 30 kyr, \( T \) is in equilibrium with the climate *and with the fixed geometry*, but not the other way around. In the next step of retuning basal drag, you further evolve the geometry and the ice temperature? Could you quantify, give an estimate how far from equilibrium you are now? Why could you not run (part of the initialisation) with freely evolving temperature?

We apologise for the confusion. Actually, there is no temperature equilibrium in the spin-up procedure used for the MAR-GRISLI experiments. However, this issue has been examined in the GMDD paper in which the 30,000-yr temperature equilibrium run appears as a sensitivity experiment. In the present paper, the temperature evolves freely at any stage of the initialisation procedure. Initial conditions inferred from the relaxation run are just restored before starting a new iteration.

What is your stopping criterium at this point and the reason for not iterating further?

Our target is to obtain the minimum ice thickness root mean square error (here \( \text{RMSE} = + 63 \) m). We stopped the iterations when the RMSE is not significantly improved (here after 8 cycles). This ensures a good compromise between the reduction of the mismatch between observed and simulated ice thickness and the rapidity of the convergence of the spin-up method. This has been clearly explained in the revised manuscript.

It is not clear why evaluation of the initial state should be based on an experiment which includes further relaxation steps. The control experiment that offers itself naturally and should be used for that purpose is just running the model after step 3 forward with constant forcing. This would give a good indication of the match with observations (at \( t=0 \) or \( t=25 \)) and the remaining model drift (after 150 years), since this is the model state actually used as initial state for the forward experiments. It anyhow seems strange to impose the observed geometry, when the model has been relaxed to a different geometry in step 3.

After the last step 2 (i.e. after the end of the 8th cycle), a 2000-yr free evolving GRISLI run is carried out under conditions identical to those used in step 2 in terms of climate forcing, initial vertical temperatures and velocity profiles. As such, the value of the basal drag coefficient is that obtained at the end of the 1st step of the 8th cycle. This has been specified in the revised manuscript.

It is a bit unusual to specify errors in ice thickness as median values, given that errors locally could be positive or negative. Why not specify the absolute error or root mean squared error augmented with the quantiles given already. A map of the mismatch with observations should be given (possibly in the appendix), but then for the model state after step 3, which is assumed as the initial state for the projections.

We agree with you that median computed from ice thickness errors with respect to observations is not always informative because of both positive and negative values. For this reason, the description of ice thickness changes has now been given as a function of different
surface elevations (see section 4.1.2) in order to aggregate regions that present similar tendencies.
As explained just above, the 2000-yr GRISLI simulation has been performed to reduce the ice volume drift. The state obtained at the end of this run is used as initial state of the forward experiments. In the revised version we mention both the new ice thickness RMSE (= + 132 m), which is different from that obtained at the end of the last step 2 (+ 63 m), the 5th and the 95th quantiles, and also the sea-level equivalent model drift (~10^-5 mm yr^-1). We added in the Supplementary Materials a figure (Fig. S1) showing the differences between the observed and the GRISLI topographies, with the GRISLI topography taken at the end of the 2000-yr relaxation run.

P8.L1 The model state that has been compared to other models in the initMIP exercise appears to be different from the state used in the forward experiments, because it includes re-imposing the observed geometry and additional relaxation for 2000 years. This should be made very clear, especially in light of the claim that the model is one of the best in the model comparison. This statement in particular requires further qualification and needs to specify what criteria to consider, since the Goelzer et al paper does not provide any explicit ranking of the models and goes into length about how different criteria for evaluating models are not independent. Please use such community efforts to improve your model, but don’t misuse them to gain credibility for your model.
To avoid confusion and misleading interpretations we removed the comparison to other models in the initMIP exercise.

P8.L5 SLR contribution as the most abstract change could be named last.
Corrected in the text following the recommendation. The new sentence reads as: “The aim of this study is to assess to what extent accounting for the atmosphere-GrIS interactions influences the GrIS evolution in terms of changes in SMB, ST, ice thickness and SLR”.

P8.L16 Is the elevation difference used for the correction calculated between Bamber (at 5 km) and Bamber (at 25 km) bi-linearly interpolated to 5 km? Please describe.
The horizontal interpolation is made using an inverse distance weighting method, as it is now specified in the revised manuscript. Moreover, to account for the differences in surface elevations between the 25 and 5 km Bamber et al. (2013) topographies, we also apply a vertical correction following Franco et al. (2012) who derived a local vertical gradient of each SMB component as a function of altitude.

P8.L21 This seems to imply that at least until 2020, NC is an appropriate approximation to the full problem. This should enter the discussion and the abstract, following an earlier comment. Is there any reason why the modification starts at 2020 and not at 2000? It would seem like a cleaner comparison to start the interaction from the moment it is possible (i.e. 2000).
For all the experiments, the “coupling” starts in 2020 when the SMB simulated by MAR at a given time is enough different from the SMB simulated at the beginning of the simulation to induce significant changes in the GrIS topography. Thus, in the PF and 2W experiments, GRISLI is forced by MAR outputs from 2005 to 2020, following the same procedure as in the NF
experiment. It would have been possible to start the coupling in 2000 or 2005, but the results would have been similar to those presented here as the SMB changes through 2005-2020 do not produce any significant topography changes in GRISLI. This has been specified in the revised abstract and in the main text.

**P8.L23** Another "coupling method" that is already discussed in the text and could be formally listed here as well is the one where MAR SMB anomalies alone are used to generate a changing ice sheet geometry (in the absence of an ice sheet model). This experiment can be performed with or without taking into account the surface elevation - SMB feedback and with or without fixed ice sheet extent.

As previously specified in our response to the General Comment”, the GrIS sea-level rise estimated from SMB integrations over fixed and time variable ice sheet masks have been discussed in Section 4.4. We decided not to present them as additional experiments since the SLR estimates are inferred from diagnostics stemming from NF and 2W experiments.

**P9.L13** This section reveals that the models are not actually fully coupled and also gives indications why a full coupling is so much more difficult to achieve. See general comment.

We agree with this comment. In the revised manuscript, we removed expressions such as “fully coupled” (see also our response to the general comment). We also explained why we have chosen the anomaly method (see section 3.2):

"Due to the topography differences between MAR and GRISLI, this approach has been chosen to avoid large inconsistencies between the SMB and ST fields computed by MAR and the ones corrected to account for the GRISLI topography."

In the discussion section (i.e. section 5), we also discussed the limits of this experiment with respect to a real fully coupling method between RCM and ISM:

“A second limitation is related to the 2000-yr relaxation GRISLI experiment, run at the end of the spin-up procedure to reduce the model drift in terms of ice volume, that produces residual differences with the observed topography (Bamber et al. 2013) used in the MAR simulations. This has important consequences on the MAR simulated climate. In particular, the steeper slopes existing in the GRISLI topography (i.e. $S_{ctrl}$) tend to produce unrealistic katabatic winds. Therefore, we choose to use an anomaly method of the surface elevation onto which the SMB and ST fields are downscaled at the 5 km resolution grid (Eq. 7). The objective of this approach was first to maintain the realism of the simulated present-day climate computed on the observed topography (Bamber et al. 2013) and, secondly, to avoid inconsistencies between the climate simulated by MAR and that used to force GRISLI. However, this implies that the forcing climate is not fully consistent with the GRISLI topography. This should be taken into consideration in a future work to improve the quality of our results. As an example, a reasonable compromise to avoid the use of anomaly method would be to use the topography obtained at the end of the spin-up iterative process (rather than $S_{ctrl}$) as initial GRISLI topography to keep the mismatch with the observed topography as low as possible, and to initialise and perform MAR simulations with this spin-up topography”
P10.L2 The mean decrease in SMB explains the shift in the ELA not the other way around. The ELA is an abstract concept, the SMB change is ‘real’.
The sentence has been modified as:
“The equilibrium line altitude (ELA, i.e. altitude for which SMB = 0) increases significantly between the beginning and the end of the 2W experiment, as a consequence of increased runoff for areas below 2000 m.”

P10.L5 I am not sure reporting the changes in ice thickness changes as mean and standard deviation makes much sense, given the bipolar nature of thickening in the centre and thinning at the margins. More useful would be for me to describe the changes for specific regions.
In order to clarify the text, the description of ice thickness changes has been given as a function of different surface elevation (see section 4.1.2). Besides the mean ice thickness anomaly values and the corresponding standard deviations, we have also reported the 5th and the 95th percentiles to indicate the range of ice thickness changes:
“The ice thickness anomaly (Fig. 4) also presents two distinct patterns. For surface elevations higher than 2000 m in the northern part, and higher than 2500 m in the central and southern parts of the ice sheet, the ice thickness increases by +5 m on average, with the increase ranging from +1.5 m (5th percentile) to +17 m (95th percentile). On the other hand, in regions whose surface elevation is lower than 2000 m, the ice thickness decreases from -248 m (5th percentile) to -3 m (95th percentile) with a mean value equal to -100 m”.

P10.L12 What exactly is the impact of ice temperature on ice dynamics? Are you implying that changes of the surface boundary conditions modify the temperature structure of the ice and its deformation?
This was a shortcut. Actually, as the model was forced by a warming scenario (i.e. the RCP8.5 and extension of year 2095 to 2150), we simplified by “warming scenario” instead of simply explaining that the ice dynamics was also impacted (in addition to ice thickness). In the revised manuscript, changes in ice velocities are related to changes in ice thickness. As a result, the new sentence has been changed in:
“The ice dynamics is also impacted by changes ice sheet geometry as illustrated by the mean surface velocity anomaly (Fig. 6a)”.

P10.L13 Are you talking about velocity or velocity anomalies here? Figure 4A shows anomalies! Please clarify.
We are talking about surface velocity anomaly. We have clarified the text (see section 4.1.3).

P10.L15 This statement calls for a figure comparing modelled and observed velocities! Add a panel to substantiate this point.
The panel showing the observed surface velocities has been added (see Fig. 7 in the revised manuscript).

P10.L18 Add "in this area" after "ice velocities" and remove it in the sentence after.
The comments related to the new figure 7 have been reorganized and the sentence you refer to has been removed from the revised manuscript. The examples of the Jakobshavn and the Kangerlussuaq glaciers are now distinguished, and the new paragraph (section 4.1.3) reads as:

“For the Jakobshavn glacier, and for altitudes above 1500 m, the vertically-averaged ice velocities increase by more than 15 m yr⁻¹ (i.e. +10 %) as a result of increasing surface slopes, and slow down by more than 200 m yr⁻¹ (i.e. +29 %) for altitudes below 1000 m due to the decreasing ice thickness (Fig. 7c). For altitudes above 500 m, the vertically-averaged velocity is mainly driven by the SIA velocity (Figs. 7c-e). On the contrary, below 500 m, basal sliding velocities are large due to low basal drag coefficient (see Fig. 3 in Le clec’h et al., 2018) and the SSA velocity component dominates the ice flow (Figs 7c-g). However, while basal drag is lower in locations below 500 m, the ice flow is limited by the strongly reduced ice thickness (Fig. 4)”.

“The Kangerlussuaq glacier is located in regions where the bedrock is characterised by a succession of valleys surrounded by mountains merging in a canyon where the deepest part is located 100 km away from the coast (Morlighem et al., 2017). The ice flow of the Kangerlussuaq is therefore divided in different branches with increasing ice velocities towards the ice sheet margin and becoming even larger when merging in the canyon (Fig. 7b). As for the Jakobshavn glacier, the ice flow accelerates at the end of the 2W experiment as a consequence of the increase in surface slope for high altitudes (~2000-2500 m, see Fig. 4). Conversely, a strong decrease of the ice flow is found in most of margin regions (Fig. 7d) directly related to the ice thinning (Fig. 4). Contrary to the case of the Jakobshavn glacier that presents large basal sliding velocities only below 500 m, the Kangerlussuaq shows low basal drag coefficients in the entire glacier (see Fig. 3 in Le clec’h et al. 2018) and thus the ice flow is mainly governed by the SSA component (Fig. 7h)”.

P10.L30 "amplification of all the changes" is a bit too general here. Better "amplification of the changes".
This has been reformulated. Following the recommendations of Reviewer 2, the Results section (Section 4) to emphasize the 2W experiments. As a consequence Section 4 has been reorganised and the changes occurring in 2100 are now discussed in Section 4.4

P10.L31 A figure showing the absolute sea-level changes for the different experiments would be in place, possible as additional panel in figure 7.
Following your suggestion, we added a figure showing the absolute sea-level changes (Fig. 12a in the revised manuscript). We also made a zoom-figure displaying the sea-level anomalies between 2000 and 2100 to better illustrate the divergence of the three experiments as soon as 2025-2030 (Fig. 12b).

P11.L3 ST is already defined
Thanks for this remark. ST is now defined once, in Section 1.

P11.L4 Replace "is strongly colder" by "sees a strong cooling" or similar.
This has been corrected

P11.L10 "Thus, the *stronger* ST decrease in 2-W compared to NC ...", assuming there is decrease in both cases. To check also in other places that you discuss differences in changes, not changes itself.
We made our best to remove all ambiguities related to changes and differences in changes. We hope the text is now clearer.

P11.L10 Not sure where "the middle of the slope is". Clarify!
We replaced “middle of the slope by “along the slope”

P11.L14 Costal regions don’t exist inland from the ice edge.
We reformulated in “in the interior of the GrIS”.

P11.L28 Replace "SMB anomalies increases by a factor of 10" by "SMB anomalies decreases by 10 cm yr$^{-1}$"
This sentence was referred to Table 1 which does no longer appears in the revised manuscript.
The new Table 1 provides values of the GrIS contribution to sea-level rise in 2050, 2100 and 2150 for the three experiments. Moreover, the section describing the SMB differences between the 2W and the NF experiments has been re-written (see Section 4.2.1)

P12.L1 Again mixing discussion of surface elevation and ice thickness here. Revise.
This has been revised and corrected in the entire revised manuscript

P12.L1 Add "difference" after "surface elevation change" and reformulate to "follow the patterns of SMB anomaly differences (Fig. 6B)".
Replaced by: “The ice thickness anomaly pattern is essentially mimicking the SMB differences between 2W and NF (Fig. 8a)”

P12.L5 Do you mean lower surface temperature in 2-W is the cause for higher SMB and therefore increasing ice thickness, or is the lower surface temperature directly impacting ice thickness (i.e. not through its effect on SMB)? In the first case, lower surface temperature and its effect on SMB should be mentioned first and higher SMB as a consequence. More precision needed here.
This part of the text has been clarified and precisions have been added:
“The main SMB differences between both experiments, averaged over the 2140-2150 period, highlight lower SMB values in 2W compared to NF for altitudes below 2000 m, with the exception of some margin locations in the eastern part (Fig. 8a). This SMB anomaly behaviour is driven by a snowfall reduction in low altitude areas (Fig. S6) and by the runoff increase in 2W with respect to NF (Fig. S7). This increased runoff results from warmer temperatures over the whole GrIS (up to 0.8°C in the western and northern parts, Fig. 8b), except in the region at the edge of the GrIS, which sees a strong cooling (as low as -10°C, Fig. 8b). The warming can be explained by the temperature-altitude feedback being active in 2W, resulting in lower altitudes
(section 4.1.2 and Fig. 8c) and therefore warmer temperatures. The cooling over the very edge of the ice sheet occurs despite the ice sheet thinning over these regions. It can be explained by changes in atmospheric circulation”.

P12.L6 "in areas of lower ST"? In my eyes, ST and ST differences (Fig. 6A) are both high and positive in regions of negative thickness anomaly. Clarify that statement.
The ST differences between the 2W and the NF experiments are positive in most of GrIS areas. However, at the very edge of the ice sheet, these differences are negative, showing that the 2W surface temperatures are lower than the NF ones as a result of the effect of katabatic winds. In the revised manuscript we changed the ST color scale, making the negative ST differences more visible.

P12.L11 I thought you are trying to describe here the impact on the ice thickness evolution of two-way coupling as opposed to no coupling. In this part, you however come to the impact on the atmospheric circulation (katabatic winds) and land model changes (albedo). From line 15 on, you go back again to ice dynamic changes. Could this material be better organised to avoid jumping between the different aspects?
Following your suggestion, this part has been reorganized. In this section, we only emphasize the effect of katabatic winds. The reduction of the ice-sheet extent simulated in the 2W experiment (and thus the effect on albedo changes) is now discussed in section 4.4.

Also, if I understand correctly, the anomalous katabatic winds created by 2-W have visible impact mainly on the narrow marginal areas of the ice sheet where anomalous cooling increases SMB. This should then be counteracted by the albedo changes described L12 and following. It is not really resolved for me how these different factors influence each other and which is the dominant mechanism in which region.
Taking into account the effect of katabatic winds leads to a cooling in 2W with respect to NF at the very edge of the ice sheet. Since the 2W temperature is lower than the NF one the predominant effect is that induced by the katabatic winds, not by albedo changes.

P12.L13 What is the difference between snow-free and snow free?
Sorry, this was a typo error

P12.L22 Melting itself does not necessarily contribute to SLR since melt water can be refrozen in the snow pack. Better replace "melting contribution" by "ice sheet contribution" or similar, also in the rest of the manuscript.
We followed your suggestion and the occurrences of “melting contribution” have been replaced by “GrIS contribution”

P12.L22 These numbers should be calculated against a control experiment to remove the contribution from remaining model drift. Has this been done?
Yes; it has. Our control experiment is the 2000-yr GRISLI relaxation run. As specified in the revised manuscript, the remaining model drift in terms of ice volume is only $10^5$ mm yr$^{-1}$, fully negligible with respect the GrIS contribution to sea-level rise.

P12.L23 I would suggest to add a panel to figure 7 with the total contributions for the three experiments and include the integrated SMB mentioned further below in this section. This has been done. Figure 7 has become Figure 12.

P12.L25 Since you discuss 2-W against NC, the surface elevation - SMB feedback which operates all over the ice sheet should also be mentioned, not just the processes at the margin. We paid attention to describe through the entire revised manuscript the processes operating in the interior of the ice sheet. In particular, the results are most often presented as a function of surface elevation: we distinguished regions of low to medium altitudes from regions of high altitudes (See in particular Section 4.1 for the description of the results inferred from the 2W method and Section 4.2 for the effects of the katabatic winds which strongly differ from central regions to margin areas.

P12.L26 The difference of 52400 km$^2$ is at the end of the experiment and then it increases with time? Reformulate
As mentionned above, absolute changes in ice-sheet extent are no longer discussed in the revised manuscript. Now, this aspect is only addressed in terms of relative changes between NF vs 2W and PF vs. 2W. This has been therefore reformulated in: “Compared to the NF and the PF experiments for which the ice-sheet mask is fixed to observations from 2000 to AD 2150, the 2W ice sheet extent is reduced by ~2.8 % in 2150 as a result of increased ablation”.

P12.L27 I think all you are saying is that the high resolution ISM mask changes are translated to partial mask changes for MAR. Clarify that the ice sheet mask (Fig S2B) is the one seen by MAR. This part has been re-written:
“Compared to the NF and the PF experiments for which the ice-sheet mask is fixed to observations from 2000 to AD 2150, the 2W ice sheet extent is reduced by ~ 2.8 % in 2150 as a result of increased ablation. As MAR sees the ice sheet retreating over time in 2W concomitantly with the increase in bare ground or tundra fractions (Fig. S5b), the albedo feedback takes place favouring further the ice melting. Although the ice sheet retreats, the extent of the ablation zone increases with time. This process is faster in 2W than in NF and PF. In 2150, the ablation zone is 14 % (resp. 11.7 %) larger in 2W than in NF (resp. PF) causing 112 Gt yr$^{-1}$ of extra ice ablation in 2W (w.r.t NF). As a consequence, the ELA is located further inland in 2W compared to NF with a maximum inland retreat of 120 km located in northeastern Greenland (Fig. 3).”

P12.L31 I think the point to make here is not about increase in uncertainty. You can show that when a fixed mask is used, you simply get the wrong result and overestimate the mass loss. Could you quantify the relative importance of this effect compared to the error that is made when not taking into account the surface elevation - SMB feedback? In Section 4.4, we quantified 1/the error made when the SMB-elevation feedbacks are ignored (i.e. 7.6%, deduced from the comparison between the SLR contributions in NF and PF
experiments) 2/ the error made when all the feedbacks are ignored (i.e. 9.3 %, deduced from the comparison between the SLR contributions from NF and 2W) and 3/ the error made when using a fixed ice-sheet mask (i.e. 6 %). To follow the suggestion of the reviewer, we added the following sentence at the end of the section:
“[...] compared to a time variable ice-sheet mask, the use of a fixed ice-sheet mask overestimates the sea-level rise by ~6 % in 2150. Though a bit lower, this number is far from being negligible compared to the errors made when the SMB-elevation feedbacks are not taken into account (i.e. 7.6 %) and when all the feedbacks are ignored (i.e. 9.3 %)”.

P13.L3 Please specify the resulting SLR.
The resulting sea-level rises obtained with the integrated-SMB methods have been explicitly mentioned in the text of the revised manuscript (Section 4.4) and reported in the new Table 2.

P13.L14 This is exactly the reason why median results are not very meaningful in this context. Mean absolute or root mean squared differences are easier to interpret.
In the revised version, we no longer mention the median values. We provide instead the ice thickness root mean square errors as well as the 5th and the 95th percentiles in ice thickness differences for regions showing similar patterns (margins vs. interior).

P13.L15 After showing figure 8, figure 9 does not add substantial information in my view. I would remove it and continue discussion about differences between 2W and 1W based on figure 8. The only reason to show figure 9 would be if you wanted to attempt modifying the parameterisation used in 1W to incorporate the katabatic wind effect, which could be a logical next step.
We acknowledge that part of the information provided in Figure 9 (Figure 11 in the revised manuscript) can be found in Figure 8 (Figure 10 in the revised manuscript). However, we believe that the new Figure 11 better illustrates the differences between the three experiments in the simulated spatial variability as a function of the altitude. This is why we finally kept this figure in the revised paper.

P13.L27 These sentences are just stating the obvious. I’d suggest to remove them.
We agree with this statement: the sentences have been removed.

P14.L5 It is not clear to me why a higher resolution should lead to increase the SLR and not the opposite. Unless there are convincing arguments to support that claim, I would leave the sign of the change open.
The reviewer is totally right. We recognize this was an overstatement. In the revised manuscript, the Discussion section has been extended and we better explained the possible influence of outlet glaciers on projected SLR. We also mention the possible decreasing influence of the outlet glacier dynamics with time:
“Regarding the ice-sheet model, a 5 km horizontal resolution does not permit to capture the complex ice flow patterns of smallest outlet glaciers, whose characteristic length scale can be less than 1 km (Aschwanden et al., 2016) and to quantify accurately the ice discharge at the marine front. This may have large implications in the sea-level rise estimates. Using a 3D ice-
sheet model with prescribed outlet glacier retreat, Goelzer et al. (2013) found an additional SLR contribution from outlet glaciers of 0.8 to 1.8 cm in 2100 and 1.3 to 3.8 cm in 2200, with the influence of their dynamics on SLR projections decreasing with time and with the increasing importance of the atmospheric forcing. This is in line with the fact that ice dynamics act to counteract ice loss from surface melting (see Section 4.2), as previously outlined by several authors (Edwards et al., 2014a, Goelzer et al., 2013, Huybrechts and de Wolde, 1999). However, despite the possible decreasing influence of marine terminating glaciers, at the centennial time scale, it seems to be preferable to evaluate more accurately the impact of ice dynamics and to better capture the complex geometry of fjords surrounding the marine-terminating glaciers”.

The same applies to the limitation of constant basal drag in the next sentence. With all the complexities surrounding the evolution of the basal conditions over time, I don’t think there is any evidence that acceleration of ice flow has to be the dominant response. Again, putting forward some convincing arguments would be appreciated. Again, we fully agree with the reviewer and this sentence has been removed from the revised manuscript. In particular, we discussed the limitations related to the time constant basal drag coefficient and to the lack of any infiltration scheme in our ice-sheet model:

“Our spin-up method adjusts the basal drag coefficient in such a way that the departure between the observed and the initial GRISLI topographies is reduced. The resulting coefficient is spatially varying but is constant in time. This assumption may likely be valid for short-term forward simulations but is probably overly simplistic. On the one hand, the basal drag tends to be smaller towards the margins with respect to the interior. As the ice sheet retreats inland, it can be expected a reduction in basal drag for a specific location, due for example to a decreasing effective pressure. On the other hand, changing basal hydrological conditions can also alter the basal drag. This can occur as a result of rainfall or surface meltwater infiltration that can re-freeze at depth or propagate all the way to the bottom of the ice sheet and increase basal lubrication (Kulessa et al., 2017). Therefore, a time constant basal drag coefficient inferred under present-day conditions may underestimate the ice flow acceleration. A few models describing the vertical inflow exist (e.g. Banwell et al., 2016, Clason et al., 2015; Koziol et al., 2017) but are generally run at the regional scale and at very high spatial resolution (a few tens to a few hundreds of meters at most). Implementing such models in large-scale ice-sheet models is currently outside the realm of possibilities. However, as there is a growing interest in performing ice-sheet projections over multi-centennial time scale, the GRISLI-like models would undoubtedly benefit from the implementation of simplified infiltration schemes (e.g. Goelzer et al., 2013) so as to account for the impact of ongoing changes in surface meltwater on ice dynamics”.

P14.L19 Additional limitations that should be discussed: - Ignoring the glacial-interglacial signature of past climate changes in this steady state spin-up of temperature typically makes the ice too warm. This needs to be compensated by other factors (likely a different set of basal drag parameters). - The steady state initialisation also ignores any influence of transients in the
observed ice sheet evolution – Mismatch of the modelled ice sheet geometry and velocity structure with observations leads to uncertainties in the projected evolution.

We added the following paragraph in the Discussion section (Section 5):

“An additional limitation related to the choice of our spin-up procedure is that the glacial-interglacial signature of past climatic changes is ignored. Neglecting the climate history of the Greenland ice sheet implies too warm ice temperatures. This may have an impact on the future GrIS evolution and on its contribution to sea-level rise. Indeed, the basal drag coefficient inferred from the inverse method may be too high so as to compensate the errors induced by the artificial warm bias. However, using a higher-order ice flow model, Seroussi et al. (2013) showed that at the centennial time scale the basal conditions and the GrIS projections are only poorly sensitive to the initial vertical temperature profile but are critically dependent on atmospheric conditions.”

P14.L28 Add "in this comparison" after "atmosphere-GrIS feedbacks". I hope you don’t think this statement is universally true.

Due to the huge changes made to the original text, this issue has been presented differently and mentioned in Section 3.3 devoted to the description of the 2W experiment: “Compared to the NF and PF approaches, this two-way coupled method is the most accurate to represent the GrIS-atmosphere feedbacks”.

P14.L30 While this statement seems true for the given results, the conclusion hinges on the change in behaviour of 2W at 2110. Unless investigated in more detail, it cannot be excluded that such change could happen at an earlier point in time, e.g. for a different model used as boundary condition to MAR.

In the new version, we showed that the results from the three experiments start to diverge from each other as soon as 2025-2030, that is a few years only after the start of the coupling. This means that the feedbacks that are accounted for in the PF (SMB-elevation feedbacks) or in the 2W experiment start to operate as early as this period. However, we also explain that the influence of the feedbacks increases over time and that they become dominant at the end of the 21st century (See in particular Section 4.4). Moreover, in the Discussion section, we clearly explain the possible dependence of our results with the GCM forcing used to force MAR: “Whatever the experimental design, the large spread in SLR projections raises the question as to whether the ice-sheet response simulated in our 2W experiment relative to that of the NF and PF experiments would be similar, amplified or mitigated with a different GCM climate forcing having a different sensitivity from MIROC5. [...]There is therefore a strong need for iterating the present study with different global climate simulations run under an extended RCP8.5 scenario and used as a MAR forcing, to assess more accurately the impact of the different GrIS-atmosphere feedbacks and to better evaluate the uncertainty associated with the projected sea-level rise contribution from GrIS”.

P15.L5 This comparison is a bit awkward. Wouldn’t it be more appropriate to compare the +0.5 cm to the total projected SLR as a relative error?

We acknowledge that this comparison was not fully appropriate and we removed it from the revised manuscript.
P15.L14 Remove repeated "respectively" after SLR.
The sentence has been changed.

P15.L19 It would be good to additionally put this number (21%) in perspective to the underestimation due to ignoring feedbacks, i.e. the difference between 2W and NC. This comparison has been done (see our response to comment P12.L31).

P15.L24 Again, I don’t see any evidence for the interpretation that higher resolution and higher order physics increase the response.
We agree with you. This comparison was not appropriate (See our response to comment P14.L15).

P15.L29 Replace "disrupt" by "modify"
OK, all occurrences of “disrupt” have been removed.

Table 1 Does "after 50 yrs" mean at year 2050? Maybe that would be a better indication. Or do you not want to assign an absolute date to your simulation? The historic and future RCP forcing is clearly linked to an absolute date, though. Since the ablation area changes so much, it may be interesting to calculate additional diagnostics for a constant region, e.g. for the observed present day ablation zone, or backwards for the area of the ablation zone area after 150 years. This way, the convolution with a changing area could be avoided.
In the revised version, most results have been discussed as a function of altitude. We mainly distinguished two type of areas: areas of high altitude (generally higher than 2000 or 2500 m) and areas of low to medium altitude (< 1000 or 1500 m). As a result it does no longer make sense to present SMB and ST values (Table 1) or GrIS thickness or ice velocities (Table 2) at the scale of the whole ice sheet, as it was done in the first version of the paper. Moreover, the results computed at the ice-sheet scale are not really informative because of the large spatial variability in the 2W-NF anomaly. In addition, the changes in ablation area and in ice-sheet extent have been discussed in the revised paper in terms of relative changes (see Section 4.4). To our opinion, they don’t need to appear in a table. We therefore removed both the former Tables 1 and 2 from this new version of the manuscript. However, we replaced these tables by new ones providing the GrIS contribution to sea-level rise inferred from NF, PW and 2W experiments (new Table 1) and from the SMB-integrated method (Table 2). We also replaced “after 50 years” by the absolute dates.

Table 2 Not sure how to interpret a velocity change of e.g. -3.0+25.0. The noise being much larger than the signal, is the valid interpretation ’no significant’ change?
We fully agree. See our previous response to comment related to Table 1.

Figures ———

The labels in the figures are upper case (A,B,C), but the panel references in the captions are all in lower case (a,b,c). Make consistent.
OK, the labels in the figures and in the captions are now identical

*Figure 2* Why are figures B and C so different? At least in the interior, one would expect a pattern very similar to the SMB anomalies in this experiment. My guess is that this is indicative of a remaining model drift. Results of a control experiment starting after step 3 with constant forcing should be shown here or in the appendix and the origin of this difference should be discussed.

To address the comments raised by Reviewer 2, the organization of the paper has been modified. We now start the result section with a thorough analysis of the GrIS evolution simulated with the most comprehensive method, i.e. the two-way (2W) coupling. The NF results are only presented in terms of differences with the 2W method. Moreover, we think that the differences you mentioned between both plots can be attributed to the choice of the color scale. The new figures (Figs 2a and 4) present similar patterns for SMB and ice thickness anomalies simulated in the 2W experiment. However, the discussion requested to explain the differences between the SMB and the ice thickness patterns has been provided for the 2W method (see Figs 2a, 4 and 5 in the revised manuscript and section 4.1.3). These differences are explained by the ice dynamics. In Section 4.1.3, we added the following paragraph to support this argument:

“The ice thickness anomaly is due to the complex combination of changes in surface atmospheric conditions (SMB, Fig. 5a), ice dynamics (ice flux divergence, Fig. 5b) and basal melting (not shown), following the continuity equation (Eq. 2). To quantify the role of ice dynamics on the GrIS geometry (Fig. 4), we plotted the ice flux divergence integrated over 150 years (2000-2150, see Fig. 5b). In particular, over the central plateau, the cumulated SMB (Fig. 5a) reaches about +50 m, 40 m of which are transported away by the ice dynamics (Fig. 5b). As a result, the ice thickness anomaly is reduced to only ~10 m in this region (Fig. 4). An opposite behaviour is found near the western coast, where the ice melting is partly compensated by ice convergence, resulting in a less negative ice thickness anomaly than that related to the SMB forcing. This shows that ice dynamics act to counteract ice loss from surface melting, as previously noticed by several authors (Huybrechts and de Wolde, 1999, Goelzer et al., 2013, Edwards et al., 2014b). As a consequence, it appears to be essential to account for ice dynamics to estimate accurately the mass balance of the whole ice sheet”.

*Figure 3* The displayed field is ice thickness, not surface elevation as written in the caption. Since the discussion is about ELA and surface elevation - SMB feedback, it may be useful to show surface elevation instead.

Figure 3 has been modified. Now it displays the surface elevation.

*Figure 4* The colour scale in A is not easy to read with small positive and small negative values sharing the exact same colour (green). This should be improved. Have you tried to plot velocity ratios instead of anomalies? Since velocity magnitudes cover several orders of magnitude, a large relative change is not visible because of the cutoff at 2 myr⁻¹, while small relative changes at the margin appear exaggerated.

The surface velocity anomalies are now plotted in the new figure 6: for the 2W experiment (instead of NF) between the end (2140-2150) and the beginning (2000-2010) of the simulation and for the anomalies between 2W and NF and between 2W and PF at the end of the
simulation. The colour scale (in log10) has been also modified to better illustrate positive and negative changes.

Figure 5 Why not use the same colour mapping here and in figure 4 for the velocity anomalies? That would make it easier to compare the two figures. Figure 5 (in the first manuscript) is now Figure 7 and uses the same colour scale as Fig. 4.

Caption: "left panel" The figure caption has been modified

Figure 6 There appears to be a slight instability in one or both of the experiments compared in figure 6A. Also Figure 8A shows signs of instability in form of a checker board pattern. While these instabilities are likely not critical for the interpretation of the large scale results presented here, they should at least be mentioned. Figures 6A and 8A are now Figures 8b and 10b. The features the reviewer refers to are related to the method used to correct for the altitude difference between the MAR and the GRISLI topographies.

Figure 7 Add a panel with absolute contributions of the three experiments. Note that results shown so far are double differences, i.e. differences in anomalous contributions since year 2000 between different experiments. Could also show sea-level contribution differences calculated from difference to a control experiment with constant forcing to remove the model drift. Same consideration holds for the absolute contributions. Figure 7 appears now as Figure 12. We added a panel showing the absolute contributions (Fig. 12a). In all the simulations presented here, the model drift has been taken into account but is fully negligible (see our response to comment P7.L32).

There seems to be a step change around 2060 and again around 2110, where the behaviour of \(2w-1w\) (yellow) changes dramatically. By comparison with \(2w-nc\) it appears to be caused by the evolution of \(2W\). What is happening at these moments in \(2w\)? Please investigate this further. The figure showing the anomalies of sea-level contributions has been replaced by a zoom-figure. Thanks to this new panel, we show that the three experiments start to diverge from each other as soon as 2025-2030 and not only around 2060. We do not observe any significant change in slope in 2060 nor in 2110.

Caption: "Differences in Greenland ice sheet sea-level contribution between the different experiments." Then explain how it is calculated. We provided further details in the figure 12 caption.

Figure 9 is not needed in my estimation. See our response to comment P13.L15. Note also that Figure 9 now appears as Figure 11.

References:

Format of many references in the text are non-standard. A few examples are given
here, but all should be re-checked.
The format of the references has been re-checked.

P3.L13 add e.g. before Gagliardini
OK, added

P3.L19 reformat list of reference and avoid double brackets
P4.L10 "(e.g. Fettweis et al., 2013
OK, reformatted

P6.L12 Author is called Fox Maule. Check reference.
OK, modified

P6.L20 Reference Le clec’h et al. (in prep) is not in the reference list.
This reference now appears as Le clec’h et al. (2018).

P8.L2 add Goelzer et al., 2017 to the reference list.
OK, added in the reference list

References:
The reference has been corrected
Assessment of the Greenland ice sheet - atmosphere feedbacks for the next century with a regional atmospheric model fully-coupled to an ice sheet model

Sebastien Le clec’h1,2, Sylvie Charbit1, Aurelien Quiquet1, Xavier Fettweis2, Christophe Dumas1, Masa Kageyama1, Coraline Wyard2, and Catherine Ritz3

1Laboratoire des sciences du climat et de l’environnement, Gif-sur-Yvette, FR
2Earth System Science and Department Geografie, Vrije Universiteit Brussel, Brussels, Belgium
3Laboratory of Climatology, Department of Geography, University of Liège, Liège, Belgium
4Institut des Géosciences de l’Environnement, Université Grenoble-Alpes, CNRS, 38000 Grenoble, France

Correspondence to: Sebastien Le clec’h (sebastien.le.clech@vub.be)

Abstract. In the context of global warming, the projected Greenland sea-level rise contribution is mainly controlled by the interactions between a growing attention is paid to the evolution of the Greenland ice sheet (GrIS) and the atmosphere, in particular through the temperature and surface mass balance—elevation feedback. In order to evaluate the importance of these feedbacks, we used three methods to represent the interactions between the GrIS model GRISLI and the polar regional atmosphere model its contribution to sea-level rise at the centennial time scale. Atmosphere-GrIS interactions, such as the temperature-elevation and the albedo feedbacks have the potential to modify the surface energy balance and thus to impact the GrIS surface mass balance (SMB). In turn, changes in the geometrical features of the ice sheet may alter both the climate and the ice dynamics governing the ice sheet evolution. However, changes in ice sheet geometry are generally not explicitly accounted for when simulating atmospheric changes over the Greenland ice sheet in the future. To account for ice sheet-climate interactions, we developed the first two-way synchronously coupled model between a regional atmospheric model (MAR) and a 3D ice-sheet model (GRISLI). Using this novel model, we simulate the ice sheet evolution from 2000 to 2150 under a prolonged RCP8.5 scenario. Changes in surface elevation and ice-sheet extent simulated by GRISLI have a direct impact on the climate simulated by MAR, under the RCP 8.5 scenario from 2020 to 2150. In the simplest method, there is no coupling: MAR computes varying atmospheric conditions using a constant GrIS geometry (topography and extent) set to observations and GRISLI is forced by these results. The second is a one-way coupling method which represents the interactions by correcting offline the MAR outputs to account for topography changes computed by GRISLI. The third method is a fully and are fed to MAR from 2020 onwards, i.e. when changes in SMB produce significant topography changes in GRISLI. We further assess the importance of the atmosphere-ice sheet feedbacks through the comparison of the two-way coupling coupled experiment with two other simulations based on simpler coupling strategies: i) a one-way coupling with no consideration of any change in ice sheet geometry; ii) an alternative one-way coupling in which the elevation changes feedbacks are parameterised in the ice-sheet model (from 2020 onwards) without taking into account the changes in ice-sheet topography in which the ice-sheet topography and extent seen by the atmospheric model evolve after each ice-sheet model time step. Due to the ice sheet elevation feedback.
The two-way coupling method amplifies the projected decrease in surface mass balance, coupled experiment simulates an important increase in surface melt below 2000 m of elevation resulting in an important SMB reduction by 2150 and a shift of the equilibrium line towards elevations as high as 2500 m despite a slight increase in SMB over the central plateau due to enhanced snowfall. In relation with these SMB changes, modifications of ice-sheet geometry tend to favour ice flux convergence towards the margins, with an increase in ice velocities in the GrIS interior due to increased surface slopes and a decrease in ice velocities at the margins due to decreasing ice thickness. This convergence tends to reduce the SMB signal in these areas. In the two-way coupling, the increase in surface temperature and the GrIS surface thinning for the coastal regions, compared to the no coupling method. Compared to both the one-way and the no coupling methods, the two-way coupling allows the changes of fine scale processes to be represented. SMB is also influenced by changes in fine scale atmospheric dynamical processes, such as the increase in katabatic winds over the coast. As a consequence, in 2150, the two-way coupling method computes a GrIS melting from central to marginal regions induced by increased surface slopes. Altogether, the GrIS contribution to sea level rise 9.3% larger than the no coupling method, and inferred from ice volume variations is equal to 20.4 cm in 2150. Our results suggest that the effect of the different feedbacks is amplified over time with the most important feedbacks being the SMB-elevation feedbacks. As a result, the experiment with parameterised SMB-elevation feedback provides a sea-level contribution from GrIS in 2150 only 2.5% larger than the one-way coupling methods. After 150 years, the GrIS extent seen by MAR in the lower than the two-way method is 52400 km$^2$ lower than with the no coupling method. Furthermore, in 2150, using a fix ice sheet mask, as in the no coupling method, overestimates by 24 coupled experiment, while the experiment with no feedback is 9.3% the SLR contribution from SMB lower. The change in the ablation area in the two-way coupled experiment is much larger than those provided by the two simplest methods, with an underestimation of 11.7% (resp. 14%) with parameterised feedbacks (resp. no feedback). In addition, we quantify that computing the GrIS contribution to sea level rise from SMB changes only over a fixed ice-sheet mask leads to an overestimation of ice loss of at least 6% compared to the use of the ice sheet mask as in the no coupling method, overestimates by 24 coupled experiment, while the experiment with no feedback is 9.3% the SMB-elevation feedback. Finally, our results suggest that ice loss estimations are diverging when using the different coupling strategies, with differences from the two-way method. Beyond the century time scale, a two-way method becomes necessary in order to avoid an underestimation of the projected ice sheet volume, topography and ice extent reduction. The one-way coupling method however seems to be sufficient to represent the interactions for projections until the becoming significant at the end of the 21st century. The no coupling method always underestimates the projected ice sheet volume loss significantly due to the lack of feedback between the GrIS and the atmosphere. In particular, even if, averaged over the whole GrIS, the climatic and ice-sheet fields are relatively similar, at the local and regional scale there are important differences, highlighting the importance of correctly representing the interactions when interested in basin scale changes.

1 Introduction

The Arctic is the region of the Earth experiencing the largest increase in temperature since the pre-industrial era (Serreze and Barry, 2011), with consequences already perceptible on the mass evolution of the polar ice caps and the Greenland ice sheet (Rignot et al., 2011). The atmospheric conditions control the variability of the near-surface temperature (ST) and the evolution
of the Greenland ice sheet (GrIS) is governed by variations of ice dynamics and surface mass balance (SMB) of the GrIS. The SMB represents, the latter being defined as the difference between snow accumulation, which is further transformed into ice, and the ablation, which are processes of ice loss. Accumulation and ablation are both sensitive to ST. Variations of ST and SMB-ablation processes (i.e., surface melting and sublimation). While surface melting strongly depends on the surface energy balance, snowfall is primarily controlled by atmospheric conditions (wind, humidity content, cloudiness...). However, various feedbacks between the atmosphere and the GrIS may lead to SMB variations that can therefore directly affect the GrIS total ice mass by impacting its characteristics such as thickness, ice volume and ice extent. In turn, variations of the GrIS characteristics affect the ST and the zones of ablation (SMB < 0) and accumulation (SMB > 0). GrIS changes can also disrupt surface characteristics, such as ice extent and thickness, with potential consequences on ice dynamics (e.g., due to changes in surface slopes). These changes may in turn alter both local and global climate. As an example, changes in near-surface temperature and surface energy balance may occur in response to changes in orography (temperature-elevation feedback) or in ice-covered area (albedo feedback; see Vizcaíno et al., 2008; Vizcaíno et al., 2015 and Lunt et al., 2004). On the other hand, topography changes may alter the atmospheric circulation over Greenland caused by changes in topography-thermal contrast between ice sheet surface and atmosphere layers, surface albedo and ice sheet area, as shown by Vizcaíno et al. (2015; Vizcaíno et al., 2008; and Lunt et al., 2004), patterns causing changes in heat and humidity transports.

Quantifying the balance between these different processes and feedbacks that regulate transient ice sheet change is required to understand and predict more confidently the evolution of the GrIS under current and future global warming. Although numerous studies highlighted the importance of correctly representing the interactions between the GrIS topography changes and the atmosphere (Vizcaíno et al., 2015; Edwards et al., 2014a; Alley and Joughin, 2012; Huybrechts et al., 2002) (Vizcaíno et al., 2015; Edwards et al., 2014a), only few global or regional models have taken the GrIS topography changes into account to compute the future evolution of the SMB-ST and energy budget over the GrIS. The climate models usually represent the ice sheet component with a fixed and constant topography, even under a warm transient climate forcing. Recently, to explore the importance of SMB-elevation feedbacks for the future GrIS evolution, Vizcaíno et al. (2015) used an atmosphere-ocean general circulation model (AOGCM) coupled with an ice sheet (ECHAM5.2) coupled to an ice-sheet model (ISM) forced under different RCP scenarios (up to 2100) and their extensions (from 2100 to 2300). Compared to a control experiment in which the ISM is forced off-line by the atmospheric model run with the fixed present-day GrIS topography, they found an amplification of ice mass loss amplification of 8–11% (by 2100) and of and 24–31% (by 2100 and 2300). Since both their ice sheet and climate models have a relatively coarse resolution (3.75° for the atmospheric component and 10 km for the ice sheet model), they focus respectively, when the elevation feedbacks are taken into account (i.e., in the coupled experiment). This results from the combination of the positive elevation-SMB feedback in low lying areas, the negative feedback related to the elevation-desertification effect in accumulation areas, and the changes of surface slopes resulting from high mass loss in ablation areas and slight snowfall increase in the accumulation zone, enhancing the ice transport from the central regions to the ice margins. Their study is focused on the added value of incorporating the coupled processes and less on exactly reproducing...
However, as specified in Vizcaíno et al. (2015), their model is not able to accurately reproduce the observed GrIS—which would require more detailed physics in both models. They explain that their study must be regarded as a necessary first step towards more advanced coupling of ice sheet and climate models at higher resolution. Indeed, a higher resolution is necessary to represent correctly—because 1/ the ice-sheet model, based on the shallow-ice approximation (Hutter, 1983), is not designed to properly represent fast ice flows in outlet glaciers and 2/ the resolution of the AGCM (∼3.75°) and the ice-sheet model (10 km) are too coarse to correctly capture the steep slopes at the ice sheet margins (typically, the altitude varies by 2000 m over distances of the order 100 km). However, the computation of atmospheric fields at a resolution similar to the ISM one margins and the atmospheric processes acting on the SMB calculation.

Using the AGCM NCAR-CAM3 run at different spatial resolutions (T21 to T85) and coupled to the SICOPOLIS ice-sheet model, proceeding how the atmospheric model resolution influences the simulated ice sheets at the Last Glacial Maximum. They found that the North American and the Eurasian ice sheets were properly reproduced with the only T85 run. According to the authors, this is likely due to the inability of the atmospheric model to properly capture the temperature and precipitation fields (used to compute the SMB) at lower horizontal resolutions, as a consequence of the poorly resolved planetary waves and smooth topography. However, running high resolution atmospheric models at the global scale requires large computing resources. Franco et al. (2012) developed an interpolation method allowing to correct each SMB. To circumvent the low resolution, some authors have used the method of elevation classes and are therefore able to offer high resolution in the direction of the slope gradient (e.g., ?).

An alternative solution consists in using regional climate models (RCM) to produce high resolution atmospheric fields and much more robust energy balance and SMB calculations. A number of RCMs have been developed for the polar regions such as MAR (Fettweis et al., 2017), RACMO2 (?), Polar MM5 (?) or HIRHAM5 (?). However, the highest resolution of the RCMs is limited by the use of the hydrostatic approximation and often remains below the resolution of Greenland ice-sheet models which are generally running at a 5-10 km scale (Buvel and Brown, 2009; ?; ?) or even below (e.g., Gagliardini et al., 2013). This means that SMB fields must be corrected for resolution (and thus for elevation) differences between the RCM and the ISM. With the aim of investigating the influence of the MAR resolution on the computed SMB fields, Franco et al. (2012) developed a method to downscale each SMB MAR component (snowfall, rainfall, runoff, sublimation and evaporation) onto a finer grid as a function of topography changes. They showed that their corrective method is able to significantly reduce the model output differences between a coarse spatial resolution model and a high resolution model. Edwards et al. (2014a) developed an alternative parametrisation of the interactions between the GrIS and the atmosphere to correct the SMB computed by a regional atmospheric model (RCM) by taking into account the GrIS topography changes computed by an ISM. This method only requires limited additional supercomputing resources. Under the SRES A1B emissions scenario, Edwards et al. (2014b) showed that a larger melting of the GrIS is obtained when the elevation feedback is taken into account (ranging from 4.4 % in 2100 to 9.6 % in 2200) of elevation changes. An alternative approach based on statistical relationships between altitude and SMB has also been derived with MAR (Edwards et al., 2014b) for regions below and above the equilibrium line altitude in the northern and southern parts of GrIS, with limited additional computing resources. However, in both parameterisations by Franco et al. (2012) and Edwards et al. (2014b), the authors only consider a strict linear relationship between topography...
and SMB changes. These previous studies show that one of the feedbacks between the GrIS and the atmosphere (Nowicki et al., 2016) is to use a high resolution atmospheric model and a detailed snow model. The higher resolution allows (Fettweis et al., 2017; ?; ?) to better represent the elevation gradients and therefore the steep topography and the extent of near the ice margins in the ablation zone. Together with the higher resolution, the use of a detailed snow model can also better estimate the GrIS surface properties, such as albedo, snow cover and surface melting. Furthermore an RCM developed for the Greenland region can represent more complex such as those implemented in MAR (Fettweis et al., 2017) or RACMO2 (?) allows a more accurate description of the surface properties (e.g., snow cover, albedo, surface melting) and therefore of the surface energy balance. RCMs developed for polar regions are also able to represent more atmospheric and land surface processes prevailing specifically over this area such as blowing snow (Gallée et al., 2001) or occurring in these regions such as bare ice albedo (Box et al., 2012) and katabatic winds (?), being also strongly dependent on topography and thus on resolution.

The second fundamental requirement to describe the interactions between atmosphere and GrIS is to represent the ice sheet topography changes in the atmospheric model by using an ISM (instead of the fixed geometry typically used) to take into account the effects of ice dynamics on the ice sheet topography changes. This can be done by using a full achieved through a numerical coupling between the RCM and the ISM. More than twenty ice sheet models exist (e.g., Ritz et al., 2001; ?; ?; Bueler and Brown, 2009; ?; ?), and are currently compared in the Ice Sheet Model Intercomparison Project (Nowicki et al., 2016)(Nowicki et al., 2016; ?). They represent thermodynamical and physical processes of the Greenland ice sheet GrIS with different levels of complexities (Gagliardini et al., 2013; Saito et al., 2016). They all compute the dynamical response of the GrIS to a given climate forcing which can be, for example, such as the SMB and the ST near surface temperature (ST) fields computed by RCMs (or global models). However, as SMB and ST from the climate models do not take into account the GrIS evolution, the climate forcing used by the ISM could be flawed.

In order to explicitly represent the feedbacks between the GrIS and the atmosphere at high resolution and to evaluate their impacts on the SMB, ST and topography changes as well as on the SLR GrIS contribution, ice-sheet evolution, we coupled the "polar" regional climate model MAR (for Modèle Atmosphérique Régional, in French, Fettweis et al., 2017)) and the high resolution ISM GRISLI (for GRenoble Ice Shelf and Land Ice Ritz et al., 2001) (Philippon et al., 2006)) and Alvarez-Solas et al. (2011a). To further investigate the (Fettweis et al., 2017) to the GRISLI ice-sheet model (Ritz et al., 2001; ?). To assess the importance of an explicit representation of the interactions between the GrIS and the atmosphere, we compared experiments using the tow way coupling method (called hereafter 2-W for two way coupling) with two other experiments using less complex methods. In the first method (referred to as NC for no coupling), the atmosphere GrIS feedbacks are not represented; in the second method (referred as 1-W for one way coupling), the feedbacks are partially represented using the Franco et al. (2012) corrective method but without any physical coupling between MAR and GRISLI.
albedo feedbacks, this coupled experiment is then compared to a one-way coupling experiment, in which the GrIS-atmosphere interactions are not taken into account, and to a second experiment where the effects of topography changes on the simulated SMB are parameterized.

A description of the atmospheric model, the ice sheet model and the climate forcing and the ice-sheet models is given in Sect. 2. Sect. 3 focuses on the description of the experimental setup of the three methods (2-W, 1-W and NC) coupling methods considered in this study. On the initialisation and the experimental set up. In Sect. 4, we describe the results of the NC experiment. Next we compare the 2-W experiment with the NC and 1-W experiments, in terms of atmospheric and ice dynamic fields and of the GrIS ice extent. Sect. 2 focuses on the limit of using the NC or the 1-W method instead of the 2-W method. We first describe the coupled experiment in detail before comparing it to the other coupling experiments. These sections are followed by the discussion and the main discussion related to the different coupling approaches (Sect. 5.1) and the conclusions of this study (Sect. 6).

2 Models and experiments

2.1 The MAR atmospheric model

MAR is a regional atmospheric model fully coupled with the land surface model SISVAT which includes a detailed snow energy balance model (Gallée and Duynkerke, 1997). It has been developed (Soil Ice Snow Vegetation Atmosphere Transfer model, see ?) which includes the detailed one-dimensional snow model Crocus (Brun et al., 1992) which simulates fluxes of mass and energy between snow layers and reproduces snow grain properties and their effect on surface albedo. MAR has been developed to simulate the GrIS SMB and has been extensively validated against in situ observations (Fettweis et al., 2017). The MAR the SMB is computed as follows:

\[
SMB = SF + RF - SU - RU
\]  

Where SF, RF, SU and RU represent snowfall, rainfall, sublimation and runoff respectively. Note that RF contributes to the SMB since liquid precipitation may percolate and refreeze at depth either in the snowpack or in the ice column.

The MAR horizontal resolution is 25 km x 25 km covering the Greenland region (6600 grid points), from 60°W to 20°W and from 85°N to 1°N. The model has 5 N, and 24 levels for the atmospheric layer from the surface to 16 km high vertical levels to describe the atmospheric column in sigma-pressure coordinates (?). SISVAT has 30 levels to represent the snowpack (with a depth of at least 200 cm over the permanent ice area) and 7 levels for the soil in the tundra area. Lateral boundary conditions can be provided either by reanalysis dataset (such as ERA-interim or NCEP) to reconstruct the recent GrIS climate (1900-2015) (Fettweis et al., 2017) or by general circulation models (GCMs) to perform future projections such as those used for the last IPCC report (e.g. Fettweis et al. (2013)).

MAR uses the solar radiation scheme of Morcrette et al. (2008). The representation of the atmospheric hydrological cycle (including a cloud microphysical model) is based on Lin et al. (1983) and Kessler (1969). The snow-ice part comes from the snowpack model Crocus (Brun et al., 1992). This 1-D model simulates fluxes of mass and energy between snow layers, and
reproduces snow grain properties and their effect on surface albedo. Each grid cell is assumed to be covered by at least 0.001% of tundra and snow. At each time step SISVAT computes the albedo of each surface type and the characteristics of the snowpack which are weighted and averaged as a function of the snow and vegetation coverage in each grid point, and then exchanged with MAR. The present work uses MAR version 3.6. The differences with previous MAR versions used in Fettweis et al. (2013, 2017) are only related to adjustments of some parameters in cloudiness: the representation of cloudiness (i.e. cloud life time) and bare ice albedo (i.e. parameterisation of the impact of melt ponds). The bare ice albedo has been improved by parametrising the melt ponds impact on the albedo.

2.1.1 Boundary conditions and climatic forcing

The topography of the GrIS as well as the surface types (ocean, tundra and permanent ice) are provided by Bamber et al. (2013). At its lateral boundaries, MAR is forced every 6 hours with 6-hourly atmospheric fields (temperature, humidity, wind and surface pressure) and surface oceanic conditions (sea surface temperature and sea ice extent) coming from reanalyses or from GCM outputs: provided either by reanalysis dataset (such as ERA-interim or NCEP) to reconstruct the recent GrIS climate (1900-2015) (Fettweis et al., 2017) or by general circulation models (GCMs) to perform future projections such as those used for the last IPCC report (e.g., Fettweis et al., 2013). As a result, the atmospheric circulation simulated by MAR over the Greenland ice sheet is strongly dependent on the quality of the climatic fields computed by GCMs or reanalyses as an input to the model. Fettweis et al. (2013) have shown that forcing GCMs which satisfactorily simulate the present-day free-atmosphere mean summer temperature at 700 hPa and the large-scale circulation over Greenland at 500 hPa–500 hPa are best suited to force MAR. For the present study we therefore choose to force MAR with the MIROC5 model output (Watanabe et al., 2010), outputs (Watanabe et al., 2010) because it has been shown by Fettweis et al. (2013), to be the best choice GCM choice from the CMIP5 database to reproduce the present-day climate with respect compared to the results of MAR forced by reanalyses compared to the other GCMs from the CMIP5 data base (Fettweis et al., 2013). The greenhouse gas forcing used in MAR (scenario RCP8.5) is the same as that used in the MIROC5 simulation (Watanabe et al., 2010). Except for the experiment presented later in this study in which MAR is coupled to an ice-sheet model, the topography of the GrIS as well as the surface types (ocean, tundra and permanent ice) are taken from the Bamber et al. (2013) dataset aggregated on the 25 km grid.

2.1.1 Model initialisation and experiment

Before starting our experiments, MAR needs to be properly initialised to limit spurious drifts, which would introduce unwanted trends in the results. The snowpack included in SISVAT requires more than 6 model implemented in SISVAT requires less than 7 years to reach an equilibrium with atmospheric fields. Here, we initialise MAR by running it: MAR is initialised with the atmospheric forcing fields from MIROC5 from 1970 until 1976 and by using the present-day GrIS geometry provided by Bamber et al. (2013). This first simulation uses a SISVAT initial state from a previous similar MAR MIROC5 simulation. In 1975 and the MAR simulations start in 1976. However, in this paper, the MAR results will be analysed for the period following year 1976. After the initialisation period, and for all experiments, MAR is forced by transient MIROC5 atmospheric fields of the
CMIP5 historical (1970-2005) and RCP8.5 scenarios (Taylor et al., 2012) until 2100. In order to extend the MAR experiment until 2150, we have repeated the MIROC5 year 2005 (representative of the years 2090s) for 50 additional years spanning from years 2000 to 2150.

2.2 The GRISLI Ice sheet model

2.2.1 Model description

The GRISLI (GRenoble Ice Shelf and Land Ice) is a coupled ISM first developed ice-sheet model was first developed to compute the dynamical evolution of the Antarctic ice sheet (Ritz et al., 2001; Philippon et al., 2006; Alvarez-Solas et al., 2011a). It has then been successfully applied to the northern hemisphere ice sheet (Peyaud et al., 2007; Alvarez-Solas et al., 2011b; Quiquet et al., 2013; Charbit et al., 2013) and the Greenland ice sheet (Quiquet et al., 2013; ?). In the present work, we use a 5 km resolution grid covering the Greenland ice sheet with 301x561 grid points. GRISLI is a three-dimensional hybrid thermo-mechanically coupled ISM computing the temporal evolution of the ice sheet, which is a function of the surface mass balance, ice flow and basal melting (Eq. 2).

\[
\frac{\partial H}{\partial t} = -\nabla(U^G H) + MSMB - b_{melt}
\]  

where \( t \) is time, \( H \) the ice thickness, \( \overline{U} \) the depth-averaged horizontal velocity, \( M \overline{U} \) the vertically-averaged velocity, SMB the surface mass balance and \( b_{melt} \) the basal melting. Basal melting occurs when the basal temperature is at the pressure melting point. The ice temperature plays a crucial role in the dynamics of the ice sheet because it also affects the viscosity, and thus the ice flow in the entire ice column (?Ritz et al., 2001). In turn, heat released by internal ice deformation and basal dragging over the bedrock modifies the temperature. The temperature field is computed by solving a time-dependent heat equation both in the ice and in the bedrock accounting for advection and vertical diffusion processes. At the surface, the boundary condition is provided by the prescribed surface temperature. At the base of the ice sheet, the boundary condition is given either by the geothermal heat flux or by the temperature melting point at the ice-bed interface.

The ice flow is governed by computed using both the shallow ice (Hutter, 1983) and shallow shelf (MacAyeal, 1989) approximations to solve the Stokes equations (Ritz et al., 2001). The SIA (shallow ice approximation) assumes that ice flow is caused only by vertical shear stress, neglecting the longitudinal stresses. This assumption is only valid for slow flowing ice. For fast flowing regions, vertical shearing becomes smaller than longitudinal shearing and the SSA assumption (shallow-shelf approximation), which neglects the vertical stresses, is used. The SIA component of the computed ice sheet velocity is mainly controlled by the ice sheet surface slope while the SSA component is mainly controlled by the ice flux and ice thickness and the ice-sheet surface slopes control the SIA and the SSA velocity components, but the SSA is also governed by basal dragging. Using both approximations in one model a hybrid model (i.e. based on both SIA and SSA approximations) allows to better represent the different rheologies deformation regimes found in an ice sheet. In GRISLI, the SSA velocity is used as a sliding velocity
(Bueler and Brown, 2009) when the basal temperature is at the ice pressure melting point. In this case, the basal drag follows a standard power law relating the basal velocity to the we assume here a power-law basal friction (Weertman, 1957) and the presence of sediments allowing for viscous deformation. The relationship between the basal shear stress with a \( \tau_b \) and the basal velocity \( u_b \) is expressed as:

\[
\tau_b = -\beta u_b
\]  

(3)

where \( \beta \) is a time constant but spatially variable basal drag coefficient. For cold base conditions, the sliding velocity is set to zero.

The resulting velocity for every model grid point is the addition of the SIA and SSA components. For floating ice points (ice shelves), we assume no basal drag. In addition, if the ice thickness of the floating ice shelves is under below 250 m and if no neighbouring points are grounded, the point is removed and the corresponding ice mass loss is considered as a calving flux. Determination of the grounding line position is based on a floating criterion. The sliding velocity is constrained by the topography and the characteristics of the Greenland bedrock. GRISLI also represents the deformation rate via the Glen flow law corrected by an enhancement factor to mimic the effect of ice anisotropy. The floatation criterion.

The isostatic adjustment in response to ice load-loading changes is governed by the flow-relaxation of the asthenosphere with a characteristic time constant of 3000 years and by the rigidity of the lithosphere deformation of an elastic lithosphere (?).

The climatic forcing is given by the mean annual SMB and the mean annual ST. Because seasonal variations of surface temperature are rapidly dampened, ST is considered as a good approximation of the bottom snowpack temperature. The model is thermo-mechanically coupled and the temperature field is computed both in the ice and in the bedrock by solving a time-dependent heat equation. The initial GrIS surface and bedrock topographies come from Bamber et al. (2013) and the geothermal heat flux is taken from ?.

### 2.2.2 Boundary conditions and climatic forcing

To compute the vertical properties of the GrIS, such as velocity and temperature, GRISLI only needs surface and bottom boundary conditions and surface climatic forcing. The annual mean near surface air temperature together with the geothermal flux (Maule et al., 2009) is used to compute the ice vertical temperature profile. To represent the GrIS initial topography and ice extent, we use boundary conditions from Bamber et al. (2013) and the Greenland Ice Mapping Project (Howat et al., 2014). The climatic forcings are from the MAR-MIROC5 experiment. At the bottom-

### 2.2.2 Spin-up procedure

Due to the long time scale response of the ice sheet we use the bedrock topography from Bamber et al. (2013). For this study, the basal drag coefficient is also a boundary condition. It is computed during the GRISLI initialisation step. This initialisation follows an optimized spin-up method (Le clec'h et al., in prep) based on data assimilation of the surface velocity field from Joughin et al. (2010) (see Sect. ??).
2.2.3 Model-initialisation and experiment

Due to the approximation and parametrisation used to solve the physical equations, it is necessary to equilibrate GRISLI with the initial boundary conditions (surface, bottom and vertical, before performing sensitivity experiments. Le clec’h et al. (in prep) optimised a data assimilation method applied to GRISLI by using the present day observed surface velocities, topography and climate forcing to obtain a GRISLI initial state. This method, similar to that used in Gillet-Chaulet et al. (2012), allows to avoid an initial to a given climate forcing, a proper initialisation of the model is required before performing forward experiments. For future sea-level projections, the aim of the initialisation is to start the simulations from a present-day equilibrated ice sheet geometry as close as possible to the observed one while ensuring consistency between internal properties of the ice-sheet (e.g. basal sliding velocities and vertical profile of temperature) with the climate forcing. Here, we use an inverse method, which is fully described in ?). Based on the same basic principles as that of ?, our method consists in the adjustment of the model to the boundary conditions during the initial years of the simulation, which could be misinterpreted as due to climate forcing. In the Greenland ice sheet, the characteristics of the ice just over the bedrock are poorly known and they are likely to change with space and time. Basal characteristics, such as basal sliding, have a crucial impact on ice sheet motion (Boulton and Hindmarsh, 1987; Weertman, 1957). As a result any error in the basal velocity computation can spread vertically in the ice and generate slowdown or acceleration of ice sheet motion. In GRISLI, the basal velocity is mostly influenced by the choice spatially-varying basal drag coefficient (and thus of the basal sliding velocities, see Eq. 3) so as to reduce the difference between the observed and the simulated ice thickness. However, while the study by ? requires long (multi-millennial) integrations for the method to converge, we use instead an iterative method of short (decadal to centennial) integrations starting from the observed ice thickness.

The basic principles of the procedure are the following: Initial vertical temperature and velocity profiles (Gillet-Chaulet et al., 2012) as well as the initial spatial distribution of the basal drag coefficient. We used an inverse method to infer the basal drag coefficient from observed surface velocities. Our computation of the basal drag coefficient is done in three main steps: The first step is a relaxation run of 200 years using initial conditions which are not necessary consistent between them (first guess of the basal drag from Edwards et al. (2014a), surface and bottom ice sheet characteristics ices from Bamber et al. (2013), vertical fields from Gillet Chaulet et al. (2012)). For this step, GRISLI is forced by the GRISLI simulations carried out within the framework of the Ice2Sea project (Edwards et al., 2014a). The climatological means of SMB and ST (1976-2005mean climate from) computed by MAR-MIROC5. Then, in order to have an ice flux as close as possible to observation, we calculate offline an ice surface velocity corrected by a factor representing the difference between the observed ice thickness of Bamber et al. (2013) and the GRISLI computed ice thickness (after 200 model years). If topography differences tend to be positive (resp. negative), the factor allows to decrease (resp. increase) the surface ice velocity in order (Fettweis et al., 2013) are used as climate forcing.

A 5-yr relaxation run is first carried out in order to avoid large inconsistencies between the different datasets used as climate forcings and boundary and initial conditions that are not necessarily mutually consistent. After this relaxation period, the method consists in an iterative adjustment of the spatially-varying basal drag coefficient (and thus of the basal sliding velocities)
so as to reduce the mismatch between the observed and the simulated ice thickness. The iterative process is divided in two main steps:

1. In the first step, we use the vertically-averaged velocity \(U^G\) simulated by GRISLI and computed from the previous time step to calculate a corrected vertically-averaged velocity field \(U^{corr}\) as a function of a corrective factor \(H^G/H^{obs}\) that represents the mismatch between the simulated \((H^G)\) and the observed \((H^{obs})\) ice thickness:

\[
U^{corr} = \frac{U^G \times H^G}{H^{obs}}
\]  

(4)

The mean velocity field is the sum of the sliding velocity and the velocity due to ice deformation. Assuming that the deformation velocity remains unchanged, the difference between \(U^{corr}\) and \(U^G\) is only due to changes in sliding velocities. This can be expressed as:

\[
\frac{U^{corr} - U^G}{U_{slid}} = \frac{U^{corr} - U_{slid}}{U_{slid}}
\]  

(5)

The adjusted basal drag coefficient \(\beta_{new}\) allowing to reduce the gap between observed and computed ice flux. In the second step, we perform another GRISLI simulation using the same initial conditions, the same forcing fields and set of parameters as in the first step. However, during the first mismatch between \(H^G\) and \(H^{obs}\) is deduced from the \(\beta_{old}\) values inferred from the previous iteration (or from initial condition for the first iteration) and from the ratio between the uncorrected and the corrected sliding velocities:

\[
\beta_{new} = \beta_{old} \times \frac{U_{slid}}{U^{corr}_{slid}}
\]  

(6)

\(\beta_{new}\) is calculated for each GRISLI grid point. \(H^G, U^G, U^{corr}_{slid},\) and \(\beta_{new}\) are updated every year during 20 years, instead of using the fixed initial basal drag coefficient, we use an iterative process to calculate each year a new basal drag value. To do that, we compare the corrected velocity fields calculated in the first step (hereafter the target velocity) to the one simulated for a specific basal drag. Depending on the ratio of the computed velocity and the target velocity, we compute a new coefficient that allows to increase or decrease the sliding for the following year. After the first 20 years, we stop calculating the coefficient and let the GRISLI simulation to evolve freely over years.

For the first iteration (after the relaxation period), the first step is skipped because there is no difference between \(H^G\) and \(H^{obs}\) and the procedure starts with the 2\(^{nd}\) step (see below).

2. Using the new basal drag coefficient \(\beta_{new}\) computed during the 1\(^{st}\) step, we let the model to freely evolve during 200 years more with the last computed basal drag. With the initial conditions used for this study, we need to repeat 8 times this second step in order to obtain a constant minimum gap between the computed and the observed GrIS topography and new corrected velocity fields \(U^{corr}\) computed from the mismatch between the simulated ice thickness (obtained at the end of the iteration) and the observed GrIS thickness. The third step allows to make the ice vertical profile of temperature consistent with the climate forcing used. Here the mean climate forcing is the same as the one used in previous steps. This 3\(^{rd}\)
The step is necessary due to the long time scale that temperature takes to be in equilibrium with the other fields. For this long experiment we perform a 30,000 years simulation where topography is kept constant during the entire experiment. All the other vertical fields (such as velocity and temperature) evolve through time. At the end of this long experiment we obtain a new set of initial conditions in which the ice sheet is in equilibrium with the climate forcing. However, because the velocity fields change (due to the evolving temperature) we compute a last basal drag coefficient over 20 years following a process similar to the second step. We use the resulting final conditions (thickness, temperature, velocity and basal drag) as conditions for all the following experiments of this study. In order to validate the final GrIS conditions obtained after the last spinup experiment, we perform a new GRISLI run using these conditions and the same mean 1976–2005 climate forcing coming from MAR. To obtain a final GrIS thickness as close as possible to observations, we use the Bamber et al. (2013) thickness as surface initial condition. After evolving freely over 2000 years, the GrIS volume drift each year by 0.0014% (0.01 mm in equivalent SLR). We then consider that the computed GrIS characteristics and dynamics reach an equilibrium with the mean climate forcing used. Despite of this equilibrium, the thickness difference between the GrIS computed at the end of the GRISLI initialization (after 2000 years of relaxation) free-evolving 200-yr simulation) and the observed one. With these new $U^\text{corr}$ values, a new cycle is started in which the 1$^{st}$ and 2$^{nd}$ steps are repeated. Each cycle uses initial conditions from the 5-year relaxed ice sheet topography. The overall process is stopped when the ice thickness root mean square error is not significantly improved. This ensures a good compromise between the reduction of the mismatch between observed and simulated ice thickness and the GrIS observed by Bamber et al. (2013) reaches a median anomaly value equal to rapidity of the convergence of the spin-up method.

In the present paper, the number of cycles that provides the best fit with observations (RMSE = +28 m and ranges on average over all the ice sheet, between -83 m (563 m) is $N_{\text{cycle}} = 8$.

To further reduce the model drift in terms of ice volume, a 2000-yr GRISLI relaxation run is performed after the last step 2 (after the end of the 8$^{th}$ quantile) and +211 m (95$^{th}$ cycle) under boundary and initial conditions identical to those of the last step 2. As such, the value of the basal drag coefficient is that obtained at the end of the 1$^{st}$ step of the 8$^{th}$ quantile. This result has been compared with other model results within the framework of an intercomparison project of initialisation methods (initMIP project). It turned out that GRISLI is one of the cycle. Over the last 150 years of this free-evolving simulation, the model drift is only $\pm 10^{-5}$ mm yr$^{-1}$ sea-level equivalent. At the end of this 2000-yr simulation, the simulated GrIS topography is slightly different from the observations (RMSE = +132 m, see Fig. S1). It will be referred hereafter to as $S_{\text{ctrl}}$ and will be used as initial topography for the transient GRISLI simulations described in the following.

For all the simulations presented in this study, including those carried out within the spin-up framework, we apply a strong negative SMB value outside the observed ice-sheet extent (Bamber et al., 2013). This avoids ice growth over actual present-day ice-free regions and allows to correct for both the potential atmospheric model biases (e.g., positive SMB values over tundra areas) and the models that compares the best with observations (Goelzer et al., 2017) spin-up procedure biases (i.e. too strong ice export towards the margins). For GrIS projections, the impact of this condition is quite negligible since GrIS will likely to keep on retreating over the next centuries.
3 Coupling methods and experiments

In this study, we compare three

3 Coupling methods

The aim of this study is to assess to what extent accounting for the atmosphere-GrIS interactions influences the GrIS evolution in terms of changes in SMB, ST, ice thickness and SLR. To achieve this goal, we designed three experiments based on coupling methods of different complexities to account for the atmosphere—GrIS interactions. We investigate how these methods affect the computed SMB, ST, SLR contribution and surface elevation changes. For the three methods, interactions between MAR and GRISLI. For all the experiments described below, the climatic forcing is designed as follows: MAR is forced every 6 hours at its lateral boundaries by the transient MIROC5 fields and run at a 25 km resolution from 1976 to 2100. The forcing fields from MIROC5 evolve in response to the historical (until 2005) and atmospheric fields from the CMIP5 historical run (1970-2005) and RCP8.5 scenarios (from 2006 to scenario (2006-2100). In order to extend the MAR simulation until AD 2150 and in the absence of a MIROC5 simulation performed under a prolonged RCP8.5 scenario (i.e. after 2100), Before forcing GRISLI, MAR is forced from 2101 to 2150 with the 25 km resolution 2095-yr MIROC5 climate. We chose the year 2095 because, averaged over the entire GrIS, the 2095 mean climate is one of the closest to the decadal 2090-2100 one. This implies that both climate changes and large-scale inter-annual variability are neglected beyond 2100.

3.1 The No Feedback experiment

The first one (referred hereafter to as NF) is based on a one-way coupling approach in which GRISLI is forced by the climatic outputs (SMB and ST outputs) obtained from the MAR model. we need to spatially interpolate them onto the 5 km GRISLI grid. However, as simulation spanning from 2000 to AD 2150. The aim of this experiment is to examine the ice-sheet response to the climatic forcing without accounting for the feedbacks related to GrIS changes. The SMB and ST are very sensitive to time series simulated by MAR from 2000 to AD 2150 are used to force the GRISLI ice-sheet model. Using an inverse distance weighting method, they are first interpolated on the GRISLI grid to account for the Greenland topography. We also need to correct them by the topography changes due to the difference of resolution between MAR both models (25 km) and GRISLI (vs 5 km). Using a linear interpolation method if the steep topography at the GrIS margin and the complex orographic features in these areas are not taken into account could lead to important biases in SMB and ST in these regions (Franco et al., 2012). In the present study, for all the experiments, we first interpolate the MAR outputs on the To account for the differences in surface elevations between the 25 km and 5 km GRISLI grid using a simple bilinear interpolation. Then, the fields are corrected for the altitude difference induced by the difference of resolution between MAR and GRISLI. These topography corrections are based on the method developed by Franco et al. (2012) who derive Bamber et al. (2013) topographies, we also apply a vertical correction following Franco et al. (2012) who derived a local vertical gradient of SMB (or ST) each SMB component as a function of altitude for each GRISLI grid point. This gradient is then used to compute the correction due to the difference
of altitude between the Bamber et al. (2013) topographies seen by MAR (resolution 25 km) and by GRISLI (resolution 5 km). Thus, this method allows to generate a 5 km resolution SMB entirely consistent with the Bamber et al. (2013) topography. While this procedure can be followed at a daily time scale (Noël et al., 2016). For our purpose, we choose to average the daily vertical gradients, in the present study, the vertical gradients are averaged at an annual time scale and to apply the altitude correction every year sale and used as corrective factors to downscale at the end of each model year the SMB and ST fields onto the 5 km grid. For all the experiments described below, the coupling between MAR and GRISLI starts in 2020 when the SMB anomalies are large enough to induce significant topography changes in GRISLI.

We investigate three levels of coupling between the ice sheet and the atmospheric models: The ”No Coupling” method (hereafter NC).

3.2 The Parameterised Feedbacks experiment

In the second experiment (referred to as NC). The present-day GrIS geometry (topography and ice extent) provided by Bamber et al. (2013) is prescribed to MAR as a boundary condition during the entire simulation duration. GRISLI is forced until 2150 (see Sect. 2.2) using the downscaled MAR outputs described above. The NC method does not allow Parameterised Feedbacks (PF in the following), the SMB and ST fields simulated by MAR are corrected each year following the method of Franco et al. (2012) to account for the ice sheet feedback on the climate. Rather, it provides the response of GrIS under a specific climate forcing – The ”One way coupling” method (hereafter called 1-W). This method goes a step further. It is based on the same principle as the NC method, but the SMB and ST fields are corrected based on an updated altitude \(H(t)\) given by (Eq. 7): evolution of the simulated GRISLI topography. This correction is made from 2020 onwards, as changes in SMB through 2006-2020 do not produce any significant topography changes in GRISLI. The new corrected SMB and ST values are computed at the altitude \(S(t)\) defined on the 5 km grid as:

\[
H(t) = H_S(t) = H_{S_{Bamber}} + \Delta H \Delta S_{GRISLI}(t)
\]  

(7)

where \(H_{S_{Bamber}}\) is the Bamber et al. (2013) topography at a given time \(t\), \(H_{S_{Bamber}}\) is the present-day observed topography defined on the 5 km and \(\Delta H_{GRISLI}\) is the topography anomaly GRISLI grid and \(\Delta S_{GRISLI}(t)\) the difference between the altitude simulated by GRISLI between the initial topography computed for year 2000 from the equilibrium state (t=0) and the ongoing time step (t). In doing so, this method artificially accounts for the elevation feedback because the at time \(t\) and \(S_{CTRL}\). Due to the topography differences between MAR and GRISLI, this approach has been chosen to avoid large inconsistencies between the SMB and ST are initially fields computed by MAR on a fixed ice sheet topography. With this method GRISLI is forced off line by the MAR atmospheric conditions already computed in the NC run, therefore allowing sensitivity experiments in GRISLI with limited additional computer time and the ones corrected to account for the GRISLI topography.

This method offers the possibility to account artificially for the elevation feedbacks when using existing RCM simulations in which the topography is kept constant. As such, it is also transferable to any ice sheet model. However, in this approach, the changes in GRISLI topography are not taken into account by MAR. The Fully Coupled method (hereafter 2-W). This
coupling method is the most accurate way to represent the interactions between the GrIS and the atmosphere but it is also more computationally expensive. GrIS geometry have no consequence on the climate as simulated by the atmospheric model.

### 3.3 The two-way coupling experiment

The third method (2W in the following) is based on a two-way coupling strategy between MAR and GRISLI. Both models use the same boundary and initial conditions as those of the NF and PF experiments. At the end of a MAR simulated model year, MAR is paused and GRISLI is forced by the 5-km interpolated downscaled SMB and ST just computed by MAR. GRISLI then fields with the method of Franco et al. (2012) as in PF (Eq. 7). Then, GRISLI computes a new GrIS topography and extent a new ice extent at 5 km which are aggregated onto the at the yearly time scale onto the 25 km MAR grid for the simulation of the next year of the MAR experiment. GRISLI and MAR are never stopped, just alternatively paused and resumed until 2150. The differences between the GRISLI equilibrium state after the initialisation step and the observed topography (Bamber et al., 2013) (cf Sect. ??) could lead to inconsistencies between the results obtained by MAR under its usual setup, i.e. calibrated with the Bamber et al. (2013) topography, and the results that would be obtained by using directly the GRISLI topography. For this reason, in both the 2-W and the 1-W experiments, we use anomalies of GrIS topography applied on Bamber et al. (2013) topography rather than the absolute topography from GRISLI. The aggregated ice extent is used to update the fraction of tundra relative to ice/snow covered surface type for the subsequent MAR run. To account for the differences between MAR and GRISLI topographies, the surface elevation which is aggregated onto MAR is computed from GRISLI surface elevation anomalies added to the present-day observed topography (Eq. 7). It is then used as the updated surface elevation in MAR. As previously mentioned, topography changes are negligible before 2020. Hence, changes in ice-sheet geometry are fed to MAR only after this date. Compared to the NF and PF approaches, this two-way coupled method is the most accurate to represent the GrIS-atmosphere feedbacks.

### 4 Results

#### 4.1 The uncoupled-simulation: Greenland ice sheet evolution the NC-2W experiment

##### 4.1.1 MARChanges in the forcing climate

The mean ST over the first two decades (2000-2020) and averaged over the whole GrIS is equal evolution of the SMB and of its different components (Eq. 1), integrated over the entire GrIS and simulated in the 2W experiment is displayed in Fig. ???. During the 2000-2040 period, the averaged SMB remains positive with a mean value equal to 280 ± 95 Gt yr\(^{-1}\) (where the notation ± represents the standard deviation computed from yearly values) but slightly decreases by 4 Gt yr\(^{-1}\). This decrease becomes substantially stronger from 2040 to -18.7°C 2100 (-17 Gt yr\(^{-1}\) on average), and the mean SMB to 434 Gt yr\(^{-1}\) (Fig. ??). After 2020, the ST increases by 0.065°C yr\(^{-1}\) until reaches strong negative values (-638 ± 271 Gt yr\(^{-1}\) over the 2090-2100 period). As the same MIROC5 year 2095 is repeatedly used to force MAR after 2100, there is no longer inter-annual variability and the integrated SMB remains quite stable between 2100 (Fig. ??). Over the same period (2020-2100), the averaged GrIS
SMB decrease by 12.3 and 2150 (-812 ± 13 Gt yr⁻¹), despite a slight increase of ~1 Gt yr⁻¹. Throughout the simulation, the evolution of the SMB signal is dominated by surface runoff whose increase rate (in absolute value) ranges from 5 Gt yr⁻¹ (Fig. ??A). The surface temperature anomaly in 2000-2040 to 19 Gt yr⁻¹ (2040-2100). After 2100 compared to 2000 shows a warming ranging from +1.5 °C in the southern part of the GrIS to more than +8 °C in the northern part (Fig. ??A). This warming in northern Greenland is a direct response to the MIROC5 forcing fields due to the polar temperature amplification. However, regional heterogeneities are observed in the annual mean GrIS SMB spatial distribution, it slightly decreases (~2 Gt yr⁻¹), explaining the slight SMB increase.

The SMB anomaly between the beginning and the end of the 2W experiment exhibits two distinct patterns (Fig. ??B). Indeed, between ??a, 65% of the grid points having surface elevations higher than 2000 and 2100 there are are characterized by a positive SMB anomaly, ranging from 0.07 m yr⁻¹ (5th percentile) to 0.2 m yr⁻¹ (i.e. more ice accumulation) in a zone located along a South North transect in the central part of the GrIS. This ice accumulation is mainly governed by the larger snowfall on the GrIS central part in winter 95th percentile) at the end of the simulation. This SMB increase is particularly pronounced in the eastern part of the GrIS between 67 and spring seasons (not shown). An opposite trend 10 times larger than ice accumulation, is simulated over the edges of the ice sheet, with a negative SMB (i.e. ice ablation). In these regions, the summer season governs this negative SMB and is characterised by larger rainfall and melting ice (meltwater and runoff increase) than for other seasons (not shown) 70 °N and in the north central part. It is due to a strong increase in snowfall (> 0.5 m yr⁻¹, Fig. ??b) which occurs mainly during the winter season in the east and during autumn in the north (Fig. S2). On the other hand, 87% of the GrIS grid points with surface elevation lower than 2000 m are dominated by an increase in surface runoff (Figs. ??c, S3) and by an increase in the fraction of rainfall over snowfall in summer and in autumn (Fig. S4).

As a consequence, the limit between the accumulation (SMB > 0) and ablation (SMB < 0) areas, also called the result, strong negative SMB anomalies are found in these regions ranging from -3.3 m yr⁻¹ (5th percentile) to -0.1 m yr⁻¹ (95th percentile) and reaching more than ~6 m yr⁻¹ along the western and the southeastern margins (Fig. ??a).

The equilibrium line altitude (ELA, a line where i.e. altitude for which SMB = 0), shifts inland through time (Fig. ??). This shift explains the mean decrease of SMB over the whole GrIS until 2100 seen in Fig. ??.

### 4.1.2 GRISLI

In 2100, a decrease of 15.8 m mean ice sheet thickness is simulated increases significantly between the beginning and the end of the 2W experiment, as a consequence of increased runoff for areas below 2000 m. As an example, at around 73.5 °N, on the eastern side of the ice sheet, the ELA moves from ~1000 m to ~2500 m (Fig. ??c) with a standard deviation of 32.7 m. In 2150, the mean ice thickness decrease reaches 28.6 ± 68.4 m ??). We can distinguish two types of regions: there is a thinning over the GrIS coastal regions and a thickening over the central GrIS regions. In other regions, at the end of the 2W experiment, the ELA is generally situated between 1500 and 2000 m high, except in the northern part where it is between 1000 and 1500 m. This shift of ELA towards higher altitudes represents an increase of 24% of the ablation area between the beginning and the end of the experiment.
The ST anomaly (Fig. ?2C). The thickness changes for the 2000-2150 period show the same patterns as the 2000-2100 period, but with a larger magnitude. Over the thinning regions, between 2000-2100 (resp. 2000-2150) the ice thickness decreases by \(40 \pm 32.5 \text{ m}^{-1}\) ranges from \(+2.2 \, ^\circ\text{C}\) (resp. \(+78.5 \pm 70 \text{ m}^{-1}\) 5th percentile) to \(+6.5 \, ^\circ\text{C}\) (95th percentile) and is characterized by a south-north gradient with the highest values found in the northern part. Beyond \(78 \, ^\circ\text{N}\) the ST anomaly reaches locally values greater than \(+11 \, ^\circ\text{C}\). This temperature increase from 2000 to 2150 contributes to the amplification of the ablation processes below the ELA. However, while the stronger temperature anomaly is found in the northeastern part of the ice sheet, ablation processes are modulated by the increasing snowfall in 2150 compared to 2000.

4.1.2 Changes in Greenland ice-sheet geometry

The ice thickness anomaly (Fig. ?) also presents two distinct patterns. For surface elevations higher than 2000 m in the northern part, and higher than 2500 m in the central and southern parts of the ice sheet, the ice thickness increases by 5 m on average, with the increase ranging from 1.5 m (5\(^{th}\) percentile) to 17 m (95\(^{th}\) percentile). On the contrary, in central regions other hand, in regions whose surface elevation is lower than 2000 m, the ice thickness increases with a median value of 4.4 decreases from \(-248 \text{ m}^{-1}\) (5\(^{th}\) percentile) to \(-3 \text{ m}^{-1}\) (95\(^{th}\) percentile) with a mean value equal to \(-100 \text{ m}\). As a result of these GrIS ice thickness changes, the surface slope between the central part of the ice sheet and the margins increases. On top of that the ice sheet mask (defined as the fraction of a MAR grid cell with permanent ice cover, Fig. S5) decreases by \(2.8 \pm 42 \text{ m}\) (resp. 17.8 \pm 0.1 \%) (mean \(13.4 \text{ m}\)). The standard deviation computed from yearly values over the 2140-2150 mean period compared to the 2000-2010 mean period, and some GrIS margin regions become ice free (red grid points in Fig. S5).

4.1.3 Changes in ice dynamics

The ice thickness anomaly is due to the complex combination of changes both in surface atmospheric conditions and ice dynamics conditions. The ice dynamics is impacted by the warming climate. Generally, the simulated ice velocities increase from the central part of the (SMB, Fig. ?a), ice dynamics (ice flux divergence, Fig. ?b) and basal melting (not shown), following the continuity equation (Eq. 2). To quantify the role of ice dynamics on the GrIS geometry (Fig. ?), we plotted the ice flux divergence integrated over 150 years (2000-2150, see Fig. ?b). In particular, over the central plateau, the cumulated SMB (Fig. ?a) reaches about \(+50 \text{ m}, 40 \text{ m}\) of which are transported away by the ice dynamics (Fig. ?b). As a result, the ice sheet to the coastal regions thickness anomaly is reduced to only \(~10 \text{ m}\) in this region (Fig. ?a). An opposite behaviour is found near the western coast, where the ice melting is partly compensated by ice convergence, resulting in a less negative ice thickness anomaly than that related to the SMB forcing. This shows that ice dynamics act to counteract ice loss from surface melting, as previously noticed by several authors (??Edwards et al., 2014a). As a consequence, it appears to be essential to account for ice dynamics to estimate accurately the mass balance of the whole ice sheet.

In turn ice dynamics is impacted by changes in ice sheet geometry as illustrated by the mean surface velocity anomaly (Fig. ?A). However, at the ?a). For regions with surface altitudes between 2000 and 2500 m, the anomaly of the ice flow increases from the inner GrIS areas towards the edges of the ice sheet. The increase in the mean ice flow for the 2140-2150 period compared to 2000-2010 period, ranges from \(0.08 \text{ m yr}^{-1}\) (5\(^{th}\) percentile) to \(17 \text{ m yr}^{-1}\) (95\(^{th}\) percentile). These faster
ice velocities at the end of the 2W experiment are mainly explained by the combination of a larger surface slope between the central and the margin regions of the ice sheet. This is consistent with information inferred from ice flux divergence as shown in Fig. ??b.

On the contrary, for the margin regions, with altitudes lower than 1500 m, the anomalies of surface ice sheet margins, the ice velocities strongly decrease. In GRISLI, the fine scale structure (Fig. ??a). Compared to the 2000-2010 period, this decrease ranges from -213 m yr$^{-1}$ (5$^{th}$ percentile) to -0.2 m yr$^{-1}$ (95$^{th}$ percentile), and is fully consistent with the decrease in ice thickness (Fig. ??).

The changes in local ice dynamics between the first and the last 10 years of the 2W experiment are also related to changes in surface slope and ice thickness, particularly at the margins. To investigate the ice dynamics changes at the local scale, we used the examples of the Jakobshavn (western coast) and the Kangerlussuaq (eastern coast) glaciers is for which the fine scale structures of the ice velocity, obtained after the GRISLI initialisation procedure, are relatively well reproduced (Fig. ??A-B). For these glaciers and their associated ice streams, within 100 years, the surface velocities slow down compared to the observations (Figs. ??a-b).

For the Jakobshavn glacier, and for altitudes above 1500 m, the vertically-averaged ice velocities increase by more than 200-15 m yr$^{-1}$ in the coastal regions, while they increase (i.e. 10 %) as a result of increasing surface slopes, and slow down by more than 60-200 m yr$^{-1}$ in the interior (i.e. 29 % for altitudes below 1000 m due to the decreasing ice thickness (Fig. ??C-D). Because the region feeding the ice stream has a frozen base, the SIA assumption is the predominant simplification used by GRISLI to compute the ice velocities. In this area, the ice velocity increase is due to ??c). For altitudes above 500 m, the vertically-averaged velocity is mainly driven by the SIA velocity component-SIA velocity (Figs. ??c-e). On the contrary, below 500 m, basal sliding velocities are large due to low basal drag coefficient (see Fig. 3 in ?) and the SSA velocity component dominates the ice flow (Figs. ??c-g). However, while basal drag is lower in locations below 500 m, the ice flow is limited by the strongly reduced ice thickness (Fig. ??E-F). As mentioned previously, by 2100, the thickness decreases at the margins and increases in the interior (Fig. ??K-L), resulting in steeper slopes and thus in larger SIA velocities. However, the coastal regions have a temperate base and the SSA component of the velocity is predominant. Thus, the ice flow velocity is mainly controlled by the ice flux coming from the inland part. As the ice flux depends on the ice thickness which decreases over the coastal areas, the SSA velocity component decreases (Fig??, ??G-H). This increase decrease ice velocity pattern has been also reported by Peano et al. (2017), using GRISLI forced by CMIP5 models under the RCP8.5 scenario. As a result of ice accumulation ablation changes and ice velocity changes, the ice mask (numbers of grid points covered by permanent ice) decreases by 3.7 % in 2100 compared to the initial one (in 2000-2005). During the first 20 years (2000-2020), the total GrIS volume remains stable with no additional contribution to SLR compared to year 2000. After 2020, the GrIS volume decreases, resulting in a global contribution of +7.6 cm in 2100. Extending the GRISLI-NC experiment until 2150, forced by the same 2005 MAR forcing-climate results in an amplification of all the changes observed in 2100 and discussed above: the extent of the ablation zone. -

The Kangerlussuaq glacier is located in regions where the bedrock is characterised by a succession of valleys surrounded by mountains merging in a canyon where the deepest part is located 100 km away from the coast (?). The ice flow of the
Kangerlussuaq is therefore divided in different branches with increasing ice velocities towards the ice sheet margin and becoming even larger when merging in the canyon (Fig. ??b). As for the Jakobshavn glacier, the ice flow accelerates at the end of the 2-W experiment as a consequence of the increase in surface slope for high altitudes (~2000-2500 m, see Fig. ??). Conversely, a strong decrease of the larger thinning and the slow down ice velocities in the coastal regions. As a consequence, the GrIS contribution to global SLR is amplified, reaching +18.7 cm in 2150. Ice flow is found in most of margin regions (Fig. ??d) directly related to the ice thinning (Fig. ??). Contrary to the case of the Jakobshavn glacier that presents large basal sliding velocities only below 500 m, the Kangerlussuaq shows low basal drag coefficients in the entire glacier (see Fig. 3 in ??) and thus the ice flow is mainly governed by the SSA component (Fig. ??h).

4.2 Differences between the 2-W and the NC experiments.

4.1.1 Impact on SMB and ST

The near-surface temperature (ST) simulated for 2150 in 2-W experiment-21st century (w.r.t 1970) in downstream regions of the Jakobshavn and Kangerlussuaq glaciers as a consequence of ice-sheet thinning at the margins.

4.2 Differences between the 2W and the NF experiments.

4.2.1 Impact on SMB and ST

To assess the importance of the atmosphere-GrIS feedbacks, we now compare the 2W and the NF experiments. The main SMB differences between both experiments, averaged over the 2140-2150 period, highlight lower SMB values in 2W compared to NF for altitudes below 2000 m, with the exception of some margin locations in the eastern part (Fig. ??a). This SMB anomaly behaviour is driven by a snowfall reduction in low altitude areas (Fig. S6) and by the runoff increase in 2W with respect to NF (Fig. S7). This increased runoff results from warmer temperatures over the whole GrIS is warmer than in the NC experiment, (up to 0.8 °C in the western and northern parts, Fig. ??b), except in the region at the edge of the GrIS, which is strongly colder (until sees a strong cooling (as low as -10 °C, Fig. ??b). The ST of this region is sensitive to the atmospheric circulation. At the edges-warming can be explained by the temperature-altitude feedback being active in 2W, resulting in lower altitudes (Sect. 4.1.2 and Fig. ??c) and therefore warmer temperatures. The cooling over the very edge of the ice sheet, there is an intensification of the strong and cold katabatic winds coming from the central part of the GrIS in 2-W compared to NC. The katabatic winds have a daily time scale resolution and are represented by the MAR model (Gallée and Pettré, 1998; Gallée et al., 1996). These stronger winds are due to the higher coastal surface slope simulated in 2-W than in NC (see Sect. ??). As a consequence, they prevent warmer and wetter air from the ice-free areas (occurs despite the ice sheet thinning over these regions and can be explained by changes in atmospheric circulation.

Indeed, unlike NF, 2W allows for an explicit computation of changes in ice sheet surface slopes due to increased melt at the margin. This has important consequences on the atmospheric circulation and in particular on the katabatic winds (Fig. ??).
Over the ice sheet, the steeper surface slopes simulated in 2W in 2150 (discussed in Sect. 4.1.2) lead to a slight increase in katabatic winds (Fig. ??). However, at the ice sheet margin, i.e. covered by tundra and ocean from reaching the GrIS margin regions (van den Broeke and Gallée, 1996). Thus, the ST decrease over the edges of the GrIS (Fig. ??A). The second where the ice mask in MAR is below 100%, there is a substantial decrease in surface winds. This is because the change in surface elevation as seen by the atmospheric model is computed from the aggregated changes in GRISLI at 5 km. As such, a non-zero fraction of tundra, which presents no change in surface elevation, results in smaller elevation changes compared to grid cell in the same region with permanent ice cover only. This induces artificially lower surface slopes at the margin with respect to the interior and a decrease in surface winds in these regions. Altogether, the slight increase in katabatic winds over the ice sheet and their reduction at the margin lead to a cold air convergence towards the ice sheet edge (Figs. ??b, ?? and Figs. S8-S9).

Another consequence of the strengthening of the katabatic winds-katabatic winds increase due to increased surface slopes in the GrIS interior, is to enhance the atmospheric exchange at the middle of the slope over the GrIS. Indeed, at the surface, the exchanges along the slope of the ice sheet. The area with lower atmospheric pressure generated by the stronger katabatic winds is filled in by the warmer air coming from higher atmospheric levels in the boundary layer. Thus the warming of the upper part of the boundary layer combined with the lower surface elevation, explains the ST increase on the coastal regions inland from the very edge of the ice sheet. These two types of colder and warmer regions simulated in 2-W with respect to NC are already present after 100 years of experiment increases in the interior of the GrIS.

### 4.2.2 Impact on ice thickness and ice dynamics

The most important ice thickness difference between the last ten years of the 2W and the NF experiments is a smaller thickness in 2W compared to NF. This is mainly explained by the positive temperature-elevation feedback in 2W that results in increased surface temperatures compared to NF, and thus increased melting, when surface elevation decreases. Areas with this type of behaviour cover most of the Greenland ice sheet slopes and reach the interior of the ice sheet from the western or the northeastern margins. The largest changes occur over the western edge of the GrIS, where the thinning between 2W and NF reaches more than -25 m (Fig. S1A-B). In 2150, the SMB difference between ??c). The ice thickness anomaly pattern is essentially mimicking the SMB differences between 2W and NF (Fig. ??a), suggesting that the two-way coupling induces only a relatively limited change in ice dynamics, as shown by the 2-W and NC experiments exhibits two distinct patterns. With the 2-W approach, the SMB increases by 0.6 m yr\(^{-1}\) over the eastern coast, the south central part and in some local regions in the northern part of the GrIS-ice flux divergence anomaly (Fig. S10), although the surface velocities (Fig. ??B). These regions are characterized by a larger snowfall in winter season compared to the NC experiment. The processes explaining these increased SMB regions are probably linked to the strengthening of the atmospheric circulation along the Greenland eastern coast coming from northern latitudes, thus bring wetter and colder mass air. Despite these regions of positive SMB anomalies in 2150, the pattern in the SMB difference between 2-W and NC is generally a negative anomaly ranging from -2.3 m yr\(^{-1}\) to -0.4 cm yr\(^{-1}\) over the coastal areas of the GrIS. Following the decrease of the ST, the surface snow melting cumulated with less snowfall and more precipitation falling as rain instead of snow drive this ??b) are slightly smaller in 2W due to smaller ice thickness (Fig. ??c).
4.3 The PF experiment

As previously described (see Sect. 3.2), the PF experiment is based on a parameterisation of the surface elevation feedbacks. In this section, we present the differences \(2W - PF\) in SMB patterns. As a result, in 2150, there is a decrease of 112 Gt yr\(^{-1}\) of ice over the ablation area in 2-W with respect to NC, and, over the accumulation area, the simulated SMB in the 2-W experiment is higher. ST and ice thickness averaged over the 2140-2150 period (Fig. ??) so as to examine the efficiency of this parameterisation. The first key feature is that the \(2W - PF\) SMB anomaly (Fig. ??a) is less negative than the \(2W - NF\) one (Fig. ??a). This results from the fact that the decreasing altitude is taken into account in PF through the altitude feedback parameterisation, leading to smaller differences with the 2W experiment. In most margin areas, the SMB simulated in PF has even become lower compared to 2W. Cumulated over the entire GrIS, the \(2W - PF\) SMB difference is -28 Gt yr\(^{-1}\) lower than the one simulated in the NC experiment. In 2100, the SMB anomaly shows the same patterns as the 2150 SMB pattern, but with lower magnitude (Fig. S1A). Over the ablation zone, \((-149 \text{ Gt yr}^{-1} \text{ for } 2W - NF)\). In the mean value of the SMB anomalies increases by a factor of 10 between 2050 and 2150 (see Table ??). The SMB changes have an impact on the extent of the ablation zone. This area increases with time and, at the end of the simulation is 14 \% larger in 2-W than in NC (Table ??). As a result, the ELA shifts more inland, in 2-W, by +12 km in the north-eastern GrIS (wrt NC) same way, the 2-W – PF differences in ST and ice thickness (Figs. ??b-c) are also less pronounced than the 2W – NF differences (Figs. ??b-c), highlighting the importance of the elevation feedbacks. These results show that over ~150 years, the topography correction used in PF, allows from an uncoupled experiment to obtain simulated fields close to those of the 2W coupled experiment.

To illustrate the spatial variability of the ice thickness response to the different coupling methods, we plotted the ice thickness anomalies between 2W and NF (red dots, Fig. ??a), PF and NF (green dots, Fig. ??)…

4.3.1 Impact on thickness and ice dynamics

The patterns of the surface elevation changes between the 2-W and the NC experiments ??a) and \(2W\) and PF (Fig. ??b) as a function of the ice sheet altitude. The \(2W\) method yields negative and positive anomalies relative to NF, while the PF method mainly yields negative values (Fig. ??c). Follow the SMB anomaly patterns. Over the eastern coast, southern central part and locally in some regions on the northern part of Greenland, the ice thickness increases, reaching more than 10 m in some locations such as East Greenland (Fig. ??C). This increasing ice thickness is explained by a larger positive SMB and a lower surface temperature over these regions. The second anomaly pattern is found all along the Greenland coast, where a decrease of ??a), illustrating the stronger spatial variability in 2W. For both 2W and PF experiments compared to NF, the regions at low to medium elevations are the most sensitive to the coupling approach with the stronger spatial variability of the ice thickness found for altitudes below 1000 m. For example, the \(2W - NF\) ice thickness anomalies range between -31.9 m (5\text{th percentile}) and +6.5 m (95\text{th percentile}), and between -27.8 m (5\text{th percentile}) and 0 m (95\text{th percentile}) for the \(PF - NF\) case. Overall, the ice thickness is found in areas of lower ST and in the ablation zone. The main changes occur on the western edge of the GrIS where the thinning between 2-W and NC reaches more than 25 m (Fig. ??C). Further inland, there is a smaller thinning (0.2 ± 3 cm in average after 150 years). As a result, averaged only over the entire ablation area, the thinning after 150 model
years is equal to $9.0 \pm 11.1$ m (see also Table ??). These ice thickness anomaly patterns are observable, with lower magnitude, when comparing 2-W and NC after 100 years (Fig. S1C). The main consequence of the increased thinning in coastal regions is the increase of the surface slope between the central part and the margin of the ice sheet. Increased surface slopes results in stronger katabatic winds. Furthermore, the thinnest parts of the GrIS become ice or snow-free or snow-free, exhibiting bare ice and modifying albedo feedbacks, with a decrease of the surface albedo which amplifies the GrIS melting. The ice dynamics computed by GRISLI are also impacted by the full representation of the interactions. Compared to the NC experiment, the ice velocities simulated with the 2-W experiment show a succession of positive and negative anomalies (Fig. ??b). The ice velocities increase from the central part of the GrIS to the coastal regions. The increase decrease velocity pattern is amplified in the 2-W compared to the NC because of the larger thickness anomaly and follows the same processes than explain in Sect. ??, anomalies decrease with increasing altitudes (Fig. ??a) and increase with time (Fig. ??b). They are also stronger in the 2W experiment than in the PF and NF simulations.

4.3.1 Impact on SLR contribution and ice sheet area

After-

4.4 Impact on GrIS contribution to sea-level rise and ice sheet mask

At the end of the simulation (i.e. after 150 model years, the melting contribution to global SLR reaches $\pm$ model years), the GrIS contribution to sea-level rise (computed from the change in GRISLI ice volume), simulated in the 2W experiment, reaches 20.4 cm in the 2-W experiment. In comparison, the melting obtained in the NC experiment is equivalent to a SLR of $\pm$ cm, against 18.5 cm. This difference (Fig. ??b) is linked to the cm and cm in the NF and PF experiments respectively (Table 1 and Fig. ??). Owing to the negligible model drift ($\sim 10^{-5}$ mm yr$^{-1}$, see Sect. 2.2.2), these differences only result from the better representation of the interactions between the GrIS and the atmosphere. The coupling allows for a better representation of the processes occurring at the margin, and in particular the ice sheet margin retreat. As GrIS-atmosphere feedbacks in 2W leading to increased runoff due to warmer temperatures (see Sect. 4.2). In 2100 (Table ??), the differences between the three experiments are smaller, with the NF and PF contributions being respectively 4.4% and 0.4% lower than the 2W contribution, against 9.3% and 2.5% in 2150. These results reflect several key aspects. First they show that the GrIS mass loss substantially accelerates from the second half of the 21st century onwards and that the effect of the different feedbacks, as simulated in 2W, is enhanced over time. Figure ??b displays the sea-level anomalies between 2000 and 2100 to better illustrate the divergence of the three experiments as soon as 2025-2030. Secondly they illustrate the effect of the feedbacks itself. As an example, accounting for the parameterised feedbacks (PF) leads to an additional SLR contribution (w.r.t NF) of 4.2% in 2100 (7.6% in 2150). This is smaller than that reported in who also used the MAR model to force the hybrid SICOPOLIS3.3 including a representation of subglacial hydrology. However, our estimate is comparable with the 4.3% additional contribution found by Edwards et al. (2014a) in 2100 who used ECHAM5 and HadCM3 to forced MAR simulations under the SRES A1B scenario and five ISM projections, and within the range of uncertainties of the 8 ± 5% additional surface mass loss reported in Fettweis et al. (2013). As for Vizcaíno et al. (2015), they also conclude that the melt-elevation feedbacks, simulated with
the ECHAM5.2-SICOPOLIS3.0 coupled model under the RCP8.5 scenario, contribute to 11 % to SMB changes and to 8 % to SLR. Compared to our own study, the larger contribution found in Vizcaino et al. (2015) may be explained by the coarser resolutions of ECHAM5.2 (~3.75° for the atmospheric component) and of SICOPOLIS3.0 (10 km) with respect to MAR and GRISLI resolutions. Our results also suggest that, at the centennial time scale, the SMB-elevation feedback is the most important since its parametrization in PF allows to reduce the mismatch between the 2W and NF GrIS SLR contributions by 73.7 % (resp. 91.4 %) in 2150 (resp. 2100), the remaining contributions being attributed to albedo and atmospheric feedbacks. However, to assess more accurately the relative importance of the elevation feedbacks, a more appropriate procedure would be to cut off the elevation feedbacks in the 2W experiment.

Compared to the NF and the PF experiments for which the ice-sheet mask is fixed to observations from 2000 to AD 2150, the 2W ice sheet extent is reduced by ~2.8 % in 2150 as a result of ice melting. GrIS coastal grid points can become ice free. The GrIS extent in the 2-W experiment is reduced by 52 400 km² compared to the NC experiment and increases exponentially with time (Table ??). Thus, the ice sheet mask field, which represents the ice coverage percentage of each grid cell of the grid used by the models, and which is therefore dependent of the ice sheet extent, decreases with time (Fig S2B). To evaluate increased ablation. As MAR sees the ice sheet retreating over time in 2W concomitantly with the increase in bare ground or tundra fractions (Fig. S5b), the albedo feedback takes place favouring further the ice melting. Although the ice sheet retreats, the projected GrIS melting contribution to SLR, the SMB integrated over the ice covered areas (i.e. the sheet mask field) is often used (Fettweis et al., 2013). However, this method could lead to strong uncertainties in the SLR contribution obtain. For example, using the NC result in extent of the ablation zone increases with time. This process is faster in 2W than in NF and PF.

In 2150, if we integrate the SMB over the no updated ice sheet mask (as in the NC method), we calculate an integrated SMB

15458 %, the ablation zone is 14 % (resp. 11.7 %) larger in 2W than in NF (resp. PF) causing 112 Gt yr−1 lower than using the updated ice sheet mask−1 of extra ice ablation in 2W (w.r.t NF). As a consequence, the ELA is located further inland in 2W compared to NF with a maximum inland retreat of 120 km located in northeastern Greenland (Fig. ??).

A widely used method to estimate the projected GrIS to global sea-level rise is to compute the GrIS mass loss as the time-integral of the SMB computed by an atmospheric model over a fixed ice-sheet mask (Fettweis et al., 2013; ?; ?). In the present study, we go a step further since the ice mass variations related to SMB changes are computed over a changing ice-sheet mask as simulated by GRISLI (as in the 2-W method). This higher integrated SMB, obtained when using no updated ice sheet mask, is only explain by taking into account the GrIS regions becoming ice free compared to the updated. However, in both the NF and the PF experiments, the atmospheric model does not account for the variations in the ice-sheet extent simulated in GRISLI and the ice-sheet mask, taken from the observations (Bamber et al., 2013) is kept constant throughout the simulation. Taking the changes in ice-sheet mask into account may have strong impacts on the computed GrIS contribution to sea-level rise. To illustrate the influence of the ice sheet mask, we show that using a fixed ice sheet mask (as in NC) leads to a large overestimation of the contribution to SLR calculated from SMB−1.

4.5 Differences between 2-W and 1-W experiments.
Anomalies of ST, SMB and surface elevation for the averaged 2145-2150 period between the 2W and the 1W experiments present similar features (Fig. ??) than those obtained between the , we used the SMB outputs from the NF experiment at the MAR resolution and applied the integrated SMB method over the fixed observed ice-sheet mask \((\text{SMB}_{\text{MSC} \times \text{NE}})\) and over the updated 2W and the NC experiments, but with lower magnitude. Over the coastal regions, a larger increase in ST is obtained with 2-W as well as a lower SMB and a larger decrease in surface elevation (Fig. ??), hence highlighting the role of the feedbacks between the ice sheet and the atmosphere that are taken into account \((\text{SMB}_{\text{MSC} \times \text{2W}})\). Results reported in Table ?? indicate differences in SMB values exceeding 23% in 2-W but not in NC. As an example, the katabatic wind feedback preventing the penetration of warm air results in colder 2-W ST compared to 1-W. In 2150, the GrIS SLR contribution obtained in the 1-W experiment reaches +19.9 cm, i.e. 0.5 cm less than in the 2-W experiment (Fig. ??). Although this difference seems quite low, the local altitudes changes are larger in 2-W than in the 1-W experiment. Indeed, even if the median value of the ice thickness anomalies (2-W vs 1-W) between 2000 and 2150 are quite similar (respectively 73.3 m and 72.4 m), some regions show stronger surface anomalies (Fig. ??A). Scatter plots of surface elevation anomalies between 2-W and NC (red dots), 1-W and NC (green dots) and 2-W and 1-W (blue dots) as a function of the ice sheet altitude show the spatial variability of the ice thickness response to warming climate (Fig. ??A). The 2-W method yields negative and positive anomalies relative to NC, while the 1-W method mainly yields negative values (relative to NC). For both experiments, the regions where the ice thickness is under 1000 m are the most impacted by the warming climate. For these lower to medium altitude points, there is a strong variability of surface elevation anomalies. Thus, for the 2-W experiment, the anomalies range between +16.4 m (98% 2150). In the same way, compared to a time variable ice-sheet mask, the use of a fixed ice-sheet mask overestimates the sea-level rise by ~6% quantile value) and 43.1 m (2 in 2150. Though a bit lower, this number is far from being negligible compared to the errors made when the SMB-elevation feedbacks are not taken into account (i.e. 7.6% quantile value). For the 1-W experiment, the surface anomalies range between -1.5 m (98% quantile value) and -45.2 m (2% quantile value). Above 1000 m, the higher the altitude, the smaller the surface anomalies (Fig. ??A). The regions at low elevations are the most sensitive to the coupling method and to the warmer climate. This sensitivity to altitude increases with time, and is stronger for the 2-W experiment than for the 1-W experiment (Fig. ??B). High altitude regions are less sensitive to climate changes and to the coupling method used (Fig. ??B). These thickness changes are correlated with changes in ice ablation and ice accumulation area. After 150 years, the ELA shifts more inland in 2-W than in 1-W (Fig. ??), and the ice ablation area is 11% greater in 2-W than 1-W after 150 years—and when all the feedbacks are ignored (i.e. 9.3%). This strongly suggests that realistic SLR projections cannot neglect the evolution of the ice-sheet extent, only accounted for through the use of an ice-sheet model.

5 Discussion

5.1 Limits of the models

The evolution of the GrIS and its contribution to sea-level rise presented in this study are the first ones inferred from a regional atmospheric model synchronously coupled to an ice-sheet model, thus accounting for the GrIS-atmosphere feedbacks. To evaluate the added value of a coupled RCM-ISM model, we explored the importance of the GrIS-atmosphere feedbacks by
comparing the results of the coupled experiment to those coming from PF (Parameterised elevation feedbacks) and NF (No feedback) and showed that the impact of taking the feedbacks into account increases over time. This study is therefore a necessary first step toward a more accurate assessment of the contribution of Greenland to future sea-level rise and of its impact on the climate system. However, future refinements could be envisaged.

One of the main uncertainty in assessing the GrIS contribution to future sea-level rise comes from the climate projections themselves. For example, using five different global climate models to force MAR at its lateral boundaries under RCP8.5 conditions, Fettweis et al. (2013) provide SMB-inferred estimates of this contribution ranging from 4.6 to 13.1 cm in 2100. This range is fully comparable to that reported by who used MAR simulations (forced by three GCMs chosen from the Fettweis et al. (2013) sample) to force the SICOPOLIS ice-sheet model. Whatever the experimental design, the large spread in SLR projections raises the question as to whether the ice-sheet response simulated in our 2W experiment relative to that of the NF and PF experiments would be similar, amplified or mitigated with a different GCM climate forcing having a different sensitivity from MIROC5. A second question concerns the impact of a constant MIROC5 climate used to force MAR beyond 2100. As outlined in Sect. 3, this results in discarding the continued change that the climate will likely undergo beyond 2100 suggesting that our SLR projections are underestimated. The second consequence is that inter-annual variability is neglected after 2100. This can lead to conservative estimates of the Greenland contribution to sea level rise in the future due to non-linearities of the SMB. On the other hand, the imprint of the 2095 MIROC5 climate may amplify regional changes of the GrIS response. There is therefore a strong need for iterating the present study with different global climate simulations run under an extended RCP8.5 scenario and used as a MAR forcing, to assess more accurately the impact of the different GrIS-atmosphere feedbacks and to better evaluate the uncertainty associated with the projected sea-level rise contribution from the GrIS.

A second limitation is related to the 2000-yr relaxation GRISLI experiment, run at the end of the spin-up procedure to reduce the model drift in terms of ice volume, that produces residual differences with the observed topography (Bamber et al., 2013) used in the MAR simulations. This has important consequences on the MAR simulated climate. In particular, the steeper slopes existing in the GRISLI topography (i.e., $S_{ctvl}$) tend to produce unrealistic katabatic winds. Therefore, we choose to use an anomaly method of the surface elevation onto which the SMB and ST fields are downscaled at the 5 km resolution grid (Eq. 7). The objective of this approach was first to maintain the realism of the simulated present-day climate computed on the observed topography (Bamber et al., 2013) and, secondly, to avoid inconsistencies between the climate simulated by MAR and that used to force GRISLI. However, this implies that the forcing climate is not fully consistent with the GRISLI topography. This should be taken into consideration in a future work to improve the quality of our results. As an example, a reasonable compromise to avoid the use of anomaly method would be to use the topography obtained at the end of the spin-up iterative process (rather than $S_{ctvl}$) as initial GRISLI topography to keep the mismatch with the observed topography as low as possible, and to initialise and perform MAR simulations with this spin-up topography.

In addition, difference of resolution between MAR (25 km) and GRISLI (5 km grid resolution of GRISLI does not allow represents the smallest peripheral GrIS glaciers to be finely represented. This could limit our results: as we have shown that the coastal GrIS regions are the...), This can cause artefacts in the results especially at the edges of the ice sheet. Indeed, in
the corresponding MAR grid cells, a fraction of permanent ice cover may coexist with a non-zero fraction of tundra. Since the surface elevation changes computed in MAR from the aggregated GRISLI topography are weighted as a function of the fraction of the different surface areas, they may be underestimated as tundra soil type is not subject to any change in altitude. This artefact has been illustrated in Sect. 4.2.1 with the example of the behaviour of katabatic winds that are artificially reduced in our simulation at the ice-sheet margin. Moreover, since the margin regions are those experiencing the strongest changes in altitude, they are also the most sensitive to climate forcing; the GrIS contribution to SLR could be enhanced by increasing the spatial resolution of the ISM in change. As a consequence, an improper estimation of the topography changes may induce improper SMB changes. This underlines the need for increasing the atmospheric model resolution as far as possible to avoid such artefacts and to better represent the fine scale atmospheric-topography feedbacks impacting the SMB. Indeed, higher spatial resolution could resolve finer scale ice sheet dynamics to better represent the ice flow in outlet glacier or better represent fine scale atmospheric-topography feedbacks impacting the SMB in these regions. Furthermore, as we hypothesise an identical basal drag over time, we underestimate the acceleration of However, a compromise must be reached between the additional computing resources and the required degree of accuracy of sea-level projections.

Regarding the ice-sheet model, a 5 km horizontal resolution does not permit to capture the complex ice flow patterns of smallest outlet glaciers, whose characteristic length scale can be less than 1 km (?) and to quantify accurately the ice discharge at the ice flow of the glaciers due to marine front. This may have large implications in the sea-level rise estimates. Using a 3D ice-sheet model with prescribed outlet glacier retreat, found an additional SLR contribution from outlet glaciers of 0.8 to 1.8 cm in 2100 and 1.3 to 3.8 cm in 2200, with the basal lubrication coming from meltwater or rainfall percolation at depth and reaching the bedrock (Kulessa et al., 2017). An other limit of the GRISLI model is its simple representation of the grounding line position and thus of the buttressing effect which could impact the ice dynamics (Gagliardini et al., 2010). Except for these aspects, GRISLI is a good tool to be coupled in future ESMs in order to take the GrIS evolution into account with a combination of a good representation of the ice dynamics and a limited impact on influence of their dynamics on SLR projections decreasing with time and with the increasing importance of the added computational resources. As for the ISM, increasing the grid resolution of MAR, would allow to better represents atmospheric topography feedbacks and more complex atmospheric processes which can have an impact on the SMB in the steep coastal regions. However, as for ice dynamic modeling, the higher is the resolution, the higher is the computational resources needed to produce results. The absence of any representation of the GrIS ocean feedbacks is also a limiting factor. Indeed, as the GrIS is an island, several glaciers are in direct contact with the ocean and feedbacks could take place between peripheral glaciers and the ocean. The warm ocean water could accelerate the melting of the glaciers and the added fresh water in the ocean could in turn, have an impact on sea surface atmospheric forcing. This is in line with the fact that ice dynamics act to counteract ice loss from surface melting (see Sect. 4.2), as previously outlined by several authors (Edwards et al., 2014a; ?, ?). However, despite the possible decreasing influence of marine terminating glaciers, at the centennial time scale, it seems to be preferable to evaluate more accurately the impact of ice dynamics and to better capture the complex geometry of fjords surrounding the marine-terminating glaciers.

There is a growing number of evidence for attributing the acceleration of outlet glaciers to the intrusion of warm waters from adjacent oceans in the fjord systems or in the cavity of floating ice tongues (e.g., ???) that can destabilise the glacier front.
and/or favour the ice-shelf breakup (Gagliardini et al., 2010), decreasing thereby the buttressing effect and increasing the ice calving. In turn, the released freshwater flux in ocean may impact sea-surface temperatures, oceanic circulation and sea-ice cover. The GrIS and the atmosphere evolution could be both modified by this added fresh water flux in the ocean system.

5.1 Limits of the methods

The 1-W coupling method neglects the spatial variability of the thickness anomaly and underestimates regional feedbacks compared to the 2-W method. These differences are explained by the linear relationship used, in the 1-W coupling method, to correct the atmospheric fields (SMB and ST) as a function of. Moreover, atmosphere-ocean feedbacks also have an impact on the GrIS. As an example, Fettweis et al. (2013) showed that the disappearance of Arctic sea ice in summer induced by ocean warming enhances surface melting in northern Greenland through a decrease of surface albedo and the surface elevation anomaly, as developed by Franco et al. (2012). Indeed, the relationship between the atmosphere and the GrIS changes is nonlinear because surface elevation changes interact not only with both SMB and ST, but also with all the other atmospheric fields which influence the SMB or the ST directly or indirectly, as for example the winds, the humidity, and the albedo. The 2-W method appears to be subsequent atmospheric warming. Thus, the best way to simulated atmospheric GrIS feedbacks. If the main objective is to compute the SLR contribution from the entire GrIS without investigating atmospheric or GrIS changes at absence of the oceanic component in our modelling setup appears as a limiting factor, although, the direct impact of ocean via sub-shelf melt at the ice sheet margin will likely be limited in the future as a result of inland retreat of GrIS.

Our spin-up method adjusts the basal drag coefficient in such a way that the departure between the observed and the initial GRISLI topographies is reduced. The resulting $\beta$ coefficient is spatially varying but is constant in time. This assumption may likely be valid for short-term forward simulations but is probably overly simplistic. On the one hand, the basal drag tends to be smaller towards the margins with respect to the interior. As the ice sheet retreats inland, it can be expected a reduction in basal drag for a specific location, due for example to a decreasing effective pressure. On the other hand, changing basal hydrological conditions can also alter the basal drag. This can occur as a result of rainfall or surface meltwater infiltration that can refreeze at depth or propagate all the way to the bottom of the ice sheet and increase basal lubrication (Kulessa et al., 2017). Therefore, a time constant basal drag coefficient inferred under present-day conditions may underestimate the ice flow acceleration. A few models describing the vertical inflow exist (e.g., ???) but are generally run at the regional scale and at very high spatial resolution (a few tens to a few hundreds of meters at most). Implementing such models in large-scale ice-sheet models is currently outside the realm of possibilities. However, as there is a growing interest in performing ice-sheet projections over multi-centennial time scale, the regional scale, the use of the 2-W coupling method with a high resolution seems avoidable until 2100. However, over longer time scales, or to study more regional processes changes, the use of a the 2-W coupling method is necessary to represent the local feedbacks between the atmosphere and the GrIS fields and ensure that the SLR contribution is not underestimated by simulating. As an example, the changes in the GrIS extent and ice surface slope have a direct impact on surface albedo and strength of katabatic winds. Although the difference in the GrIS melting contributions to SLR between 1-W and 2-W seems low, the use of the 2-W method to compute the ice sheet evolution for 50 additional years (from 2100 to 2150) with the constant forcing of the year 2095 contributes to increase the ice mass loss contribution to SLR by +0.5 cm compared
In this study, we have improved the representation of the interactions between the GrIS and the atmosphere by developing a full coupling between the Greenland ice sheet model GRISLI and an additional limitation related to the atmospheric model MAR (2 W experiment). To assess the importance of this improvement, we have investigated the atmosphere and ice sheet responses to the RCP 8.5 warming climate scenario, and we have compared the 150 years of our fully coupled experiment choice of our spin-up procedure is that the glacial-interglacial signature of past climatic changes is ignored. Neglecting the climate history of the Greenland ice sheet implies too warm ice temperatures. This may have an impact on the future GrIS evolution and on its contribution to sea-level rise. Indeed, the basal drag coefficient inferred from the inverse method may be too high so as to compensate the errors induced by the artificial warming. However, using a higher-order ice flow model, showed that at the centennial time scale the basal conditions and the GrIS projections are only poorly sensitive to the initial vertical temperature profile but are critically dependent on atmospheric conditions.

Despite these limitations, the sea-level projections performed with GRISLI compare well with those conducted with more sophisticated ice-sheet models Edwards et al. (2014a) and the simulated surface ice velocities present a good agreement with the observed ones (Fig. S11). It appears thus as a good numerical tool to be coupled with a regional climate model with a reasonably good representation of the ice dynamics and a limited computational resources.

6 Conclusions

This study is based on the first regional atmospheric – ice-sheet coupled model allowing the GrIS-atmosphere feedbacks to be accounted for. Using this new model, we investigated the GrIS evolution and its contribution to sea-level rise from 2000 to AD 2150 under a prolonged RCP8.5 scenario (2 W) 2 W experiment). The importance of the GrIS-atmosphere feedbacks has been assessed through the comparison of the two-way coupled experiment with two other experiments using a less complex coupling method (1 W) and a no coupling at all (NC). The fully coupled approach under the RCP 8.5 scenario produces a GrIS melting contribution to SLR of simulations based on simpler coupling strategies: the NF experiment in which the MAR outputs are directly used as GRISLI forcing and the PF experiment in which the elevation feedbacks are parameterized. In both NF and PF experiments, changes in topography simulated by GRISLI are not updated in the atmospheric model. The main conclusions drawn from this study are the following:

- Accounting for the GrIS-atmosphere feedbacks amplifies the ice mass loss and changes in ice-sheet geometry with increased surface slopes from the central regions to margin areas. In turn, changes in the shape of Greenland modify the ice velocities.

- The effect of accounting the feedbacks between GrIS and the atmosphere increases over time and becomes significant at the end of the 21st century, as illustrated by a GrIS contribution to sea-level rise of 20.4 cm in 2150, while the 1 W and
NC methods produce a GrIS contribution to SLR respectively of +19.9 cm and +18.4 cm, respectively. The difference of 0.5 cm between the 2-W and the 1-W methods represents at least 25% of the contribution of the peripheral Greenland glaciers melting estimated for the next 100 years using the same RCP 8.5 scenario (Radić et al., 2014; Machguth et al., 2013). This difference, increasing with time, is mainly explained by representation of local interactions between the GrIs and the atmosphere, only possible with the 2-W method. Furthermore, even if the difference is not perceptible in 2100 and it is low in 2150, we have shown that the ice loss computed from the integration of the SMB against 7.9 cm only in 2100.

- Accounting for the parameterized elevation feedbacks in the PF experiment leads to an additional SLR contribution of \( \sim 7.6\% \) in 2150 compared to NF. On the other hand, the parametrization used in PF allows to reduce the mismatch (in terms of SLR projections) between the one-way and the two-way coupled approaches by 73.7% in 2150, showing that at this time scale, changes in ice-sheet geometry appear to be dominated by the SMB-elevation feedback.

- Finally, we showed that estimating the GrIS contribution over a fixed ice sheet mask is 21% (as in PF and NF experiments) overestimates the SLR contribution by \( \sim 6\% \) higher than that obtained with the use of an evolving ice sheet mask. This means that, suggesting that most of RCM-based studies have probably overestimated the ice loss computed from a change in SMB. However, with the 5 km grid-resolution of GRISLI, we cannot reproduce the fine scale structure of the Greenland coast and glaciers. Using an ice sheet model with higher resolution and more complex physics (i.e. Full-Stokes models) and a fully coupled method would probably amplify the sensitivity of these coastal regions. This argument is also valid for the atmospheric model for which a higher resolution would be beneficial for the representation of fine scale atmospheric processes over the ice sheet. We showed that it is at small spatial scales that the coupling method makes most difference. It would therefore be very interesting to find the optimal resolution of the ice sheet and the atmospheric model, for ISM-RCM coupling. Furthermore, since the Greenland ice sheet and glaciers are in contact with the oceanic component, changes in oceanic characteristics, due to the input of freshwater from GrIS melting or due to the warming climate scenario, could in turn disrupt the GrIS and atmosphere evolution. The next step of this study will be therefore to improve the representation of the oceanic component by developing a fully coupled method between an ISM, an RCM and an oceanic model to evaluate the impacts on the Greenland polar region but also on remote regions.

7 Data availability

The model outputs from the simulations described in this paper are freely available from the authors upon request without conditions. The source code of MAR version 3.7 is available on the MAR website: http://mar.cnrs.fr.

Author contributions. The implementation of the three coupling methods as well as the simulations were done by X. Fettweis and C. Wyard. S Le clec’h, S. Charbit and A. Quiquet analysed the results and wrote the manuscript with contributions from M. Kageyama, C. Dumas and X. Fettweis. The GRISLI model was developed by C. Ritz.
Competing interests. The authors declare that they have no conflict of interest.

Acknowledgements. The authors are very grateful to J. Fyke and two anonymous reviewers for their numerous and fruitful comments that helped to improve the writing of the manuscript. S. Le clec'h, M. Kageyama, S. Charbit and C. Dumas acknowledge the financial support from the French-Swedish GIWA project and the ANR AC-AHC2, as well as the CEA for the S. Le clec'h PhD funding. A Quiquet is funded by the European Research Council grant ACCLIMATE no 339108. Computational resources (MAR and GRISLI) have been provided by the Consortium des Équipements de Calcul Intensif (CÉCI), funded by the Fonds de la Recherche Scientifique de Belgique (F.R.S.–FNRS) under grant no. 2.5020.11 and the Tier-1 supercomputer (Zenobe) of the Fédération Wallonie Bruxelles infrastructure funded by the Walloon Region under the grant agreement no. 1117545.
References


