



Supplement of

An ice-sheet-wide framework for englacial attenuation from ice-penetrating radar data

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Supplementary material for: An ice sheet wide framework for englacial attenuation from ice penetrating radar data

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In this supplementary material we firstly discuss the temperature sensitivity of the algorithm sample region (Sect. S1), and the tuning of the sample region tolerance (Sect. S2), which both accompany Sect 2.5 in the main article. We then illustrate the ice sheet wide properties of the algorithm using the SICOPOLIS temperature field as an input (Sect. S3), as was shown for the GISM temperature field

5 in Sect. 3.1. The convergence of the attenuation solution in drainage basins other than the SE GrIS (Sect. 3.2 in the main article) is then discussed (Sect. S4). Finally, we provide evidence of algorithm repeatability for independently analysed field campaign data (Sect. S5), which expands upon Sect. 3.4 in the main article.

S1 Temperature sensitivity of the algorithm sample regions

- 10 Maps for the target window vectors, R_1 - R_4 , using the SICOPOLIS temperature field are shown in Fig. S1. These are equivalent to Fig. 7 in the main text for the GISM temperature field, and a target window vector difference plot for the two temperature fields is shown in Fig. S2. It is evident that both temperature fields share the same ice sheet wide trend that the target window vectors are larger in the interior of the ice sheet. Generally, the GISM temperature field has larger window vectors than
- 15 SICOPOLIS toward the ice margins. An explanation for this trend is that the horizontal gradient in the absolute values $\langle \hat{B} \rangle$ (which is approximately equivalent to depth-averaged temperature) are larger toward the margins for the SICOPOLIS temperature field (also refer to Fig. 5. in main article). The target windows vectors are more sensitive to temperature field, (in terms of absolute differences in target window vector length), in the northern interior of the GrIS. This is thought to arise because

20 the segment approximation that is used in the target windowing procedure becomes more sensitive in colder ice, where there is a smaller horizontal gradient in $\langle \hat{B} \rangle$. Future refinements to the algorithm could therefore consider developing a more refined target windowing procedure for the sample region in colder ice

S2 Tuning the tolerance of the algorithm sample region

threshold of 20 data points.

A free parameter in our sample region windowing method is the RMS tolerance metric that is used 25 to define the target window vectors and boundaries, (refer to Appendix C). If the RMS tolerance is too small, then the ice thickness range is insufficient to discriminate an attenuation trend and the solution accuracy and related coverage regions are reduced. This point is illustrated by a comparison of algorithm coverage regions for RMS=1 dB km⁻¹ and RMS=0.5 dB km⁻¹ in Fig. S3. The lower coverage for RMS=0.5 dB km⁻¹ arises primarily due to lower $r_{[P^C]}^2$ values. Additionally, due to the 30 smaller sample region windows for RMS=0.5 dB km⁻¹ less of the grid points meet the minimum

As the RMS tolerance is increased, the sample regions will contain individual ice columns with a greater range of depth-averaged attenuation rates. This is undesirable for a bed-returned power

- attenuation algorithm, which either assumes: (i) local constancy or, (ii) makes a local attenuation 35 correction. For algorithms that assume local constancy, the degree of systematic underestimation the attenuation rate (described in Sect. 2.6 and Fig. 8) is predicted to increase with window size/RMS tolerance. For the local attenuation correction variant of the bed-returned power method described here, the local attenuation correction, Eq. (9), increases with the RMS tolerance. It therefore follows
- that greater RMS tolerance values are more susceptible to bias from the Arrhenius temperature model 40 input. Subsequently, we suggest that the desired level of solution accuracy should inform the choice the RMS tolerance. As discussed in the main article, the desired attenuation loss accuracy for basal melt discrimination is ~ 5 dB. If this is rescaled by ice thickness for a sample region in the interior of the ice sheet (mean ice thickness ~ 2500 m), this results in a desired attenuation rate accuracy ~ 1
- dB km⁻¹. This choice is consolidated by our estimates of attenuation rate solution accuracy in Sect. 45 3.4, and Sect. S5 which are similar to the predicted range of values over which the ice columns are sampled.

Ice sheet wide properties using the SICOPOLIS temperature field **S3**

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The ice sheet wide properties of the radar algorithm using the SICOPOLIS temperature field are shown in Fig. S4, and are equivalent to Fig. 9 in the main article for GISM. The map for $r_{[P^C]}^2$ for SICOPLOLIS is very similar to GISM, and supports the view that the algorithm has greater utility for identifying an attenuation trend in regions of rougher basal topography and where the absolute values of the attenuation rate are higher. The regions of high $r_{[\hat{R}]}^2$ /low r_{ratio}^2 for the two temperature

fields are less well correlated. For the reasons discussed in Sect. 2.7 it is difficult to establish whether,

- 55 for either temperature field, the majority of these regions correspond to: (i) true variation in the basal reflector with ice thickness, (ii) a strong bias in the ice sheet model temperature field. However, for regions where there is ice core temperature data available it is possible to predict which scenario is more likely. Notably, for the region surrounding the NEEM and Camp Century ice cores in North West Greenland, the GISM temperature field is very close to the core data (Fig. S5). This is therefore
- 60 suggestive that the regions of high $r_{[\hat{R}]}^2$ in Fig. 9 correspond to a true variation in the basal reflector with ice thickness.

As described in Sect. 2.6, a minimum of 20 IPR (along-track averaged) measurements within each target window is set as a threshold for including a grid cell in the linear regression procedure. Due to the generally smaller target window vectors for the SICOPOLIS temperature field toward the ice

65 margins, Fig. S2, there are more grid cells that fall below this threshold if this temperature field is used as in input. This effect is particularly apparent along the western ice margin where the maps for SICOPOLIS input in Fig. S4 are non-continuous.

S4 Demonstration of solution convergence in other drainage basins

- We observe qualitatively similar attenuation solution behavior in drainage basins 3,5,6 (southern and eastern Greenland) to that described for drainage basin 4 in the main article. In all these basins we observe algorithm solution convergence and an associated reduction in the solution bias from the Arrhenius model input. An example of the attenuation rate solution differences for drainage basin 6 (the SW GrIS) is shown in Fig. S6a, with the radar-inferred distribution being approximately normally distributed about zero, (μ ± σ= -0.18 ± 1.56 dB km⁻¹), and the Arrhenius model input
 75 having a mean systematic bias of μ=-1.66 dB km⁻¹. In drainage basins 1,2,7,8 (northern and western
- Greenland) we do not observe analogous solution convergence, (in the sense of normally distributed difference distribution for the radar-inferred values). We do, however, typically see a reduction in the mean systematic bias of the attenuation rate solution, relative to the Arrhenius model input. For example, in drainage basin 8 (the NW GrIS), the Arrhenius model mean attenuation rate bias is μ =-
- 3.62 dB km⁻¹, whereas the radar-inferred value is μ =-2.62 dB km⁻¹ (Fig. S6b). An explanation for why we observe solution convergence in some basins but not others is thought to relate to the temperature dependence of the target windows (Sect. S1), which is generally higher in the northern interior. In turn, this results in different IPR data being sampled for the different temperature field inputs.

85 S5 Repeatability and uncertainty for independently analysed field seasons

In Sect. 3.4. we proposed that the attenuation solution uncertainty can be measured by: (i) solution variation for different input temperature fields, (ii) solution variation for independently analysed field

seasons (with a fixed temperature field input). Since, for equivalent algorithm settings, the solution variation for different input temperature fields is slightly higher, we focused upon this in the main article as it gives a better overall indicator of solution uncertainty.

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A full summary of attenuation rate solution variation for different field season pairs is shown in Table S1 and an example of the distributions and mutual coverage regions is shown in Fig. S7. The mutual coverage regions occur where there are regions of intersecting flight tracks (Fig. 1a). As these overlap regions are quite small, we considered the GrIS as a whole when measuring solution

- variation. The general relationship between solution variation (as measured by the standard deviation of the solution difference distribution) and algorithm coverage (as measure by the frequency of grid cells that contain IPR data), is similar to the case for different input temperature fields. For $(\alpha, \beta) =$ (0.60, 0.80), the mean standard deviation aggregated over the different field season combinations is 1.29 dB km⁻¹. This supports the conclusion in the main paper that this algorithm quality setting
- 100 represents sufficient accuracy for defining a coverage region for discrimination of basal melt from the reflection values.

References

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Figure S1. Maps for target window radi vector length using the SICOPOLIS temperature field. (a) Vector R_1 , (b) Vector R_2 , (c) Vector R_3 , (d) Vector R_4 . The orientation of each radi vector is shown in each subplot.



Figure S2. Difference maps for the length of target window radi vectors (GISM-SICOPOLIS). (a) Vector R_1 , (b) Vector R_2 , (c) Vector R_3 , (d) Vector R_4 . The orientation of each radi vector is shown in each subplot.



Figure S3. Example of algorithm coverage regions for different window tolerances. (**a**) RMS=1 dB km¹, (**b**) RMS=0.5 dB km¹. Both plots are for the 2011 P3 field season data using the GISM temperature field.



Figure S4. Ice sheet wide properties of the radar algorithm using the SICOPOLIS temperature field. (a) Powerthickness correlation, $r_{[P^C]}^2$. (b) Arrhenius reflection coefficient-thickness correlation, $r_{[\hat{R}]}^2$. (c) Correlation ratio, r_{ratio}^2 . (d) Coverage for three thresholds (green is a subset of red and red is a subset of blue). (e) Radarinferred attenuation rate, $\langle B(T_{SIC}) \rangle$, for $(\alpha, \beta) = (0.60, 0.80)$. (f) Radar-inferred attenuation loss, $[L(T_{SIC})]$, for $(\alpha, \beta) = (0.60, 0.80)$.



Figure S5. Vertical temperature profiles. (a) Camp Century (Weertman, 1968), (b) NEEM (Macgregor et al., 2015).



Figure S6. Attenuation rate difference distributions. (a) SW Greenland (drainage basin 6 in Fig. 1), (b) NW Greenland (drainage basin 8 in Fig. 1).

Season pair	Algorithm settings	Mean, μ , (dB km ⁻¹)	Standard deviation, σ , (dB km ⁻¹)	Frequency, N.
2012-2011	T_{GISM} : $(\alpha, \beta) = (0.60, 0.80)$	-0.28	1.34	9869
	T_{GISM} : $(\alpha, \beta) = (0.70, 0.85)$	-0.11	1.08	5522
	T_{GISM} : $(\alpha, \beta) = (0.80, 0.90)$	0.26	0.65	953
	T_{SIC} : $(\alpha, \beta) = (0.60, 0.80)$	-0.71	1.71	6339
	T_{SIC} : $(\alpha, \beta) = (0.70, 0.85)$	-0.81	1.36	4164
	T_{SIC} : $(\alpha, \beta) = (0.80, 0.90)$	-0.50	1.01	1371
2013-2011	$T_{\text{GISM}}: (\alpha, \beta) = (0.60, 0.80)$	-0.13	1.46	6698
	T_{GISM} : $(\alpha, \beta) = (0.70, 0.85)$	-0.26	1.03	3584
	T_{GISM} : $(\alpha, \beta) = (0.80, 0.90)$	0.181	0.62	521
	T_{SIC} : $(\alpha, \beta) = (0.60, 0.80)$	-0.73	1.42	5330
	T_{SIC} : $(\alpha, \beta) = (0.70, 0.85)$	-0.84	1.14	3509
	T_{SIC} : $(\alpha, \beta) = (0.80, 0.90)$	-1.09	0.85	953
2014-2011	T_{GISM} : $(\alpha, \beta) = (0.60, 0.80)$	0.05	1.55	10036
	$T_{\text{GISM}}: (\alpha, \beta) = (0.70, 0.85)$	- 0.18	1.12	5045
	T_{GISM} : $(\alpha, \beta) = (0.80, 0.90)$	0.57	0.85	533
	T_{SIC} : $(\alpha, \beta) = (0.60, 0.80)$	-0.22	1.48	7934
	T_{SIC} : $(\alpha, \beta) = (0.70, 0.85)$	-0.38	1.18	5614
	T_{SIC} : $(\alpha, \beta) = (0.80, 0.90)$	-0.16	1.02	1765
2013-2012	$T_{\text{GISM}}: (\alpha, \beta) = (0.60, 0.80)$	0.33	1.06	6557
	T_{GISM} : $(\alpha, \beta) = (0.70, 0.85)$	0.33	0.92	4913
	T_{GISM} : $(\alpha, \beta) = (0.80, 0.90)$	0.38	0.44	343
	T_{SIC} : $(\alpha, \beta) = (0.60, 0.80)$	0.10	1.02	5272
	T_{SIC} : $(\alpha, \beta) = (0.70, 0.85)$	0.16	0.85	3319
	T_{SIC} : $(\alpha, \beta) = (0.80, 0.90)$	0.27	0.59	1202
2014-2012	T_{GISM} : $(\alpha, \beta) = (0.60, 0.80)$	0.53	1.26	10195
	$T_{\text{GISM}}: (\alpha, \beta) = (0.70, 0.85)$	0.33	0.92	4913
	T_{GISM} : $(\alpha, \beta) = (0.80, 0.90)$	0.38	0.65	343
	T_{SIC} : $(\alpha, \beta) = (0.60, 0.80)$	0.32	1.23	7015
	T_{SIC} : $(\alpha, \beta) = (0.70, 0.85)$	0.35	1.12	4838
	T_{SIC} : $(\alpha, \beta) = (0.80, 0.90)$	0.32	1.23	586
2014-2013	$T_{\text{GISM}}: (\alpha, \beta) = (0.60, 0.80)$	0.02	0.98	6350
	T_{GISM} : $(\alpha, \beta) = (0.70, 0.85)$	0.11	0.59	3952
	T_{GISM} : $(\alpha, \beta) = (0.80, 0.90)$	-0.03	0.32	540
	T_{SIC} : $(\alpha, \beta) = (0.60, 0.80)$	0.25	1.00	5371
	T_{SIC} : $(\alpha, \beta) = (0.70, 0.85),$	0.21	0.79	3889
	T_{SIC} : $(\alpha, \beta) = (0.80, 0.90),$	0.40	0.32	1353

Table S1. Summary of attenuation rate solution variation for independently analysed field season pairs.



Figure S7. (a) Example of attenuation rate difference distributions for independently analysed field seasons (2012 P3-2011 P3). (b) Coverage overlap region.