

Climatic controls and climate proxy potential of Lewis Glacier, Mt Kenya

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Abstract

The Lewis Glacier on Mt Kenya is one of the best studied tropical glaciers and has experienced considerable retreat since a maximum extent in the late 19th century (L19). From distributed mass and energy balance modelling, this study evaluates the current sensitivity of the surface mass and energy balance to climatic drivers, explores climate conditions under which the L19 maximum extent might have sustained, and discusses the potential for using the glacier retreat to quantify climate change. Multiyear meteorological measurements at 4828 m provide data for input, optimization and evaluation of a spatially distributed glacier mass balance model to quantify the exchanges of energy and mass at the glacier–atmosphere interface. Currently the glacier loses mass due to the imbalance between insufficient accumulation and enhanced melt, because radiative energy gains cannot be compensated by turbulent energy sinks. Exchanging model input data with synthetic climate scenarios, which were sampled from the meteorological measurements and account for coupled climatic variable perturbations, reveal that the current mass balance is most sensitive to changes in atmospheric moisture (via its impact on solid precipitation, cloudiness and surface albedo). Positive mass balances result from scenarios with an increase of annual (seasonal) accumulation of 30% (100%), compared to values observed today, without significant changes in air temperature required. Scenarios with lower air temperatures are drier and associated with lower accumulation and increased net radiation due to reduced cloudiness and albedo. If the scenarios currently producing positive mass balances are applied to the L19 extent, negative mass balances are the result, meaning that the conditions required to sustain the glacier in its L19 extent are not reflected in today’s observations. Alternatively, a balanced mass budget for the L19 extent can be explained by changing model parameters that imply a distinctly different coupling between the glacier’s local surface-air layer and its surrounding boundary-layer. This result underlines the difficulty of deriving paleoclimates for larger glacier extents on the basis of modern measurements of small glaciers.

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1 Introduction

Glaciological observations in East Africa over the last century show a pronounced decrease in glacier length, area and mass (Cullen et al., 2013; Hastenrath, 1984, 2005b, 2008; Mölg et al., 2013; Prinz et al., 2011, 2012). Immediate changes in glacier mass are governed by concurrent weather and, consequently by climate through energy and mass exchanges between the glacier surface and the atmosphere. Integrated over time, and filtered by glacier dynamics, these exchanges result in changes in glacier extent. Thus, an accurate understanding of the glacier–climate interaction can be used to reveal the main atmospheric drivers of observed glacier extent changes.

The identification of climate signals from glaciers in East Africa is of particular interest, because they exist in elevations between approximately 5 and 6 km a.s.l. and therefore capture climate signals from the mid-troposphere (Mölg et al., 2009a), where our knowledge of climate change is scarce and controversial (e.g. Hartmann et al., 2013; Karl et al., 2006; Pepin and Lundquist, 2008). The temperature regime of the tropical East African climate – in particular at high elevations – is dominated by a pronounced diurnal cycle, driven by high incoming global radiation during daytime and strong long-wave energy loss during the night (Hastenrath, 1983), resulting in diurnal temperature variations being larger than seasonal temperature variations. The annual cycle is dominated by the hygric seasonality, expressed in the “long rains” (March to May) and the “short rains” (October to December), driven by the passage of the Intertropical Convergence Zone (e.g. Mutai and Ward, 2000) and modulated by Indian Ocean sea surface temperatures (e.g. Yang et al., 2014).

While some initially speculated that glacier recession on Kilimanjaro is caused by rising local air temperature, subsequent physical, process-resolving studies revealed that these glaciers are most sensitive to changes in atmospheric moisture and precipitation (Mölg and Hardy, 2004; Mölg et al., 2008, 2009a) due to their location far above the mean annual freezing level. Although there are additional controls on

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elevation, they offer a different climatic proxy than that offered by the glaciers of Kilimanjaro.

Nicholson et al. (2013) investigated the recent micrometeorological conditions and energy fluxes on LG at the point scale, in the context of other tropical glaciers.

Conditions at the summit of Mt Kenya were found to be much warmer and more humid than on Kilimanjaro, allowing convective clouds to converge over the summit of Mt Kenya much more frequently than over the summit of Kilimanjaro, even though both summits are influenced by the same air masses. The point modelling undertaken by Nicholson et al. (2013) suggests that, unlike on Kilimanjaro, the glacier mass balance variability is not dominated by a single variable or season. Building on that work, this paper aims to (i) extend the point surface energy and mass balance from Nicholson et al. (2013) to glacier wide-values for LG, (ii) evaluate the climate sensitivity of the glacier-wide surface mass and energy balance, (iii) explore climate conditions under which the late 19th century maximum extent of LG might have been sustained and (iv) discuss the potential for using shrinkage of LG to quantify climate change for a time period not covered by instrumental records.

2 Data and methods

2.1 Study site and in situ meteorological and mass balance observations

Lewis Glacier (0.1 km² in 2010) lies ~ 370 m below the summit of Mt Kenya in a south-westerly exposed, quasi-cirque location between the true summit and a secondary peak (Fig. 1). Several authors surveyed LG since 1934 (Prinz et al., 2011 and references therein) and reconstructed the late 19th century maximum extent (L19) from moraines and sketches (Patzelt et al., 1984). The most recent mapping was performed in 2010 (Prinz et al., 2011, 2012) and is used as topographic reference in this paper.

An automatic weather station (AWS) was installed on the glacier surface at an elevation of 4828 m which is ~ 30 m below the upper limit of the glacier. Meteorological

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The model used in this study originates from an energy balance model (Mölg and Hardy, 2004), that was developed into a mass balance model suitable for single point or glacier-wide applications (Mölg et al., 2008, 2009a). The model structure used in this study is explained in detail by Mölg et al. (2012) with the latest model version (2.4) used by Mölg (2015). The model treats the surface and near-subsurface mass and energy fluxes and has been successfully applied to address various questions in a range of climatic conditions (e.g. Collier et al., 2013; Conway and Cullen, 2015; Cullen et al., 2007, 2014; Gurgiser et al., 2013a, b; MacDonell et al., 2013; Mölg et al., 2014; Nicholson et al., 2013). The model computes the mass balance as the sum of solid precipitation, surface deposition, internal accumulation (refreezing of liquid water in snow), change in englacial liquid water storage, subsurface and surface melt, and sublimation. This approach is based on the surface energy balance of a glacier in the following form:

$$SWI(1 - \alpha) + LWI + LWO + QS + QL + QPRC + QC + QPS = F \quad (1)$$

where SWI is incoming shortwave radiation (global radiation corrected for aspect/slope), α is surface albedo, LWI and LWO are incoming and emitted longwave radiation fluxes, QS and QL are the turbulent fluxes of sensible and latent heat, respectively, QPRC is the heat flux from precipitation, QC the conductive heat flux in the subsurface and QPS the energy flux from shortwave radiation penetrating into the subsurface. The sum of these fluxes yields a residual flux F which, if the glacier surface temperature (TS) reaches 273.15 K, represents the latent energy for melting. If TS is below 273.15 K, energy conservation is achieved by solving TS to balance the fluxes (Mölg et al., 2009a; van den Broeke et al., 2006). Input data for mass accumulation is provided by a surface height change record or as water equivalent, and SWI, α , LWI, LWO, air temperature, atmospheric humidity, air pressure and wind speed are required in order to solve the energy balance for the remaining mass fluxes. The model can be validated on the basis of measured surface height changes at the AWS and nearby reference stakes, and/or with TS measured directly or derived from measured

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LWO. In order to enable glacier-wide mass balance runs and sensitivity tests some of these input variables must be replaced by parameterizations based on key atmospheric properties that can then be varied in the sensitivity study. Thus, SWI, α , LWI are parameterized and TS (LWO) is computed from the energy balance as noted above, to capture the feedback effects on the mass balance in case of a climate perturbation (e.g. Mölg et al., 2009a). The required MBM inputs are therefore reduced to air temperature and humidity, air pressure, a cloud cover factor, wind speed and accumulation rate. Additionally, for the spatially distributed case, a digital elevation model is compulsory as lower boundary condition, from which the grid cell sky view and shading parameters are computed. Vertical gradients in precipitation and air temperature are essential parameters to distribute the meteorological data from the AWS to the whole glacier surface.

2.3 Parameterization of radiative fluxes

The MBM's parameterizations for SWI employs the approach from Budyko (1974), which was applied for Equatorial East Africa in earlier studies (Hastenrath, 1984; Mölg et al., 2009b),

$$G = (S_{CS} + D_{CS})(1 - kn_{eff}), \quad (2)$$

where G is the global radiation, separated into direct and diffuse clear sky components (S_{CS} and D_{CS}), and n_{eff} is the effective cloud cover fraction (0–1). The constant k controls the global radiation under cloudy conditions. First, clear sky G is modelled using concepts of Iqbal (1983) and Meyers and Dale (1983) and optimized against measured clear sky days (Mölg et al., 2009b), defined from the meteorological record at the AWS when G was $> 77\%$ of top of atmosphere radiation and mean daily longwave net radiation lower than its 5th percentile as a proxy for absent cloudiness (van den Broeke et al., 2006). From 773 days, only 17 clear sky days were found, all occurring in late December to late February. All-sky G is modelled by optimizing the product kn_{eff} over the total measurement period (773 days) and yielding a k of 0.72,

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meteorological conditions. In this study model uncertainty at the AWS location was obtained by a combination of Monte Carlo optimization in conjunction with a k -fold cross-validation that provides a robust assessment of model performance outside the training period. The 773 days of meteorological input were divided into five periods for which MBM initialization (snow depth, TS and an estimate of the last snow fall event) was possible from independent field observations (Table 1). For each of the five (k) periods 1000 model realizations were performed, applying the Monte Carlo simulation with the same quasi-random parameter matrix, where each parameter spans a physically sound range based on measurements or previously published data (Table 2). Minimizing the combined RMSE between measured and modelled daily surface height change, daily means of TS and α , results in five optimal model parameter sets. To test their performance outside their optimization periods, these sets were cross-validated against the k th period excluded from the optimization. Table 3 gives the obtained errors for surface height change, TS, and α for the whole range of 773 days and for the 2011/12 mass balance year represented by the last 358 days of the record, as this period is the focus of the subsequent sensitivity analysis. The modelled mass balance uncertainty for the mass balance year 2011/12 is $\pm 154 \text{ kg m}^{-2}$, derived from the RMSE in daily surface height change together with a conservatively assumed density of 900 kg m^{-3} . If the mass balance, modelled at the AWS location, is compared against the available independent measurements at the nearest ablation stake, which is 2 m away, the model uncertainty is $\pm 82 \text{ kg m}^{-2}$, but the stake readings offer only two points in time for evaluation (the biannual ablation stake readings for 2011/12), and therefore the larger error is conservatively taken to be more representative of the model uncertainty.

2.5 Distributing model input variables over the glacier surface

After determining the MBM error at the point of the meteorological measurements, the MBM is applied in its distributed mode to the whole LG surface, represented in a digital elevation model (DEM) of 25 m grid point spacing. Vertical gradients are used to

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do not allow for a quantified partitioning of potential dynamic processes causing strong daytime gradients observed at the margins of glaciers (Ayala et al., 2015; Carturan et al., 2015; Petersen and Pellicciotti, 2011), the daytime gradient of $-0.015^{\circ}\text{C m}^{-1}$ was found through optimization using the spatially distributed stake mass balances of 2011/12. As the component of LWI that is parameterized for terrestrial irradiance from the surroundings is not constrained by measurements, the optimized daytime vertical air temperature gradient will also compensate for shortcomings in the LWI parameterization. LG spans only 220 m in altitude, so the mean nocturnal and diurnal air temperature differences from the top (measured) to the terminus of the glacier (from the optimized vertical air temperature gradient extrapolations) are 1.4 and 3.3 $^{\circ}\text{C}$ respectively.

2.6 Sampling synthetic climate scenarios

As described by Nicholson et al. (2013) it is unlikely that a single climatic variable dominates the mass balance variability on LG. Thus, exploring mass balance sensitivity of single climatic variable perturbation was rejected in favour of coupled perturbations reflecting the variability in the climate more comprehensively (Mölg et al., 2009a). Consequently, in order to assess the sensitivity of the mass balance to a perturbation in its forcing climate, alternative climate scenarios were constructed by reassembling the AWS records on a diurnal basis in a very simple, weather generator like concept (e.g. Hutchinson, 1987). The period of AWS data has been shown to be representative for the recent decades in terms of monthly ERA-interim air temperature (1979–2012) and TRMM precipitation (1998–2012) time series (Fig. 3 in Nicholson et al., 2013). Out of the 773 days with meteorological data from LG AWS, 365 days (representing one arbitrary mass balance year from 1 March to 28 February) were sampled with replacement to construct four differently perturbed, synthetic climate scenarios meeting the following characteristics compared to the 2011/12 REF year: +1 K air temperature (warm and wet, WW), -1 K air temperature (cold and dry, CD), +50 % accumulation (WET), and -20 % accumulation (DRY). To maintain the actual hygric seasonality, the

3 Results and discussion

3.1 The mass and energy balance for the year 2011/12

Modelled annual mass balance at the 23 ablation stakes available for 2011/12 is significantly correlated to the measured mass balance with $r = 0.86$ at an RSME of 320 kg m^{-2} (Fig. 4). Propagating the combined errors of the cross validation and comparison to the measured mass balance at all stakes, the modelled glacier mass balance for 2011/12 is $-911 \pm 355 \text{ kg m}^{-2}$, which agrees well with the measured value¹ of -961 kg m^{-2} . The model sensitivity of the mass balance to the vertical air temperature gradient is -20 kg m^{-2} for each increase of $0.001 \text{ }^\circ\text{C m}^{-1}$ (between -0.015 and $-0.0065 \text{ }^\circ\text{C m}^{-1}$) and 6 kg m^{-2} for each increase of the vertical precipitation gradient of $0.001 \text{ } \%$ m^{-1} (between -0.005 and $-0.01 \text{ } \%$ m^{-1}).

Glacier-wide mean monthly energy and mass flux densities for 2011/12 are shown in Fig. 5. The governing role of the net short-wave flux on the energy and mass balance is clear. When α is high, due to abundant snow accumulation, the energy available for ablation is reduced, enabling conditions for net accumulation. In contrast, during the January/February dry season increased net radiation and turbulent fluxes cause high ablation rates. Both long-wave fluxes are almost constant throughout the year, as a potential lower LWI in colder and drier conditions is compensated by increased emission from surrounding terrain due to its enhanced solar heating, and TS reaches regularly $0 \text{ }^\circ\text{C}$. The seasonal cycle in 2011/12 is attenuated in the first half of the record and amplified in the latter, due to accumulation amounts below normal in the “long rains” and above normal in the “short rains”, respectively (Nicholson et al., 2013).

¹For the mass balance year 2011/12 Prinz et al. (2012) reported an annual mass balance of -1030 kg m^{-2} . This value changed to -961 kg m^{-2} as the spatial interpolation of observed mass change from the ablation stakes to the glacier area was homogenized by using contours with constant equidistance.

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the air above is maintained due to effective long-wave surface cooling and energy consumption by QL, and higher wind speeds enhance both turbulent fluxes. Thus, the physically-based modelling study presented here indicates that for climate scenarios that are within the limits of observations, only a change in the mountain's precipitation climatology to substantially more accumulation is able to sustain LG in its current extent.

3.3 Modelling and interpreting L19 mass balance

The mass balance model was applied to the LG in its L19 extent with the two most positive mass balance scenarios synthesized from the range of observed modern climate conditions (scenarios WET and AMPC) to see if these perturbed climates are sufficient to sustain the L19 glacier extent. The modelling produced negative mass balances in both cases of -233 kg m^{-2} (WET) and -338 kg m^{-2} (AMPC), respectively. Again, in these simulations, the impact of the strong air temperature gradient on the phase change of precipitation over the L19 extent is minor and confined to the lowest parts ($< 4550 \text{ m}$), where the maximum reduction of accumulation in favour of rain is 13% for the lowest grid point. Over the total glacier area the fraction of rain is less than 1% for both scenarios WET and AMPC.

The negative mass balances for these scenarios, being most favourable to glaciation, imply that even the extremes of the present day climate are incapable of reproducing the L19 conditions. One interpretation of this finding is that the range of modern-day meteorological conditions in the summit region of Mt Kenya no longer overlaps with the L19 range, and/or the covariance of meteorological conditions was substantially different in L19 than today.

3.4 The impact of glacier extent on the proxy potential of Lewis Glacier

Given that LG is now 83% smaller than its L19 extent, the relative importance of the glacier microclimate relative to the surrounding terrain is likely to have changed

significantly between these two glacier geometries (Fig. 1). The limited aerial and vertical extent of the modern glacier favours the steep vertical air temperature gradient along the glacier surface, but on larger glaciers such as LG during its L19 extent, the air temperature gradient over the glacier surface is strongly modified by the katabatic wind field (Ayala et al., 2015; Greuell and Böhm, 1998; Shea and Moore, 2010), and the influence of longwave emissions from surrounding terrain is drastically reduced as the glacier fills the cirque (Fig. 1).

Given the small area of LG in the modern day, it could be that the glacier is too small to form a substantial katabatic layer and modify its own microclimate and is instead more strongly influenced by the surroundings with the off-glacier boundary layer conditions dominating over much of the lower glacier. If this is the case, the air temperature distribution optimized for the modern-day LG extent cannot capture the dynamic processes that play a part in governing the air temperature distribution over the larger L19 glacier extent. Repeating the modelling for the L19 glacier extent using the moist adiabatic lapse rate for all hours of the day gives mass balances of 190 and 68 kg m⁻² for WET and AMPC respectively, and quasi-zero mass balances (17, -2 kg m⁻², for WET and AMPC respectively) can be achieved by using daytime vertical air temperature gradients of -0.010 and -0.008 °C m⁻¹, respectively. This supports the idea that a larger glacier can develop a deeper katabatic boundary layer that is more difficult to entrain through advection and turbulent mixing of warm air. Thus, reconstructions of former or future climates that are based on model optimizations for a modern-day glacier extent may not be applicable to substantially different glacier extents. This might have particular relevance for reconstructions based on very small glaciers such as LG, especially when glacier geometries changed significantly. Ongoing retreat of remaining mountain glaciers suggests that scale effects such as this might also become increasingly important for paleoclimate reconstructions from mountain glaciers in the future.

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4 Conclusion

Distributed surface energy and mass balance modelling indicates that the energy and mass balance of present-day Lewis Glacier is most sensitive to atmospheric moisture (solid precipitation, cloudiness and albedo). In the tropical atmosphere of Mount Kenya, air temperature changes are always coexistent with changes in atmospheric moisture; consequently, air temperature variation cannot be isolated as a single driver of glacier mass change. Although it has been proposed that a reduction in air temperature of 0.7°C would be sufficient to bring this small glacier to a zero mass balance state (Hastenrath, 2010), a colder climate scenario results in a less negative mass balance, but without additional accumulation it is insufficient to achieve equilibrium.

Two scenarios suggest that higher accumulation ($+30$ to $+100\%$ year $^{-1}$ or wet season, respectively), higher relative humidity (4 to 8% units per year or dry season), a change of fractional cloud cover (-5 to $+1\%$ units per dry season or year) and higher wind speed (0.3 to 0.6 ms^{-1} year $^{-1}$ or dry season), which are all mutually linked, allow a zero mass balance for the present day Lewis Glacier, without significant changes in the air temperature.

Using the mass balance model as optimized for the modern-day glacier, driven by climate perturbations reflecting the observed variability in precipitation and air temperature, indicates that L19 conditions at LG were distinctly different to the present day, and it is not possible to fully quantify the climatic conditions that could sustain LG at its maximum L19 extent. Additionally, the modelling suggests that extracting proxy climate conditions from a particular glacier geometry using a modelling system optimized on a dramatically different geometry may invalidate the approach, particularly if changes in boundary layer dynamics are substantial and not resolved in the model. This issue might warrant further investigation given that paleoclimate reconstructions based on mountain glacier fluctuations inherently involve these scale contrasts, yet they are rarely considered in the tools used.

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Table 1. Periods of available meteorological input for cross-validation.

| period | from to | number of days |
|--------|------------------------------------|----------------|
| 1 | 26 September 2009–24 January 2010 | 121 |
| 2 | 2 March 2010–19 July 2010 | 140 |
| 3 | 29 September 2010–1 March 2011 | 154 |
| 4 | 2 March 2011–14 September 2011 | 197 |
| 5 | 15 September 2011–22 February 2012 | 161 |

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Table 3. Model uncertainties expressed as root mean squared error (RSME) and correlation coefficient (r) for surface height change (sfc), daily mean surface temperature (TS) and albedo (α) at the location of the AWS.

| | RSME | | | r | | |
|---|----------|---------|----------|------|------|----------|
| | sfc [cm] | TS [°C] | α | sfc | TS | α |
| all 773 days | 29.7 | 1.23 | 0.15 | 0.81 | 0.72 | 0.54 |
| last 358 days (= mass balance year 2011/12) | 17.2 | 1.03 | 0.09 | 0.78 | 0.78 | 0.83 |

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Table 4. Mean air temperature (T), relative humidity (RH), effective cloud cover fraction (n_{eff}), wind speed (v), and accumulation sum for different scenarios and the 2011/12 mass balance year for different seasons: annual (wet season/dry season). Mass balance modelled under each scenario for the 2010 glacier extent is denoted by B ($\pm 355 \text{ kg m}^{-2}$).

| Variable (unit) | Scenarios | | | | | | | 2011/12 |
|-------------------------------|------------------------|------------------------|------------------------|------------------------|------------------------|------------------------|------------------------|------------------------|
| | WW | CD | WET | DRY | AMPW | ATT | AMPC | REF |
| T (°C) | -0.11 (-0.15/-0.06) | -2.11 (-2.11/-2.10) | -1.05 (-1.07/-1.03) | -0.93 (-0.92/-0.94) | -1.13 (-0.15/-2.10) | -1.09 (-2.11/-0.06) | -1.59 (-1.07/-2.10) | -1.11 (-0.96/-1.27) |
| RH (%) | 78 (81/76) | 67 (71/64) | 77 (79/75) | 75 (78/71) | 72 (81/64) | 74 (71/76) | 71 (79/64) | 73 (79/67) |
| n_{eff} (%) | 28 (29/27) | 23 (26/21) | 28 (28/27) | 26 (28/24) | 25 (29/21) | 26 (26/27) | 25 (28/21) | 26 (27/26) |
| v (ms^{-1}) | 2.6 (2.6/2.7) | 3.1 (2.8/3.4) | 2.7 (2.7/2.7) | 2.8 (2.6/2.9) | 3.0 (2.6/3.4) | 2.7 (2.8/2.7) | 3.1 (2.7/3.4) | 2.8 (2.8/2.7) |
| acc (cm) | 355 (300/55) | 152 (140/12) | 349 (296/54) | 188 (166/21) | 312 (300/12) | 195 (140/55) | 308 (296/12) | 232 (142/88) |
| B (kg m^{-2}) | -527 (+580/-1107) | -578 (+115/-693) | +447 (+696/-249) | -1384 (-186/-1197) | +66 (+592/-526) | -1242 (+112/-1354) | +260 (+688/-427) | -966 (-414/-552) |

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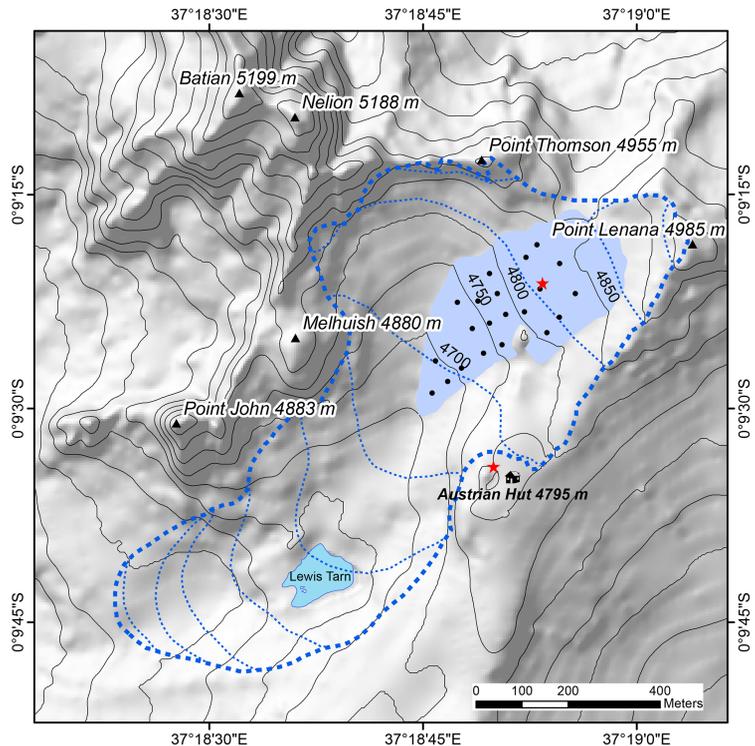



Figure 1. Overview map of LG for 2010 (0.1 km²) and the late 19th century (L19, 0.6 km², dashed). Red stars denote AWS locations, black dots the ablation stakes. L19 outline from Patzelt et al. (1984) with reconstructed contour lines. Off-glacier contours were taken from Schneider (1964) and updated for LG basin 2010 (Prinz et al., 2012).

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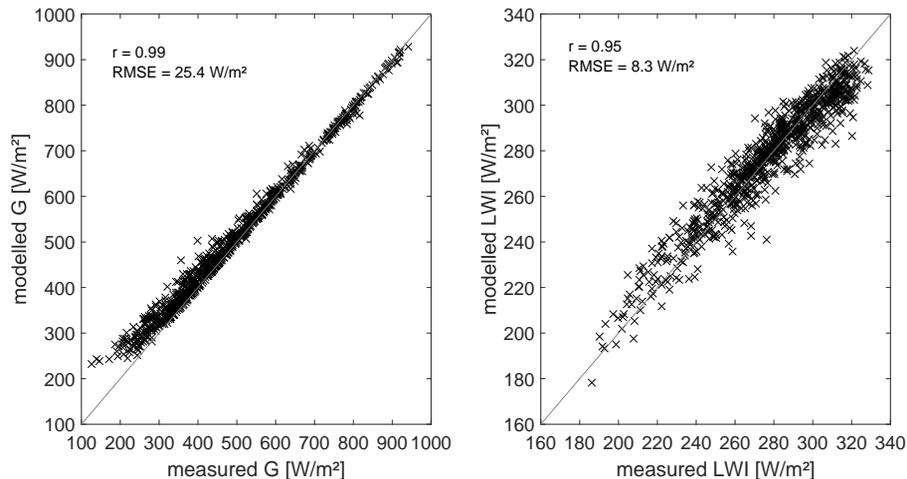


Figure 2. Parameterization performance of downward radiative fluxes for 772 days at the location of the AWS. Scatterplot of mean daily measured vs. modelled G (left) and LWI (right). Hourly values show similar correlation of $r = 0.99$ (0.93) and $RSME = 26.4$ (33.8) $W m^{-2}$ for G (LWI), respectively.

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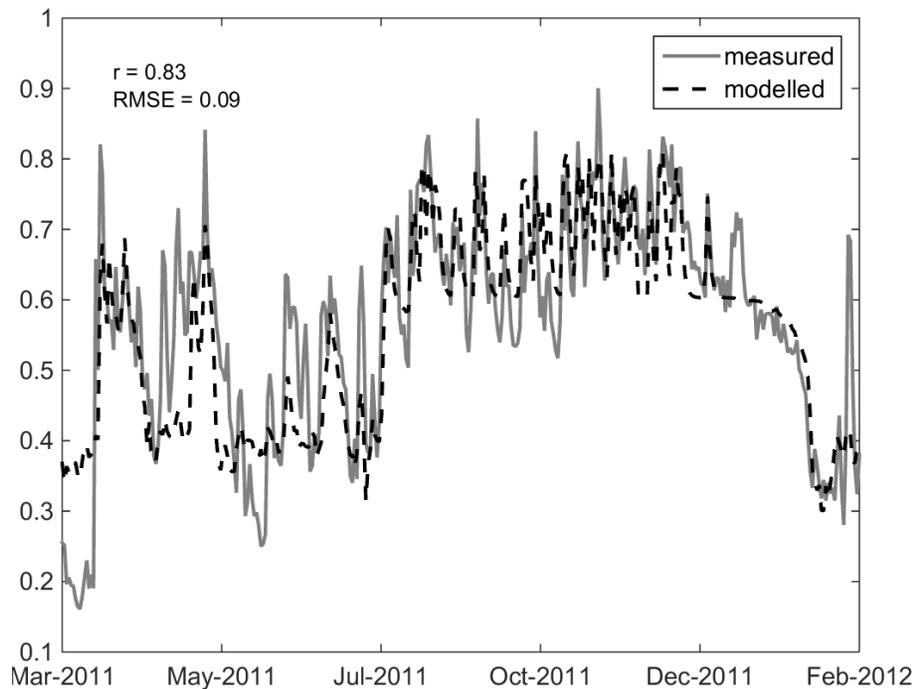


Figure 3. Daily mean values of measured and modelled α for the 358 day mass balance year 2 March 2011 until 22 February 2012. Varying ice albedo is computed as a function of the dew point temperature (DPT): $0.0056\text{ }^{\circ}\text{C}^{-1}\text{ DPT} + 0.4179$.

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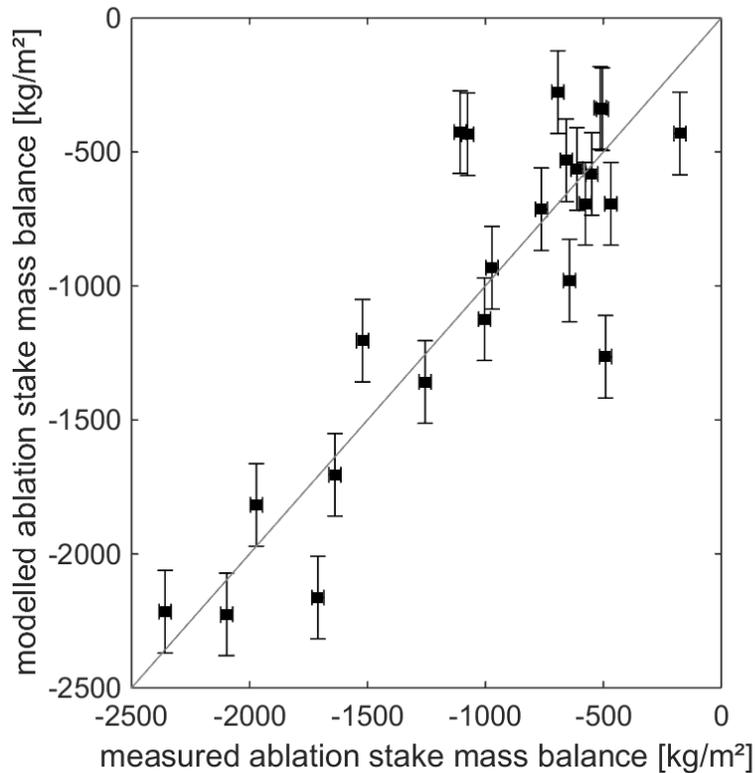


Figure 4. Model performance at the ablation stakes for the mass balance year 2011/12.

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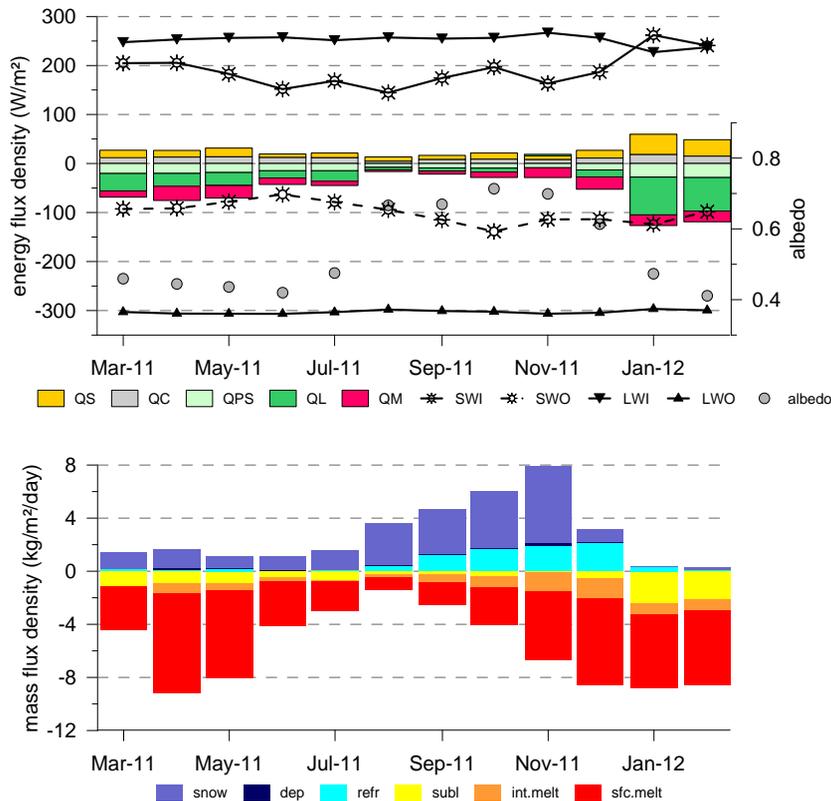


Figure 5. Glacier-wide mean monthly energy (upper panel) and mass flux densities (lower panel) for the mass balance year 2011/12 (REF): abbreviations as defined in Sect. 2.2 – the heat flux from precipitation is not shown due to its very low values; sfc.melt (surface melt), int.melt (internal melt in the subsurface), subl (sublimation), refr (refreezing), dep (deposition).

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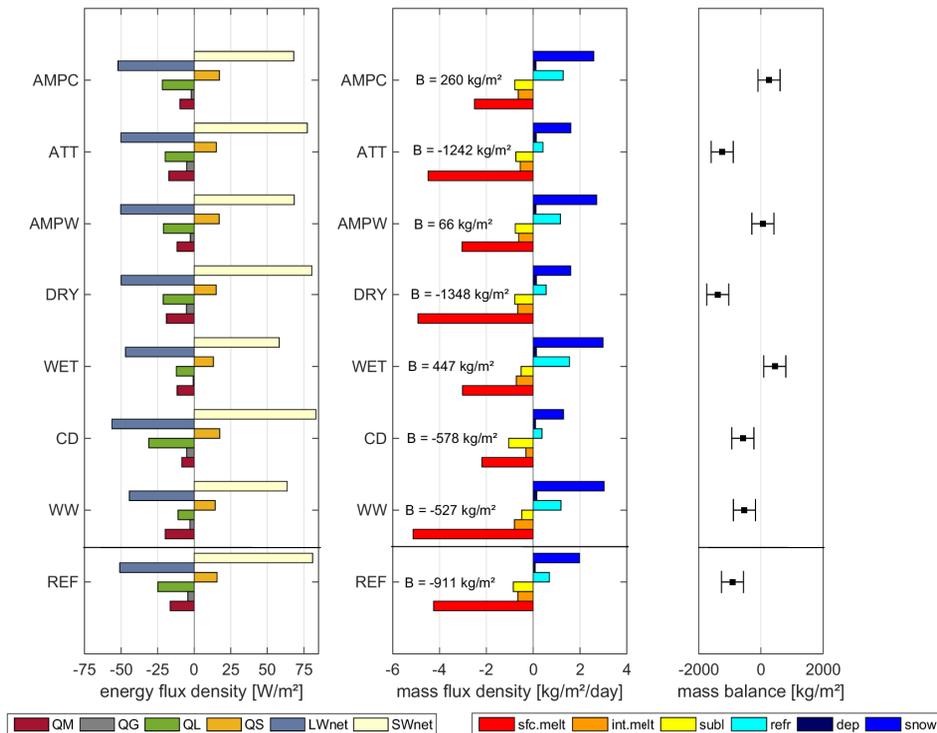


Figure 6. Glacier-wide mean energy (left) and mass flux densities (middle) for the eight different scenarios. The annual mass balances (B) are shown in the middle panel and their error ranges in the right panel. Abbreviations as defined in Sect. 2.2 and Fig. 5, except QG (ground heat flux as the sum of QC and QPS), LWnet (net long-wave radiative flux), and SWnet (net short-wave radiative flux).

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