

Brief communication: Two well marked cases of aerodynamic adjustment of sastrugi

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Abstract

In polar regions, sastrugi are a direct manifestation of drifting snow and form the main surface roughness elements. In turn, sastrugi alter the generation of atmospheric turbulence and thus modify wind profiles and aeolian snow mass fluxes. Little attention has been paid to these feedback processes, mainly because of experimental difficulties, and, as a result most polar atmospheric models currently ignore sastrugi. This paper focuses on two cases during which sastrugi responses to shifts in wind direction were evidenced by variations in aeolian snow mass fluxes and neutral stability 10-m air-snow drag coefficients C_{DN10} computed from wind profiles collected during austral winter 2013 in coastal Adélie Land, East Antarctica. Using this dataset, it was shown that (i) C_{DN10} values were in the range of $1.3\text{--}1.5 \times 10^{-3}$ when the wind was well aligned with the sastrugi and could increase to nearly 3.3×10^{-3} with wind shifts of only 20–30°, (ii) as C_{DN10} increases, the aeolian snow mass flux can decrease (to 80%) in response to a shift in wind direction, (iii) the timescale of sastrugi aerodynamic adjustment can be as short as 3 h for friction velocities of 1 m s^{-1} or above and during strong drifting snow conditions, and (iv) knowing C_{DN10} is not sufficient to estimate the erosion flux that results from drag partitioning at the surface because C_{DN10} includes the contribution of the sastrugi form drag.

1. Introduction

In polar regions, sastrugi are a direct manifestation of drifting snow. They are generally regarded as elongated ridges of wind-packed snow 1 to 2 meters in length whose longitudinal

axis is parallel to the prevailing wind at the time of their formation. These erosional surface roughness features are very widespread over the Antarctic ice sheet (Kotlyakov 1961) where they can be major determinants of surface roughness (Jackson and Carroll 1978; Inoue 1989; Andreas and Claffey 1995). Sastrugi orientations have been recognized as useful indicators of the Antarctic near-surface wind direction (Mather 1962, 1969; Mather and Miller 1966; Long and Drinkwater 2000) in agreement with continent-scale modeling studies (Parish and Bromwich 1987, 2007).

The development of sastrugi depends on the ability of snow to be eroded and thus on the threshold velocity needed to lift snow particles from the surface. In the literature, aeolian erosion thresholds have been reported to vary depending on temperature and diverse properties of surface snow. From observations in Antarctica, Mellor (1965) reported that 10-m wind speeds of 3 to 8 m s⁻¹ are strong enough to cause aerodynamic entrainment of loose, unbounded snow, whereas winds exceeding 30 m s⁻¹ are needed to erode snow consolidated by the freeze-thaw process. Budd et al. (1966) suggested a high threshold wind speed (14 m s⁻¹) was needed to trigger snow transport in the cold environment of Byrd station. Schmidt (1980) reported that the cohesion of the snow surface determines the threshold speed required for snow erosion to occur. Schmidt (1980) also showed that the threshold wind speed increases with time since snow deposition, and that this increase slows with time and is slower at lower temperatures. Pomeroy et al. (1993) identified significantly lower thresholds for fresh, loose, dry snow than for older, wind hardened, dense or wet snow. Ôura et al. (1967) and, later on, Li and Pomeroy (1997) discussed the major role of temperature in surface erodibility (i.e. the potential of a surface to be eroded; Shao 2008) through metamorphism of snow (i.e. changes in snow structure over time), and showed an empirical but generally positive correlation between threshold wind speed and air temperature. All studies suggest that the physical properties of the snow play a major role in the formation of sastrugi.

Sastrugi contribute to the drag exerted on the atmosphere over the snow surface and enhance interactions at the air-snow interface compared to over a smooth snow surface. Rougher snow surfaces favor the generation of turbulence in the near surface air stream that is likely to further increase the aeolian snow mass flux (Das et al. 2013). On the other hand, sastrugi are responsible for a loss of wind momentum through pressure fluctuation gradients in their immediate vicinity (sastrugi form drag) that directly reduces the energy budget available for erosion of snow. This attenuating effect on snow erosion is taken into account in the coupled atmosphere-snowpack-aeolian snow transport model MAR (Gallée et al. 2013) and was parametrized as in Marticorena and Bergametti (1995). By comparing observed and simulated

1 aeolian snow mass fluxes over Adélie Land using MAR, Amory et al. (2015) showed that in
2 the model, erosion efficiency is highly sensitive to the parameterization of surface roughness,
3 and underlined the need for observational characterization of interactions between wind-
4 induced roughness features and aeolian transport of snow.

5 Some authors have shown that the sastrugi form drag actually depends on how the wind is
6 oriented with respect to the main sastrugi axis. Based on measurements of wind speed and
7 temperature profile in the atmospheric surface layer at the South Pole, Jackson and Carroll
8 (1978) reported that sastrugi form drag was essentially absent when the wind was perfectly
9 aligned with the sastrugi up to a height of 50 cm. As the wind rotated, sastrugi form drag
10 increased, to reach maximum when the wind direction was perpendicular to the prior sastrugi
11 pattern. These authors developed an idealized single sastruga model from Lettau's (1969)
12 findings to reproduce their observations. Using another analytical sastruga model adapted from
13 Raupach (1992), Andreas (1995) also found a minimum and a maximum drag for wind
14 directions respectively parallel and perpendicular to the sastruga longitudinal axis. However,
15 these modeling efforts were undertaken without accounting for the erodible character of
16 sastrugi or for their possible reorganization when realigning with persistent (erosive) winds
17 blowing transversally to their elongated sidewalls. If the crosswise flow continues from a
18 relatively constant direction thereby allowing sufficient shear stress to dislodge snow surface
19 particles, sastrugi can adjust aerodynamically; transversal sastrugi are eroded, and new
20 streamlined sastrugi form parallel to the mean wind (Andreas and Claffey 1995). This results
21 in a gradual decrease in the contribution of the sastrugi to the total surface drag, and hence in
22 an increase in erosion efficiency. Andreas and Claffey (1995) reported that the timescale for
23 this streamlining process on Weddell Sea ice in winter was about half a day with $6\text{--}8\text{ m s}^{-1}$
24 winds, but might be shorter if the winds are stronger. To date, no observational study has
25 provided quantitative insight into the potential effect of erodible roughness elements of the
26 snow surface on snow erosion.

27 Quantifying the variable influence of sastrugi on the local wind field and associated surface
28 drag could improve parameterization of surface roughness and erosion in polar atmospheric
29 models that currently ignore sastrugi. The present paper focuses on two erosion events during
30 which sastrugi responses to shifts in wind direction were interpreted from temporal variations
31 in both measured drag and aeolian snow mass flux in coastal Adélie Land during austral winter
32 2013.

33 2. Data and Method

2.1. Field area

Site D17 (66.7°S, 139.9°E; ~450 m asl.) is located about 10 km inland in a coastal accumulation zone of Adélie Land (Agosta et al. 2012), roughly 15 km southwest of the permanent French station Dumont d'Urville (Fig. 1). An annual temperature of -10.8 °C and a mean wind of around 10 m s⁻¹ have been reported at Dumont d'Urville station (König-Langlo et al. 1998). The measurement area consists in a gently sloping snowfield with a long unobstructed upstream fetch several hundred kilometers over a uniform snow surface. Local topographic channeling acts together with the Coriolis force to produce southeasterly flows all year round that result either from pure katabatic or combined katabatic-synoptic forcings (Parish et al. 1993). Site D17 is visited only during summer (December to February), when the presence of sastrugi is often reported. Frequent strong winds combined with the permanent snow surface lead to frequent aeolian snow transport events (Trouvilliez et al. 2014), thereby favoring aerodynamic adjustment of the snow surface. This results in a net south-southeast orientation of the sastrugi (Fig. 2).

2.2. Instrumentation

The measurement structure deployed at site D17 is a 7-m high meteorological mast. Wind speed, relative humidity and air temperature are recorded along the mast at 6 logarithmically spaced intervals between 0.8 and 7 m above the snow surface using Vector A100LK cup anemometers and HMP45A thermo-hygrometers installed in naturally ventilated MET21 radiation shields. The anemometers are mounted on roughly 1-m long booms pointing southeastward. Wind direction is only sampled at the upper level by a Vector W200P wind vane. Surface level variations are measured by a Campbell SR50A acoustic depth gauge. Information on drifting snow is obtained from a second-generation acoustic FlowCapt™ device that was set up vertically close to the ground to allow detection of the beginning of aeolian snow transport events. The sensor is a 1-m long tube that converts the acoustic pressure caused by snow particles impacting the tube into an aeolian snow mass flux integrated over the length of the tube. The second-generation FlowCapt™ was evaluated in the French Alps by Trouvilliez et al. (2015). The authors reported that the instrument underestimates the aeolian snow mass flux compared to a reference optical sensor (Snow Particle Counter S7; Sato et al. 1993), especially during snowfalls. Nevertheless, the equivocal behavior of the second-generation FlowCapt™ does not affect its ability to accurately detect the occurrence of aeolian snow

1 transport. Data were sampled at 15 s intervals, averaged to half-hourly means and stored in a
2 Campbell CR3000 datalogger.

3 **2.3. The 10-m drag coefficient in near-neutral conditions**

4 Computing the drag coefficient (C_D) is a convenient way to estimate the local drag exerted by
5 the surface on the overlying air. C_D can be computed by measuring the vertical wind speed
6 gradient (profile method) under near-neutral conditions following the Monin-Obukhov
7 similarity theory. Assuming stationarity and horizontal homogeneity when the atmospheric
8 surface layer is statically neutral, the wind speed profile is logarithmic and can be written as

$$U(z) = \frac{u_*}{\kappa} \ln\left(\frac{z}{z_0}\right), \quad (1)$$

9 where $U(z)$ is the average wind speed as a function of height z , κ is the von Kármán constant
10 (taken as 0.4), z_0 is the aerodynamic roughness length, and u_* the friction velocity describing
11 the wind shear at the surface and being related to the vertical momentum flux at the surface (τ ;
12 also known as Reynolds shear stress)

$$\tau = \rho u_*^2 = -\rho \overline{uw} = \rho C_{DNz} U_z^2, \quad (2)$$

13 where ρ is the air density, u and w are fluctuations in the longitudinal and vertical turbulent
14 velocity, respectively, and C_{DNz} and U_z are the neutral-stability drag coefficient and the average
15 wind speed at height z , respectively. The overbar stands for a time average. C_{DN} is usually
16 discussed at a standard reference height of 10 m (C_{DN10}). From (2) and (3), it follows that

$$C_{DN10} = \left[\kappa / \ln\left(\frac{10}{z_0}\right) \right]^2, \quad (3)$$

17 with z_0 expressed in meters. Here C_{DN10} and z_0 are two equivalent quantities for evaluating the
18 momentum exchange at the air-snow interface that results from the integrated (in space and
19 time) turbulent drag caused by the roughness elements.

20 The wind profiles used to compute C_{DN10} were selected following a strict procedure. After
21 discarding icing or malfunctioning cases and half-hourly runs for which a rare (northwesterly)
22 flow was likely to be disturbed by the measurement structure, stationary conditions were
23 selected by requiring that temperature changes between two consecutive half-hourly runs not
24 exceed 0.3 K, following Joffre (1982)'s recommendations. Near-neutral conditions were then
25 selected requiring $U > 5 \text{ m s}^{-1}$ and an absolute value of the bulk Richardson number below 10^{-1}

². The last selection criterion was applied following a suggestion by Andreas and Claffey (1995) that demands

$$\frac{\sum_{i=1}^6 [U(z_i) - (u_*/\kappa) \ln(z_i/z_0)]^2}{u_*^2} \leq \varepsilon, \quad (4)$$

where ε is an empirical constant determined from visual inspection of the observed wind speed profiles. Here it was set to 0.15. Wind profiles that survived this filtering process were fitted (1) using a least-square log-linear regression technique, and u_* and z_0 deduced from the regression coefficients. All of them yielded a correlation coefficient (r^2) larger than 0.99. The 80% confidence limits of each calculated C_{DN10} value were determined following the statistical method proposed by Wilkinson (1984). The highest uncertainty bounds deduced from these confidence limits reached $\pm 14\%$.

3. Results

The two erosion events depicted in Figure 3 occurred respectively in March (left panels) and October (right panels), 2013, during particularly constant wind direction conditions, which persisted after a wind shift of a few tens of degrees. Such a constancy in wind direction, necessary for the following demonstration, is very rare. Combined to the strict selection procedure, only two cases were exploitable in this context. The 2-m wind speed, wind direction, profiled-derived C_{DN10} values and aeolian snow mass flux recovered by the second-generation FlowCapt™ sensor are shown in Figure 3. As the friction velocity is the actual dynamic quantity involved in aerodynamic entrainment of surface snow particles, it is also plotted on the graph. The two events are split into three parts, before (A_i), during (B_i) and after (C_i) the shift in wind direction. The occurrence of precipitation may affect the detection of erosion events because the FlowCapt™ sensor does not distinguish between eroded (saltating particles and/or suspended particles of snow) and precipitating snow particles. No visual observation of precipitation from the nearby Dumont d'Urville station were available for the period concerned. Moreover, as Adélie Land is very prone to aeolian transport of snow, these observations, if performed, are limited by the inability to discriminate between actual precipitation and pure drifting snow. Here we used the operational analyses of the European Center for Medium-Range Weather Forecasts (horizontal resolution of ~ 16 km) to evaluate the occurrence of precipitation at our measurement site. We assumed that both events were pure erosion events

1 after finding negligible precipitation rates for the fully continental grid point including D17.
 2 At the beginning of Julian day (JD) 87 (part A₁), the wind direction was around 140°, the
 3 friction velocity was above the erosion threshold with a related aeolian snow mass flux of 100
 4 g m⁻² s⁻¹, and C_{DN10} was near 1.5 x 10⁻³. At the end of JD 87 (part B₁), the wind rotated toward
 5 160° while C_{DN10} increased to nearly 3.3 x 10⁻³, i.e. by 120%, in response to a wind shift of
 6 only 20°. As assumed in Jackson and Carroll (1978) and Andreas and Claffey (1995), it is likely
 7 that as the wind turned, it was deflected from the mean sastrugi axis, thereby encountering a
 8 rougher surface. As a result, C_{DN10} soared, reflecting the growing contribution of the sastrugi
 9 form drag to the vertical momentum flux at the surface, and hence to the total surface drag.
 10 Within the same time frame, the measured aeolian snow mass flux fell by ~30% from 365 to
 11 260 g m⁻² s⁻¹, despite increasing friction velocity (wind speed) from 0.7 to 1.6 (18 to 24) m s⁻¹.
 12 Then, until the end of the event (part C₁), the wind direction remained centered about 160°.

13 From 0330 UT to 0630 UT on JD 88, C_{DN10} fell back to 1.5 x 10⁻³ as high winds presumably
 14 streamlined the surface. In other words, C_{DN10} was reduced by ~50% in only 3 hours. As C_{DN10}
 15 decreased, the aeolian snow mass flux again rose above 400 g m⁻² s⁻¹. The erosion event lasted
 16 through JD 90 when u_{*} (wind speed) dropped to 0.7 (15) m s⁻¹, causing a significant decrease
 17 in the aeolian snow mass flux. After nearly 48 hours of persistent erosive winds, C_{DN10} was as
 18 low as 1.3 x 10⁻².

19 During the two days that preceded the second erosion event (part A₂), the wind direction was
 20 within ± 10° of 150°, the friction velocity was generally not strong enough to erode the snow
 21 surface, and C_{DN10} was between 1.3 – 1.6 x 10⁻³. C_{DN10} and wind direction were strongly
 22 correlated during this period, with the lowest drag coefficients occurring for a wind direction
 23 of around 140°, suggesting that this was the sastrugi alignment before erosion started and the
 24 wind changed direction. Then, the same situation depicted in the left panels of Figure 3 occurred
 25 again. At mid-JD 286 (part B₂), u_{*} increased beyond the erosion threshold as the wind rotated
 26 from 150° to 180°. Consequently, C_{DN10} increased to 1.9 x 10⁻³. The aeolian snow mass flux
 27 dropped simultaneously from 320 to 55 g m⁻² s⁻¹ under increasing friction velocity. That is, for
 28 a ~30% increase in C_{DN10} as the result of a wind deflection of 30°, the aeolian snow mass flux
 29 decreased by ~80%. Together with the first case of erosion, this illustrates how the form drag
 30 exerted by sastrugi can significantly affect snow erosion when the wind and sastrugi are not
 31 aligned (this effect is discussed later in the paper; see Section 4). Then (part C₂), the wind
 32 direction remained roughly unchanged until erosion ceased. Again, the rise in aeolian snow

mass flux coincided with a decrease in C_{DN10} . After nearly 3 hours of winds above 20 m s^{-1} ($u_* > 0.9 \text{ m s}^{-1}$) from 180° , C_{DN10} fell from 1.9×10^{-3} to 1.4×10^{-3} , i.e. decreased by $\sim 30\%$. In summary, for friction velocities (wind speeds) around 1 (20) m s^{-1} and above, the sastrugi streamlining timescale can be as fast as 3 hours. For a windflow initially aligned with the sastrugi, a deviation of $20\text{-}30^\circ$ from the streamlining direction has the potential to both increase C_{DN10} by $30\text{-}120\%$ and to significantly reduce (up to 80%) the aeolian snow mass flux, even under increasing friction velocity.

4. Discussion

At Ice Station Weddell, Andreas and Claffey (1995) measured a decrease in C_{DN10} of $20\text{-}30\%$ in 12 hours with considerably weaker winds ($< 12 \text{ m s}^{-1}$) than those reported here. The observations reported in this paper show that this timescale can be 4 times faster for winds exceeding 20 m s^{-1} ($u_* > 1 \text{ m s}^{-1}$), and the associated decrease in C_{DN10} can reach 50% . Andreas and Claffey (1995) also proposed generic C_{DN10} values in the range $1.5\text{-}1.7 \times 10^{-3}$ when the wind is well aligned with the sastrugi, and around 2.5×10^{-3} when the wind is at an angle of 20° to the dominant orientation of the sastrugi. The present results differ slightly from these values: C_{DN10} was more in the range $1.3\text{-}1.5 \times 10^{-3}$ for sastrugi-parallel winds, and could increase to more nearly 3.4×10^{-3} with a wind shift of similar amplitude. For a given erosion threshold, the quantity of windborne snow increased with wind strength according to a power law (Mann 2000). As sastrugi mainly form through snow erosion/deposition processes (Filhol and Sturm 2015), it is likely that under the strong wind (shear) conditions in Adélie Land, rougher snow surfaces develop, whose aerodynamic adjustment ability is greater than at the less windy Ice Station Weddell.

During both erosion events, significant aeolian snow mass fluxes were measured for 2-m wind speeds (u_*) of 10 (0.6) m s^{-1} or above. As the wind (friction) velocity likely frequently exceeds this threshold on the coastal slopes of Adélie Land, the sastrugi alignment process might be also frequently active, depending on persistence of the wind. As explained in Section 1, this mechanism is probably also strongly controlled by the properties of the snow surface that determine the threshold shear stress required for erosion to begin rather than only the characteristics of the wind. Since the erosion flux is the integrated result of both the capacity of the wind to erode and carry snow, and snow surface erodibility, the sastrugi streamlining timescale presumably mostly depends on this specific quantity. The implication is that the drag coefficient must be strongly related to other factors including the current wind orientation and

1 the history of the wind's interactions with the snow surface as well as past timescales and past
2 temperatures of the snowpack.

3 On the other hand, the sastrugi streamlining timescale also appears to control snow erosion in
4 the form of feedback by fixing the time during which the sastrugi form drag mainly contributes
5 to total surface drag. With friction velocities above the snow erosion threshold, increasing u_*
6 could be expected to result in an increase in erosion efficiency. However, in both cases, the
7 observations showed a significant decrease in the aeolian snow mass flux in phase with an
8 increase in the drag coefficient (Figure 3, parts B). By analogy with measurements made in a
9 water flume (Wiberg and Nelson 1992; Le Bouteiller and Venditti 2015), it can be considered
10 that the flow and turbulence in the sastrugi region are the result of interaction between flow
11 separation and wake formation, which can lead to a local Reynolds shear stress peak
12 corresponding to flow separation. Above the region of influence of the wake, named outer
13 region, the flow has adjusted to increased roughness and exhibited a logarithmic profile, as
14 shown by the relative continuous time series of C_{DN10} and u_* (Fig. 3) despite the strict selection
15 procedure. Even if the shear stress of the outer flow (τ) is relatively easy to measure, it cannot
16 be extrapolated to the snow bed. The averaged snow bed shear stress (also referred to as skin
17 friction in the literature), which is the ultimate parameter for aeolian erosion (Li and Shao
18 2003), varies depending on its position along the sastrugi field. In absence of direct
19 measurements, it is necessary to link outer shear stress, sastrugi geometry, and skin friction to
20 be able to estimate aeolian snow mass fluxes. This is quite important since the reduction of
21 shear stress near the surface is crucial in limiting the growth of the mass flux (Groot Zwaafink
22 et al., 2014). For erodible forms in riverbeds such as ripples, Smith and McLean (1977) and
23 later Wiberg and Nelson (1992) developed a method for partitioning the outer shear stress.
24 These authors considered that the averaged bed shear stress is equal to the difference between
25 the outer shear stress and the drag-related stress produced as the flow is forced around the
26 bedform – i.e., in the present case, the form drag induced by the sastrugi. As mentioned above,
27 an increasing form drag can be expected, and hence a decrease in skin friction and in aeolian
28 snow mass flux, when the wind direction gradually shifts away from the longitudinal axis of
29 the sastrugi. Because C_{DN10} reflects the contribution of the sastrugi form drag, knowing the
30 drag coefficient is not sufficient to estimate skin friction. A better knowledge of skin friction
31 over a sastrugi field is also needed to improve aeolian snow mass flux parameterizations in
32 aeolian erosion models. The measurements made in the present study showed that a
33 considerable decrease (even 80%) of the aeolian snow mass flux can occur during the
34 transitional regime during which the wind and sastrugi are not aligned (Figure 3, parts B). But

it should be also noted that the rapid aerodynamic adjustment of sastrugi (3 hours) will limit errors if the erosion event considered is strong and sufficiently long.

5. Conclusion

An experimental meteorological dataset collected in coastal Adélie Land during austral winter 2013 was exploited to document surface turbulent fluxes of momentum and snow over an Antarctic sastrugi field. The main results of the analysis of two erosion events can be summarized as follows:

- C_{DN10} values were in the range of $1.3-1.5 \times 10^{-3}$ when the wind was well aligned with sastrugi and could increase to nearly 3.3×10^{-3} with wind shifts of only $20-30^\circ$,
- As C_{DN10} increases, the aeolian snow mass flux may decrease (to 80%) in response to the wind shift in direction,
- the timescale for the aerodynamic adjustment of sastrugi can be as low as three hours for friction velocities of 1 m s^{-1} or above and during strong drifting snow conditions,
- because C_{DN10} includes the contribution of the sastrugi form drag, knowing C_{DN10} is not sufficient to estimate the erosion flux that results from drag partitioning at the surface.

These results support the existence of feedback mechanisms linking aeolian erosion and surface drag properties over (Antarctic) snow surfaces, as already demonstrated for other erodible natural surfaces (Marticorena and Bergametti 1995). In contrast with non-erodible roughness elements such as rocks or vegetation, these mechanisms involve the time needed for sastrugi to adjust to the main wind (3 hours in both erosion events), during which both the drag coefficient and the aeolian snow mass flux can be greatly modified. In comparison, Andreas and Claffey (1995) reported a longer timescale (12 hours) for the sastrugi to realign with weaker winds. Because lighter winds are supposed to be associated with lower erosion fluxes, it is suggested that the sastrugi streamlining timescale most likely depends on the snow erosion flux.

Real-time observations of the distribution (size, abundance, orientation) of the sastrugi would further advance understanding of the physical processes involved in the development of sastrugi and enable better characterization of sastrugi aerodynamic adjustment timescales. In addition, having a more accurate representation of the distribution of sastrugi would make small-scale modeling in a wind tunnel possible, in which case, it would be possible to realistically estimate shear stress partitioning. One possible way to monitor sastrugi would be to set up an automatic

mini laser-scan. Such a device was developed in the framework of the MONISNOW research project (Picard and Arnaud, LGGE, personal communication) and has been operating daily at Dome C in Antarctica since the beginning of 2015. These complementary approaches are vital to improve parameterization schemes for aeolian snow transport models and general drag parameterizations for weather, climate and earth system models.

References

Amory, C., Trouvilliez, A., Gallée, H., Favier, V., Naaïm-Bouvet, F., Genthon, C., Agosta, C., Piard, L., and Bellot, H., 2015. Comparison between observed and simulated aeolian snow mass fluxes in Adélie Land, East Antarctica. *The Cryosphere*, 9, 1–12.

Andreas, E. L., 1995. Air-ice drag coefficients in the western weddell sea. 2. A model based on form drag and drifting snow. *J. Geophys. Res.*, 100(C3), 4833–4843.

Andreas, E. L. and Claffey, K. J., 1995. Air-ice drag coefficients in the western weddell sea. 1. Values deduced from profile measurements. *J. Geophys. Res.*, 100(C3), 4821–4831.

Budd, W. F., 1966. The drifting of non-uniform snow particles. *Studies in Antarctic Meteorology*, M. J. Rubin, Ed., Antarctic Research Series, Vol. 9, American Geophysical Union, 59-70.

Das, I., Bell, R. E., Scambos, T. A., Wolovick, M., Creyts, T. T., Studinger, M., Frearson, N., Nicolas, J. P., Lenaerts, J. T. M., and van den Broeke, M. R., 2013. Influence of persistent wind scour on the surface mass balance of Antarctica, *Nat. Geosci.*, 6, 367–371.

Filhol, S., and Sturm, M., 2015. Snow bedforms: A review, new data and a formation model. *J. Geophys. Res.*, 120, 9, 1645-1669. doi: 10.1002/2015JF003529.

Gallée, H., Trouvilliez, A., Agosta, C., Genthon, C., Favier, V., and Naaïm-Bouvet, F., 2013. Transport of Snow by the Wind: A Comparison Between Observations in Adélie Land, Antarctica, and Simulations Made with the Regional Climate Model MAR. *Boundary-Layer Meteorol.*, 146, 133–147.

Groot Zwaaftink, C. D., Diebold, M., Horender, S., Overney, J., Lieberherr, G., Parlange, M., and Lehning, M., 2014. Modelling small-scale drifting snow with a Lagrangian stochastic model based on large-eddy simulations. *Boundary-Layer Meteorol.*, 153, 117–139,

- 1 doi:10.1007/s10546-014-9934-2.
- 2 Inoue, J., 1989. Surface drag over the snow surface of the Antarctic Plateau. 1: factors
3 controlling surface drag over the katabatic wind region. *J. Geophys. Res.*, 94, 2207-2217.
- 4 Joffre, S. M., 1982. Momentum and heat transfers in the surface layer over a frozen sea.
5 *Boundary-Layer Meteorol.*, 24, 211–229.
- 6 König-Langlo, G., King, J. C., and Pettré, P., 1998. Climatology of the three coastal Antarctic
7 stations Dumont d’Urville, Neumayer, and Halley. *J. Geophys. Res.*, 103(D9), 10,935–10,946,
8 doi:10.1029/97JD00527.
- 9 Kotlyakov, V. M., 1961. The Snow Cover of the Antarctic and its role in the Present-Day
10 Glaciation of the Continent (Snezhni pokrov antarktity i ego rol’ v somvremennom oledenении
11 materika), Translated from Russian 1966, Israel Program for Scientific Translation, Jerusalem,
12 256 p.
- 13 Lettau, H. H., 1969. Note on aerodynamic roughness-parameter estimation on the basis of
14 roughness element description. *Journal of Applied Meteorology*, 8, 828-832.
- 15 Li, L., and Pomeroy, J. W., 1997. Estimates of Threshold Wind Speeds for Snow Transport
16 Using Meteorological Data. *J. Appl. Meteorol.*, 36, 205–213.
- 17 Li, A., and Shao, Y., 2003. Numerical simulation of drag partition over rough surfaces.
18 *Boundary-Layer Meteorol.*, 108, 317–342.
- 19 Le Bouteiller, C., and Venditti, J. G., 2015. Sediment transport and shear stress partitioning in
20 a vegetated flow. *Water Resour. Res.*, 51, 2901–2922, doi:10.1002/2014WR015825.
- 21 Long, D. G., and Drinkwater, M. R., 2000. Azimuth variation in microwave scatterometer and
22 radiometer data over Antarctica. *Transactions on Geoscience and Remote Sensing*, 38,
23 1857–1870.
- 24 Marticorena, B., and Bergametti, G., 1995. Modeling the atmospheric dust cycle: 1. Design of
25 a soil-derived dust emission scheme. *J. Geophys. Res.*, 100, 16415–16430.
- 26 Mather, K. B., 1962. Further observations on sastrugi, snow dunes and the pattern of surface
27 winds in Antarctica. *The Polar Record*, 11, 158-171.

- 1 Mather, K. B., 1969. The pattern of surface wind flow in Antarctica. *Pure and Applied*
2 *Geophysics*, 75, 332-354.
- 3 Mather, K. B., and Miller, G. S., 1966. Wind drainage off the high plateau of Eastern Antarctica.
4 *Nature*, 209, 281-284.
- 5 Mellor, M., 1965. Blowing snow. *CRREL Monograph III*, A3c, Cold Regions Res. And Eng.
6 Lab., Hanover, N.H.
- 7 Naaim-Bouvet, F., Bellot, H., and Naaim, M., 2010. Back analysis of drifting-snow
8 measurements over an instrumented mountainous site. *Ann. Glaciol.*, 51(54), 207–217.
9 doi:10.3189/172756410791386661.
- 10 Ôura, H., Ishida, T., Kobayashi, D., Kobayashi, S., and Yamada, T., 1967. Studies on blowing
11 snow: Part II. *Physics of Snow and Ice: Int. Conf. on Low Temperature Science*, Sapporo, Japan,
12 Institute of Low Temperature Science, Hokkaido University, 1099–1117.
- 13 Parish, T. R., and Bromwich D. H., 1987. The surface wind field over the Antarctic ice sheet.
14 *Nature*, 328, 51-54.
- 15 Parish, T. R., and Bromwich, D. H., 2007. Reexamination of the near-surface airflow over the
16 Antarctic continent and implications on atmospheric circulations at high southern latitudes.
17 *Monthly Weather Review*, 135, 1961-1973.
- 18 Parish, T. R., Pettré, P., and Wendler, G., 1993. The influence of large scale forcing on the
19 katabatic wind regime of Adélie Land, Antarctica. *Meteor. Atmos. Phys.*, 51, 165-176.
- 20 Pomeroy, J. W., Gray, D. M., and Landine, P. G., 1993. The prairie blowing snow model:
21 Characteristics, validation, operation. *J. Hydrol.*, 144, 165–192.
- 22 Raupach, M. R., 1992. Drag and drag partition on rough surfaces. *Boundary-Layer Meteorol.*,
23 60, 375-395.
- 24 Sato, T., Kimura, T., Ishimaru, T., and Maruyama, T., 1993. Field test of a new snow-particle
25 counter (SPC) system. *Ann. Glaciol.*, 18, 149–154.
- 26 Schmidt, R. A., 1980. Threshold wind-speeds and elastic impact in snow transport. *J. Glaciol.*,
27 26-94, 453–467.

- 1 Shao, Y., 2008. Physics and modelling of wind erosion. Springer Verlag.
- 2 Smith, J. D., and McLean, S. R., 1977. Spatially averaged flow over a wavy surface. *J. Geophys.*
3 *Res.*, 82, 1735-1746.
- 4 Trouvilliez, A., Naaïm-Bouvet, F., Genthon, C., Piard, L., Favier, V., Bellot, H., Agosta, C.,
5 Trouvilliez, A., Naaïm-Bouvet, F., Bellot, H., Genthon, C., and Gallée, H., 2015. Evaluation of
6 FlowCapt acoustic sensor for snowdrift measurements. *J. Atmos. Ocean. Technol.*, 32(9), 1630-
7 1641.
- 8 Wiberg, P. L., and Nelson, J. M., 1992. Unidirectional flow over asymmetric and symmetric
9 ripples, *J. Geophys. Res.*, 97(C8), 12745-12761.
- 10 Wilkinson, R. H., 1984. A method for evaluating statistical errors associated with logarithmic
11 velocity profiles. *Geo-Mar. Lett.*, 3, 49–52.
- 12

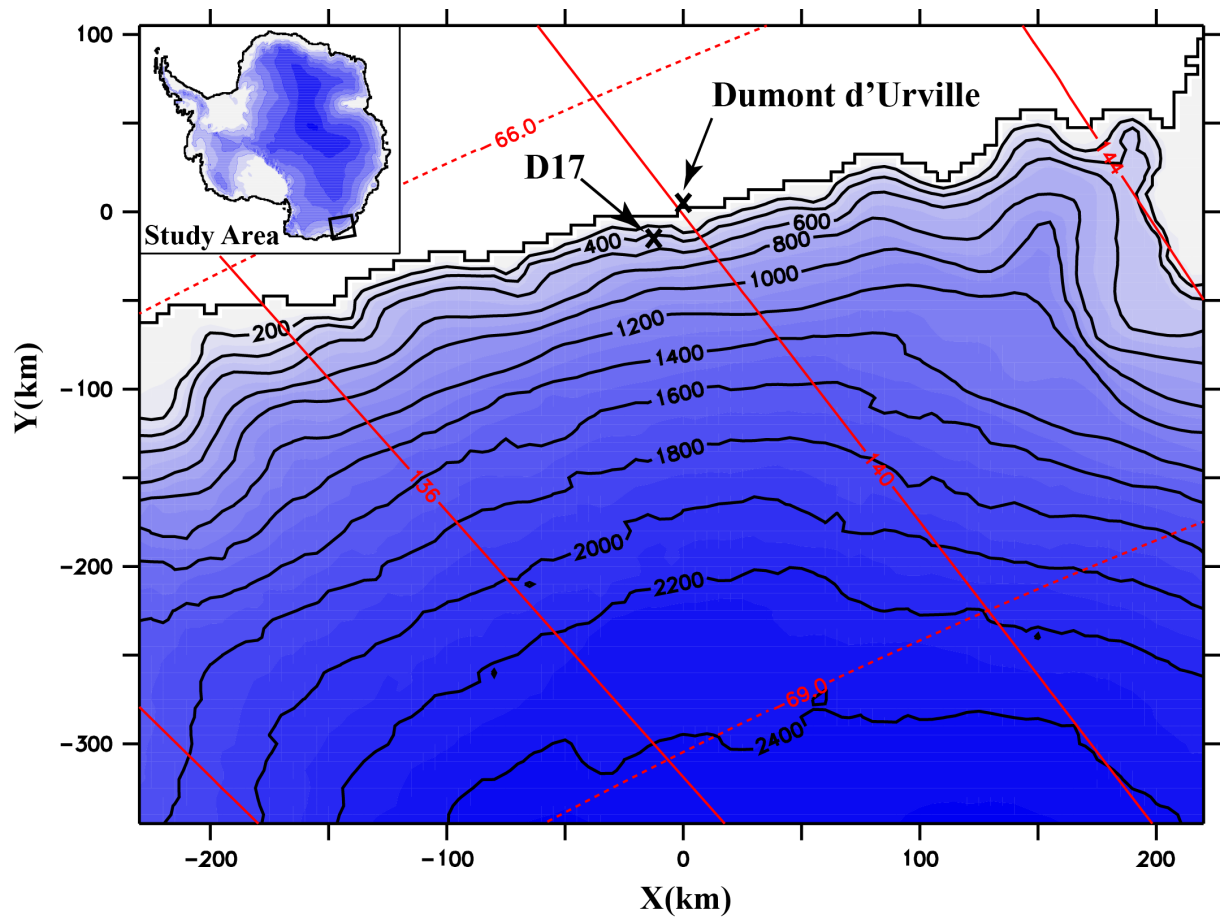
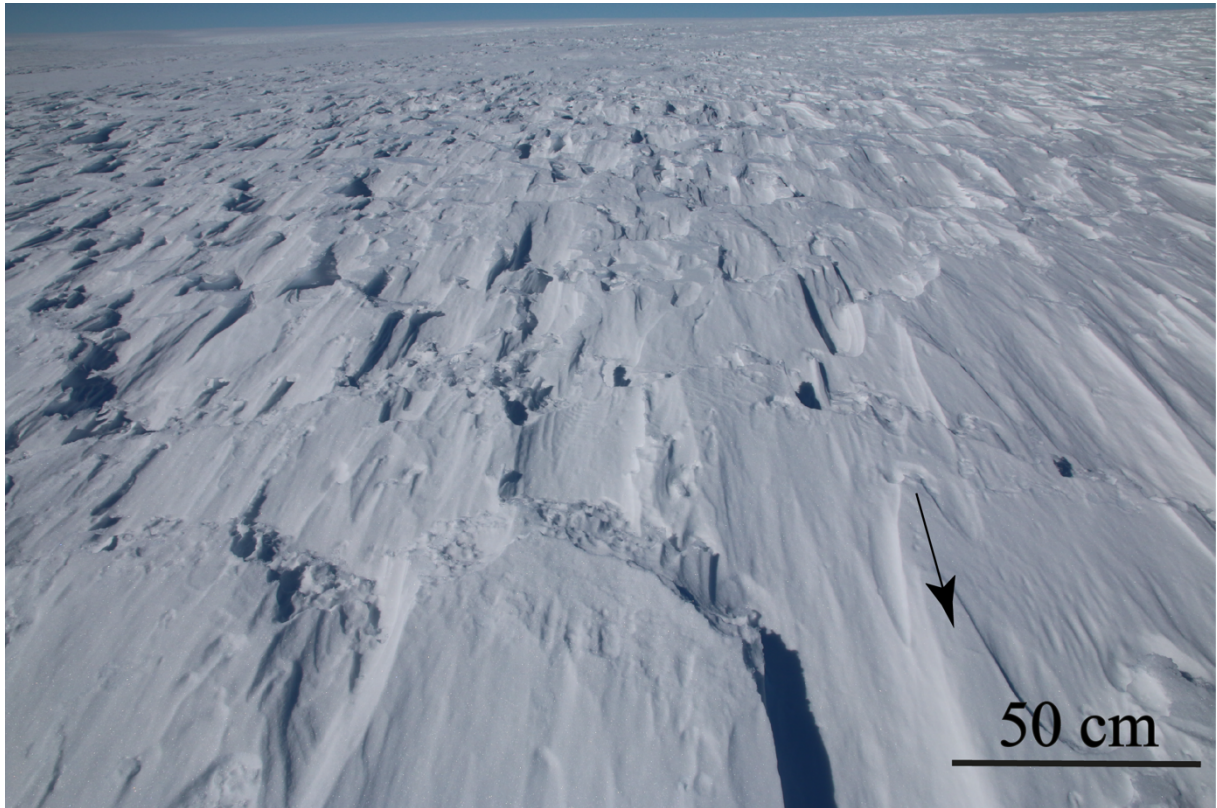


Fig. 1. Map of Adélie Land showing the location of Dumont d'Urville station and measurement site D17. Contour lines are in meters.



1
2 **Fig. 2.** Photograph of the snow surface at D17 in January 2014. The arrow indicates the mean
3 direction of the wind episode that led to the formation of the sastrugi.

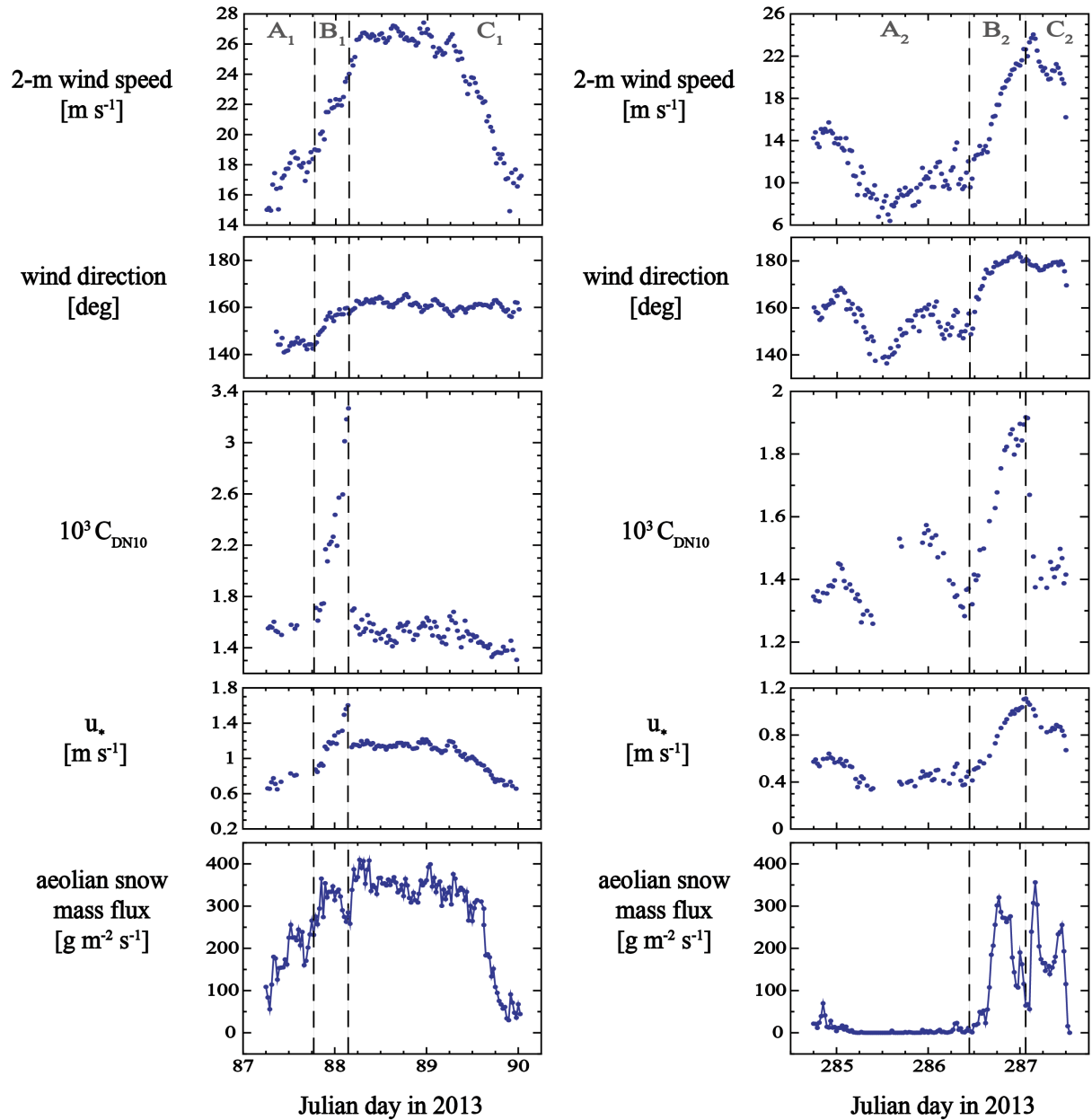


Fig. 3. Two erosion events showing sastrugi responses to shifts in wind direction. Note the different vertical scales between right and left panels concerning measured 2-m wind speed and profile-derived C_{DN10} and u_* values. The aeolian snow mass fluxes come from the second-generation FlowCapt™ sensor set up from 0 to 1 m above the snow surface. In both cases, the event is split into three parts, respectively before (A_i), during (B_i) and after (C_i) the wind shift.