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Atmospheric and oceanic forcing of Larsen C Ice Shelf thinning

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Abstract

The catastrophic collapses of Larsen A and B ice shelves on the eastern Antarctic Peninsula have caused their tributary glaciers to accelerate, contributing to sea-level rise and freshening the Antarctic Bottom Water formed nearby. The surface of Larsen C Ice Shelf (LCIS), the largest ice shelf on the peninsula, is lowering. This could be caused by unbalanced ocean melting (ice loss) or enhanced firn melting and compaction (englacial air loss). Using a novel method to analyse eight radar surveys, this study derives separate estimates of ice and air thickness changes during a 15 year period. The uncertainties are considerable, but the primary estimate is that the surveyed lowering ($0.066 \pm 0.017 \text{ myr}^{-1}$) is caused by both ice loss ($0.28 \pm 0.18 \text{ myr}^{-1}$) and firn air loss ($0.037 \pm 0.026 \text{ myr}^{-1}$). Though the ice loss is much larger, ice and air loss contribute approximately equally to the lowering. The ice loss could be explained by high basal melting and/or ice divergence, and the air loss by low surface accumulation or high surface melting and/or compaction. The primary estimate therefore requires that at least two forcings caused the surveyed lowering. Mechanisms are discussed by which LCIS stability could be compromised in future, suggesting destabilisation timescales of a few centuries. The most rapid pathways to collapse are offered by a flow perturbation arising from the ungrounding of LCIS from Bawden Ice Rise, or ice-front retreat past a “compressive arch” in strain rates.

1 Introduction

The ice shelves of the Antarctic Peninsula (AP) have shown a progressive decline in extent over the last five decades, including the catastrophic collapses of Larsen A Ice Shelf (LAIS) in 1995 and Larsen B Ice Shelf (LBIS) in 2002 (Scambos et al., 2003; Cook and Vaughan, 2010). The collapse of LBIS was unprecedented in at least the last 12 000 years (Domack et al., 2005). These collapses have reduced the restraint of the ice shelves on the flow of grounded tributary glaciers, causing them to accelerate

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(Rignot et al., 2004; Berthier et al., 2012) and thereby contributing to sea-level rise (Shepherd et al., 2012). Increased freshwater input to the ocean from the collapses and subsequent excess ice discharge is also thought to have freshened the Antarctic Bottom Water formed nearby (Hellmer et al., 2011; Jullion et al., 2013).

The final collapses have been attributed to meltwater-induced ice fracture following years of extreme atmospheric melting (Scambos et al., 2003; van den Broeke, 2005; Banwell et al., 2013), but longer-term processes such as ice thinning and firn compaction could have first weakened these ice shelves towards a state liable to collapse. Apparently following the southward progression of ice-shelf instability on the AP, satellite altimetry shows that the surface of Larsen C Ice Shelf (LCIS) has lowered in recent decades (Shepherd et al., 2003; Fricker and Padman, 2012; Pritchard et al., 2012). The lowering is more rapid in the north of LCIS (Fig. 1; updated from Fricker and Padman, 2012, as described in Sect. 2). Ice flow in this northern region has also accelerated slightly, which may be related to a decrease in back-stress from Bawden Ice Rise following an iceberg calving in 2004/05 (Haug et al., 2010; Khazendar et al., 2011). However, the origin of the lowering remains uncertain. Since the ice shelf is floating, the lowering could be caused by a loss of firn air of nearly the same magnitude, a loss of solid ice approximately 10 times larger, or a combination of the two. There is an urgent need to understand the cause of this lowering in order to project the possible future collapse of LCIS and the impacts of its many glacier catchments upon sea-level rise and ocean freshening.

The LCIS lowering was initially attributed to increased oceanic basal melting (i.e. ice loss) on the basis that firn compaction from derived surface melting trends was insufficient to account for the signal (Shepherd et al., 2003). However, sparse observations of the ocean beneath LCIS found the ocean to be at or below the sea-surface freezing temperature, suggesting that it is not capable of rapid melting (Nicholls et al., 2012). Observations of the meltwater emanating from the cavity (Nicholls et al., 2004) and widespread marine ice in LCIS (Holland et al., 2009; Jansen et al., 2013; McGrath et al., 2014) suggest that these temperatures are spatially and temporally prevalent. Ocean

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waters entering the LCIS cavity appear to be constrained to the freezing temperature by nearby sea-ice formation. Since the Weddell Sea has consistently high rates of sea-ice production it has been regarded as hard to conceive of an ocean warming sufficient to increase melting enough to explain the lowering (Nicholls et al., 2004). However, year-round sonar measurements at a single location in the south of LCIS yield a mean melt rate of $\sim 0.8 \text{ myr}^{-1}$ (with a range of $0\text{--}1.5 \text{ myr}^{-1}$), which is significantly higher and more variable than expected (K. W. Nicholls, personal communication, 2014; Nicholls et al., 2012). Furthermore, ocean data collected in January 1993 from the LCIS ice front (Bathmann et al., 1994) show anomalous waters that are considerably warmer than any subsequently observed in the cavity or inferred as sources for melting (Nicholls et al., 2004, 2012). If they entered the cavity, such warm waters could produce a melting anomaly large enough to significantly perturb the LCIS ice mass budget. Given our incomplete understanding of ocean processes and melting beneath LCIS, oceanic thinning of LCIS remains a credible explanation for the lowering.

On the other hand, there is some evidence supporting a hypothesis that the lowering results from an atmosphere-driven increase in firn compaction (i.e. air loss), either through dry compaction or through firn melting and refreezing. The AP has experienced strong atmospheric warming since the 1950s (Marshall et al., 2006; Turner et al., 2014). A spatial correspondence between ice-shelf collapses and mean atmospheric temperature suggests that atmospheric warming may have pushed some ice shelves beyond a thermal limit of viability (Morris and Vaughan, 2003), and the northern edge of LCIS is at this limit. Observations of LCIS firn-air thickness confirm that there is sufficient firn air available for compaction, that lower firn air spatially corresponds with high melting, and that the northward-intensified surface lowering spatially corresponds to areas of high melting and firn compaction (Holland et al., 2011; Trusel et al., 2013; Luckman et al., 2014). Modelled firn compaction entirely offset the lowering in one study of 2003–2008 (Pritchard et al., 2012). A temporal correspondence between high annual melting and ice shelf collapse (van den Broeke, 2005) would be expected to hold also for firn compaction before collapse. However, attributing the lowering to simple

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atmospheric temperature trends is not straightforward. Observed AP surface melt days and modelled meltwater fluxes both lack significant trends during 1979–2010 and have trends that are strongly negative during 1989–2010 (Kuipers Munneke et al., 2012a). An Automatic Weather Station on LCIS lacks any significant 1985–2011 trend in air temperature in any season (Valisuo et al., 2014), and there is no convincing evidence of trends in melting derived from reanalysis models during recent decades (Valisuo et al., 2014). Even without a trend in atmospheric forcing within recent decades, the period could still be anomalous relative to the long-term mean, and so an atmosphere-driven lowering remains viable.

In summary there is a wealth of circumstantial evidence related to the lowering, but no direct test of its origin. In this study we analyse repeated radio-echo sounding surveys of LCIS, applying a novel method to separate changes in ice thickness from changes in firn-air thickness (Holland et al., 2011). The method is presented in Sect. 2 and its results in Sect. 3. We then consider whether the uncertainties in these ice and air trends are sufficiently well-constrained to isolate the origin of the LCIS lowering (Sect. 4), and speculate upon the prognosis for the ice shelf's future stability (Sect. 5).

2 Method

Radar sounding provides the two-way travel time (TWTT) of a radar wave between the ice-shelf surface and base. This can be combined with accurate measurements of surface elevation to derive separate thicknesses of the solid ice and englacial firn air that comprise an ice shelf (Holland et al., 2011). With multiple surveys it is therefore possible to determine differences in ice and air thickness over time. There have been many radar surveys of LCIS, but we find that a very large number of observations are needed to sufficiently reduce the random error in the ice and air differences. Therefore, only repeated survey lines provide usable data; inter-survey cross-overs are not sufficient. Fortunately, a nearly meridional (across-ice flow) survey line sampling the centre of LCIS has been occupied eight times between 1998 and 2012 by airborne

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and ground-based radar surveys (Fig. 1b, Table 1), offering the opportunity to derive interannual trends in ice and air thickness from these data. The survey line also passes through five satellite cross-overs of European Space Agency radar altimeter missions, allowing direct comparison to the known lowering.

2.1 Theory

We separate the total ice-shelf thickness into its constituent thicknesses of solid ice and firn air by following the method of Holland et al. (2011), with a few modifications. The presence of liquid meltwater is neglected on the basis that most surveys were undertaken early in the austral spring (see Sect. 4). If the ice is freely floating then the hydrostatic ice and ocean forces must balance at the ice base, so the total mass of the shelf ice and firn air equals that of the atmosphere and ocean displaced

$$\rho_i l + \rho_a A = \rho_A S + \rho_o (l + A - S). \quad (1)$$

Here l is the total solid ice thickness, A is the total firn air thickness, S is the ice freeboard (surface elevation above sea level), and $\rho_i = 918 \text{ kg m}^{-3}$, $\rho_a = 2 \text{ kg m}^{-3}$, $\rho_A = 1.3 \text{ kg m}^{-3}$, and $\rho_o = 1028 \text{ kg m}^{-3}$, are densities of solid ice, englacial air (partly pressurised), atmospheric air, and ocean respectively, which are all assumed constant. Adopting a similar approach and separating the radar delay of ice from that of air, the TWTT of a radar wave through the ice shelf is

$$T = \frac{2}{c} (n_i l + n_a A), \quad (2)$$

where T is the TWTT, $c = 3 \times 10^8 \text{ m s}^{-1}$ is the speed of light in a vacuum and $n_i = 1.78$ and $n_a = 1.0$ are refractive indices of pure ice and air. Combining Eqs. (1) and (2) and eliminating variables as appropriate, we obtain expressions for the constituent ice and air thicknesses (and hence total thickness, $l + A$) as functions of known quantities and

the measured TWTT and surface elevation:

$$A = \left[\frac{c(\rho_o - \rho_i)}{2n_i} T + (\rho_A - \rho_o) S \right] / \left[(\rho_a - \rho_o) + \frac{n_a(\rho_o - \rho_i)}{n_i} \right] \quad (3)$$

$$I = \left[\frac{c(\rho_o - \rho_a)}{2n_a} T + (\rho_A - \rho_o) S \right] / \left[(\rho_i - \rho_o) + \frac{n_i(\rho_o - \rho_a)}{n_a} \right]. \quad (4)$$

Taking the temporal derivative of these expressions, we obtain the trends in ice and air thickness as a function of the trends in elevation and TWTT:

$$\frac{\partial A}{\partial t} = \left[\frac{c(\rho_o - \rho_i)}{2n_i} \frac{\partial T}{\partial t} + (\rho_A - \rho_o) \frac{\partial S}{\partial t} \right] / \left[(\rho_a - \rho_o) + \frac{n_a(\rho_o - \rho_i)}{n_i} \right] \quad (5)$$

$$\frac{\partial I}{\partial t} = \left[\frac{c(\rho_o - \rho_a)}{2n_a} \frac{\partial T}{\partial t} + (\rho_A - \rho_o) \frac{\partial S}{\partial t} \right] / \left[(\rho_i - \rho_o) + \frac{n_i(\rho_o - \rho_a)}{n_a} \right]. \quad (6)$$

Evaluating the known quantities in these terms, we find that

$$\frac{\partial A}{\partial t} = 1.06 \frac{\partial S}{\partial t} - 0.114 \frac{c}{2n_i} \frac{\partial T}{\partial t} \quad (7)$$

$$\frac{\partial I}{\partial t} = -0.598 \frac{\partial S}{\partial t} + 1.06 \frac{c}{2n_i} \frac{\partial T}{\partial t}, \quad (8)$$

where the TWTT is expressed as a solid ice equivalent for clarity.

Note that the derivation of Eqs. (5) and (6) from Eqs. (3) and (4) neglects temporal derivatives of all densities, of which the most variable is the ocean density. Repeating the derivation and retaining ocean density terms provides an expression in which 0.3 myr^{-1} ice loss would require a $\sim 2 \text{ kg m}^{-3} \text{ yr}^{-1}$ reduction in ocean density, and 0.03 myr^{-1} air loss would require a $\sim 0.1 \text{ kg m}^{-3} \text{ yr}^{-1}$ increase in ocean density. Such changes persisting over 15 years are clearly implausible, and we conclude that ocean density changes have negligible effect.

2.2 Application to Larsen C Ice Shelf

We apply the method to eight radar surveys between February 1998 and December 2012 along a line traversing the centre of LCIS (Fig. 1b, red line). The surveys were carried out by ground-based field parties and a variety of aircraft flying at different heights and speeds, and many different radar instruments and methods for measuring elevation were used (Table 1). The processed elevation and TWTT data are shown in Fig. 2a and b. The most densely-spaced TWTT data were gathered during the 2004 NASA-CECS airborne survey, so this is chosen as a baseline dataset. For each elevation and TWTT measurement in the other surveys, we find the difference from the nearest corresponding measurement in the 2004 survey, discarding all observations that do not have a 2004 analogue within 1000 m. These elevation and TWTT differences are shown in Fig. 2c and d. There is a great deal of scatter in the differences, which could result from several factors, including the advection of ice topography across the survey line at $\sim 400 \text{ m yr}^{-1}$ (Rignot et al., 2011). The differences are therefore binned spatially to extract the overall signals by averaging random noise, and linear trends in surface elevation and TWTT are calculated for the bins. Equations (5) and (6) are then used to determine the trends in ice and firn-air thickness from trends in surface elevation and TWTT. We apply this methodology in two ways, first considering the overall trends for the entire survey line, and then dividing the survey into five bins, surrounding each of the five satellite crossover points (Fig. 1).

The 2012 British Antarctic Survey (BAS) ground-based survey was a mission of opportunity within a wider seismic season (Brisbourne et al., 2014) and deviated from the rest of the surveys, heading due south (Fig. 1b, yellow line). However, it did repeat a flight line from the 1998 BAS airborne survey, so to include the data we first calculate the mean difference between the 2012 and 1998 surveys along the meridional line, and then the mean difference between the 2004 and 1998 surveys along the primary line, and then use these to obtain the 2012 – 2004 difference. The results are only in-

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cluded in the northernmost bin when we consider along-survey variability; they are not included in the whole-survey results.

2.3 Radio-echo sounding survey data

Different techniques are available for picking radar return echoes from echograms, and so to ensure that our inter-survey trends are as robust as possible the ice surface and base echoes from all surveys were re-picked in a consistent manner. Automatic first-break picks on time-windowed and scaled traces were manually edited to remove or correct mis-picks. For airborne surveys, TWTT was calculated as the difference between ice surface and basal returns, thus minimising inter-survey biases by removing any error associated with the absolute accuracy of the radar. Basal return TWTTs from the ground-based survey data were corrected for the radar antenna separation. In the NASA IceBridge 2009 and 2010 and BAS 2011 airborne surveys, the altitude of the aircraft in specific sections caused the surface multiple return to appear at a TWTT similar to that of the basal return, significantly contaminating the picks. Therefore, the radar-grams were overlain with an estimate of the surface multiple return calculated from the aircraft altitude and also an estimate of the basal return derived from the aircraft altitude, surface elevation and hydrostatic assumption. Wherever the TWTT of these two signals was indistinguishable in the radargram, no basal return pick was recorded. Significant marine ice bands were omitted from all surveys, because basal returns become indistinct and the meteoric-marine transition may be visible instead.

TWTTs from the 2009 IceBridge survey were found to contain consistently shorter radar-wave delays than the 2009 McGrath ground-based survey despite being collected only two weeks earlier, with a mean ice equivalent thickness approximately 10 m thinner and therefore a significant outlier relative to the other surveys. The data were investigated and re-picked, but the problem seems to result from transmit/receive switches not meeting their switching-time specification in the survey (<https://data.cresis.ku.edu/#RDS>), so the 2009 IceBridge TWTT data are neglected.

2.4 Surface elevation survey data

Surveyed ice elevation data have several corrections applied to make them directly comparable. Corrections for the steady geoid and mean dynamic ocean topography are not required because the method employs only temporal differences in elevation, as shown by Eqs. (5) and (6). All data are de-tided using the CATS2008a_opt model (L. Padman, personal communication, 2014) and have a local sea-level rise of 4 mm yr^{-1} removed (Rye et al., 2014).

Most of the instruments used to derive elevation were well-calibrated in the field (e.g. <http://nsidc.org/data/docs/daac/icebridge/ilatm2/index.html>), but the two BAS airborne surveys in 1998 and 2011 were not calibrated to the centimetre-scale accuracy required here. The 1998 survey passed over the open ocean in many locations, so these elevations were corrected for tides, EIGEN-6C geoid (<http://icgem.gfz-potsdam.de/ICGEM/ICGEM.html>), and DTU12 mean dynamic topography (http://www.space.dtu.dk/english/Research/Scientific_data_and_models/downloaddata), and then the mean difference from zero (sea surface) of 1.01 m was removed from the entire dataset. Repeating this procedure for the 2011 survey produced a 1.33 m offset, but from only a small area of open-ocean data. Fortunately, it was possible to correct the 2011 elevations to match the well-calibrated 2010 NASA IceBridge survey that took place 10 weeks earlier. However, this was complicated by a progressive southward decrease in the difference between surveys (Fig. 3). Radar altimetry penetrates the surface and reflects from within the firn layer, whereas laser altimetry reflects from the surface, so we ascribe this southward decrease to firn penetration by the radar altimeter in the 2011 survey, consistent with the known southward increase in firn air content (Holland et al., 2011). North of 67.85° S there is no broad-scale spatial variation in the offset between datasets, implying either uniform or no radar penetration. We subtract the mean offset in this area, 1.59 m, from the 2011 data and then treat the variable radar penetration to the south separately.

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The elevation estimates derived from the two BAS radar altimeter surveys need a firn-penetration correction to make them comparable to those derived from the laser altimeter and GPS. This correction consists of a linear fit south of 67.85° S to the difference between the IceBridge 2010 and BAS 2011 surveys (Fig. 3). Thus, out of necessity, the correction includes implicit assumptions that there is no firn penetration north of this during either radar survey, and that penetration to the south did not change between February 1998 and January 2011.

2.5 Satellite radar altimeter elevation data

Satellite radar altimeter data are used to corroborate the surveyed elevation data and provide a context for the lowering. The satellite elevation timeseries combine radar altimeter data from the ERS-1, ERS-2, and Envisat satellites using an existing methodology (Fricker and Padman, 2012) but including new data to the end of 2011. These data consist of repeat measurements of ice-shelf surface elevation at satellite orbit crossing points, available approximately every 35 days during Austral winters (April–November) during 1992–2011.

When analysing the data we found a strong correlation between changes in elevation and changes in surface backscatter for the period 1992–1993 (the first two years of ERS-1). This anomalous behaviour in the altimeter backscatter, which alters the shape of the waveform from which the elevation is deduced, occurs throughout Antarctica. This leads us to believe that these data may not be reliable, so we only use data from 1994 onwards in this study. Shepherd et al. (2010) also neglected data prior to 1994 in their analysis. This is important because other studies of LCIS that include these early data have very rapid lowering in the 1990s (Shepherd et al., 2003; Fricker and Padman, 2012) that is not supported by the remaining data. To illustrate the lowering of LCIS we first consider the period 1994–2011 (Fig. 1a), though our main analysis focuses upon the 1998–2011 period covered by the radar surveys (Fig. 1b). During the latter period the LCIS lowering has the same general pattern, but the trends at the five crossovers

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covered by the survey line are slightly different. Importantly, the survey line does not sample the northern section of LCIS in which the fastest lowering occurs.

To compare elevation trends derived from the survey data to those derived from satellite radar altimeter data (Fig. 4 and Table 2) a single satellite elevation trend was derived that represents all five independent satellite crossovers. First, the mean elevation for the Austral winter of 1998 was calculated for each independent crossover and subtracted from each crossover's time series. The resulting temporal elevation anomaly data were then treated as individual data points in a single merged time series, and from that a linear trend was calculated to compare to the surveyed trends. Linear trends were also calculated at each crossover, as presented in Figs. 1 and 5–7.

2.6 Ice and air mass balances

We consider the derived ice and air losses in the context of the ice and air mass balances of LCIS. The mass balance of the ice fraction of the ice shelf (i.e. excluding firn air) yields an equation governing the depth-integrated ice thickness

$$\frac{\partial I}{\partial t} + I \nabla \cdot \mathbf{u} + \mathbf{u} \cdot \nabla I = a_i - m_b \quad (9)$$

where \mathbf{u} is the two-dimensional horizontal ice velocity vector, a_i is net surface ice accumulation, and m_b is basal melting. The mass balance of the air fraction of the ice shelf yields a similar equation for depth-integrated air thickness

$$\frac{\partial A}{\partial t} + A \nabla \cdot \mathbf{u} + \mathbf{u} \cdot \nabla A = a_A - m_s - d \quad (10)$$

where a_A is the air trapped in the firn by accumulation, m_s is the loss of air by surface melting, percolation, and refreezing, and d is the loss of air by dry compaction. The terms on the left-hand side of both equations are the unsteady term, divergence, and advection.

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When analysing the results we map the terms in the ice mass balance Eq. (9) following a previous study (McGrath et al., 2014) that combined data from several sources. Divergence, advection, and mass input terms can be mapped from satellite-based observations of ice velocity (Rignot et al., 2011) and ice-shelf elevation (Griggs and Bamber, 2009), firn-air thickness derived from airborne radar measurements (Holland et al., 2011), and model estimates of net surface accumulation (Lenaerts et al., 2012). Though we also possess a spatial map of ice surface elevation change, an unknown fraction of this is caused by firn-air changes and so we cannot derive ice thickness change outside the temporal and spatial range of our survey data. Neglecting the unsteady term, we can derive a map of steady-state melting from the other terms. Prior to these calculations the ice thickness and velocity fields are smoothed over a 20 km footprint (masked outside the ice shelf) to remove small-scale noise that is amplified in the spatial derivatives. The firn-air mass balance Eq. (10) contains so many unknown quantities that we do not attempt to derive its terms.

3 Results

We first present the main results of the study, before a full analysis of the uncertainties in Sect. 4.

3.1 Trends over the whole survey line

Figure 4 shows the elevation and TWTT for each survey, as mean differences from 2004 over the entire survey line, and Table 2 gives the “primary” derived trends for these “reference” data and also a variety of alternatives. The surveyed elevation differences show a lowering trend ($-0.066 \pm 0.017 \text{ m yr}^{-1}$) that is very similar to that obtained from the satellite altimeter data (-0.062 m yr^{-1}); the trends are not expected to be identical due to method uncertainties and spatial and temporal differences in sampling. Crucially, there is also a decreasing trend in surveyed TWTT ($-0.296 \pm 0.17 \text{ m yr}^{-1}$)

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ice equivalent), though there is considerably more inter-survey scatter in this quantity and uncertainty in the resulting trend (see Sect. 4). Combining these observed trends using Eqs. (5) and (6) reveals that the lowering is caused partly by air loss ($-0.0367 \pm 0.026 \text{ m yr}^{-1}$) and partly by ice loss ($-0.274 \pm 0.18 \text{ m yr}^{-1}$). The ice loss has a much greater magnitude, but the surface lowering is approximately ten times more sensitive to air loss, so that ice loss and air loss contribute approximately equally to the surface lowering. There is considerable scatter in the data and several sources of uncertainty in the methodology, but our conclusion that ice and air loss both contribute to the lowering is robust when several different combinations of data are used in the calculations (see Sect. 4).

3.2 Variation within survey line

We now consider spatial variability by binning the survey data around each satellite crossover (Fig. 5a). The derived ice loss is reasonably uniform along the line, while the derived air loss is noticeably higher towards the southern end of the survey line. However, the surveyed elevation trends at the southern end of the line show considerably more lowering than the satellite elevation trends. Inspection of the data underlying the timeseries in each bin (Fig. 6) reveals that the surveyed elevations are reasonable apart from the 1998 data in the southernmost bin (centred on 68.3° S), which exceed the range of the figure. We consider the satellite altimeter data to be a more reliable measure of lowering because the 1998 surveyed elevation data are subject to calibration and firn-penetration corrections that are uncertain in this area. Replacing the surveyed elevation trends with the satellite elevation trends (Fig. 7a) has virtually no effect on the derived ice loss, but removes the air loss completely from the southernmost bin, so that the air loss is concentrated on the centre of the survey line.

The air and ice losses shown in Figs. 5a and 7a are scaled so that their resultant surface lowering can be read on the left-hand axis. From Fig. 7a we conclude that air loss contributes the majority of the lowering in the centre of the survey line, while ice loss also contributes to this lowering and is responsible for the lowering at both ends.

It is unsurprising that the ice and air loss have different spatial patterns, given their different (oceanic, ice-dynamic, and atmospheric) forcings.

3.3 Ice and air budgets

Figure 8 shows the maps of each term in the LCIS ice mass balance Eq. (9). Thinning along flowlines causes a sink of ice through divergence (Fig. 8a), advection is generally a source of ice where the ice shelf flows from thick to thin (Fig. 8b), and modelled surface accumulation is almost uniform (Fig. 8c). Their sum, the steady melting map (Fig. 8d), contains obvious artefacts but also many features that match our existing knowledge of ocean melting beneath LCIS. For example, the results are in agreement with a simple ocean-layer model (Holland et al., 2009) that predicts strong melting along the grounding line and freezing in the thinner ice immediately offshore of islands and peninsulas on the western coast (also visible as negative values in the advection term). A more sophisticated three-dimensional ocean model (Mueller et al., 2012), forced only by tides, predicts large values of tidally-driven melting next to Bawden Ice Rise and Kenyon Peninsula, which also seem apparent in Fig. 8d, though other areas of high melting near the ice front and south of Kenyon Peninsula are not consistent with the model.

Combining the estimated mean terms in the ice mass budget (Fig. 8) with the ice loss derived along the survey line (Figs. 5a and 7a) allows us to consider the full unsteady ice budget (Fig. 5b and 7b). The basic ice balance is between accumulation and divergence, with advection becoming important at the southern end of the line. If the ice shelf were in steady state the derived oceanic melt rate would be an order of magnitude smaller than accumulation and divergence (0.06 m yr^{-1}). In fact, our derived ice loss profiles suggest a mean oceanic melt rate over the survey line of 0.26 m yr^{-1} , peaking at 0.5 m yr^{-1} in the southernmost bin. These estimates are consistent with modelled patterns of melting (Holland et al., 2009; Mueller et al., 2012) and observations in a higher-melting region nearby (K. W. Nicholls, personal communication, 2014; Nicholls et al., 2012). Crucially, without basal melting the components of the

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mass budget are approximately balanced, so the majority of the melting is causing net ice loss. This emphasizes that for ice shelves melted by cold ocean waters, relatively small absolute changes in melting can have a significant influence on the ice shelf mass balance. In comparison, warm-water ice shelves such as Pine Island Glacier can have much larger melting perturbations (e.g. 5 myr^{-1} ; Wingham et al., 2009), causing equally large thinning rates, but these perturbations are a much smaller fraction of the mean melt rate (e.g. 100 myr^{-1} ; Dutrieux et al., 2013).

The terms in the analogous firn-air budget are extremely uncertain. To put the derived air loss of 0.04 myr^{-1} into context, we simply note that there was 10–15 m of air in the surveyed section during the 1997/98 survey (Holland et al., 2011), and if fresh snow is deposited at a density of $350\text{--}450 \text{ kg m}^{-3}$ (Kuipers Munneke et al., 2012b) then the accumulation of 0.5 myr^{-1} ice implies the addition of $0.5\text{--}1 \text{ myr}^{-1}$ firn air each year before compaction is taken into account. Therefore, our best estimate is that the net air loss is only 5–10 % of the annual air input.

4 Error estimation

The data contain a considerable amount of scatter and their interpretation relies upon a clear understanding of the uncertainties inherent in the derived trends. For this reason, we present a thorough error analysis before proceeding to discuss the implications of our findings. This analysis starts with a simple technique for visually assessing the reliability of the results, before proceeding to more formal methods.

4.1 Visual Assessment

It is possible to visually assess the reliability of ice and air trends from appropriately-plotted trends in elevation and TWTT. If the TWTT trend is expressed as a solid-ice

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surface-elevation equivalent, i.e.

$$\frac{\partial T_s}{\partial t} = \frac{c}{2n_i} \frac{\rho_o - \rho_i}{\rho_o - \rho_A} \frac{\partial T}{\partial t}, \quad (11)$$

then comparing $\partial T_s / \partial t$ to the elevation trend $\partial S / \partial t$ allows us to determine the value of $\partial A / \partial t$ from Eq. (5). Any elevation trend that is more negative than $\partial T_s / \partial t$ implies a loss of air, with the air loss equal to 1.06 times the difference between $\partial S / \partial t$ and $\partial T_s / \partial t$. For this purpose, the two y axes of Figs. 4, 5a, 6, and 7a are scaled such that the left-hand axis shows both ice surface elevation ($\partial S / \partial t$) and TWTT expressed as solid-ice surface equivalent ($\partial T_s / \partial t$). Consideration of the numerator of Eq. (6) shows that $\partial T_s / \partial t$ merely has to be more negative than $-0.107 \times \partial S / \partial t$ to imply a loss of ice; any $\partial T_s / \partial t$ that is negative enough to be distinguished in the figures implies some ice loss. In plain terms, Fig. 4 is scaled such that if the red line (scaled TWTT trend) is parallel to the green line (elevation trend) then the lowering is due solely to ice loss, and if the red line is flat then all of the lowering is due to air loss.

These criteria allow a simple visual assessment of the signal present in the available data. Our assessment of Fig. 4 is that the scaled TWTT is decreasing, but that this result is not robust in the sense that it is dependent upon all datasets and removing certain surveys would remove the decrease. This reduces confidence in the conclusion that ice loss has occurred. On the other hand, we do not believe that the scaled TWTT data could support a trend that is more negative than the elevation trend, and therefore we are confident in our conclusion that air loss has occurred.

A formal analysis revisits these conclusions below, but this requires many assumptions about the nature of the errors and so is not necessarily superior. There are many sources of error in our surveys, which we divide into two classes. The first class of errors produces random intra-survey scatter, which affects the extent to which the data from each survey estimate the mean signal within that survey. The second class of errors create a systematic signal across a whole survey, directly affecting inter-survey

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differences. The latter are of greatest concern because they have the largest effect on trends.

4.2 Intra-survey errors

Predominantly intra-survey errors include:

- instrument and processing error (including radar picking error, assumed intra-survey because all surveys were re-picked consistently);
- spatial offset from the reference line (there is no systematic spatial difference between surveys, but the data deviate from a straight line within surveys);
- advection of complex ice topography through the survey line (assumed intra-survey because ice features are smaller than both the along-survey distance and the advection lengthscale in the across-survey direction: $15 \text{ years} \times 400 \text{ myr}^{-1}$).

We can easily quantify these random errors by considering, for each survey, the statistics of each population of differences of data points from their 2004 analogues (Table 3). Standard deviations (SD) are relatively large, 1–2 m for elevation and ~ 10 m ice equivalent for TWTT, as expected from previous analyses of the error in individual point measurements (Holland et al., 2009). However, when all data are considered the standard errors are small due to the large sample sizes. Assuming that the differences are independent and normally distributed, 95 % confidence interval bounds for the survey mean are given by multiplying the standard error by 1.96, as shown by the error bars in Fig. 4. We estimate overall 95 % confidence interval bounds as ± 0.04 m for elevation and ± 0.5 m ice equivalent for TWTT. Thus, from a random error perspective, we are confident that all surveys differ significantly from 2004 apart from the elevation differences in the 2011 and McGrath 2009 surveys and both elevation and TWTT datasets in 2012. Simple examination of the error bars in Fig. 4 shows that variation within these random error bounds will have negligible effect on the computed trends.

4.3 Inter-survey errors

Predominantly inter-survey errors include:

- differences between survey instruments, calibration, and processing (radar altimeter penetration, ice-penetrating radar power and frequency, speed and altitude of acquisition platform);
- time-variable presence of liquid meltwater in the firn column;
- time-variable firn penetration in the ice-penetrating radar surface pick;
- the time-variable part of dynamic ocean topography (inter-survey because most surveys are rapid compared to the relevant variations in ocean flow; affects elevation only);
- error in the tidal model correction (inter-survey because most surveys are rapid compared to tides; affects elevation only);
- the inverse barometer effect (inter-survey because most surveys are rapid compared to the relevant variations in atmospheric pressure; affects elevation only).

An initial concern is that the NASA IceBridge and NASA-CECS surveys (high-altitude, high-speed, consistent radar systems, laser altimeter) differ from the BAS airborne surveys (lower-altitude, slower, different radar, radar altimeter) and both differ from the ground-based surveys (low-frequency radar, GPS elevation). However, the three types of survey are interleaved in time, so such differences do not necessarily cause systematic trends. The issue is assessed by re-calculating the trends using different combinations of data (Table 2). Considering only the two BAS surveys produces broadly the same results. However, considering only NASA IceBridge and NASA-CECS surveys produces a much weaker surface lowering and no decrease in TWTT, so that the ice loss disappears. Systematically removing the surveys from the calculation reveals that it is neglecting the BAS 1998 survey that removes these trends (Table 2). We know of

no reason to neglect this survey, but this suggests that we treat TWTT and ice trends with additional caution.

The presence of meltwater in the firn would require us to adapt the methodology because it affects both the hydrostatic floatation and radar-wave delay of the ice shelf, as described by Holland et al. (2011), leading to different ice and air thicknesses being derived from the same TWTT and elevation. This potentially confounding issue is neglected because most surveys were sampled in November, before the onset of melt (Barrand et al., 2013). However, the two BAS surveys were in summer and could be contaminated by the presence of meltwater. Repeating the derivation of Eqs. (3) and (4) but including the effects of meltwater produces new equations from which 0.57 m more air and 5.6 m less ice would be derived for every 1 m of meltwater present (Holland et al., 2011). A maximum LCIS meltwater content of 0.4 m (Holland et al., 2011) therefore implies a maximum underestimate of 0.23 m air and overestimate of 2.24 m ice. The summer of 1997/98 was a high melting year (Tedesco, 2009), and if meltwater was present during the 1998 survey the derived air content should be higher and ice content lower, enhancing the air loss trend and reducing the ice loss trend. A linear regression to 0.23 m air error and -2.24 m ice error in 1998 and no meltwater-derived error in the other surveys yields maximum trend errors of $-0.0137 \text{ m yr}^{-1}$ air and $+0.134 \text{ m yr}^{-1}$ ice. Melt estimates for 2010/11 are not available, but any 2011 meltwater would have the opposite effect on the inter-survey trends to 1998 meltwater, and thus mitigate this issue.

For the airborne surveys, surface penetration could affect both radar altimeters and the surface pick of ice-penetrating radars. We have used a penetration correction in radar altimeter data (see above), and their agreement with the satellite elevation trend implies that deviation from this correction is not important. Our strategy of finding the ice TWTT by picking the surface and basal returns and differencing the result means that surface penetration could affect the TWTT. We examine this by comparing the radar surface picks with altimeter data. This test is imperfect because it introduces errors from the aircraft altitude and surface elevation data, and requires absolute accuracy in

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the radar data that is not needed of the TWTT differences used. The test cannot even be performed for the NASA IceBridge and NASA-CECS surveys because the absolute timing of the radar pulse transmission is not known to the required accuracy. The mean difference between altimeter-derived surface elevations and radar-derived surface elevations is 2.14 m for the BAS 1998 survey and 2.38 m for the BAS 2011 survey. The altimeter-derived elevation is higher than the radar-derived elevation in both cases, so the difference may be caused by surface penetration. This very limited dataset suggests that radar firn penetration is of order 2 m, with an interannual variability of order 0.2 m.

These differences between radar surface picks and altimeter data are also the only independent information we have to quantify overall inter-survey error in TWTT differences. They are again imperfect in this role because they include error in aircraft altitude and surface elevation data that does not appear in the TWTT differences used in Eqs. (5) and (6). Also, if the error in basal and surface picks is identical (e.g. from an absolute calibration error) then the error in their difference is zero. On the other hand, if the surface and basal errors are uncorrelated and of the same magnitude then the TWTT difference error is the surface pick error multiplied by $\sqrt{2}$. We believe that an inter-survey error of 2 m ice equivalent for TWTT is a reasonable compromise, and this value is in good agreement with the deviation of the TWTT points from the trend line in Fig. 4.

The effects of unsteady dynamic ocean topography, error in the tidal correction, and inverse barometer effect should each contribute an inter-survey error of order 0.1 m to the surface elevation differences (L. Padman, personal communication, 2014; Padman et al., 2003; King and Padman, 2005). If these errors are uncorrelated, this would create a total error of about 0.2 m, and this estimate is consistent with both the deviation of the surveys from the linear trend and the difference in elevation between the two 2009 surveys (Fig. 4). In any case, the surface lowering from the satellite crossovers provides an independent test of the surveyed elevation trend, and the two trends are only slightly

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different (Table 2), as might be expected from the difference in spatial and temporal sampling.

Given these overall inter-survey error estimates (0.2 m elevation and 2 m TWTT ice equivalent), we used a Monte Carlo approach to estimate the resultant uncertainty in the elevation and TWTT trends. The trends were recalculated 500 000 times with all data points subject to a perturbation drawn from a normal distribution with 95 % confidence interval bounds equal to the error estimates. This yields a population of trends with 95 % confidence interval bounds of $\pm 0.017 \text{ m yr}^{-1}$ for elevation trends and $\pm 0.17 \text{ m yr}^{-1}$ ice equivalent for TWTT trends. Evaluating the terms as in Eqs. (7) and (8) and combining the errors in quadrature yields

$$\varepsilon_{At} = \sqrt{0.013\varepsilon_{Tt}^2 + 1.13\varepsilon_{St}^2} \quad (12)$$

$$\varepsilon_{It} = \sqrt{1.13\varepsilon_{Tt}^2 + 0.36\varepsilon_{St}^2}. \quad (13)$$

Where ε_{At} , ε_{It} and ε_{St} are errors in $\partial A/\partial t$, $\partial I/\partial t$, and $\partial S/\partial t$ respectively. The symbol ε_{Tt} represents the error in $c/2\eta_i\partial T/\partial t$, TWTT converted to solid ice thickness. These formulae yield uncertainties of $\pm 0.026 \text{ m yr}^{-1}$ for $\partial A/\partial t$ and $\pm 0.18 \text{ m yr}^{-1}$ for $\partial I/\partial t$.

4.4 Error summary

In summary, formal error estimates suggest that both the ice and air loss derived in our reference calculation are robust. However, visual assessment of Fig. 4 suggests that the data support air loss more strongly than ice loss. Recalculating the trends with different combinations of the data (Table 2) shows that almost all possible calculations have significant air loss; the only way to obtain insignificant air loss is to include 2009 IceBridge TWTT data known to be erroneous. On the other hand, removing either the 1998 BAS or 2009 McGrath surveys is sufficient to render the ice loss insignificant. Any meltwater that were present during the 1998 BAS survey would further strengthen the

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air loss and weaken the ice loss. Our best estimate is that the lowering is a result of both air loss and ice loss, but there remains a possibility that air loss is solely responsible.

The preceding calculations apply to the whole-survey comparisons shown in Fig. 4. The latitude bins shown in Figs. 5–7 contain fewer data, so the intra-survey standard error should increase. Standard errors scale with the reciprocal square root of the number of datapoints, so the 95 % confidence interval bounds approximately double (± 0.08 m for elevation and ± 1 m ice equivalent for TWTT) when the data sample size are reduced by a factor of 5. Inter-survey systematic error should in principle remain similar, but on the shorter length scale of an individual bin, several intra-survey errors become inter-survey in character (differences in radar picking, survey path, and advection of ice features, which can be a significant fraction of a bin length in the along-survey direction). Scrutinising the time series in Fig. 6 suggests a reasonable confidence in the binned trends. In most cases a downward trend of the TWTT is apparent, suggesting some ice loss has occurred, and the scaled TWTT data would not support a downwards trend steeper than the satellite elevation, suggesting air loss has occurred. The steepest elevation trends and shallowest TWTT trends are in the centre of the survey line, implying greatest air loss.

5 Discussion

The uncertainties are considerable, but our primary estimate is that the lowering ($0.066 \pm 0.017 \text{ m yr}^{-1}$, or 0.99 ± 0.26 m) is caused by both ice loss ($0.28 \pm 0.18 \text{ m yr}^{-1}$, or 4.2 ± 2.7 m) and firn air loss ($0.037 \pm 0.026 \text{ m yr}^{-1}$, or 0.56 ± 0.39 m). It is notable that though their effect on the lowering is approximately equal, ice loss is an order of magnitude larger than air loss. The derivation of these values allows us to speculate upon the possible sources of the changes, and their future implications.

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5.1 Sources of change

The existence of mean rates of change in ice and air over our 15 year period imply an imbalance in the other terms of Eqs. (9) and (10) during this time. We consider the ability of each of these terms to cause the imbalance and therefore the ice and air losses.

Whether the budget was ever balanced in the past, with the observed imbalance then implying that changes have occurred, is a separate question that we cannot answer.

We start with sources and sinks. Above-balance basal melting will cause ice loss but not air loss, and can easily account for our ice loss signal. Any melting greater than a few centimetres per year can cause an imbalance (Fig. 7), and observations and models easily support the rates of $\sim 0.26 \text{ myr}^{-1}$ needed to explain the ice loss (Holland et al., 2009; Mueller et al., 2012; Nicholls et al., 2012). Above-balance surface melting and refreezing or dry compaction (through atmospheric warming) will cause only air loss, and it is again easy for these processes to account for the air loss signal observed here. Below-balance surface accumulation will cause air and ice loss at a ratio of 2 : 1–1 : 1 if snow is initially deposited at a density of $350\text{--}450 \text{ kg m}^{-3}$ (Kuipers Munneke et al., 2012b) and compensating compaction changes are ignored. Below-balance accumulation of approximately half of the modelled value (Fig. 7) would be required to solely explain our ice loss, and the fact that our ice loss is an order of magnitude larger than the air loss suggests that below-balance accumulation alone cannot account for both. A small below-balance accumulation could, however, explain the air loss. Since the total input of air into the firn is $0.5\text{--}1 \text{ myr}^{-1}$, relatively small anomalies in surface melting, dry compaction, or accumulation are required to yield the observed 0.04 myr^{-1} air loss.

We now turn to dynamic mechanisms. Above-balance ice flow advection will affect air and ice thicknesses in proportion to their relative gradients along-flow. According to the results of Holland et al. (2011), increased advection would enhance the flow of thicker ice with less firn air across the survey line. The air thickness increases along-flow by approximately 1 m for every 10 m decrease in along-flow ice thickness. Above-balance

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advection would therefore cause air loss, but accompanied by ice gain approximately ten times faster, which entirely contradicts our observed signals. Above-balance ice flow divergence will cause air and ice losses in proportion to their relative thicknesses, approximately 1 : 30 for characteristic ice and air thicknesses of 10 and 300 m. The largest velocity change in the literature is an acceleration of 80 myr^{-1} between 2000 and 2006 surveys of northern LCIS (Haug et al., 2010; Khazendar et al., 2011). If this acceleration caused unbalanced divergence over a length scale of 100 km, it would cause ice loss of $\sim 0.24 \text{ myr}^{-1}$ and air loss of $\sim 0.008 \text{ myr}^{-1}$. Above-balance divergence could explain the ice loss, but not the air loss, if maintained at this level and not accompanied by above-balance advection.

In summary, the ice loss we observe could be explained by above-balance basal melting and/or ice divergence, and the air loss could be explained by below-balance accumulation or above-balance surface melting and/or compaction. Our results therefore suggest that at least two different forcings caused the lowering of LCIS during our survey period. Elsewhere around Antarctica, rapid ice-shelf thinning is thought to be driven by unbalanced ocean melting (e.g. Shepherd et al., 2004; Holland et al., 2010; Padman et al., 2012; Khazendar et al., 2013), and our robust evidence of a firn-air loss from LCIS in response to surface processes is the first direct evidence of an exception to this. The existence of at least two different mechanisms underlying the change is also consistent with our observation that the ice and air loss signals have different spatial variation along the survey line.

The surveys do not encompass all of the known ice-shelf lowering (Fig. 1), and it is likely that the balance of ice and air losses, and their driving mechanisms, varies in different regions and periods. In particular, our surveys do not capture the rapid lowering in northern LCIS. Ice divergence may play a part in this, since the known acceleration of LCIS is northward-intensified (Haug et al., 2010; Khazendar et al., 2011), but there are also good reasons to expect changes in surface melting to be largest in the north (Holland et al., 2011; Trusel et al., 2013; Luckman et al., 2014). The pattern of changes in basal melting is unknown.

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5.2 Ice-shelf stability

Our results have important implications for the future stability of LCIS and thus the AP Ice Sheet. Previous ice-shelf collapses are thought to have been accomplished by surface meltwater-driven crevassing (Scambos et al., 2003; van den Broeke, 2005; Banwell et al., 2013) or ice-front retreat past a “compressive arch” in strain rates (Doake et al., 1998). The northeastern part of LCIS is likely to be least stable, since it has high surface melting and low firm air (Holland et al., 2011), is showing the most rapid lowering (Shepherd et al., 2003) and acceleration (Khazendar et al., 2011), is highly crevassed (McGrath et al., 2012), is slow-moving and largely sustained by accumulation, and has a stress field conducive to instability (Kulesa et al., 2014). We conceive several interconnected mechanisms by which LCIS stability could be compromised: (1) ice-front retreats past a compressive arch, (2) increased surface melting causes firm depletion and meltwater-driven crevassing, (3) decreased ocean freezing or increased melting depletes marine ice, permitting the propagation of crevasses, (4) collapse of Scar Inlet opens a new ice front at the northern margin of LCIS, (5) ungrounding from Bawden Ice Rise removes an ice-front pinning point, (6) ice thinning and acceleration enhances the propagation of crevasses and weakens shear zones.

5.2.1 Retreat past compressive arch

Doake et al. (1998) suggested that LBIS was in a stable configuration when the second principal strain rate was compressive everywhere inshore of a “compressive arch” near the ice front. Once this arch was breached by calving, a significant collapse followed. Kulesa et al. (2014) showed that LCIS has a large region near the ice front in which the second principal stress is tensile and thus offshore of a compressive arch. Kulesa et al. (2014) also considered the angle between the flow and first principal stress under the assumption that rifts strike perpendicular to the flow; if the first principal stress is aligned with the flow it will therefore tend to open rifts, rendering the ice shelf unstable. LCIS has a large region near the ice front in which the first principal stress is oriented

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across-flow, thus stabilising the ice shelf. It is argued that this region is secured by marine ice (Kulesa et al., 2014), but there is clearly a risk that calving in this region will remove ice that both stabilises rifts and shields the compressive arch, leading to a progressive collapse of LCIS. We are unable to assess a timescale for this possibility.

5.2.2 Meltwater-driven crevassing

The final collapse of many AP ice shelves has been linked to the availability of surface meltwater to enhance the downward propagation of surface crevasses (Scambos et al., 2003; van den Broeke, 2005; Banwell et al., 2013). There are significant crevasse fields on the surface of LCIS, so we hypothesise that meltwater ponding is sufficient to drive collapse. Meltwater is already ponding in limited areas near the LCIS grounding line (Holland et al., 2011; Luckman et al., 2014), but these do not pose an imminent risk of collapse. Before ponding can occur it is necessary for the firn to be largely depleted of its air content, since otherwise meltwater will simply percolate and refreeze. Holland et al. (2011) showed that the northern part of LCIS had approximately 10 m of firn air remaining in 1998, while the retreating LBIS had very little. Our derived air loss of 0.04 myr^{-1} would require 250 years to deplete 10 m of air and threaten LCIS stability. However, the lowest air content and highest lowering are north of the survey line, and it is likely that surface melting will increase over the coming centuries (Kuipers Munneke et al., 2014), so this timescale is probably an upper bound.

5.2.3 Depletion of marine ice

There is plenty of evidence that LCIS is stabilised by marine ice (Holland et al., 2009; Khazendar et al., 2011; Jansen et al., 2013; Kulesa et al., 2014; McGrath et al., 2014), and this implies that decreased marine ice deposition or increased melting could allow LCIS to collapse under its existing stress field. The marine ice at the ice front can form a very small fraction of the ice column, implying that the stability of basal crevassing and ice-front calving is controlled by only tens of metres of marine ice (McGrath et al.,

2014). Elsewhere the marine ice can be hundreds of metres thick (Jansen et al., 2013; Kulesa et al., 2014; McGrath et al., 2014). If our ice loss estimate of 0.3 myr^{-1} is caused by unbalanced basal melting, this suggests a timescale of 170 years to remove the bottom 50 m of ice, destabilising the ice front, and 500 years to remove the lowest 150 m of ice, destabilising the eastern half of LCIS. These timescales are extremely uncertain because the ocean processes driving melting and freezing are unknown and thus impossible to project. Counter-intuitively, increased ice-shelf melting could actually increase the meltwater-driven ocean currents and increase their marine ice deposition downstream (Holland et al., 2009). If marine ice deposition were to cease altogether, it would take 400–500 years to remove the existing marine ice from LCIS solely by lateral ice advection and iceberg calving.

5.2.4 Collapse of Scar Inlet

Albrecht and Levermann (2014) propose that the collapse of any ice shelf can destabilise neighbouring ice shelves by changing their stress regime. In the context of LCIS, this translates into the risk that the collapse of LBIS could remove buttressing across Jason Peninsula. Scar Inlet, the last remaining part of LBIS, is presumably at risk of disintegration, so we assess the impact of this possibility on LCIS. Jason Peninsula anchors a large area of stagnant ice that is a significant stabilising influence on both LCIS and Scar Inlet (Fig. 9a). The ice between LCIS and Scar Inlet, Phillipi Rise, is poorly surveyed but appears to be well-grounded at present, with ice 150 m above floatation (calculated using 5 m firn air from Holland et al., 2011, EIGEN-6C geoid, and mean dynamic ocean topography of -1 m ; Fig. 9b). However, the ice base is hundreds of metres below sea level in places (Fig. 9c), so there is certainly the possibility that ice thinning could open a new oceanographic pathway through Jason Peninsula. However, given the stagnant nature of this ice, it is unclear to what extent this would influence LCIS stability.

5.2.5 Ungrounding from Bawden Ice Rise

Of far greater concern is the stability of Bawden Ice Rise. An ungrounding from Bawden Ice Rise would prompt significant acceleration of the ice shelf (Borstad et al., 2013) and re-organisation of its strain rate field, probably destabilising the ice front (Kulesa et al., 2014). The ice rise is only a few kilometres across, but has a noticeable effect upon the flow and structure of the ice shelf (Fig. 10a). Three radar survey lines show that the ice rise is very lightly grounded in the north, but approximately 40 m above floatation at its summit in the south (Fig. 10b), where the ice base is about 150 m below sea level (Fig. 10c). (Height above floatation is calculated using a 10 m firn air content derived from nearby surveyed floating ice and finding elevation relative to sea level using nearby surveyed open water.) Our ice loss estimate of 0.3 myr^{-1} would take 130 years to unground the ice rise entirely, but this timescale is subject to great uncertainty, including the ice loss estimate itself, its applicability to this region, and the projection of this rate into the future. It is almost certainly an upper bound because lowering is rapid in the region and Bawden would cease to provide a significant stabilising influence, and may even destabilise the ice front, long before the ice actually ungrounds through thinning. For example, Doake and Vaughan (1991) showed that ice rises destabilised Wordie Ice Shelf by acting as an “indenting wedge” during the retreat of its ice front. A large calving south of Bawden between late December 2004 and early January 2005 and the ongoing thinning and acceleration in the region might even indicate that ungrounding from Bawden is already underway.

5.2.6 Crevassing weakens shear zones

Whatever its source, the ongoing thinning and acceleration of LCIS could ultimately cause its demise by weakening the structural integrity of the ice shelf. LAIS and LBIS both accelerated before collapsing (Bindshadler et al., 1994; Rignot et al., 2004), and LBIS apparently collapsed after weakening of the shear zones between ice flow units (Khazendar et al., 2007; Vieli et al., 2007; Glasser and Scambos, 2008). The shear

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zones in the north of LCIS are slower-moving and not so strongly sheared (Khazendar et al., 2011) and hence more stable, but the ice is already quite damaged (Jansen et al., 2010; McGrath et al., 2012; Borstad et al., 2013). The uncertainties in this interaction are large and we are unable to assess a timescale for this risk.

6 Conclusions

We analyse eight repeated radar surveys between 1998 and 2012 along a nearly meridional line that traverses the centre of Larsen C Ice Shelf (LCIS), applying a novel method to derive the separate ice and air losses along this line contributing to the known lowering of the ice shelf. The uncertainties are considerable, but our primary estimate is that the lowering ($0.066 \pm 0.017 \text{ m yr}^{-1}$, or $0.99 \pm 0.26 \text{ m}$) is caused by both ice loss ($0.28 \pm 0.18 \text{ m yr}^{-1}$, or $4.2 \pm 2.7 \text{ m}$) and firn air loss ($0.037 \pm 0.026 \text{ m yr}^{-1}$, or $0.56 \pm 0.39 \text{ m}$). Though their effect on the surface lowering is approximately equal, ice loss is an order of magnitude larger than air loss. The derivation of these values allows us to speculate upon the possible sources of the changes, and their future implications.

The ice loss we observe could be explained by above-balance basal melting and/or ice divergence, and the air loss could be explained by below-balance accumulation or above-balance surface melting and/or compaction. We conclude that at least two different forcings caused the lowering of LCIS during our survey period. The surveys do not sample the most rapid ice-shelf lowering in northern LCIS and it is likely that the balance of ice and air losses, and their driving mechanisms, varies for different regions and periods.

We conceive several interconnected mechanisms by which LCIS stability could be compromised, and our ice and air loss rates suggest typical timescales for LCIS collapse of a few centuries. The two mechanisms that offer the earliest possibility of collapse are a flow perturbation arising from the ungrounding of LCIS from Bawden Ice Rise, and ice-front retreat past a “compressive arch” in strain rates, suggesting that the stability of Bawden Ice Rise and calving from the ice front should be monitored closely.

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Table 1. Details of the radio-echo sounding and altimeter surveys used in this analysis.

Date	Origin	Platform	Ice-Sounding Radar	Ice Elevation
20 Feb 1998	BAS-Argentina	Twin Otter	150 MHz ^a	radar altimeter ^a
26 Nov 2002	NASA-CECS	P-3	ICORDS2 140–160 MHz ^b	laser ATM ^c
29 Nov 2004	NASA-CECS	P-3	ACORDS 140–160 MHz ^b	laser ATM ^c
4 Nov 2009	NASA IceBridge	DC-8	MCoRDS 190–200 MHz ^{b,d,f}	laser ATM ^c
19–21 Nov 2009	McGrath	Sledge	25 MHz ^e	GPS ^e
13 Nov 2010	NASA IceBridge	DC-8	MCoRDS 190–200 MHz ^{b, d}	laser ATM ^c
27 Jan 2011	BAS	Twin Otter	150 MHz	radar altimeter
13–14 Dec 2012	Brisbourne	Sledge	50 MHz	GPS

^a Holland et al. (2009);

^b <https://data.cresis.ku.edu/#RDS>;

^c <http://nsidc.org/data/ilatm2>;

^d <http://nsidc.org/data/irmcr2.html>;

^e McGrath et al. (2014).

^f Data neglected due to transmit/receive switch problem; see Sect. 2.4.

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Table 2. Elevation and TWTT trends and their derived ice and air trends from calculations performed using different combinations of data. TWTT trends are expressed as solid-ice thickness equivalent. Trends in bold are smaller than the derived uncertainty (see main text).

case	elevation (myr^{-1})	TWTT (m ice yr^{-1})	ice (myr^{-1})	air (myr^{-1})
Reference	−0.0660	−0.296	−0.274	−0.0367
Using satellite altimetry	−0.0616	−0.296	−0.277	−0.0320
BAS only ^a	−0.0752	−0.264	−0.235	−0.0500
NASA only ^b	−0.0303	0.087	0.110	−0.0421
Without 1998	−0.0311	−0.041	−0.025	−0.0285
Without 2002	−0.0694	−0.389	−0.371	−0.0297
Without 2004	−0.0713	−0.281	−0.256	−0.0439
Without 2009 MG	−0.0654	−0.195	−0.168	−0.0474
Without 2010	−0.0648	−0.351	−0.334	−0.0290
Without 2011	−0.0695	−0.394	−0.377	−0.0292
With 2009 IB TWTT	−0.0660	−0.482	−0.471	−0.0155
With 2012	−0.0670	−0.212	−0.185	−0.0473
Uncertainty (see text)	0.017	0.17	0.18	0.026

^a All 1998 and 2011 data.

^b All 2002, 2004, 2010 data and elevation for IceBridge 2009.

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Table 3. Statistics of the differences between all data from each survey and their nearest 2004 analogue, as shown in Fig. 4. TWTT is expressed as solid-ice thickness equivalent.

survey	elevation differences from 2004 (m)				TWTT differences from 2004 (m ice)			
	count	mean	SD	stderr	count	mean	SD	stderr
1998	2213	0.993	1.365	0.029	1382	2.320	7.507	0.202
2002	5092	0.376	1.329	0.019	952	−3.384	9.365	0.304
2004	6097	0	0	0	18 385	0	0	0
2009 MG	8731	−0.013	1.726	0.019	4385	−5.441	9.501	0.144
2009 IB	4779	0.215	1.139	0.017	4444	−11.62	11.91	0.179
2010	4461	−0.088	1.836	0.028	5317	−1.784	9.847	0.135
2011	12 126	0.020	1.573	0.014	9190	−1.097	9.802	0.102
2012	303	−0.225	2.401	0.138	187	−0.976	9.651	0.706

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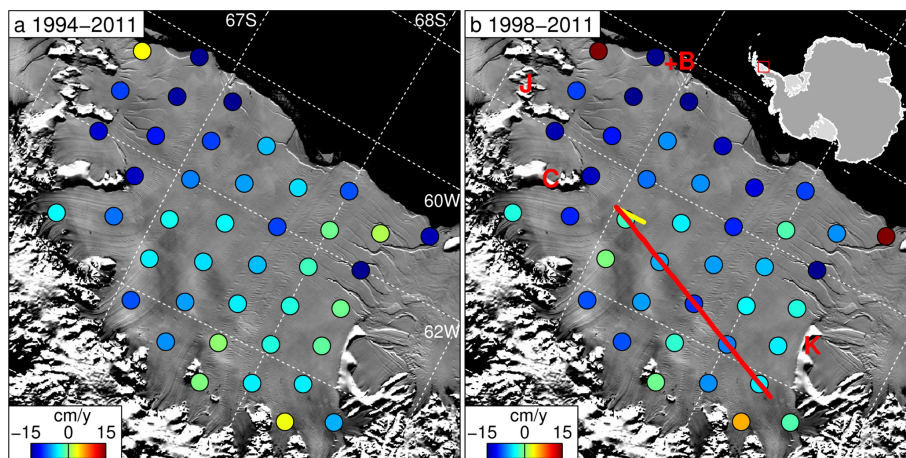


Figure 1. MODIS Mosaic of Antarctica imagery of LCIS (Scambos et al., 2007) showing the location of satellite radar altimeter crossovers and estimated surface lowering rates (updated from Fricker and Padman, 2012, as described in Sect. 2) for two periods. **(a)** 1994–2011, the full period for which ERS-1/2 and Envisat data are reliable; **(b)** 1998–2011, the period for which we have radar surveys. The main survey line is shown in red, with the 2012 survey shown in yellow. **(b)** shows geographical features referred to in the text: B: Bawden Ice Rise; C: Churchill Peninsula; J: Jason Peninsula; K: Kenyon Peninsula.

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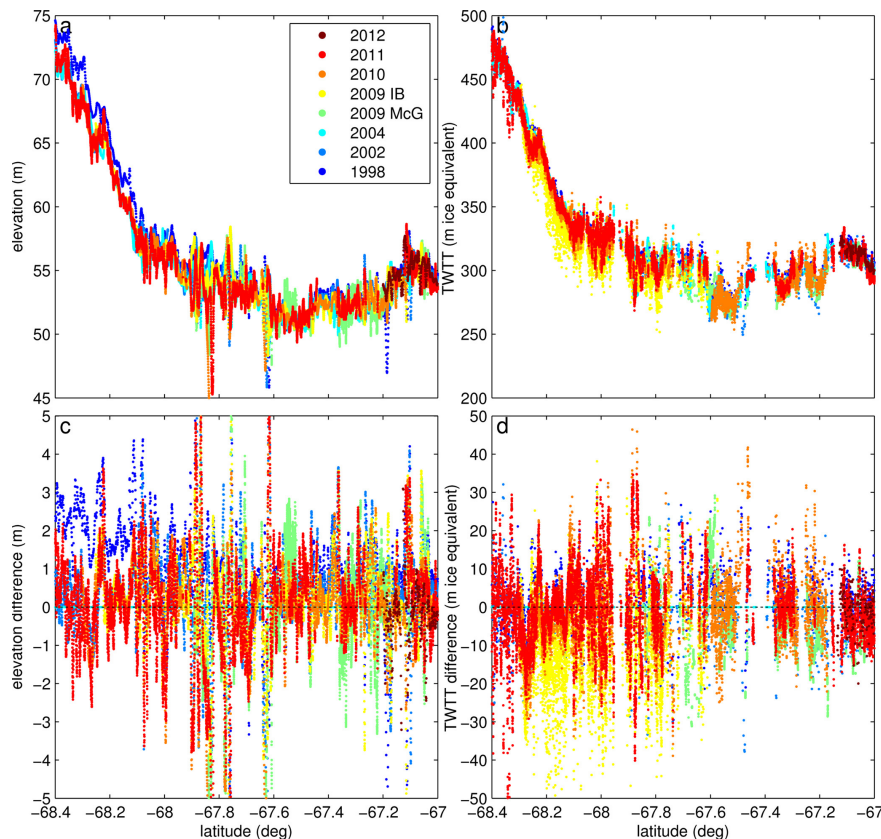


Figure 2. Processed data from the eight surveys, from which the air and ice thickness changes are derived. **(a)** Surface elevation relative to WGS84 ellipsoid. **(b)** Radar two-way travel time (TWTT), expressed as an equivalent thickness of solid ice. **(c)** Difference between each elevation observation and nearest 2004 analogue. **(d)** Difference between each TWTT observation and nearest 2004 analogue.

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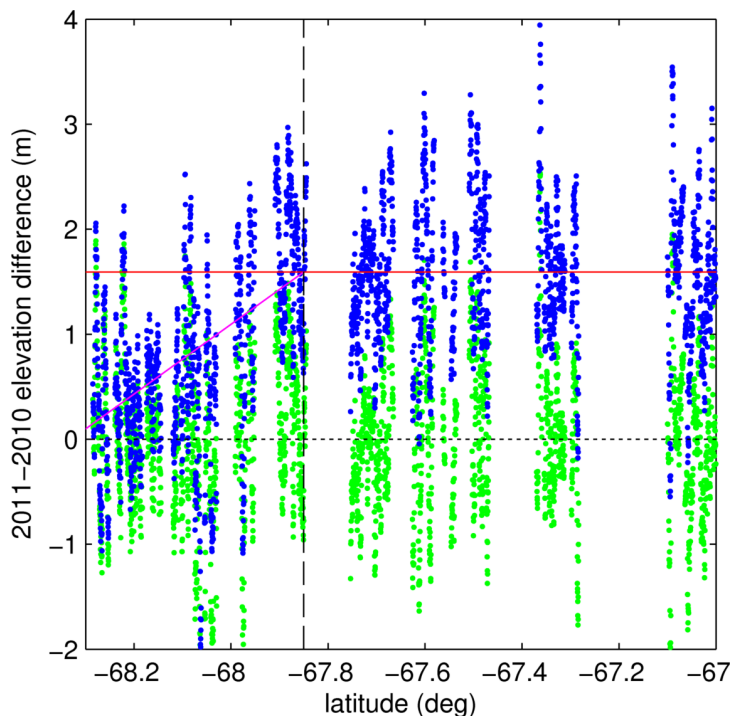



Figure 3. Calibration of the elevation data in the 2011 BAS airborne survey. Blue dots indicate the difference between elevations derived from BAS radar altimetry on 27 January 2011 and from IceBridge laser altimetry on 13 November 2010. Assuming negligible elevation change over the intervening ~ 10 weeks, the 2011 data need to be corrected by subtracting a constant offset of 1.59 m (red line; the mean difference from 2010 for all data north of 67.85°S). South of 67.85°S , the 2011 data become progressively lower, which is attributed to increasing radar penetration of the firn (Holland et al., 2011). In this region we add an additional penetration correction equal to the difference between the red and magenta lines. This penetration correction is also applied to the 1998 BAS radar altimeter data. Green dots show the difference between the 2010 data and the corrected 2011 data.

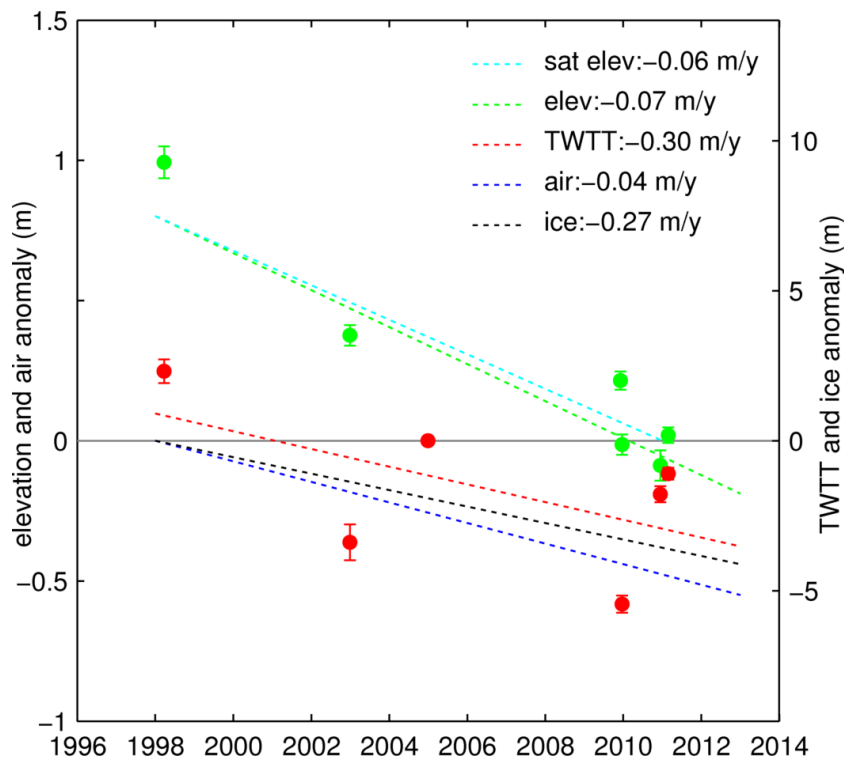


Figure 4. Inter-survey differences in elevation, TWTT, ice and air. Mean differences between each survey and 2004 for elevation are shown in green and for radar two-way travel time (TWTT; ice equivalent) in red. Error bars represent 95 % confidence intervals, and dashed lines represent linear trend lines. The elevation trend derived from satellite radar altimetry is also shown in cyan. Trends in ice (black) and air (blue) thickness are derived directly from the trends in TWTT and elevation, revealing that LCIS has lost both ice and air over the period surveyed. Elevation and air thickness use the left axis, while TWTT and ice thickness are plotted with absolute values on the right axis and equivalent surface elevation on the left axis (see Sect. 4).

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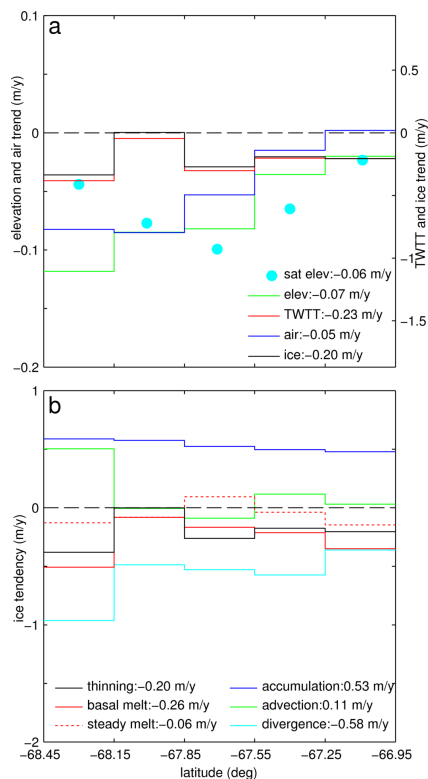


Figure 5. Spatial variation in derived quantities along the survey line within latitude bins centred upon the locations of the satellite cross-over points (see Fig. 1b). **(a)** Trends in elevation (green), TWTT (red; ice equivalent), and air (black) and ice (blue) thickness, showing significant ice and air loss. Elevation trends derived from satellite radar altimetry at the crossovers are cyan. Elevation and air thickness use the left axis, while TWTT and ice thickness are plotted with absolute values on the right axis and equivalent surface elevation on the left axis. **(b)** Spatial variation in ice mass budget. Divergence balances accumulation, and ice thinning implies unbalanced basal melting. Values in the legends represent means over all bins.

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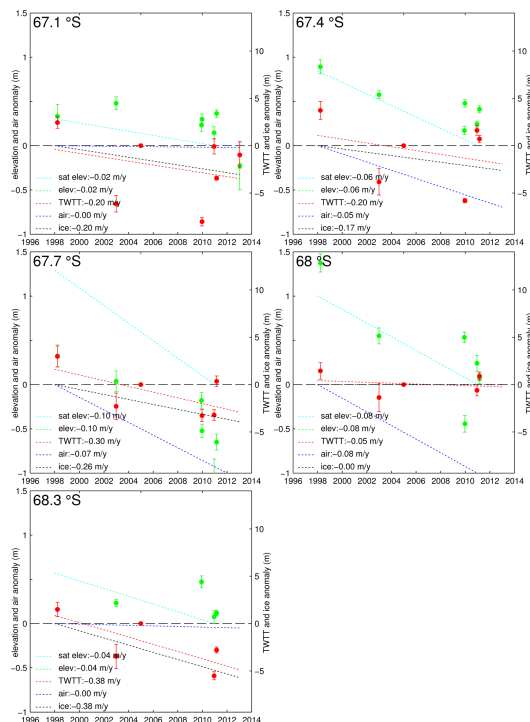


Figure 6. Data and trends for the five latitude bins defined by the satellite altimetry crossovers, labelled with the latitude of the accompanying crossover. Data points show the mean and 95 % confidence intervals of the differences between each survey and the 2004 baseline for surface elevation (green) and TWTT (red, expressed as solid-ice equivalent). The satellite-altimeter derived elevation trend for the crossover at the centre of each bin is also shown (cyan). Surveyed trends in elevation and TWTT are converted to trends in ice (black) and air (blue) thickness. Elevation and air thickness are plotted on the left-hand axis, while TWTT and ice thickness are plotted such that the right-hand axis shows absolute values and the left-hand axis shows the equivalent surface elevation.

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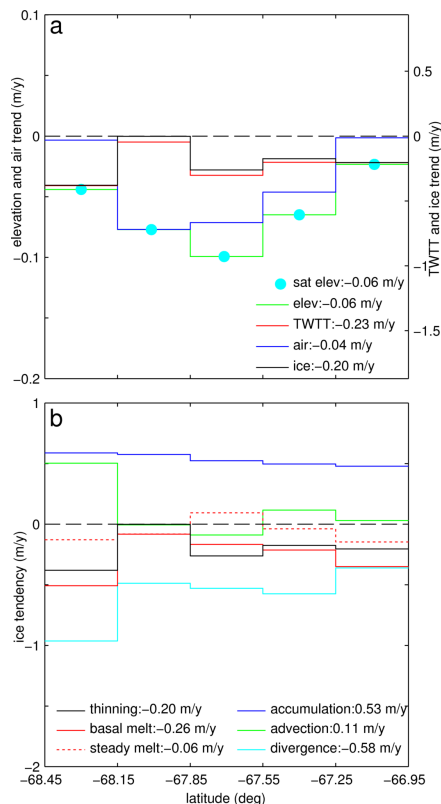


Figure 7. Version of Fig. 5 in which the binned survey elevation trends are replaced by satellite crossover elevation trends. **(a)** Spatial variation of trends in elevation (green), TWTT (red, ice equivalent), and air (black) and ice (blue) thickness. Satellite crossover trends are cyan. Elevation and air thickness use the left axis, while TWTT and ice thickness are plotted with absolute values on the right axis and equivalent surface elevation on the left axis. **(b)** Meridional variation in ice mass budget.

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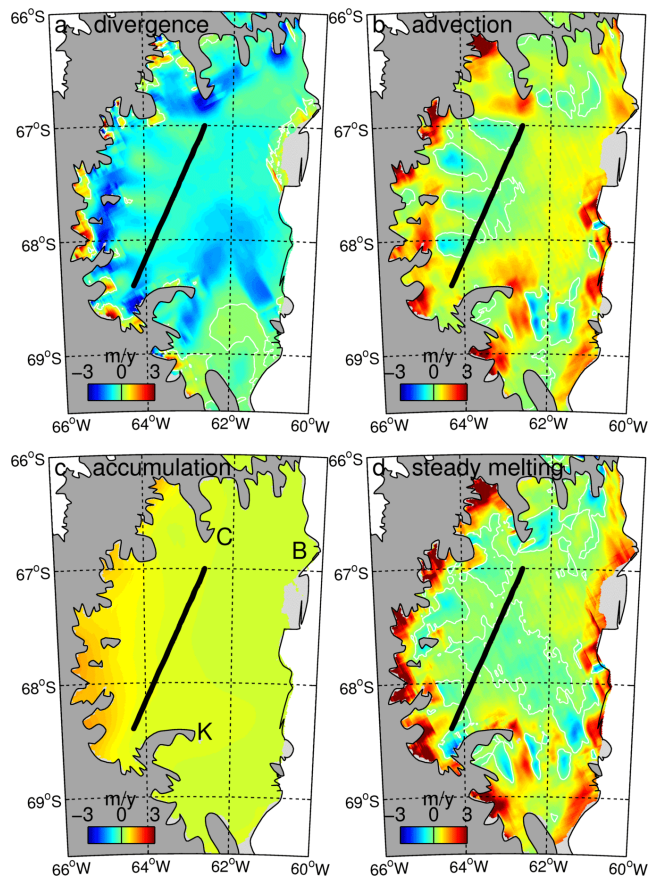


Figure 8. Fields of derived values for the terms in the ice-only mass balance (positive implies melting). **(a)** Ice divergence ($-1/\nabla \cdot \mathbf{u}$); **(b)** ice advection ($-\mathbf{u} \cdot \nabla I$); **(c)** ice surface accumulation; **(d)** derived steady-state basal melting. **(c)** shows geographical features referred to in the text: B: Bawden Ice Rise; C: Churchill Peninsula; K: Kenyon Peninsula.

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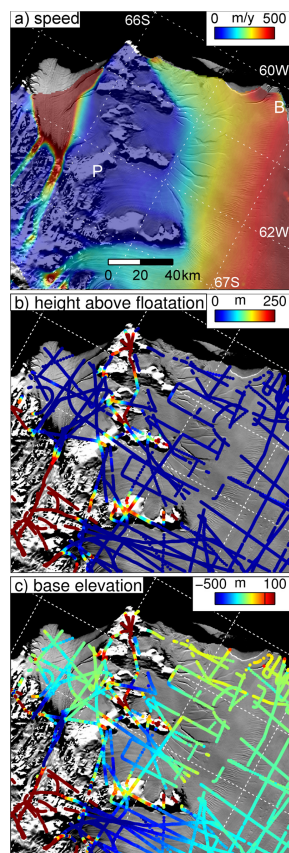


Figure 9. Northern LCIS and Jason Peninsula, showing various quantities overlain on MODIS Mosaic of Antarctica (Scambos et al., 2007). **(a)** Ice flow speed (Rignot et al., 2011); **(b)** height of ice surface above hydrostatic flotation; **(c)** elevation of ice base relative to sea level. Bawden Ice Rise is labelled B and Phillipi Ice Rise is labelled P.

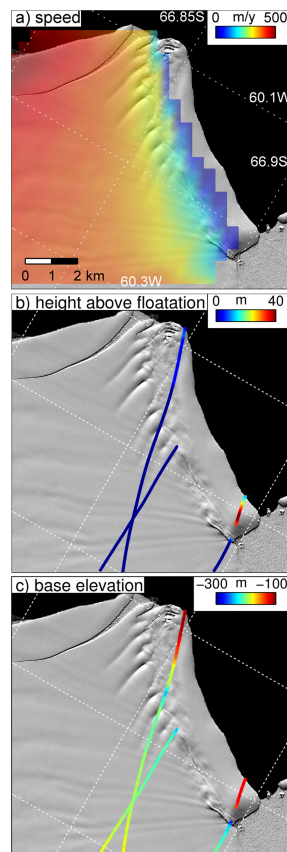


Figure 10. High-resolution WorldView2 satellite imagery of Bawden Ice Rise acquired 15 October 2012 (copyright Digital Globe) with various quantities overlain. **(a)** Ice flow speed (Rignot et al., 2011); **(b)** height of ice surface above hydrostatic floatation; **(c)** elevation of ice base relative to sea level.