Dear Prof. dr. Hanna,

We hereby submit our revised manuscript, initially submitted to The Cryosphere Discussions on 19 March 2015. The reviewers, while embracing our manuscript in its initial form, provided very valuable input towards its improvement. In response, we have performed changes, the most important of which we outline in the following:

We have tried to simplify the text throughout the manuscript in order to enhance the flow, and facilitate the reader towards our conclusions. The most significant changes are the exemplified paragraph in the review by Xavier Fettweis, and the last paragraph of our Discussion, where we removed the unnecessary quantifications and simply referred to the respective subsections in the Results. Also, an incorrect analysis at the “Melt-albedo feedback” subsection was identified and corrected (p. 2883, l. 20 of Discussion Paper), as the total amount of meltwater equals the sum of runoff and refreezing minus rainfall.

We have now clarified that the snow density in the simulations is the rounded measured density of snow-pit measurements in 2013. Also, in response to a very sharp comment by Xavier Fettweis, we have modified Table 4, and included the surface height changes. In the accumulation area, the permeable surface is a source of uncertainty for surface height monitoring, since during the melt season all objects tend to sink into the warmed snow and firn, following the percolation of meltwater. It is rather important for studies in the accumulation area to provide both surface height and mass budget estimates, and we are glad we were prompted to do so.
We have removed all discussion of the North Atlantic Oscillation (NAO), because it is indeed a climatic aspect more relevant to studies of higher spatiotemporal scales. Our study reserves its conclusions even without the NAO, which indicates that its mention was unnecessarily high level of detail.

As suggested, we have also introduced a paragraph in the Discussion on the MODIS instrument sensitivity while referencing specialized studies on that matter (Wang et al., 2012; Lyapustin et al., 2014).

Further changes include an alternative reference for the PROMICE project (Ahlstrom et al., 2008), as well as the citation of a paper describing our weather station setup (Citterio et al., 2015). A short comment about the uncertainty of refreezing estimates was also inserted in the introduction with the inclusion of a relevant citation (Vernon et al., 2013).

We should mention that two of our references are currently under review (Charalampidis et al., under review; Machguth et al., under review). However, we expect the editorial decision from their respective journals no later than the end of October this year.

Lastly – but more importantly – as you might have already noticed, we have decided to include Alun Hubbard and Samuel Doyle, as well as their current affiliations in the revised manuscript. Their contribution to our efforts is undeniable since throughout the years they have provided for the maintenance and logistical support of the KAN_U weather station within the framework of the Greenland Analogue Project. We felt that the simple mention in the Acknowledgements of the Discussion Paper did not do justice to their contribution. They have welcomed our decision to include them as authors, and have kindly offered to proofread the study.

We look forward to your response. Don’t hesitate to contact me if you have further questions.

Yours sincerely,

Charalampos Charalampidis
References


Response to both reviewers

Anonymous referee

This is a very thorough analysis and interpretation of 5 years of automatic weather station data from the western flank of the Greenland Ice Sheet. The MS contains a wealth of meteorological data that will be of interest to other Greenland workers, and as such is worth publishing. There is also useful consideration of ice surface albedo feedback and sub-surface refreezing. I found the paper to be well-structured, well-written and -illustrated and easy to follow. I concur with many of the comments of Reviewer 1.

Referee Xavier Fettweis

This paper presents recent changes in the energy balance and the snowpack behaviours at KAN_U situated near the equilibrium line in the accumulation zone of the Greenland ice sheet. It is not the first time that energy balance and melt from in situ observations is discussed in this south-western part of the GrIS (van den Broeke et al., 2011) but KAN_U is situated in the accumulation (while measurements from the ablation zone only was presented in van den Broeke et al. (2011)) and the discussion about the snowpack changes in 2012 is interesting, innovative and deserves to be published in TC with some minor revisions only.

We thank Xavier Fettweis and an anonymous reviewer for their enthusiastic response, encouragement and constructive comments on our discussion paper. In the following, we address all points:

The paper is clear and fits well with TC. The text is well written but sometimes it is hard to read due to the abundance of numbers and statistics in the text. Some simplifications when nothing important is told (e.g.: lines 1-10, pg 2883) could be made in the text by simply referencing to the corresponding tables.

Taking also into account that Reviewer #2 found our discussion paper easy to follow, we have attempted some minor simplifications throughout the text, in order to make our quantifications more straightforward. The paragraph exemplified, we reformulated as: “Figure 9a, which depicts total monthly surface energy exchanges throughout the study period, illustrates that $E_{\text{S}}^{\text{net}}$ and $E_{\text{L}}^{\text{net}}$ dominate the SEB from May to September, while $E_{\text{L}}^{\text{net}}$ and $E_{\text{H}}$ dominate the SEB during the remainder of the year. During the years exclusive of 2012 considered here (2009, 2010, 2011 and 2013), the total summer energy input to the ice sheet surface was 620–650 MJ m$^{-2}$. This energy reaches a peak in July. In July 2010, for example, the total energy input reached 246 MJ m$^{-2}$. By contrast, in 2012, the total summer energy input exceeded 770 MJ m$^{-2}$, and in July it reached 304 MJ m$^{-2}$. The 2012 total energy used for melt was 414 MJ m$^{-2}$ (65 % higher than in 2010), of which 183 MJ m$^{-2}$ was used for melt in July.”

Line 25, pg 2875 vs line 11 pg 2878: 360 or 400 kg/m3 for the snow density?

Admittedly, it is an unclear wording. The measured average snow density is 360 kg m$^{-3}$, used in Table 4 in the discussion paper. The snow density used in the simulations was rounded to 400 kg m$^{-3}$. We have
reformulated as: “Solid precipitation is added in the model based on KAN_U sonic ranger measurements, assuming the rounded average snow density found in snow-pit measurements, i.e. 400 kg m$^{-3}$."

**Table 4:** I am a bit surprised that we use here a mean density of 360 Kg/m$^{3}$ for estimating the mean ablation rate. As snow is melting, the snowpack density should be higher.

This is a fair point. Of course, there is density increase after melt percolates, as illustrated in Figures 9b and 11c of the discussion paper. We used this density in an attempt to estimate the mass flux melted for the first time from the initial snowpack each melt season. Arguably, this can be confusing, and not necessarily correct. Therefore, we decided to remove the estimated ablation rates from Table 4, as they are also not a crucial point in the study.

**Where does the density uncertainty of 40 kg/m$^{3}$ come from?**

This uncertainty is based on the standard deviation among the measurements from snow pit measurements conducted in spring 2013, which we now clarified in the caption.

**Just giving the difference in snow height is for me more reliable.**

Indeed, we also think that heights are more reliable and generally should be also provided, as SMBs are sometimes a matter of interpretation, especially in the accumulation area of the ice sheet. We have now inserted the winter/summer heights in Table 4. We did, however, keep also our SMB estimates for reference. The updated Table 4 of the discussion paper is shown here as Table I.

**Line 16, pg 2880:** these low albedo values are for me more likely the result of the snowpack erosion by the wind (making apparent old firn) than reduced winter precipitation. The regional model MAR does not suggest particular low winter accumulation at KAN_U in 2012-2013.

We agree, and MAR is accurately suggesting substantial accumulation during winter 2012–2013 at KAN_U. We also agree that there might have been snowpack erosion at our study site (Leanerts et al., 2014). However, this erosion is unlikely to have caused exposure of old firn. By November 2012, the snow thickness was 0.6 m (Fig. 2). From that point onward, and until spring 2013, the sonic ranger measurements suggest that there was limited accumulated snow on top of that initial snowpack in autumn; hence wind might have had an effect on surface, but certainly not erode the whole snowpack, thereby exposing firn. The snow at the surface probably lost part of its reflectivity after the prolonged exposure to the atmosphere. The area received substantial accumulation in spring. This was also verified by the Arctic Circle Traverse 2013 (ACT-13) in late April 2013, when we were at the location and the snow cover was ~ 0.9 m. After two weeks, the snow cover had increased by ~ 0.3 m.

**Lines 15-20, pg 2886:** I do not see the interest of discussing NAO here. The role of NAO over Greenland is well known for explaining the recent melt increase and for me, Fig 10a as well as these 5 lines should be removed.

Initially, our aim was to provide with a description as complete as possible, but we agree that information on the NAO, as well as its connection to recent climatic changes are readily available in the literature, and
perhaps more relevant to larger-scale studies. We have now removed all discussion of the NAO, and also Figure 10a. The new version of Figure 10 of the discussion paper is shown here as Figure I.

_lines 5-12, pg 2887:_ The comparison with MODIS is interesting but a part of the MODIS based albedo decrease could be the result of the declining instrument sensitivity of the MODIS sensors. This issue should be discussed. However the same albedo trend is also simulated by MAR (forced by NCEP-NCARv1) which also simulates the exceptional low albedo in summer 2012 (see Fig. 1 next page)!

It is, indeed, an issue that should be mentioned when discussing remotely-sensed albedo. We have now included the following in the manuscript: “Part of the MODIS based albedo decrease could be the result of the declining instrument sensitivity of the Terra MODIS sensor (Wang et al. 2012; Lyapustin et al. 2014) though updated (through 2014) comparisons between MOD10A1 and ground observations from GC-Net data (Box et al. 2012; not shown) do not indicate an obvious nor statistically significant difference.”

According to MAR, it is the first time in summer 2012 since 1950 that significant ice lenses appear but in 1960, MAR also simulates high runoff rates due to snowpack meltwater saturation suggesting that it is not the first time that significant melt events occur at Kan_U.

We have reasons to believe that this is partly incorrect. In detail, two firn cores were retrieved in May–June 1989 from Site J (66° 51.9' N, 46° 15.9' W, 2030 m a.s.l.; Kameda et al., 1995), ~36.2 km east-southeast from KAN_U (Fig. IIa). According to the deduced Melt Feature Percentage (MFP) shown in Figure IIb (Kameda et al., 2004), 1960 has a higher melt feature percentage (MFP) than usual, but assuming that strong melt in 1960 would mostly percolate into previous year's firn layers, there is no indication that 1960 stands out much.

Another observational study analyses 10 m firn temperatures based on measurements prior to 1965 (Mock and Weeks, 1966). In their analysis, the closest site to KAN_U was ~60 km south-southwest, i.e. at a lower elevation. According to this study, the estimated 10 m firn temperature at KAN_U was at that time around −14 °C, suggesting that there was no significant refreezing in that period (and therefore ice content within the firn).

The above observations come from different settings, but they provide evidence from both higher and lower elevations than KAN_U, thus bracketing out location. From these observations, we cannot explain the MAR-simulated meltwater at KAN_U in 1959–1964 (Fig. 1 of the review) saturating the snowpack and not percolating to available pore space below. We believe this to be corroborated by the concurrent summer (JJA) albedo from MAR, the high values of which do not imply meltwater presence at the surface, but rather increased snow metamorphosis.

Finally, while some runoff still occurs in 2013 (while the summer was cold) as a result of the 2012 summer induced snowpack compaction, runoff disappears in summer 2014 suggesting that we need several successive summers as 2012 to have a significant snowpack degradation.

Correct, and a very good point. Hopefully, with the present study we communicated effectively that the years 2010–2011–2012 were three consecutive years of unusual atmospheric conditions that resulted in
negative net mass budgets and increased melt. The exceptional condition of the firn in 2012 was in part preconditioned by the two previous years. This can be understood by the analysis of the subsurface temperature measurements (Fig. 11a), and is the subject of another study (Charalampidis et al., under review). We should add at this point, that we are very pleased to see agreement on runoff between MAR and our results for the years 2012 and 2013.

Some RACMO (or eventually MAR) outputs could be added in the manuscript to put the 2012 summer in a longer term perspective instead of using Kangerlussuaq measurements.

A very good point, indeed. While in the present study we tried to base our argumentation exclusively on observations, it is our aim to increase the spatiotemporal perspective of our investigations using RCM output, which will be in fact the subject of a forthcoming study.

References


Table I. Surface height changes and mass budgets (measured in winter and calculated in summer) at KAN_U in meters and m w.e., respectively, and ablation duration. The uncertainty of the surface height change is estimated at 0.2 m. The mass budgets are calculated assuming snow density of 360 kg m\(^{-3}\) (the average density of the uppermost 0.9 m measured on 26 April 2013), with uncertainty estimated at 40 kg m\(^{-3}\) (standard deviation among the snow-pit measurements). The snow density assumption was not needed in 2012 and 2013 when actual density measurements were conducted.

<table>
<thead>
<tr>
<th></th>
<th>winter height change</th>
<th>winter budget</th>
<th>summer height change</th>
<th>summer budget</th>
<th>net budget</th>
<th>ablation period</th>
</tr>
</thead>
<tbody>
<tr>
<td>2008–2009</td>
<td>+1.64*</td>
<td>+0.59±0.15</td>
<td>−0.71</td>
<td>−0.26±0.08</td>
<td>+0.34±0.12</td>
<td>01/06–19/08</td>
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<td>2009–2010</td>
<td>+0.70</td>
<td>+0.25±0.08</td>
<td>−1.22</td>
<td>−0.44±0.09</td>
<td>−0.19±0.12</td>
<td>30/04–05/09</td>
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<tr>
<td>2010–2011</td>
<td>+1.02</td>
<td>+0.37±0.08</td>
<td>−1.13</td>
<td>−0.41±0.09</td>
<td>−0.04±0.12</td>
<td>28/05–13/08</td>
</tr>
<tr>
<td>2011–2012**</td>
<td>+0.70</td>
<td>+0.25±0.08</td>
<td>−1.80</td>
<td>−0.86±0.14</td>
<td>−0.61±0.16</td>
<td>27/05–24/08</td>
</tr>
<tr>
<td>2012–2013***</td>
<td>+1.24</td>
<td>+0.45±0.09</td>
<td>−0.75</td>
<td>−0.27±0.08</td>
<td>+0.18±0.12</td>
<td>29/05–17/08</td>
</tr>
</tbody>
</table>

* value inferred from Van de Wal et al. (2012)

** estimate based on snow-pit densities from May 2012

*** estimate based on snow-pit densities from May 2013
Figure I. (a) Monthly air temperature from Kangerlussuaq and at KAN_U. Correlation coefficients: 0.97 for the extent of the KAN_U data, 0.66–0.99 for the months individually, minimum being January. (b) Monthly reference period (1976–1999) air temperature at Kangerlussuaq. (c) Monthly (May to September) and summer (June-July-August average) air temperature anomalies at Kangerlussuaq for the years 2000–2013. Error bars indicate two standard deviations.

Figure II. (a) Location of Site J (Kameda et al., 1994) with respect to KAN_U. (b) Melt percentage data from the top part of a firn core retrieved in 1989 at Site J by the Japanese Arctic Glaciological Expedition (JAGE89; Kameda et al., 2004).
Changing surface-atmosphere energy exchange and refreezing capacity of the lower accumulation area, West Greenland

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Abstract

We present five years (2009–2013) of automatic weather station measurements from the
lower accumulation area (1840 m above sea level) of the Greenland ice sheet in the Kangerlussuaq region, western Greenland. Here, the summers of 2010 and 2012 were both exceptionally warm, but only 2012 resulted in a strongly negative surface mass budget (SMB) and surface meltwater runoff. The observed runoff was due to a large ice fraction in the upper 10 m of firn that prevented meltwater from percolating to available pore volume below. Analysis reveals an anomalously low 2012 summer albedo of ~0.7, as meltwater was present at the ice sheet surface. Consequently, during the 2012 melt season, the ice sheet surface absorbed 28 % (213 MJ m$^{-2}$) more solar radiation than the average of all other years.

A surface energy balance model is used to evaluate the seasonal and interannual variability of all surface energy fluxes. The model reproduces the observed melt rates as well as the SMB for each season. A sensitivity test analysis reveals that 71 % of the additional solar radiation in 2012 was used for melt, corresponding to 36 % (0.64 m) of the 2012 surface lowering. The remaining 1.14 m of surface lowering resulted from was primarily due to the high atmospheric temperatures, up to +2.6 °C daily average, indicating that 2012 would have been a negative SMB year at this site even without the melt-albedo feedback.

Longer time series of SMB, regional temperature and remotely sensed albedo (MODIS) show that 2012 was the first strongly negative SMB year, with the lowest albedo$_2$ at this elevation on record. The warm conditions of recent the last years has resulted in enhanced melt and reduction of the refreezing capacity at in the lower accumulation area. If high temperatures continue, the current lower accumulation area will turn into a region with superimposed ice in coming years.

1 Introduction

Glaciers and ice caps have dominated the cryospheric component to global average sea level rise during the past century (0.5 mm yr$^{-1}$, i.e. or about 70 % of the total cryospheric component for the period 1961–2003; Solomon et al., 2007) due to their relatively short response times to climate variability. However, the largest freshwater reservoir in the Northern Hemisphere is the Greenland ice sheet, which would cause a sea level rise of 7.4 m if completely melted (Bamber et al., 2013). The average sea level rise contribution from the ice sheet has increased from 0.09 mm yr$^{-1}$ over the period 1992–2001 to 0.6 mm yr$^{-1}$ over the period 2002–2011, according to the latest IPCC report (Vaughan et al., 2013). The sheer
volume of the ice sheet and the relatively large warming of the polar regions may yield an increasingly dominant contribution to cryospheric mass loss in coming decades.

An increasingly important driver of this accelerating mass loss is surface melt and subsequent runoff (Shepherd et al., 2012). Between 2009 and 2012, roughly 84% of the increased Greenland ice sheet’s increased mass loss was due to enhanced surface runoff and reduced SMB (Enderlin et al., 2014), causing a reduction of the surface mass budget SMB (SMB; Ettema et al., 2009; 2010; Enderlin et al., 2014). Increased melt is primarily the result of atmospheric warming (Huybrechts and de Wolde, 1999; Huybrechts et al., 2011; Huybrechts and de Wolde, 1999) and the darkening of the ice sheet (Bøggild et al., 2010; Wientjes and Oerlemans, 2010; Box et al., 2012; Van As et al., 2013; Wientjes and Oerlemans, 2010). It has been postulated that the sea level rise associated with an increase in meltwater production can be substantially buffered is commonly expected to be largely compensated for by water refreezing in snow and firn (Harper et al., 2012). However, it has also been suggested that under a moderate warming scenario the ice sheet will lose 50% of its capacity to retain water by the end of the century (Van Angelen et al., 2013), although there is considerable uncertainty involved in retention estimates, based on SMB reconstructions (Vernon et al., 2013).

In-situ measurements are essential for understanding the impact of the changing atmospheric conditions on the ice sheet. In the Kangerlussuaq region, West Greenland, seven automatic weather stations (AWS) and nine SMB stakes constitute a relatively dense network of in-situ measurements (Van de Wal et al., 1995; Greuell et al., 2001; Van den Broeke et al., 2008a; Van As et al., 2012). The uppermost AWS, KAN_U, was established on 4 April 2009 (67° 0' 0" N, 47° 1' 1" W; Fig. 1). Located approximately 140 km inland from the ice margin and at about 1840 m above sea level (a.s.l.), KAN_U is one of the few AWS in Greenland located in the lower accumulation area, where small changes in climate forcing will likely have the largest impact on the ice sheet near-surface stratigraphy.

In the Kangerlussuaq region, approximately 150 km of mountainous tundra separates the ice sheet from the ocean. Characteristic for the ice sheet in this region is a relatively wide (approx. 100 km) ablation area. The equilibrium line altitude (ELA), where annual accumulation and ablation are equal, was estimated to be 1535 m a.s.l. for the period of 1990–2003 (Van de Wal et al., 2005), but is reported to have increased to 1553 m a.s.l. for the
period of 1990–2011 (Van de Wal et al., 2012). At 1520 m a.s.l., superimposed ice becomes apparent at the ice sheet surface at the end of every ablation season, and its upglacier extent is estimated until to reach about 1750 m a.s.l. (Van den Broeke et al., 2008a). The percolation area is found at higher elevations, up to about 2500 m a.s.l., which is the lower limit of the dry snow area.

The ablation area in this region has been studied extensively. Van den Broeke et al. (2008a) presented four years of radiation measurements below the ELA. The lowest albedo values are found at the intermediate AWS S6 (1020 m a.s.l.) due to a surface melt-water induced "dark band" induced by surface meltwater (Greuell, 2000; Wientjes and Oerlemans, 2010). Melt modelling revealed an increase of in summer melt toward the margin, and the a decrease in decreasing role of sensible heat flux with increasing elevation, but also the an increasing in the importance dominance of shortwave radiation in the surface energy balance (SEB) during melt at higher elevations (Van den Broeke et al., 2008b, 2011). An annual cycle in surface roughness length was found to exist over a large part of the ablation area (Smeets and van den Broeke, 2008). This determines part of the variability in the turbulent heat fluxes during the summer months as presented by Van den Broeke et al., (2009). This latter study showed that the regional katabatic nature of the winds over the region, in combination with the variable surface roughness provides significant year-round turbulent heat transfer in a stable surface layer. An increasing wind speed with surface elevation was identified, contrary to what would be expected from katabatically forced wind over an ice surface flattening with elevation. This is due to the larger surface roughness near the margin (Smeets and van den Broeke, 2008), the increasing influence of the large scale pressure gradient force (Van Angelen et al., 2011) and the proximity of pooled cold air over the tundra that sets up an opposing pressure gradient force in the boundary layer during winter. Van As et al. (2012) quantified the extreme surface melt in the Kangerlussuaq region in 2010, validated by river discharge measurements.

At elevations above the superimposed ice area and below the dry snow area (i.e. ~1750–2500 m a.s.l.), sufficient melt occurs to impact snow/firn properties, but not enough to reveal bare ice. In a warming climate with melt at increasingly occurring at higher elevations, this area would comprise an increasingly large portion of the ice sheet due to the ice sheet’s flattening with increasing elevation (McGrath et al., 2013). A rare event in July 2012 caused melt at all elevations of the ice sheet (Nghiem et al., 2012). Bennartz et al. (2013) partially attributed this
Greenland-wide event of increased near-surface temperatures partially to thin, low-level liquid clouds. These clouds, while optically thick and low enough to enhance downward longwave radiation, were thin enough for solar radiation to reach the ice sheet surface. They were present at Summit station about 30% of the time during the 2012 summer months.

A large difference between the ablation and accumulation areas is that in the accumulation area, processes within the snow/firm layers, such as meltwater percolation and refreezing, significantly impact the mass budget (Harper et al., 2012). Additionally, an important process is the melt-albedo feedback (Box et al., 2012). Our aim is to assess the SMB sensitivity to atmospheric forcing in the lower accumulation area using AWS measurements that serve as input for a SEB model. The five-year period with AWS measurements (2009–2013) spans a wide range of melting conditions, including the record melting years of 2010 and 2012 (Tedesco et al., 2011; Van As et al., 2012; Nghiem et al., 2012; Hanna et al. 2014) and years with limited melting such as 2009 and 2013. We add a temporal perspective by discussing Kangerlussuaq temperatures since 1976 and Moderate Resolution Imaging Spectroradiometer (MODIS) albedo values since 2000. In the following, we first describe the observations and SEB calculations, after which we will present the atmospheric conditions and surface energy fluxes at KAN_U, and the changes therein due to recent years with extreme melt. Finally, we investigate the importance of the melt-albedo feedback on the SMB of the lower accumulation area and discuss how changes in the firn can yield SMB variability on an short, interannual time scale.

2 Methods

2.1 AWS measurements

KAN_U is part of the network of ~20 AWSs comprising in the Programme for Monitoring of the Greenland Ice Sheet (PROMICE) network (Ahlstrøm et al., 2014–2008). Measurements include ambient air pressure, relative humidity and aspirated temperature \( T_a \) at 2.7 m height above the ice sheet surface, wind speed and direction at 3.1 m height, as well as incoming and reflected solar/shortwave \( (E_S^\uparrow, E_S^\downarrow) \) and downward and emitted terrestrial/longwave \( (E_L^\downarrow, E_L^\uparrow) \) radiation components at 10-minute intervals. Accumulation and ablation are measured by two sonic rangers, one attached to the AWS and one on a separate stake assembly (Fig. 1b).
Sensor specifications are listed in Table 1. The AWS transmits hourly measurements during the summer period and daily during winter (Citterio et al., 2015).

AWSs installed on glaciers are prone to tilt due to the transient changes in the ice or firn surface. The importance of accounting for the pyranometer tilt has been discussed by e.g., MacWhorter and Weller (1991). AWSs located in accumulation areas are comparatively stable due to the accumulated snow on the base of the tripod. The maximum tilt registered by KAN_U is 3.0 degrees. A tilt correction for the solar radiation measurements is made after Van As (2011).

Two gaps in (sub)hourly measurements exist due to a malfunctioning memory card, from 27 October 2010 until 22 April 2011 and from 26 October 2011 until 21 January 2012. During these periods, when only transmitted values are available, measurements from a second AWS, S10, located in 17 August 2010 at ~50 m distance from KAN_U, were used and adjusted by linear regression to eliminate systematic errors due to a different tilt in measurement heights. The overlapping time series of the two time series revealed high cross-correlations and low Root Mean Squared Deviations (RMSD) for every measured parameter (Table 2). Due to technical issues with S10, $E_L^\uparrow$, $E_L^\downarrow$, and $T_a$ measurement gaps from 9 February 2011 until 30 April 2012 were filled with a similar approach, using measurements from AWS S9 located 53 km closer to the ice sheet margin. Any added uncertainty from using adjusted wintertime measurements will have minimal impact on the summertime results presented below.

The broadband albedo is the fraction of the incoming shortwave radiation reflected at the ice surface and an important parameter in studying the changes in the accumulation area:

$$\alpha = \frac{|E_L^\uparrow|}{|E_L^\downarrow|} \quad (1)$$

To verify its accuracy, albedo was compared for both AWSs KAN_U and S10 for the warm seasons (May–September) of 2010, 2011 and 2012 (Table 2). For hourly values, the RMSD for 2010 and 2011 was only ~0.03. The RMSD for 2012 was 0.07, due to the larger higher spatial variability in surface reflectance after substantial melt.
2.2 Surface radiation budget

The radiation budget at the ice sheet surface is given by the sum of solar and terrestrial radiation components:

\[ E_R = E_{S↓} + E_{S↑} + E_{L↓} + E_{L↑} = E_{S\text{Net}} + E_{L\text{Net}} \quad (2) \]

Fluxes are here taken as positive when directed toward the ice sheet surface. By the inclusion of albedo and utilizing the Stefan-Boltzmann law, this can be rewritten as:

\[ E_R = (1 - \alpha) E_{S↓} + \varepsilon E_{L↓} - \varepsilon \sigma T_s^4 \quad (3) \]

with \( \sigma \) being the Stefan-Boltzmann constant \((5.67 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4})\) and \( T_s \) the surface temperature. The longwave emissivity \( \varepsilon \) for snow/firn is assumed equal to 1 (black-body assumption).

2.3 SEB model

Various studies have applied SEB models in glaciated areas under different climatic conditions, such as a high Antarctic plateau \((\text{-} \text{Van As et al., 2005})\) and the Greenland ablation area \((\text{-} \text{Van den Broeke et al., 2008b-\text{z} 2011})\). The energy balance at the atmosphere-surface interface is:

\[ E_M = E_R + E_H + E_E + E_G + E_P \quad (4) \]

where \( E_H, E_E, E_G \) and \( E_P \) are the turbulent sensible, turbulent latent, subsurface conductive and rain-induced energy fluxes, respectively.

Rainfall is assumed to be at melting-point temperature \((T_0 = 273.15 \text{ K})\), thus \( E_P \) is non-zero when \( T_s \) is below freezing:

\[ E_P = \rho_w c_w \dot{r}(T_0 - T_s) \quad (5) \]

where \( c_w \) is the specific heat of water \((4.21 \text{ kJ kg}^{-1} \text{ K}^{-1} \text{ at } 0 \text{ °C} \text{ and } 999.84 \text{ kg m}^{-3})\) and \( \dot{r} \) is the rainfall rate. The latter is assumed to be non-zero when under conditions of heavy cloud cover during periods with non-freezing air temperatures are met (see below).

The energy balance is solved for the one unknown variable \( T_s \). If \( T_s \) is limited to the melting-
point temperature (273.15 K), the imbalance in Eq. 4 is attributed to melt \( (E_M) \). For sub-
freezing \( T_s \) values all other SEB components are in balance and surface melt does not occur.
\( E_H \) and \( E_E \) are calculated using the “bulk method” as described by Van As et al. (2005). This
method uses atmospheric stability, and thus depends on \( T_s \), implying that Eq. 4 has to be
solved iteratively.

The average surface roughness length for momentum \( z_0 \) at this elevation would realistically be
\(~10^{-4} \) m (Smeets and van den Broeke, 2008). During summer, the ice sheet surface melts
occasionally, smoothing it and thus attaining a smaller \( z_0 \) \((~10^{-5} \) m). Slightly increased
roughness is expected during wintertime due to sastrugi and drifting snow (Lenaerts et al.,
2014) can increase \( z_0 \) in cases up to \( 10^{-3} \) m (Lenaerts et al., 2014). In the present study, \( z_0 \) is
kept constant at \( 10^{-4} \) m. A series of test runs showed that the results of this study were not
very sensitive to in the range of plausible \( z_0 \) values. The scalar roughness lengths for heat and
moisture are calculated according to Andreas (1987).

Subsurface heat transfer is calculated on a 200-layer grid with 0.1 m spacing (20 m total) and
is forced by temperature changes at the surface and latent heat release when water refreezes
within the firm. Heat conduction is calculated using effective conductivity as a function of firm
density (Sturm et al., 1997) and specific heat of firm as a function of temperature (Yen, 1981).
The calculations are initialized using thermistor string temperatures from April 2009 and
depth-adjusted firm core densities measured on 2 May 2012. The subsurface part of the model
includes a percolation/refreezing scheme based on Illangasekare et al. (1990), assuming active
percolation within snow/firm. Provided that there is production of meltwater at the surface, the
amount of refreezing is limited either by the available pore volume or by the available cold
content at each level. The scheme simulates water transport and subsequent refreezing as the
progression of a uniform warming front into the snow/firm and is active for all melt seasons
except for 2012. In 2012, surface runoff dominated water movement after 14 July, as clearly
visible on Landsat imagery (not shown). This coincided with the surfacing of a 6 m thick ice
layer in the model, which was also found in firm cores (Machguth et al., under review).

**Consistent with these observations. Here** we use 6 m of ice (density of 900 kg m\(^{-3}\)) as a
threshold that causes meltwater to run off horizontally, shutting down vertical percolation.

**Solid precipitation is added in the model based on KAN_U sonic ranger measurements,**
assuming thea rounded average snow density of 400 kg m\(^{-3}\) observed found in snow-pit
measurements, i.e. 400 kg m⁻³. Solid precipitation is added in the model based on KAN-U sonic ranger measurements, assuming snow density of 400 kg m⁻³ in accordance to snow pit measurements. Although rain occurs only sporadically–infrequently at 1840 m a.s.l., a rain estimate is incorporated with prescribed precipitation rates for each year during hours with a thick cloud cover producing \( E_L \) values that exceed black-body radiation using the air temperature \( (E_L > \sigma T^4_a) \) and \( T_a \) is above freezing. Testing–Evaluating this against winter accumulation, the following precipitation rates were found derived and applied prescribed to the rain calculation: \( 2.0 \times 10^{-3} \) m h⁻¹ for 2009–2010 and 2012–2013, \( 3.5 \times 10^{-3} \) m h⁻¹ for 2010–2011 and \( 0.5 \times 10^{-3} \) m h⁻¹ for 2011–2012. Using this approach, the model produces liquid precipitation during the summer months only; by the end of the five-year period it amounts to a total of 0.26 m water equivalent (m w.e.), 15 % of the total precipitation over the five years. The contribution of rain in the energy balance is minor; the total energy added to the surface for the whole study period is approximately 1.15 MJ m⁻², which could yield a total of 9 mm of melted snow.

The performance of the model in terms of ablation is illustrated by comparing the simulated surface changes with the measured surface height changes due to ablation and accumulation in (Fig. 2a). The model accurately reproduces the melt rates during every melt season. Yet this validation does not cover the whole melt season. We found that the AWS tripod and stake assembly are prone to sinking somewhat into warm, melting firm during the second part of the melt season (note the measurement gaps). In a second model validation method–exercise, we compare the simulated and with the measured \( T_s \) (inferred from the \( E_L \)) in Fig. 2b, and find they correlate well \( (R^2 = 0.98) \) with an average difference of 0.11 °C and Root mean square–squared Error error (RMSE) of 1.43 °C. Part of this difference can be attributed to the seemingly overestimated 10 % \( E_L \) measurement uncertainty as reported by the sensor manufacturer, which would yield a RMSE of 6.2 °C of \( T_s \) values.

2.4 Additional measurements

For a study with a five-year time span, it is useful to provide a longer temporal perspective. For this, we use North Atlantic Oscillation (NAO) index measurements from the National Oceanic and Atmospheric Administration (NOAA). We use the air temperature record from Kangerlussuaq airport observed by the Danish Meteorological Institute (DMI) since initiated in—1973 in support of to facilitate aircraft flight operations (Cappelen, 2013).
observational suite coverage is available from since 1976. Monthly $T_d$ from the airport correlate well with the KAN_U time series ($R = 0.97$), indicating that Kangerlussuaq measurements can be used for providing temporal perspective, despite in spite of the 160 km distance that separates the two measurement sites. Finally, from the 5x5 km regridded MODIS albedo measurements (MOD10A1) we use the pixel nearest to KAN_U in 5 by 5 km regridded MODIS albedo product (MOD10A1) to investigate albedo variability over in a-the 2000–2013 perspective period.

3 Results

3.1 Meteorological observations

The importance of katabatic and synoptic forcing on near-surface wind direction is are roughly equivalent equally important; at the elevation of KAN_U (Van Angelen et al., 2011). The average wind direction is south-southeast (~150°; Fig. 3a). Yet However, in a case study of the 2012/2013 winter (Van As et al., 2014), the prevailing wind direction was ~135° (southeast), suggesting an influential katabatic regime in which air drains down-slope and is deflected by the Coriolis forcing effect. Wind speed is higher during winter (Fig. 3b); annual average values are 6–7 m s$^{-1}$ whereas summer average values are around 5 m s$^{-1}$ (Table 3). Winds exceeding 15 m s$^{-1}$ occur primarily during the winter period and rarely exceed 20 m s$^{-1}$ when averaged over 24 hours. The barometric pressure of about 800 hPa exhibits an annual cycle with relatively high pressure in summer (Fig. 3c), favoring more stable, clear-sky conditions. Also the specific humidity varies annually, peaking in summer. Annual values are about 1.5 g kg$^{-1}$.

The year 2010 was the warmest year of the record (Table 3), with the winter (December-January-February) of 2009–2010 being 4.0 °C warmer than the 2009–2013 average, and the summer (June-July-August) only being equaled by 2012 (~1.8 °C; Table 3). Especially May 2010 was especially warm, at ~6.2 °C, or 5.1 °C above the 2009–2013 average. Positive $T_d$ persisted during the end of the melt season resulting in a ~1.1 °C monthly average for August. The high 2010 temperatures influenced impacted the surface ablation by inducing early onset. In 2010, ablation which lasted from late April until early September, whereas for instance the 2009 melt season at KAN_U occurred spanned early June until mid-August.
The average SMB over the period 1994–2010 at KAN_U is +0.27 m w.e. (Van de Wal et al., 2012). Melt at this elevation occurs frequently during each melt season. The winter 2009/2010 accumulation was relatively low, amounting to 65% of the 2009–2013 average (0.25 m w.e.; Table 4). During the 2010 melt season, all the snow that had accumulated since the end of the previous melt season ablated, including part of the underlying firn, resulting in the first negative SMB year on record (Table 4). The stake measurements from Van de Wal et al. (2012) document a two-year surface height change of +0.42 m on average for 2008–2010 at the same location (S10), corresponding to +0.15 m w.e. assuming a snow pit density of 360 kg m⁻³. From this estimate, we infer the winter and net SMB for 2009 to be +0.59 and +0.34 m w.e., respectively.

During winter 2011–2012, accumulation was the same as in winter 2009–2010. In spring 2012, positive $T_a$ was first recorded during April (with −12.8 °C the warmest April on record), followed by a relatively warm May (−8.6 °C). Already in late May 2012 ablation rates were high (7.2 mm w.e. d⁻¹; Charalampidis and van As, 2015 in press). June and July were the warmest of the five-year record with −1.5 °C and −0.6 °C monthly average $T_a$, respectively. With the summer of 2012 on average as warm as that of 2010, but the ablation period shorter by 398 days (Table 4), the summer SMB was −0.86 m w.e., making 2012 the most strongly negative SMB year (−0.61 m w.e.) to be recorded at this location (Van de Wal et al., 2005; 2012).

### 3.2 Surface energy fluxes

Solar radiation exhibits a strong annual cycle at this location above the Arctic Circle (Fig. 4a). In the absence of topographic shading or a significant surface slope (~0.37°) the day-to-day variability in incoming shortwave radiation at this elevation is dominated by cloudiness and the solar zenith angle. The highest daily $E_S^\uparrow$ values occur in June and exceed 400 W m⁻², while at the ELA they are just below 400 W m⁻² (Van den Broeke et al., 2008a) due to more frequent cloud cover and a thicker overlying atmosphere. Whereas $E_S^\uparrow$ increases with elevation from the ELA to KAN_U, $E_S^\text{Net}$ obtains values of up to 100 W m⁻² both at the ELA and at KAN_U, implying regulated solar energy input that is regulated by surface reflectance.

Terrestrial radiation exhibits an annual cycle of smaller amplitude (Fig. 4a). The annual
variations of the downward and emitted longwave radiation are governed by the temperature and emissivity variations of the atmosphere and the ice sheet surface, respectively. Hence, the absolute magnitudes of both components are larger during the summer period. $E_L^{↓}$ fluctuations depend primarily on cloud cover. $E_L^{↑}$ is a sink to the SEB and during summer is limited by the melting surface with the maximum energy loss being 316 W m$^{-2}$. This results in predominantly negative $E_L^{\text{Net}}$ values throughout the year. The energy loss peaks during June and July.

The $E_R$ annual cycle displays an energy gain at the ice sheet surface during May to August and energy loss the rest of the year (Fig. 4b). This winter energy loss is primarily compensated by downward sensible heat flux. Calculated $E_H$ is typically positive throughout the year, with highest values in winter when $E_R$ is most negative, heating the ice sheet surface while cooling the atmospheric boundary layer (Fig 4b). This facilitates the katabatic forcing and thus enhances wind speed and further turbulent energy exchange between the atmosphere and the ice sheet surface. The contribution of $E_H$ to melt is smaller than at lower elevations (Van den Broeke et al., 2011). The dominant melt energy source at KAN_U is $E_R$.

$E_E$ changes sign from winter to summer, and is on average a small contributor to the annual SEB. During the summer period, $E_E$ is comparable to $E_H$ but with opposite sign, enabling surface cooling by sublimation and/or evaporation. In the winter, $E_E$ is directed mostly toward the cold ice sheet surface, resulting in heating from deposition.

The annually averaged $E_G$ is mostly negative and of the same magnitude as $E_E$ (3–4 W m$^{-2}$), but with no distinct annual cycle. Melt seasons with substantial refreezing exhibit increased positive summer-averaged $E_G$ since the near-surface firn temperature is on average higher than $T_s$, leading to conductive heat transport toward the ice sheet surface. Low $E_G$ values in summer indicate limited refreezing in the firn just below the ice sheet surface.

$E_P$ is non-zero, but still negligible in summer, when positive air temperatures occur and thus precipitation is liquid.

### 3.3 Interannual variability of the SEB and implications for melt

With the exception of August 2009, when predominantly clear skies caused $E_S^{↓}$ being 40 W m$^{-2}$ larger, and $E_L^{↓}$ 36 W m$^{-2}$ smaller, than in the other years, monthly average values of $E_S^{↓}$
at this site are fairly invariant (difference < 25 W m$^{-2}$; Fig. 5a). Often $E_R$ increases when clouds are present over an ice sheet; a so-called radiation paradox (Ambach, 1974), as it was observed in April 2012.

Figure 5b illustrates the annual cycle of monthly average albedo, not including the winter months October–February, when shortwave radiation values are too low for accurate albedo estimation, yet it is expected to be characteristic of attain fresh dry snow values (0.8–0.9). High albedo persists until May due to fresh snow deposited on the ice sheet surface. An exception occurred during March and April 2013, when the monthly albedo of 0.78 suggests reduced precipitation input for a prolonged period and the presence of aging old, clean, dry snow on the ice sheet surface (Cuffey and Paterson, 2010). In the years 2009–2011 and 2013 the albedo gradually decreased beginning late May and during the summer due to the effects of relatively high temperatures and melt on snow metamorphism. During summer, albedo still exceeded 0.75. While in August melt at KAN_U can still occur, this does not counteract the effect of snowfall events that increase the surface albedo.

The anomalously warm period in June and July 2012 (Fig. 3d) coincided with a larger decrease of surface albedo than in the other years. The combination of enhanced melting, heat-induced metamorphism and firn saturation, reduced the albedo from 0.85 in April to 0.67 in July reaching a value that is characteristic of corresponds to the lower soaked snow facies close to the lower elevation snow/firn line (Cuffey and Paterson, 2010). As a consequence, $E_{S_{Net}}$ increased by approximately 25 W m$^{-2}$ in June and July (32%; Fig. 5c). This darkening thus functions as an amplifier of melt (Box et al., 2012; Van As et al., 2013), contributing to the large observed ablation rates (Table 4).

The longest longwave radiation surface emissions occurred during August 2010 and June–July 2012, approaching the theoretical limit of −316 W m$^{-2}$ for a continuously melting ice sheet surface (Fig. 6b). The concurrent high $E_{L_{Net}}$ (Fig. 6a; Table 5) was related to high atmospheric temperatures. This caused summer $E_{L_{Net}}$ in the summers 2010 and 2012 to exceed its value in other years (Table 5; Van As et al., 2012). While summer $E_{S_{Net}}$ was similar in 2009 and 2010, summer $E_R$ was about 690% larger in 2010 than in 2009, primarily due to the high atmospheric temperatures alone. During summer 2012, summer $E_{L_{Net}}$ was similar as in 2010. The large summer $E_{S_{Net}}$, resulted in summer $E_R$ 67% higher than in 2010 (Table 5). The highest daily $E_R$ attained 100 W m$^{-2}$ on 9 July and coincided with the start of a Greenland-
wide warm event. On 12 July, nearly the entire ice sheet surface was reported to melt (Nghiem et al., 2012), followed shortly after by the highest meltwater discharge in 56 years on 12 July 2012, as inferred can be concluded by the partial destruction of a bridge constructed over the the 1956 Watson river bridge in Kangerlussuaq in 1956, on 12 July 2012. At KAN_U, WWell above the long-term ELA at KAN_U, not only a strongly negative SMB was recorded in 2012, but it was the only year with a positive annual radiation budget ($E_R = +4 \text{ W m}^{-2}$; Table 5).

$E_H$ was largest during 2010 and smallest during 2011 (Table 5), the years of highest and lowest annual $T_o$, respectively (Table 3). Sensible heat transfer toward the ice sheet surface was also low on average also low—in 2012, owing to the cold winter months. The high July 2011 $E_H$ was due to warm air advection that occurred over a cold surface, yielding large near-surface temperature gradients and thus sensible heat exchange (Fig. 7a). In the During summer of 2013, when the air temperatures remained relatively low, the ice sheet surface exhibited the lowest sensible heat gain compared to the other melt seasons. In all, $E_H$ did not contribute to SEB interannual variability as much as the radiative components.

Summer $E_E$ values are correlated with the atmospheric pressure ($R = 0.96$), which that influences the gradients in near-surface specific humidity and wind speed. In the During summer of 2012, pressure and specific humidity were relatively high (811 hPa and 3.7 g kg$^{-1}$, respectively; Table 3), while the wind speed was reduced, thus contributing to the lowest absolute summer $E_E$ with the lowest cooling rates due to evaporation/sublimation. The maximum latent heat loss that year occurred in May. Thereafter, the moisture content in the near-surface air became relatively large, with $E_E$ decreasing in absolute value until July. Summer 2013, was conversely characterized by relatively low pressure and specific humidity (804 hPa and 2.8 g kg$^{-1}$, respectively) resulting in high evaporation/sublimation rates especially in June and July (Fig. 7b).

Monthly $E_G$ values were small and displayed small interannual variability especially in summer. The summers of 2010 and 2011 exhibited the most positive $E_G$ as a consequence of substantial refreezing (Fig. 7c), which influenced impacting the near-surface firm temperature gradients in the firm. The summer average Summer $E_G$ values for 2009 and 2013 (Table 5) were was lower due to the moderate melt seasons of smaller duration. Summer $E_G$ was lower in summer 2012 due to a warm ice sheet surface conducting heat into the firm and the absence
of refreezing.

The melt rates in 2009 and 2013 were similar. In both years, with the largest $E_M$ occurring in July and did not exceeding 30 W m$^{-2}$ (Fig. 8a). $E_M$ peaked similarly in 2010 and 2011, in June reaching about 20 W m$^{-2}$ and in July exceeding 35 W m$^{-2}$. May and August 2010 sustained exhibited significant melt in response to the warm atmospheric conditions (Van As et al., 2012). Both 2010 and 2012 exhibited significant melt in May (10 W m$^{-2}$). During summer 2012, $E_M$ far exceeded all other years, with e.g., a July value of 68 W m$^{-2}$, leading to the largest ablation reported in Table 4.

There is a large dominance of the radiative fluxes dominate over the interannual variability of melt at KAN_U, with the variations in $E_L^{Net}$ being most influential over the amount of available $E_M$ in the years 2009–2011 and 2013. In 2012, it was the large $E_S^{Net}$ that mainly contributed to the melt anomaly.

### 3.4 Melt-albedo feedback

Figure 9a, which depicts total monthly surface energy exchanges throughout the study period, illustrates that $E_L^{Net}$ and $E_L^{Net}$ dominate the SEB from May to September, while $E_L^{Net}$ and $E_H$ dominate the SEB during the remainder of the year. During the years exclusive of 2012 considered here (2009, 2010, 2011 and 2013), the total summer energy input to the ice sheet surface was 620–650 MJ m$^{-2}$. This energy reaches a peak in July. In July 2010, for example, the total energy input reached 246 MJ m$^{-2}$. By contrast, in 2012, the total summer energy input exceeded 770 MJ m$^{-2}$, and in July it reached 304 MJ m$^{-2}$. The 2012 total energy used for melt was 414 MJ m$^{-2}$ (65 % higher than in 2010), of which 183 MJ m$^{-2}$ was used for melt in July. Figure 9a shows the total surface energy exchange for each month throughout the study period. It illustrates that $E_L^{Net}$ and $E_L^{Net}$ dominate the energy balance during the months May to September, while $E_L^{Net}$ and $E_H$ govern the SEB during wintertime (Fig. 9a). For the years 2009–2011 and 2013 the total energy in the summer months was 1250–1300 MJ m$^{-2}$. In July, when the energy input is largest, the years 2009 and 2013 exhibited a total amount of energy of 464 MJ m$^{-2}$ while in the years 2010 and 2011 it reached 500 MJ m$^{-2}$. In 2012, the summer total energy was significantly higher exceeding 1500 MJ m$^{-2}$, while July alone reached 608 MJ m$^{-2}$ with 183 MJ m$^{-2}$ invested in melting. The total melt energy in 2012 amounted to 414 MJ m$^{-2}$.
Figure 9b illustrates the simulated mass fluxes at the ice sheet surface (note the different y scales for positive and negative values). A total of 40 kg m$^{-2}$ of mass loss occurs on average by the sum of sublimation and evaporation during spring and summer. Conversely, deposition amounts to 10 kg m$^{-2}$ each winter season. The total snowfall from April 2009 until September 2013 amounted ~1500 kg m$^{-2}$ (also Table 4). Up to the end of May 2012, all meltwater had accumulated internally through percolation into the firn, adding a mass of 1158 kg m$^{-2}$ (1020 kg m$^{-2}$ from snowfall and 138 kg m$^{-2}$ from rainfall). Due to ice layer blocking vertical percolation in summer 2012, 444 kg m$^{-2}$ ran off, removing approximately 38% of accumulated mass since April 2009.

The total amount of meltwater generated at the ice sheet surface, equaling equivalent to the sum of runoff and refreezing minus rainfall, amounted 1307–1232 kg m$^{-2}$ in 2012. As the calculated surface ablation was 860 kg m$^{-2}$ (Table 4), 30.66% (372477 kg m$^{-2}$) of the produced meltwater was melted more than once during the ablation season. This suggests that essentially, 48% (416 kg m$^{-2}$) (48% of the total ablation or 34% of the produced meltwater) was effectively retained in near surface firn layers.

2010 was the first year on record during which surface ablation exceeded accumulation from the preceding winter at KAN_U (Table 4; Van de Wal et al., 2012). Even though atmospheric temperatures were high and the impact on ablation was large in 2010, the response of the snow surface was much larger in 2012, when ablation was more than three times larger than the accumulation. Albedo-Jin 2012, albedo dropped decreased to -0.7 already by in mid-June (Charalampidis and van As, in press 2015), implying substantial metamorphosis of the snow surface, while in all other years this value albedo was reached only in July or August. The albedo reduced even more on 10 July (-0.6), signifying the saturation of the ice sheet surface and the exposure of thick firn. Until 6 August, the albedo value corresponded to that of soaked facies close to the snow/firn line (Cuffey and Paterson, 2010). It should be noted that snowfall events increased the albedo during several periods in the summer season (Charalampidis and van As, in press 2015).

To quantify the impact of a relatively dark ice sheet surface on the SEB, the average annual cycle in albedo of all years excluding 2012 was used to replace the low 2012 albedo in dedicated a model-sensitivity analysis run to quantify the impact of a relatively dark ice sheet surface on the SEB. Figure 10a shows the albedo anomaly of 2012, which resulted in started
invoking enhanced ablation in late May/early June (Fig. 10b). At the end of August, the ice sheet surface lowered an additional 0.64 m due to 58% more melt energy compared to a situation with average albedo. The excess $E_M$ from the melt-albedo feedback amounted to 152 MJ m$^{-2}$, while the excess $E_{\text{eff}}^{\text{Net}}$ supplied was 213 MJ m$^{-2}$ (Fig. 10c). The remaining $E_{\text{eff}}^{\text{Net}}$ was consumed by other fluxes, primarily $E_H$. As the total surface ablation of 2012 was 1.78 m of surface height change (Fig. 2) the remaining 1.14 m was primarily due to the warm atmospheric conditions and similar to 2010 (1.21 m). This sensitivity analysis implies that the location would have experienced a negative SMB in 2012 even without the melt-albedo feedback.

4 Discussion

4.1 Uncertainties

Model performance is limited by the accuracy of the instruments of KAN_U as given in Table 1. The radiometer uncertainties are the largest, based on what is reported by the manufacturer (10% for daily totals, although this is likely to be an overestimate (Van den Broeke et al., 2004)). Nevertheless, the accurate simulation of surface temperature and snow ablation rates (Fig. 2) throughout the period of observations builds confidence in both the measurements and the model.

The model exhibits considerable sensitivity to the subsurface calculations, suggesting importance of pore volume and firn temperature, and how much more complicated SEB calculations are in the lower accumulation area than for bare ice in the ablation area. The model is able to capture the seasonal variations of temperature in the firn and calculated temperatures are commonly within 3.6 °C of those measured with an average of −0.3 °C (Fig. 11a). The shallow percolation of a wetting front in the firn is estimated at depths of 1–3 m in the years 2009 and 2013 (Fig. 11b), while in the years of larger melt, pore volume until 10 m depth is filled, possibly overestimating the percolation depth given the relative temperature buildup in the simulated firn below roughly 5 m depth (Fig. 11a; Charalampidis et al., under review). In particular for 2012, available simulated pore volume at ~6 m is significantly affected by meltwater that percolates below the formed thick ice layer, which may indicate that the runoff threshold of a 6 m ice layer is an overestimate, highlighting the need for a
better runoff criterion. Further investigation on this criterion and inclusion of water content held in the firm by capillary forces, saturation of the surface and proximity of impermeable ice to the surface is necessary.

The fact that the subsurface calculations are initialized in 2009 by use of vertically shifted firm densities from a 2012 core does not influence the calculation of the surface energy fluxes and thus the outcome of this paper. Importantly, the timing that simulated runoff occurred in July 2012 is in agreement with satellite observations due to the runoff criterion, thereby providing confidence in realistic calculation of \( E_G \).

Although the subsurface calculations are on a vertical grid of 10 cm (see also 2.3), there is loss of detail in the density profile with time due to the interpolation scheme that shifts the column vertically when it needs to account for surface height variations (Fig. 11c). Increased spatial resolution requires a finer temporal resolution to avoid model instability. Since this would also increase calculation time severely, while the calculated SEB would be unaffected, we accepted the loss of detail in density in this study. Nevertheless, during each melt season when refreezing is important to be accounted for, no loss of detail is expected near the surface since the column is shifted almost continuously upward.

Rainfall is known to occur during summer on the higher elevations of the ice sheet. The exact amount is unknown as in situ measurements for precipitation are rare and difficult. Therefore, the rainfall calculated by our model should be considered a first-order estimate. Nevertheless, the amount of rain is expected to be small and its effect on the SEB is negligible, as shown by the model results.

It is possible that other factors than heat-induced snow metamorphism and the presence of surface water contributed to the 2012 albedo anomaly. Such could be aerosol particles or impurities at the snow surface, effectively reducing its albedo (Doherty et al., 2013). Also, in cases of extreme melt, microbial activity can develop at the ice sheet surface with the subsequent production of a dark-colored pigment (Benning et al., 2014).

4.2 Long-term perspective in temperature and albedo

The regional barometric pressure is linked to the NAO with KAN_U measurements exhibiting a negative correlation (\( R = -0.72 \) for monthly averages) with the NAO index. The correlation
between the NAO index and pressure in west Greenland implies that the variability of the regional climate is partly controlled by the North Atlantic climate (Figs. 12a,b). For instance, warm summers on the ice sheet coincide with extended periods of negative NAO index. The Kangerlussuaq airport temperature record since 1976 was used to provide a temporal perspective to the KAN_U temperature in recent years (Fig. 12c). The standard deviations reveal variability during the winter period of more than 10 °C while for the months of July and August standard deviations are ~2.0 °C. The temperature measurements reveal that the region has been warming on average starting in 1996 (not shown). Figure 12d illustrates this for 2000–2013; e.g. the summers (JJA) were 1.2 °C warmer than in the reference period 1976–1999. The warm 2010 and 2012 summer have an anomaly value of +1.9 °C and +1.8 °C, respectively. The high temperatures in recent years are most apparent for June when in 10 out of 14 years the 1976–1999 standard deviation is exceeded. A further increase of the regional temperatures, as anticipated by climate models, will likely further increase the frequency of large melt events and the extent of each melt season, leading to conditions similar to or more extreme than in 2012 (McGrath et al., 2013).

The MOD10A1 time series from the years 2000–2013 shows an albedo decrease of 0.05–0.10 during the 14 years of measurements in response to the increased temperatures (Fig. 13). In particular, May albedo reached record low values in 2010 and 2012. July albedo is considerably lower in the years 2007–2013 than it was in the first half of the record. The exceptional surface conditions in July 2012 were also captured by MODIS with the lowest monthly albedo (~0.6) of the time series. The albedo in August is generally higher than in July due to snowfall, but the values remain sufficiently low to enhance melt. We note that part of the MODIS based albedo decrease could be the result of the declining instrument sensitivity of the Terra MODIS sensor (Wang et al. 2012; Lyapustin et al. 2014) though updated (through 2014) comparisons between MOD10A1 and ground observations from GC-Net data (Box et al. 2012; not shown) do not indicate an obvious nor statistically significant difference.

Increased meltwater infiltration into the firn during events of increased melt has led to the formation of thick, near-surface ice lenses between 2 to 7 m, judging from the 2012 firn core. This contrasts the aquifers (i.e. liquid water storage) that are observed in the firn in southeast Greenland (Forster et al., 2013; Koening et al., 2014; Kuipers Munneke et al., 2014). The southwestern ice sheet receives about one third of the annual precipitation amount in the
southeast (Ettema et al., 2009), resulting in differences in thermal insulation of the infiltrated water and available pore volume. Persistent shallow refreezing on an interannual scale has led to the formation of thick impermeable ice enabling runoff in 2012 (Machguth et al., under review).

The DMI measurements indicate that 2009 is representative of the reference period 1976–1999 (Fig. 11c; Van As et al., 2012). In summer with respect to summer 2009, 2010 the radiation budget in summer 2010 was increased due to reduced lower $E_L^{\text{Net}}$ by 10 W m$^{-2}$ (Table 5; Sect. 3.3). The concurrent increase in melt was 13 W m$^{-2}$. In summer 2012, $E_L^{\text{Net}}$ was the same, while $E_S^{\text{Net}}$ was 47 W m$^{-2}$ larger than in 2010, approximately 70% Most of which this $E_S^{\text{Net}}$ excess was consumed by melting (42 W m$^{-2}$, Sect. 3.4). The melt-albedo feedback (Box et al., 2012) will contribute to the rise of the ELA in a warming climate (Fettweis, 2007; Van de Wal et al., 2012) and might transform the lower accumulation area into superimposed ice if warming prevails. We have shown that the melt-albedo feedback makes that warm summers can have great impact on melt and runoff in the lower accumulation area. Our results suggest that if warm atmospheric conditions persist in the future, the additional input of solar radiation at the ice sheet surface will be of higher importance to surface changes than atmospheric warming.

5 Conclusions

We used five years of automatic weather station measurements to characterize the prevailing meteorology and surface energy fluxes at a location in the lower accumulation area of the southwest Greenland ice sheet. The analysis showed large control of the radiative components over the interannual variability, and mostly from net longwave radiation. The main contributor to melt is absorbed solar radiation. In all but one year of observations, however, this solar radiation did not induce surface mass balance variability. This was not the case during the 2012 melt season, when the area attained unusually low albedo values (<0.7) owing to large melt and the subsequent exposure of water-saturated, high-density firn. The consequential enhanced solar absorption along with warm atmospheric conditions resulted in intensified melt during the strongest most negative surface mass budget year since 1994, and presumably since at least 1976 given the Kangerlussuaq temperature record. A sensitivity test with our energy balance model indicates that the melt-
albedo feedback contributed an additional 58% (152 MJ m$^{-2}$) to melt energy in 2012, though increased atmospheric temperatures alone would have yielded a negative surface mass budget as well.

Percolation of meltwater within snow and firn is generally considered to refreeze in firn at 1840 m a.s.l. on the ice sheet, which prevents runoff and therefore limits Greenland’s contribution to sea level rise. This concept is applicable to all the higher elevation regions of the ice sheet where there is that experience moderate melt and where deep percolation is possible. However, the lower accumulation area of the southwestern ice sheet showed high sensitivity to the warm atmosphere in 2012, primarily due to because of the relatively low precipitation in the region, which prevents the timely replenishment of saturated enables the immediate loss of-pore volume in the near surface upper meters of-firn under extreme melting conditions. Water retained in the firn can lead to substantial density increase due to refreezing, which in warm years not only may function as a mechanism to block percolation, but also lowers the surface albedo and enhances melt, accelerating the transformation of the lower accumulation area underlain by firn into an ablation-dominated region area underlain by superimposed ice. This highlights the importance of accurately modelling of percolation and refreezing within the firn, in order to best be able to estimate the sea level rise contribution associated with from the Greenland ice sheet to sea level rise meltwater production.

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We are grateful to Xavier Fettweis and an anonymous reviewer for constructive comments. We also thank Andreas Mikkelsen for snow density measurements, and Robert Fausto, Filippo Cali Quaglia and Daniel Binder for valuable discussions. The KAN_U weather station was funded by the nuclear waste management organizations in Sweden (Svensk Kärnbränslehantering AB), Finland (Posiva Oy) and Canada (NWMO) through the Greenland Analogue Project (GAP Sub-Project A). It is operated by the Geological Survey of Denmark and Greenland (GEUS) with 2009–2012 logistical, technical and manpower support from Aberystwyth University funded through the UK Natural Environmental Research Council (NERC grant NE/G005796/1), a Royal Geographical Society (RGS) Gilchrist Fieldwork Award, and an Aberystwyth University doctoral scholarship. Technical and salary support
were received from the Programme for Monitoring of the Greenland Ice Sheet (PROMICE), launched and funded by the Danish Energy Agency (Energistyrelsen) under the Danish Ministry of Energy, Utilities and Climate, and within the Danish Cooperation for Environment in the Arctic (DANCEA). We further acknowledge support from the Netherlands Polar Programme of the Netherlands Organization for Scientific Research (NWO). This is a PROMICE publication and contribution number 52 of the Nordic Centre of Excellence SVALI, “Stability and Variations of Arctic Land Ice”, funded by the Nordic Top-level Research Initiative (TRI). We thank Andreas Mikkelsen for collection of snow densities, Alun Hubbard and Samuel Doyle for retrieval of KAN_U weather station measurements and Robert Fausto, Filippo Calì Quaglia and Daniel Binder for valuable discussions. The KAN weather stations are financed by the Greenland Analogue Project (GAP) and operated by the Geological Survey of Denmark and Greenland (GEUS). We also acknowledge support from the Netherlands Polar Programme of the Netherlands Organization for Scientific Research, section Earth and Life Sciences (NWO/ALW). This is a publication in the framework of the Programme for Monitoring of the Greenland Ice Sheet (PROMICE) and contribution number 52 of the Nordic Centre of Excellence SVALI, “Stability and Variations of Arctic Land Ice”, funded by the Nordic Top-level Research Initiative (TRI).
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<table>
<thead>
<tr>
<th>parameter</th>
<th>sensor</th>
<th>accuracy</th>
</tr>
</thead>
<tbody>
<tr>
<td>air pressure</td>
<td>Campbell CS100</td>
<td>2 hPa at –40 °C to 60 °C</td>
</tr>
<tr>
<td>aspirated air temperature</td>
<td>Rotronic MP100H aspirated</td>
<td>0.03 °C at 0 °C</td>
</tr>
<tr>
<td></td>
<td>(Pt100)</td>
<td></td>
</tr>
<tr>
<td>relative humidity</td>
<td>Rotronic MP100H aspirated</td>
<td>1.5% at 23 °C</td>
</tr>
<tr>
<td></td>
<td>(HygroClip R3)</td>
<td></td>
</tr>
<tr>
<td>shortwave radiation</td>
<td>Kipp &amp; Zonen CNR4</td>
<td>10% for daily totals</td>
</tr>
<tr>
<td>(incoming and reflected)</td>
<td>(Pyranometer)</td>
<td></td>
</tr>
<tr>
<td>longwave radiation</td>
<td>Kipp &amp; Zonen CNR4</td>
<td>10% for daily totals</td>
</tr>
<tr>
<td>(incoming and emitted)</td>
<td>(Pyrgeometer)</td>
<td></td>
</tr>
<tr>
<td>wind speed and direction</td>
<td>Young 05103-5</td>
<td>0.3 m s⁻¹; 3°</td>
</tr>
<tr>
<td>surface height</td>
<td>Campbell SR50A</td>
<td>10⁻² m or 0.4%</td>
</tr>
</tbody>
</table>
Table 2. Linear regression parameters for hourly values of KAN_U and S10 AWSs: Slope (χ), intercept (ψ), correlation coefficients (R) and root-mean-square deviations (RMSD).

<table>
<thead>
<tr>
<th>S10-KAN_U</th>
<th>χ</th>
<th>ψ</th>
<th>R</th>
<th>RMSD</th>
</tr>
</thead>
<tbody>
<tr>
<td>E_s↓*</td>
<td>1.010</td>
<td>-</td>
<td>0.99</td>
<td>37.25 (W m⁻²)</td>
</tr>
<tr>
<td>E_s↑*</td>
<td>0.987</td>
<td>-</td>
<td>0.99</td>
<td>24.71 (W m⁻²)</td>
</tr>
<tr>
<td>E_L↓</td>
<td>1.003</td>
<td>-6.06</td>
<td>0.99</td>
<td>8.92 (W m⁻²)</td>
</tr>
<tr>
<td>E_L↑</td>
<td>0.990</td>
<td>-0.25</td>
<td>1.00</td>
<td>3.62 (W m⁻²)</td>
</tr>
<tr>
<td>T_a</td>
<td>0.995</td>
<td>-0.25</td>
<td>1.00</td>
<td>0.50 (°C)</td>
</tr>
<tr>
<td>ambient air</td>
<td>0.990</td>
<td>7.77</td>
<td>1.00</td>
<td>0.45 (hPa)</td>
</tr>
<tr>
<td>relative humidity</td>
<td>0.899</td>
<td>10.31</td>
<td>0.91</td>
<td>3.78 (%)</td>
</tr>
<tr>
<td>wind speed*</td>
<td>0.928</td>
<td>-</td>
<td>0.99</td>
<td>0.66 (m s⁻¹)</td>
</tr>
<tr>
<td>α_{2010}**</td>
<td>-</td>
<td>-</td>
<td>0.93</td>
<td>0.032 (-)</td>
</tr>
<tr>
<td>α_{2011}**</td>
<td>-</td>
<td>-</td>
<td>0.94</td>
<td>0.028 (-)</td>
</tr>
<tr>
<td>α_{2012}**</td>
<td>-</td>
<td>-</td>
<td>0.91</td>
<td>0.066 (-)</td>
</tr>
</tbody>
</table>

* regression line forced through zero  
** 24-hour running averages for the months May until September
Table 3. Annual and summer (June-July-August) average meteorological parameters at KAN_U.

<table>
<thead>
<tr>
<th>KAN_U</th>
<th>2009*</th>
<th>2010</th>
<th>2011</th>
<th>2012</th>
<th>2013**</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td><strong>annual averages</strong></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$T_a$ (°C)</td>
<td>−15.5</td>
<td>−11.6</td>
<td>−18.0</td>
<td>−14.3</td>
<td>−15.4</td>
</tr>
<tr>
<td>ambient air pressure (hPa)</td>
<td>799</td>
<td>804</td>
<td>797</td>
<td>800</td>
<td>799</td>
</tr>
<tr>
<td>specific humidity (g kg$^{-1}$)</td>
<td>1.5</td>
<td>2.0</td>
<td>1.4</td>
<td>1.9</td>
<td>1.5</td>
</tr>
<tr>
<td>wind speed (m s$^{-1}$)</td>
<td>7.0</td>
<td>7.0</td>
<td>6.2</td>
<td>6.5</td>
<td>7.0</td>
</tr>
<tr>
<td>albedo</td>
<td>0.85</td>
<td>0.82</td>
<td>0.82</td>
<td>0.79</td>
<td>0.80</td>
</tr>
<tr>
<td><strong>summer (JJA) averages</strong></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$T_a$ (°C)</td>
<td>−4.3</td>
<td>−1.8</td>
<td>−2.9</td>
<td>−1.8</td>
<td>−4.5</td>
</tr>
<tr>
<td>ambient air pressure (hPa)</td>
<td>809</td>
<td>808</td>
<td>811</td>
<td>811</td>
<td>804</td>
</tr>
<tr>
<td>specific humidity (g kg$^{-1}$)</td>
<td>2.9</td>
<td>3.6</td>
<td>3.3</td>
<td>3.7</td>
<td>2.8</td>
</tr>
<tr>
<td>wind speed (m s$^{-1}$)</td>
<td>5.3</td>
<td>5.2</td>
<td>5.0</td>
<td>4.6</td>
<td>5.2</td>
</tr>
<tr>
<td>albedo</td>
<td>0.78</td>
<td>0.77</td>
<td>0.78</td>
<td>0.71</td>
<td>0.78</td>
</tr>
</tbody>
</table>

* average 2010–2013 January, February and March

** average 2009–2012 October, November and December
Table 4. Surface height changes and mass budgets (measured in winter and calculated in summer) at KAN_U in meters and m w.e., respectively, and ablation duration. The uncertainty associated with the surface height change is estimated at 0.2 m. The mass budgets are calculated with an assumed snow density of 360 kg m\(^{-3}\) (the average density of the uppermost 0.9 m measured on 26 April 2013), with uncertainty estimated at 40 kg m\(^{-3}\) (standard deviation among the snow-pit measurements). The snow density assumption was not needed in 2012 and 2013, when actual density measurements were conducted.

<table>
<thead>
<tr>
<th>Winter height change</th>
<th>Winter budget</th>
<th>Summer height change</th>
<th>Summer budget</th>
<th>Net budget</th>
<th>Ablation period</th>
</tr>
</thead>
<tbody>
<tr>
<td>2008–2009</td>
<td>+1.6*</td>
<td>+0.59±0.15</td>
<td>-0.7</td>
<td>-0.26±0.08</td>
<td>+0.34±0.12</td>
</tr>
<tr>
<td>2009–2010</td>
<td>+0.7</td>
<td>+0.25±0.08</td>
<td>-1.2</td>
<td>-0.44±0.09</td>
<td>-0.19±0.12</td>
</tr>
<tr>
<td>2010–2011</td>
<td>+1.0</td>
<td>+0.37±0.08</td>
<td>-1.1</td>
<td>-0.41±0.09</td>
<td>-0.04±0.12</td>
</tr>
<tr>
<td>2011–2012**</td>
<td>+0.7</td>
<td>+0.25±0.08</td>
<td>-1.8</td>
<td>-0.86±0.14</td>
<td>-0.61±0.16</td>
</tr>
<tr>
<td>2012–2013***</td>
<td>+1.2</td>
<td>+0.45±0.09</td>
<td>-0.8</td>
<td>-0.27±0.08</td>
<td>+0.18±0.12</td>
</tr>
</tbody>
</table>

* value inferred from Van de Wal et al. (2012)
** estimate based on snow-pit densities from May 2012
*** estimate based on snow-pit densities from May 2013

Surface mass budgets (measured in winter and calculated in summer) at KAN_U in m w.e., ablation duration and average ablation rates in mm w.e. day\(^{-1}\) assuming snow density of 360 kg m\(^{-3}\) (the average density of the uppermost 0.9 m measured on 26 April 2013). This assumption was not needed in 2012 and 2013 when actual density measurements were conducted. The uncertainties of the surface height change and snow density are estimated at 0.2 m and 40 kg m\(^{-3}\), respectively.
<table>
<thead>
<tr>
<th></th>
<th>winter-budget</th>
<th>summer budget</th>
<th>net-budget</th>
<th>ablation period</th>
<th>average ablation rate</th>
</tr>
</thead>
<tbody>
<tr>
<td>2008–2009</td>
<td>+0.59±0.15</td>
<td>-0.26±0.08</td>
<td>+0.34±0.12</td>
<td>01/06–19/08</td>
<td>3.25</td>
</tr>
<tr>
<td>2009–2010</td>
<td>+0.25±0.08</td>
<td>-0.44±0.09</td>
<td>-0.19±0.12</td>
<td>30/04–05/09</td>
<td>3.41</td>
</tr>
<tr>
<td>2010–2011</td>
<td>+0.37±0.08</td>
<td>-0.41±0.09</td>
<td>-0.04±0.12</td>
<td>28/05–13/08</td>
<td>5.26</td>
</tr>
<tr>
<td>2011–2012**</td>
<td>+0.25±0.08</td>
<td>-0.86±0.14</td>
<td>-0.61±0.16</td>
<td>26/05–24/08</td>
<td>9.05</td>
</tr>
<tr>
<td>2012–2013***</td>
<td>+0.45±0.09</td>
<td>-0.27±0.08</td>
<td>+0.18±0.12</td>
<td>29/05–17/08</td>
<td>3.33</td>
</tr>
</tbody>
</table>

* value inferred from Van de Wal et al. (2012)

** estimate based on snow-pit densities from May 2012

*** estimate based on snow-pit densities from May 2013
Table 5. Annual and summer (June-July-August) average energy fluxes at KAN_U (W m\(^{-2}\)).

<table>
<thead>
<tr>
<th></th>
<th>2009*</th>
<th>2010</th>
<th>2011</th>
<th>2012</th>
<th>2013**</th>
</tr>
</thead>
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<tr>
<td><strong>annual averages</strong></td>
<td></td>
<td></td>
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<td></td>
<td></td>
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<tr>
<td>(E_S^\downarrow)</td>
<td>155</td>
<td>153</td>
<td>150</td>
<td>145</td>
<td>151</td>
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<tr>
<td>(E_S^\uparrow)</td>
<td>-125</td>
<td>-121</td>
<td>-121</td>
<td>-110</td>
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<tr>
<td>(E_S^{\text{Net}})</td>
<td>30</td>
<td>32</td>
<td>29</td>
<td>35</td>
<td>32</td>
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<tr>
<td>(E_L^\downarrow)</td>
<td>207</td>
<td>224</td>
<td>205</td>
<td>223</td>
<td>212</td>
</tr>
<tr>
<td>(E_L^\uparrow)</td>
<td>-246</td>
<td>-262</td>
<td>-239</td>
<td>-254</td>
<td>-248</td>
</tr>
<tr>
<td>(E_L^{\text{Net}})</td>
<td>-39</td>
<td>-38</td>
<td>-34</td>
<td>-31</td>
<td>-36</td>
</tr>
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<td>(E_R)</td>
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<td>4</td>
<td>-4</td>
</tr>
<tr>
<td>(E_H)</td>
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<td>12</td>
<td>14</td>
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<tr>
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<td>-2</td>
<td>-1</td>
<td>-3</td>
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<tr>
<td>(E_G)</td>
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<td>-3</td>
<td>1</td>
<td>-2</td>
<td>-2</td>
</tr>
<tr>
<td>(E_P)</td>
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<td>0.006</td>
<td>0.009</td>
<td>0.012</td>
<td>0.005</td>
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<tr>
<td>(E_M)</td>
<td>4</td>
<td>8</td>
<td>6</td>
<td>13</td>
<td>5</td>
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<tr>
<td><strong>summer (JJA) averages</strong></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>(E_S^\downarrow)</td>
<td>322</td>
<td>305</td>
<td>302</td>
<td>296</td>
<td>313</td>
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<td>(E_S^\uparrow)</td>
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<td>-236</td>
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<td>-242</td>
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<td>(E_S^{\text{Net}})</td>
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<td>71</td>
<td>66</td>
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<td>71</td>
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<td>(E_L^\downarrow)</td>
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<td>252</td>
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<td>5</td>
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<td>(E_E)</td>
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<td></td>
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<tr>
<td>$E_G$</td>
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<td>$E_M$</td>
<td>15</td>
<td>28</td>
<td>24</td>
<td>49</td>
<td>17</td>
</tr>
</tbody>
</table>

* used average 2010–2013 values for January, February and March

** used average 2009–2012 values for October, November and December
Figure 1. (a) Map of Greenland and the location of KAN_U. (b) Picture taken after the installation of KAN_U (April 2009).
Figure 2. SEB model validation: (a) Observed and simulated relative surface height for the period of observations. (b) Observed against simulated $T_s$ ($R^2 = 0.98$; $(ΔT_s)_{avg} = 0.11^\circ$C; RMSE = 1.43 $^\circ$C.)
Figure 3. Average values of: (a) wind direction, (b) wind speed, (c) air pressure and (d) air temperature at KAN_U.
Figure 4. (a) 31-day running average values of all radiation budget components at KAN_U. Solid lines indicate the net solar and terrestrial radiation components. (b) Same, as (a), but for all surface energy balance components.
Figure 5. Seasonal cycles for the years 2009–2013 based on monthly averages of: (a) incoming shortwave energy flux, (b) surface albedo and (c) net shortwave energy flux.
Figure 6. Seasonal cycles for the years 2009–2013 based on monthly averages of: (a) incoming, (b) emitted and (c) net longwave energy flux.
Figure 7. Seasonal cycles for the years 2009–2013 based on monthly averages of: (a) sensible heat flux, (b) latent heat flux and (c) subsurface heat flux.
Figure 8. Seasonal cycle for the years 2009–2013 based on monthly averages of energy consumed by melt.
Figure 9. (a) Total energy per unit surface area. (b) Cumulative fluxes of all mass components. Note the different y-scales in (b).
Figure 10. (a) 2012 albedo anomaly measured by KAN_U for the months May–September, (b) simulated relative surface height anomaly and (c) simulated cumulative energy anomalies for all contributing fluxes.
Figure 11. (a) Difference between firm temperature measured by the KAN_U thermistor string and simulated firn temperature. The blue lines indicate the position of the thermistors below the surface. The white areas near the surface are due to surface thermistors. Note that the thermistor string was replaced by a new one drilled on 28 April 2013. (b) Simulated refreezing rates. (c) Simulated firn density.
Figure 12. (a) Monthly average NAO index. (ab) Monthly air temperature from Kangerlussuaq and at KAN U. Correlation coefficients (R): 0.97 for the extent of the KAN_U data, 0.66–0.99 for the months individually, minimum being January. (bc) Monthly reference period (1976–1999) air temperature at Kangerlussuaq. (cd) Monthly (May to September) and summer (June-July-August average) air temperature anomalies at Kangerlussuaq for the years 2000–2013. Error bars indicate two standard deviations.
Figure 13. 11-day Gaussian filtered nearest neighbor 5x5 km MOD10A1 albedo (2000–2013) and KAN_U (2009–2013) albedo for the months: (a) May (R = 0.91, \((\Delta\alpha)_{\text{avg}} = -0.02\ m\), RMSD = 0.02 \(m\)), (b) June (R = 0.77, \((\Delta\alpha)_{\text{avg}} = -0.05\ m\), RMSD = 0.05 \(m\)), (c) July (R = 0.95, \((\Delta\alpha)_{\text{avg}} = -0.05\ m\), RMSD = 0.05 \(m\)), (d) August (R = 0.60, \((\Delta\alpha)_{\text{avg}} = -0.03\ m\), RMSD = 0.04 \(m\)) and (e) September (R = -0.19, \((\Delta\alpha)_{\text{avg}} = 0.02\ m\), RMSD = 0.04 \(m\)).